A blast from the infant Universe: the very high-$z$ GRB 210905A


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ABSTRACT

We present a detailed follow-up of the very energetic GRB 210905A at a high redshift of $z = 6.312$ and its luminous X-ray and optical afterglow. Following the detection by Swift and Konus-Wind, we obtained a photometric and spectroscopic follow-up in the optical and near-infrared (NIR), covering both the prompt and afterglow emission from a few minutes up to 20 Ms after burst. With an isotropic gamma-ray energy release of $E_{\text{iso}} = 1.27^{+0.28}_{-0.15} \times 10^{54}$ erg, GRB 210905A lies in the top ~7% of gamma-ray bursts (GRBs) in the Konus-Wind catalogue in terms of energy released. Its afterglow is among the most luminous ever observed, and, in particular, it is one of the most luminous in the optical at $t \gtrsim 0.5$ d in the rest frame. The afterglow starts with a shallow evolution that can be explained by energy injection, and it is followed by a steeper decay, which is the highest ever measured for a $z \gtrsim 6$ burst, but within the range covered by closer events. The result is that the central engine of this burst was a newly formed black hole. Despite the outstanding energetics and luminosity of both GRB 210905A and its afterglow, we demonstrate that they are consistent within 2σ and it suggests that the central engine of this burst was a newly formed black hole. Despite the outstanding energetics and luminosity of both GRB 210905A and its afterglow, we demonstrate that they are consistent within 2σ and it suggests that the central engine of this burst was a newly formed black hole. Despite the outstanding energetics and luminosity of both GRB 210905A and its afterglow, we demonstrate that they are consistent within 2σ and it suggests that the central engine of this burst was a newly formed black hole. Despite the outstanding energetics and luminosity of both GRB 210905A and its afterglow, we demonstrate that they are consistent within 2σ and it suggests that the central engine of this burst was a newly formed black hole.

Key words. gamma-ray burst: general – Gamma-ray burst: individual: GRB 210905A

1. Introduction

The discovery of a $z > 6$ gamma-ray burst (GRB) is a rare occurrence that, thanks to the extreme luminosity of these sources, offers a window into the infant Universe, which is otherwise difficult to observe. Long GRBs, with gamma-ray emission generally longer than 2 s (Kouveliotou et al. 1993), originate from the explosions of very massive stars (Hjorth et al. 2003; Stanek et al. 2003; Woosley & Bloom 2006). Under the assumptions that the stellar initial mass function (IMF) in distant galaxies is not broadly different from that of closer objects and that the opening angles do not evolve strongly with redshift, the rate of GRBs can be used both to estimate the star-formation rate (SFR) (Kistler et al. 2009; Robertson & Ellis 2012) and to study the effects of metallicity on supernovae (SNe)-Ibc and GRB progenitors (Greco et al. 2012). The SFR is expected to change at very high redshift with the transition from the first massive population III (pop-III) stars in the remote Universe to pop-II and pop-I stars (Salvaterra 2015; Fryer et al. 2022). How this happens remains an open question that may be addressed through GRB studies. We note that the prompt emission is not affected by dust extinction, and thus GRBs can provide a census of obscured star formation at all redshifts (Blain & Natarajan 2000). Due to their immense brightness, GRBs can also act as beacons illuminating the local circumburst medium (e.g. Savaglio et al. 2003; Prochaska et al. 2008; Schady et al. 2011; Watson et al. 2013; Heintz et al. 2018), the interstellar medium (ISM) of their hosts (e.g. Fynbo et al. 2006; Savaglio et al. 2012; Cucchiara et al. 2015; Bolmer et al. 2019; Heintz et al. 2019), and the surrounding intergalactic medium (IGM) in the line of sight (Totani et al. 2006; Hartoog et al. 2015). They are therefore powerful probes of the ionisation and chemical enrichment history of the early universe. To shed light on these open issues through very high-redshift GRBs, several mission concepts have been studied and proposed (e.g. Amati et al. 2018; Tanvir et al. 2021b; White et al. 2021).

So far, out of the ~555 GRBs with a well-constrained spectroscopic redshift (as of 20 July 2022), only five have been de-
Here we present a follow-up of the bright GRB 210905A, the tenth burst with redshift $z \gtrsim 6$ detected in the last 16 years. It was detected by the Neil Gehrels Swift Observatory (Gehrels et al. 2004, Swift hereafter) and Konus-Wind (Aptekar et al. 1995). X-ray as well as optical and near-infrared (NIR) follow-up observations of its bright afterglow led us to determine a spectroscopic redshift of $z = 6.312$ (refined with respect to Tanvir et al. 2021a). The burst was also detected by the Cadmium Zinc Telluride Imager (CZTI) on-board Astrosat (Prasad et al. 2021) and, following the detection by the Swift Burst Alert Telescope (BAT, Barthelmy et al. 2005a), it was also found via a targeted search in data of the Gamma-ray Burst Monitor (GBM) on-board Fermi (Veres et al. 2021).

In §2 we describe the observations of both the GRB and the afterglow, and in §3 we present the analysis of the data. In §4 we discuss the results and compare them with other bursts at low and high redshift, and we draw our conclusions in §5. Throughout this work, the flux density of the afterglow is described as $F_\nu(t) \propto t^{-\alpha}$. A $\Lambda$CDM cosmological model with $\Omega_M = 0.308$, $\Omega_\Lambda = 0.692$, and $H_0 = 67.8$ km s$^{-1}$ Mpc$^{-1}$ (Planck Collaboration et al. 2016) has been assumed for calculations. All data are in the observer frame and $1\sigma$ errors are used throughout the paper, unless otherwise specified.

2. Observations

2.1. Gamma-ray and X-ray observations.

GRB 210905A was discovered by BAT on-board Swift at $T_0 = 00:12:41.3$ UT on 5 September 2021 (Sonbas et al. 2021). The BAT light curve shows a complex structure with three pulses, detected until ~800 s after the burst trigger.
Since GRB 210905A was too weak to trigger Konus-Wind (KW), the burst data are available only from the instrument’s waiting mode, as first reported by Frederiks et al. (2021). In this mode, count rates with a coarse time resolution of 2.944 s are recorded continuously in three energy bands: G1 (20 − 100 keV), G2 (100 − 400 keV), and G3 (400 − 1500 keV). A Bayesian block analysis of the KW waiting mode data in S1 (one of the two NaI(Tl) detectors) reveals three (separated in time) emission episodes, each featuring a statistically significant count rate increase in the combined G1+G2 band (Figure 1), while no statistically significant emission was detected in the G3 band throughout the burst.

The first episode, which triggered Swift/BAT, started at \(T_0 - 30\) s and ends at \(T_0 + 11\) s (hereafter Pulse 1). The weaker second episode (\(\sim T_0 + 344\) s to \(T_0 + 426\) s; Pulse 2) coincided in time with the bright flare in the XRT windowed-timing (WT) mode light curve around \(T_0 + 400\) s (Figure 1). The onset of the final emission episode, observed by KW from \(\sim T_0 + 747\) s to \(T_0 + 862\) s (Pulse 3), is clearly visible in the BAT mask-weighted data, which are available up to \(\sim 800\) s after the trigger. The \(T_0\) duration of the GRB 210905A prompt emission derived from the KW observation is \(\sim 870\) s.

Swift/XRT started observing the BAT error circle 91.7 s after the trigger and found an unknown X-ray source at the UVOT-enhanced position coordinates RA (J2000) = 20°36′11.64, Dec. (J2000) = −44°26′24.3 with a final uncertainty of 1′5 (Beardmore et al. 2021, Swift/XRT catalogue). Pointed Swift observations continued until 3.8 Ms after the GRB, when the source became too faint to be detected. Light curves and spectra, as well as the result of their modelling, have been obtained from the Swift/XRT repository (Evans et al. 2007, 2009). However, to build more accurate multi-wavelength spectral energy distributions (SEDs), given that some data available in the Swift/XRT repository suffer from bad centroid determination, we have processed the Swift data corresponding to the epochs of our SED analysis (obs. IDs 01071993001/002/003, Fig. 4). To reduce the data, the software package HeaSoft 6.29 was used with the latest calibration file available. For the data processing, we used standard procedures, consisting of the use of the packagexrtpipeline, available within the FTOOLS distribution, with standard-grade filtering. Using the most refined position provided by the Swift team, the selection of the GRB position in the X-ray data and the extraction of both source and background spectra, were done with the xselect package, while for the construction of the corresponding ancillary response file (.arf) we usedxrtsmrkarf on each corresponding exposure file. In the following, a Galactic equivalent hydrogen column density of \(N_H = 3.38 \times 10^{22}\) cm\(^{-2}\) is adopted (Willingale et al. 2013).

### 2.2. Optical/NIR Imaging and photometry

Swift/UVOT started observing about 156 s after the trigger but no credible afterglow candidate was found (Siegel et al. 2021). The MASTER Global Robotic Net (Lipunov et al. 2010) was also pointed at GRB 210905A 6 s after notice time and 414 s after trigger time but could not detect any afterglow candidate (Lipunov et al. 2021). We obtained optical/NIR observations with the 0.6m robotic Rapid Eye Mount telescope (REM, Zerbi et al. 2001), starting 428 s after the burst. A transient source was detected immediately in the \(H\) band and later in \(i′\) and \(ZJK\) bands (i.e. all except...)
Observations continued for about 3 hr before the declining afterglow brightness fell below the instrument detection limits in all filters (D’Avanzo et al. 2021). Images were automatically reduced using the jitter script of the ec1lpsae package (Devillard 1997) which aligns and stacks the images to obtain one average image for each sequence. A combination of IRAF (Tody 1993) and SExtractor packages (Bertin & Arnouts 2010) were then used to perform aperture photometry.

We triggered Bessel R- and I-band observations with the 1m telescope of the Las Cumbres Observatory Global Telescope (LCOGT) network, equipped with the Sinistro instrument, at the Cerro Tololo Inter-American Observatory (CTIO), Chile. The midpoints of the first epoch are $t_i = 1.06$ hr and $t_R = 1.29$ hr, in the $I$ and $R$ bands respectively. The data provided by the LCO are reduced using the BANZAI pipeline (McCully et al. 2018) that performs bias and dark subtraction, flat-fielding, bad-pixel masking, and astrometric calibration. Afterwards, we use our own pipeline, which aligns and stacks the images using the astrolign Python package (Beroiz et al. 2020), and afterwards uses SExtractor to perform the photometry and calibration against a sample of USNO-B catalogue stars (Monet et al. 2003). Using the data-reduction pipeline from LCO, and our relative photometry pipeline, we calculate a magnitude of $I = 19.46 \pm 0.15$ mag and a $3\sigma$ upper limit of $R > 22.44$ mag. The lack of an $R$-band detection alerted us to the possibility that this burst may lie at very high redshift ($z > 5$, first reported by Strausbaugh & Cucchiara 2021a,b).

GRB 210905A was observed simultaneously in $g'r'i'z'JHK$ with GROND (Gamma-Ray Burst Optical-Near-Infrared Detector; Greiner et al. 2008; Greiner 2019) mounted on the 2.2m MPG telescope at ESO La Silla Observatory in Chile (Nicuesa Guelbenzu et al. 2021). The first epoch observations were done around 23 hr after the GRB trigger. The afterglow was detected only in the $z'JHK$ bands. A second set of observations obtained 7 hr later was shallower and yielded only upper limits. Subsequent follow-up observations were obtained on 7 and 8 September 2021, but the afterglow was also not detected in the latter epochs. We continued our ground-based follow-up using both the VLT/HAWK-I (High Acuity Widefield K-band Imager, Pirard et al. 2004) NIR imager on Paranal, as well as the Dark Energy Camera (DECam) mounted on the 4m Victor Blanco telescope at CTIO. Also we used the acquisition camera of the ESO VLT/X-shooter spectrograph to obtain $g'r'i'J$-band imaging before moving on to spectroscopy. We obtained a last ground-based observation 87 d after the GRB with VLT/FORS2 in the $I_{Bessel}$ band.

Finally, the field was observed with the Hubble Space Telescope (HST) on 24 April 2022. At this epoch four dithered observations with a total duration of 4797 s were obtained in the F140W filter. The data were obtained from the MAST archive and processed with astrodrrizzle to create a final combined charge transfer efficiency corrected image with a pixel scale of $0.07''$/pixel. Aperture photometry was performed with a radius of $0.4''$ to minimise any contribution from the nearby sources (see Figure 5).

X-shooter and GROND optical/NIR images were reduced in a standard manner using PyRAF/IRAF (Tody 1993). In particular, GROND data reduction was done with a customised pipeline (Krühler et al. 2008) that is based on standard routines in IRAF. FORS I-band and HAWK-I $JHK_s$-band data have been reduced using the ESO Reflex environment (Freudling et al. 2013). We obtained PSF photometry with the DAOPHOT and ALLSTAR tasks of IRAF. PSF-fitting was used to measure the magnitudes of the GRB afterglow. Only for the late-time FORS2 observation in $I_{Bessel}$ at 87 days and HST F140W at 232 days did we use aperture photometry.

All optical photometry except $I_{Bessel}$-band data were calibrated against the SkyMapper catalogue (Wolf et al. 2018), while the ground-based NIR photometric calibration was performed against the 2MASS catalogue (Krutskii et al. 2006). This procedure results in a typical systematic accuracy of $0.04$ mag in $g'r'i'z'$. 0.06 mag in $JH$ and 0.08 mag in $K_s$. To cross-calibrate all the $I$-band imaging we applied the Lupton formulae to a set of local standard stars from the SkyMapper catalogue.

The $I$ filters used by X-shooter and LCO extend beyond 10000 Å. Therefore, we expect that not all the flux is dimmed by the Lyman-$\alpha$ dropout at around 8900 Å in these filters. On the contrary, the FORS2 I-band filter has negligible transmission above Lyman-$\alpha$ (at the redshift of GRB 210905A). Therefore, we speculate that the possible (note the large error) late I-band emission (see § 4.5) does not originate from the afterglow, but instead from a foreground source.

The optical/NIR afterglow lies at coordinates RA (J2000) = $20^\mathrm{h}36^\mathrm{m}11.57^\mathrm{s}$, Dec. (J2000) = $-44^\circ26'24.7''$ as measured in our first HAWK-I image and calibrated against field stars in the GAIA DR2 catalogue (Gaia Collaboration et al. 2018) with the astrometric precision being 0.15. This refines the position reported by LCO (Strausbaugh & Cucchiara 2021a) and is in agreement with the more precise localisation provided by ALMA (Laskar et al. 2021a). Table 1 provides a summary of all photometry of the transient (non-relevant upper limits are not reported). All reported magnitudes are in the AB photometric system and not corrected for the Galactic foreground extinction of $E(B-V) = 0.029$ mag (Schlafly & Finkbeiner 2011).

### 2.3. X-shooter spectroscopy and redshift

Starting ~2.53 hr after the GRB detection, we obtained UV to NIR spectroscopy of the afterglow with the X-shooter instrument (Vernet et al. 2011) mounted on the VLT on Cerro Paranal (ESO, Chile), via the Stargate Large Programme for GRB studies.

The afterglow is well detected in the red part of the visible arm. A clear break is detected around 9000 Å, which we interpret as the Lyman-$\alpha$ break (first reported in Tanvir et al. 2021a). Other lines such as Fe ii, Al ii, C iv and Si ii and fine structure lines are visible and display two velocity components, which belong to the ISM of the same galaxy. All these lines allow us to determine $z = 6.312$ as the redshift of the GRB. A very strong foreground system at $z = 2.8296$ (Mg ii, Fe ii lines) and another at $z = 5.7390$ (C ii, Fe ii, C iv, Si iv lines) are also present. The details concerning the reduction and analysis of the absorption lines in the X-shooter spectra are given in Saccardi et al. (2022). This high redshift explains the non-detection by UVOT and MASTER and the red $R_C - I_C$ and $r' - z'$ colours found with LCO and X-shooter as due to Lyman dropout. In Fausey et al. (in prep.) we will study the IGM neutral fraction in light of the GRB 210905A afterglow spectrum.

### 3. Modelling and results

#### 3.1. Joint BAT-KW modelling

To derive the broad-band spectral parameters of the prompt emission of this burst, we performed a joint spectral analysis of the BAT data ($15 - 150$ keV) and the KW waiting-mode data.

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The photometry was confirmed after the cross-calibration mentioned below.

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Table 2. Fits to the prompt emission spectra.

<table>
<thead>
<tr>
<th>Spectrum</th>
<th>Instruments</th>
<th>Model</th>
<th>Time interval (relative to $T_0$, s)</th>
<th>$\alpha$ (photon index)</th>
<th>$E_{\text{peak}}$</th>
<th>$E_{\text{break}}$</th>
<th>Flux (15–1500 keV)</th>
<th>$\chi^2$/d.o.f.</th>
</tr>
</thead>
<tbody>
<tr>
<td>'peak'</td>
<td>BAT+KW</td>
<td>CPL</td>
<td>[−0.465, 2.479]</td>
<td>−0.66$^{+0.35}_{-0.30}$</td>
<td>144$^{+26}_{-28}$</td>
<td>2.83$^{+0.56}_{-0.40}$</td>
<td>40.2 (58)</td>
<td></td>
</tr>
<tr>
<td>Pulse 1</td>
<td>BAT+KW</td>
<td>CPL</td>
<td>[−29.905, 11.311]</td>
<td>−0.99$^{+0.18}_{-0.17}$</td>
<td>127$^{+31}_{-19}$</td>
<td>1.14$^{+0.14}_{-0.11}$</td>
<td>36.5 (58)</td>
<td></td>
</tr>
<tr>
<td>Pulse 1</td>
<td>BAT+KW</td>
<td>DBPL</td>
<td>[−29.905, 11.311]</td>
<td></td>
<td>127</td>
<td>27.09$^{+1.41}_{-1.34}$</td>
<td>35.2 (56)</td>
<td></td>
</tr>
<tr>
<td>X-ray flare</td>
<td>XRT+KW</td>
<td>DBPL</td>
<td>[80.0, 120.0]</td>
<td>1.46$^{+0.21}_{-0.26}$</td>
<td>unconstrained</td>
<td>225.9 (281)</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Pulse 2</td>
<td>XRT+BAT+KW</td>
<td>CPL</td>
<td>[343.983, 426.415]</td>
<td>−0.80$^{+0.58}_{-0.47}$</td>
<td>70$^{22}_{11}$</td>
<td>1.13$^{+0.11}_{-0.10}$</td>
<td>48.6 (58)</td>
<td></td>
</tr>
<tr>
<td>Pulse 2</td>
<td>XRT+BAT+KW</td>
<td>DBPL</td>
<td>[343.983, 426.415]</td>
<td></td>
<td>50</td>
<td>1.13$^{+0.11}_{-0.10}$</td>
<td>714.2 (714)</td>
<td></td>
</tr>
<tr>
<td>Pulse 3</td>
<td>BAT</td>
<td>CPL</td>
<td>[747.311, 797.359]</td>
<td>−0.89$^{+0.37}_{-0.30}$</td>
<td>154$^{77}_{38}$</td>
<td>0.76$^{+0.17}_{-0.11}$</td>
<td>55.2 (58)</td>
<td></td>
</tr>
<tr>
<td>Pulse 3</td>
<td>KW</td>
<td>CPL</td>
<td>[747.311, 862.127]</td>
<td>−0.88$^{+0.30}_{-0.30}$</td>
<td>167$^{48}_{30}$</td>
<td>0.97$^{+0.11}_{-0.23}$</td>
<td>0 (0)</td>
<td></td>
</tr>
<tr>
<td>Pulse 3</td>
<td>REM+BAT+KW</td>
<td>DBPL</td>
<td>[747.311, 797.359]</td>
<td></td>
<td>154</td>
<td>0.006$^{a}$</td>
<td>59.3 (56)</td>
<td></td>
</tr>
</tbody>
</table>

$^a$ All spectra, except the first, are time-averaged.
$^b$ This spectrum was used to calculate the peak energy flux.
$^c$ The interval covered by XRT.
$^d$ The interval covered by BAT.
$^e$ KW-only fit: for the CPL model, the 3-channel fit has 0 degrees of freedom.
$^f$ CPL stands for cut-off power-law. DBPL stands for double-broken power-law, used for the synchrotron model. In this last case, the power-law indices were fixed as described in §3.2.
$^g$ This break was fixed to match the H-band data.

The 15 – 1500 keV energy fluences of Pulses 1, 2, and 3, derived from our time-averaged fits, are summarised in Table 5, together with the fluence integrated over all three emission episodes. We use these results to calculate the isotropic energy (see also §4.2). The spectrum in the interval ($T_0 = 0.465$ s, $T_0 + 2.479$ s) inside Pulse 1, which corresponds to the peak count rate, is characterised by $\alpha \sim −0.66$ and $E_{\text{peak}} \sim 144$ keV. Using this spectrum and the BAT light curve, we estimate the 1 s peak energy flux of GRB 210905A to be $3.83^{+0.73}_{-0.54} \times 10^{-7}$ erg cm$^{-2}$ s$^{-1}$ (15 – 1500 keV).

3.2. Joint modelling of the prompt emission from gamma-rays to the optical

In the previous section we have analysed the gamma-ray spectra during the three pulses and found that they can be modelled almost equally well with a CPL or a Band function with very similar (within errors) low-energy photon index $\beta = −2.3$ and $E_{\text{peak}}$. The high-energy index of the Band function is $\beta < −2.3$, poorly constrained by the sparse KW data. Values of $\beta = −2.3$ are very typical low- and high-energy photon indices for GRBs (e.g. Preece et al. 1998; Nava et al. 2011). Follow-up studies (Ravasio et al. 2018, 2019) and has been studied in detail (Rastinejad et al. 2022) in the temporally long merger event GRB 211211A (Rastinejad et al. 2022). It has been suggested to be a common feature of GRB prompt emission spectra (Tofano et al. 2021). Therefore, the low-energy part of the spectrum, with photon index $−1$, into two power-law photon indices describing...
the spectrum below and above the low-energy break, and have distributions centred around $-2/3$ and $-3/2$ (or $1/3$ and $-1/2$ for the flux density spectrum $F_\nu$), respectively. These indices are the same as those below and above the cooling break $\nu_c$ and expected by the synchrotron theory in the fast-cooling regime (see also Ravasio et al. 2018, 2019). Further confirmation of these empirical fits was obtained by direct fitting of prompt GRB spectra with a synchrotron model (Ronchi et al. 2020; Burgess et al. 2020) and the synchrotron interpretation is discussed for example in Ghisellini et al. (2020).

To determine if the prompt emission of GRB 210905A is in agreement with these theoretical expectations, we have modelled the NIR and X- to gamma-ray SEDs of five epochs during the whole prompt emission with a double broken power-law with photon indices fixed to the synchrotron model predictions. That the optical-to-gamma emission is the result of a common radiative process is justified by the simultaneous evolution of the optical, X-ray and gamma-ray prompt emission. The selected epochs are the three gamma-ray pulses, the first X-ray flare at $\sim 120$ s and an additional epoch at $\sim 630 - 690$ s simultaneous to an $H$-band observation. This is the only epoch before the last pulse with few counts in the BAT spectrum. We have fixed the high-energy break ($\nu_H$, the frequency corresponding to the minimum injection energy in a fast-cooling synchrotron model) to the break energy in the Band modelling above. The high-energy photon index above this break has been fixed to $-2.4$, that is also consistent with the Band fit. The results are shown in Table 2 and in Figure 2. The analysis of the X-ray flare alone shows that it is well modelled by $\nu_m$ at $\sim 1$ keV, a photon index $^{10}$ $-2.4$ and intrinsic absorption $N_H = 7.7^{+3.2}_{-1.6} \times 10^{21}$ cm$^{-2}$. In the following, we fixed the intrinsic hydrogen column density to this value. During the first two pulses the data are consistent with a broken power-law with photon index 0.5.

In the last two SEDs, we also include the $H$-band follow-up obtained with REM (Figure 1). Note that in the fourth SED we have simply scaled the solution from the last epoch, because there are basically just two measurements for three possible free parameters$^{11}$, not enough to constrain all breaks. Therefore, this is not shown in Table 2.

For these last SEDs (i.e. before and during Pulse 3) the $H$-band observation is below the extrapolation of the photon index from the gamma-rays, and thus $\nu_c$ must be in between the $H$ and X-ray bands. We further discuss the implications of this finding in §4.1. Unfortunately, for both SEDs the lack of any colour information and possible contribution from the emerging afterglow in the observed optical/NIR prevents us from affirming without doubt that the low-energy photon index is $-2/3$. However, we can confirm that for both SEDs the synchrotron model is in agreement with the observations.

### 3.3. Joint afterglow light curve and SED

Figure 3 shows both optical/NIR and X-ray light curves of the afterglow. Regions in grey have not been considered in this section because of: i) the presence of flares likely due to long-lasting activity from the central engine and ii) a possible late break when compared to the earlier evolution (Figure 3) that we discuss below in §3.5.

A complete understanding of the afterglow behaviour would require a full numerical simulation. Nevertheless, we can derive some conclusions by modelling the SEDs and light curves of the

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**Table 3. Optical/NIR to X-ray modelling of the afterglow (see Figure 4).**

<table>
<thead>
<tr>
<th>Model</th>
<th>Time</th>
<th>$\beta_{opt}^a$</th>
<th>Break (keV)</th>
<th>$\chi^2$/d.o.f.</th>
</tr>
</thead>
<tbody>
<tr>
<td>BPL</td>
<td>0.1</td>
<td>0.62 ± 0.04</td>
<td>1.7$^{+2.6}_{-1.6}$</td>
<td>1.4 ± 1.2</td>
</tr>
<tr>
<td>SPL</td>
<td>0.1</td>
<td>0.63 ± 0.03</td>
<td>--</td>
<td>--</td>
</tr>
<tr>
<td>BPL</td>
<td>1.0</td>
<td>0.60 ± 0.04</td>
<td>0.35$^{+0.28}_{-0.13}$</td>
<td>0.35</td>
</tr>
<tr>
<td>SPL</td>
<td>1.0</td>
<td>0.71 ± 0.02</td>
<td>--</td>
<td>--</td>
</tr>
<tr>
<td>BPL</td>
<td>2.2</td>
<td>0.56 ± 0.16</td>
<td>0.18$^{+0.06}_{-0.03}$</td>
<td>0.22 ± 0.19</td>
</tr>
<tr>
<td>SPL</td>
<td>2.2</td>
<td>0.80 ± 0.03</td>
<td>--</td>
<td>--</td>
</tr>
</tbody>
</table>

$^a$ In the broken power-law we assumed $\beta = \beta_{opt} + 0.5$.
$^b$ Obtained from the best-fit value at 1 d with $\nu(t) = \nu(1d)(t/1d)^{\beta_{opt}}$, with $k = 0.5$ after 1 d (ISM, slow cooling scenario) and $k = 0.6$ before 1 d assuming energy injection (see §4.4).
afterglow. We modelled the afterglow SED from NIR to X-ray frequencies at three different epochs, 0.1, 1.0, and 2.18 days, using Xspec v12.12.0 (Arnaud 1996). We have not considered the optical data (z-band and bluer bands) because they are affected by the Lyman-α break and thus do not add anything useful to this modelling. The redshift was fixed to 6.312 and we fixed the Galactic and intrinsic hydrogen column density (see §2.1). To avoid being affected by the uncertain gas absorption, we have not considered data below 0.5 keV in the modelling. We have modelled the NIR-to-X-ray SED both with a single and a broken power-law with $\beta_X = \beta_{opt} + 0.5$, at all three epochs. The best-fits are shown in Figure 4 and their parameters are shown in Table 3. All fits give negligible dust extinction $A_V \lesssim 0.03$ mag, independent of the extinction law\(^\text{12}\), which is not unusual for high-z GRB afterglows (see §4.8). It is not straightforward to decide between the single and broken power-law models as the SEDs are fit comparably well in both cases. However, we note that in the first epoch they give basically the same value for the low-energy spectral index. Therefore, we conclude that $\nu_c$ is within or above the X-ray band at 0.1 d. That $\nu_c$ is then in between the two bands is even more clear in the second SED at 1 d whose best-fit gives $\beta_{opt} = 0.60 \pm 0.04$, and thus an electron index $p = 2.20 \pm 0.08$. To confirm these findings, we need to also consider the light-curve evolution.

\(^{12}\) In the zdust model.
We have modelled the optical and NIR light curves simultaneously with a smoothly broken power-law (Beuermann et al. 1999): \( F = (F_0 + F_\infty)^{1/\kappa} \), where \( F_0 = f_{\text{break}}/(t/f_{\text{break}})^{-\alpha_{\text{break}}} \) being the flux density at break time \( t_{\text{break}} \), \( \kappa \) the break smoothness parameter, and the subscripts 1, 2 indicate pre- and post-break, respectively. We find a shallow break with large uncertainty at \( t_{\text{break, opt}} = 0.99 \pm 0.73 \) d (85.9 \pm 67.2 ks) and decay indices \( \alpha_{1, \text{opt}} = 0.69 \pm 0.04 \) and \( \alpha_{2, \text{opt}} = 0.94 \pm 0.04 \), with break smoothness \( \kappa = 10 \) fixed (Zeh et al. 2006)\(^{13}\). With respect to a simple power-law, the \( x^2/d.o.f. \) decreases from 1.36 to 0.92.

The X-ray light curve shows an initial peak at 97 s followed by the typical steep decay (Tagliaferri et al. 2005b; Barthelmy et al. 2005b) with \( \alpha = 2.3^{+0.6}_{-0.16} \) until \( \sim 270 \) s after the burst, when it is interrupted by a flare also visible in gamma-ray data. After \( \sim 3000 \) s, it is best modelled by a broken power-law with a shallow break at \( \sim 1 \) d: from \( \alpha_{1,X} = 0.74^{+0.03}_{-0.01} \) to \( \alpha_{2,X} = 1.10 \pm \) 0.04, with \( t_{\text{break,X}} = \sim 60 \pm 30 \) ks\(^{14}\) (1σ errors). Finally, we note that modelling simultaneously the X-ray and optical bands with the best-fit indices found above, the shallow break is seen at a common time of \( 0.70 \pm 0.26 \) days (60.5 \pm 22.5 ks).

In Table 4, we compare the observed evolution with the predicted values of the temporal slopes in the optical/NIR and the X-ray bands for various slow-cooling afterglow scenarios (see, e.g. Zhang et al. 2006; Schulze et al. 2011) and the electron index, \( p = 2.20 \pm 0.08 \). We cannot find a good solution for the data before 0.7 d, however, after the first modest break the data are best modelled within a scenario where the jet is expanding into a constant-density medium (hereafter referred to as the interstellar medium or ISM environment). A single power-law SED solution cannot explain the observed temporal decay index in X-rays, \( \alpha_{2,X} = 1.1 \), with emission below the cooling frequency, \( \nu_c \). Moreover, within this solution \( \beta_{\text{opt}} \) should be constant, but instead it evolves with time. These results indicate that \( \nu_c \) should lie between the optical and X-ray bands (see also Figure 4). A \( \nu_c \) that has moved out of the X-ray band can explain the difference in the temporal decay index between optical and X-rays after the shallow break, therefore, we consider a broken power-law as the best description for the optical-to-X-ray SED. For an upper branch\(^{15}\) \( \beta_X = 1.10 \pm 0.04 \), obtained at 1d (the epoch with the best statistics), the electron index is \( p = 2.20 \pm 0.08 \). The large errors on the cooling frequency do not permit to test whether the break shifts in time as \( r^{-1/2} \), although the results seem consistent with such a relation (see Table 3).

### 3.4. The late NIR imaging

In Figure 5 we show the most recent observation of the field obtained with \( \text{HST} \) in the F140W band. At 0"09\pm0"02 from the NIR afterglow position we clearly detect an extended source \((\text{F140}W = 25.66 \pm 0.05 \) mag\). The relative offset is measured comparing the centroids in the first HAWK-I image and the \( \text{HST} \) image, after aligning these two images using a common set of sources. It is slightly elongated in the NNE-SSW direction and has a FWHM of 0"4 and 0"3, larger than the FWHM of field stars (0"25 \pm 0"02). Therefore, we conclude that the \( \text{HST} \) detection is dominated by a constant source. The statistical prob-

\(^{13}\) We have also evaluated smaller fixed \( \kappa \) values (5, 2, 1) and find that \( x^2/d.o.f. \) increases, \( t_b \) remains similar, but even at \( \kappa = 5 \), the error exceeds the value of the break time, and increases further.

\(^{14}\) As shown by the Swift/XRT light curve repository (Evans et al. 2007, 2009).

\(^{15}\) The spectral index \( \beta_X = 0.90 \pm 0.15 \) at 22.8 ks reported in the XRT pages is well in agreement with this result.

### Table 4. Closure relations.

<table>
<thead>
<tr>
<th>Afterglow model</th>
<th>Theoretical</th>
<th>Observed</th>
</tr>
</thead>
<tbody>
<tr>
<td>( \alpha )</td>
<td>( \alpha_{1, \text{opt}} = 0.69 \pm 0.04 )</td>
<td>( \alpha_{1,X} = 0.74 \pm 0.03 )</td>
</tr>
<tr>
<td>( \alpha_{2, \text{opt}} = 0.93 \pm 0.04 )</td>
<td>( \alpha_{2,X} = 0.93 \pm 0.04 )</td>
<td></td>
</tr>
</tbody>
</table>

\(^{\alpha}\) The \( \sigma \)-level is the difference of the predicted and the observed temporal slope, normalised to the square root of the sum of their quadratic errors.

\(^{\beta}\) The solution that matches the closure relations within 1 \( \sigma \) is highlighted in bold (see §3.3).

\(^{\gamma}\) We follow the common use and refer to the constant-density medium as ISM.

### 3.5. Constraints on the jet break

A sizeable number of GRB afterglow light curves break to steeper power-law decays, usually within a few days after the trigger. These breaks have generally been interpreted as due to the outflow being collimated in a jet, where the break occurs when the relativistic beaming angle becomes wider than the jet’s half-opening angle \( \theta_{\text{jet}} \) (Rhoads 1997; Sari et al. 1999). In the forward-shock model the jet breaks have to be achromatic, thus to have the same slope (and slope change) simultaneously in all bands\(^{16}\).

In §3.3 we have shown that a moderate break is present in both optical and X-rays at a common time of \( \sim 0.74 \) d. However, the post-break slope for both X-rays and optical is only \( \sim 1 \), that is too shallow for a jet break, both observationally (Wang et al. 2015) and theoretically (Sari & Piran 1999; Zhang et al. 2006; Panaitescu 2007). Instead, the last XRT detection, together with the late observation by the \( \text{Chandra} \) X-ray Observatory (Laskar et al. 2009).  

\(^{16}\) The value of the light-curve post-jet-break slope depends on \( \nu_{\text{obs}} \) being above or below \( \nu_m \) and \( \nu_V \), where \( \nu_m \) is the synchrotron self-absorption frequency. Optical and X-ray afterglow SEDs are usually observed to be above \( \nu_m \) (e.g. Greiner et al. 2011).
Using the spectral slope \( \beta_{\text{opt}} = 0.6 \) obtained from the SED fitting of the afterglow, the colour correction is just \( H - F140W = -0.10 \) mag, and thus will not make an appreciable difference in our analysis.

To better constrain the break time, we modelled jointly the \( H \)-band and X-ray light curves after the early break at 0.7 d with a smoothly broken power-law \( F = (F^2 + F^3)^{1/\kappa} \), following the definition in §3.3 but with the subscripts \( \alpha_{\text{opt}} \) and \( \alpha_{\text{X}} \) indicating the pre- and post-jet break respectively. We fixed the pre-jet-break index to the model values \( \alpha_{\text{opt}} = 0.9 \) and \( \alpha_{\text{X}} = 1.15 \) (see Table 4). In our analysis we adopt the jet model (with sideways expansion) and slow cooling (e.g. Sari et al. 1998; Zhang & Mészáros 2004). Therefore, we assume that the post-jet-break index is \( \alpha_{\text{opt}} = \alpha_{\text{X}} \approx p = 2.2 \) (see §3.3). Note that the sparse data after the break prevent us from constraining the \( \kappa \) parameter.

4. Discussion

4.1. The nature of the prompt emission

GRB 210905A is among the few exceptional cases where optical data could be obtained during a gamma-ray pulse (Figure 1).
the past, in less than a dozen cases has modelling of the prompt emission been possible from optical/NIR to gamma-rays, such as in the cases of GRBs 990123, 041219A, 060526, 080319B, 080603A, 080928, 090727, 091024, 110205A, 111209A, 130427A, and the more recent GRBs 160625B and 180325A (e.g. Sari & Piran 1999; Vestrand et al. 2005; Thöne et al. 2010; Racusin et al. 2008; Guidorzi et al. 2011; Rossi et al. 2011; Kopač et al. 2013; Virgili et al. 2013; Stratta et al. 2013; Kann et al. 2011; Zheng et al. 2012; Gendre et al. 2013; Vestrand et al. 2014; Troja et al. 2017; Becerra et al. 2021).

At \( z > 6 \), this analysis was possible only for GRB 050904 (Boër et al. 2006). In all these cases, modelling of the data with a broken power-law shows that the X-to-gamma-ray SED of the prompt pulses is in agreement with synchrotron emission, and in particular with fast cooling. This is in agreement with studies on large samples as we have mentioned in § 3.2 (see Ghisellini et al. 2020), for a discussion on the possible implications.

However, when including the optical data the situation can be more complex: for example, the main and earlier pulses of GRBs 990123 (Sari & Piran 1999; Galama et al. 1999; Corsi et al. 2005; Maiorano et al. 2005), 080319B (Racusin et al. 2008; Bloom et al. 2009), 111205A (Zheng et al. 2012), 130427A (Vestrand et al. 2014), 160625B (Troja et al. 2017), and 180325A (Becerra et al. 2021) show a convex spectrum between optical and X/gamma-rays (e.g. Guiriec et al. 2016), this feature can be explained by synchrotron emission from internal forward shocks dominating the gamma-ray and X-ray prompt emission, while the early optical flashes are generated by a reverse shock.

The analysis of Pulse 3 of GRB 210905A is however clearly in disagreement with this latter scenario, with the \( H \)-band emission being fainter than the extrapolation of the power-law modelling the gamma-rays (Figure 2). Therefore, although simultaneous optical-to-gamma coverage of the first prompt pulses is missing in the case of GRB 210905A, we show that at least during the last pulse there is no indication that the NIR data have an origin different from the X/gamma-ray emission, and all the observed epochs during the prompt phase can be explained by synchrotron emission from internal shocks. This is not surprising, as in several events (e.g. GRBs 990123, 130427A, 160525B) the optical-to-gamma SED later evolves and can be entirely explained as emission from the forward shock. Oganesyan et al. (2019) have shown that the later SEDs are consistent with being produced through synchrotron emission in the moderately fast-cooling regime from the same emission region.

4.2. Prompt emission in context

Using \( z = 6.312 \), we estimate the rest-frame properties of the burst prompt emission. Isotropically-equivalent energy release (\( E_{\text{iso}} \)) and rest-frame spectral peak energies \( E_{\text{peak,rest}} = (1 + z)E_{\text{peak}} \) for the individual emission episodes were calculated from the CPL spectral fits (§3.1); they are listed in Table 5. Integrated over the three intervals, the total energy release of GRB 210905A in \( \gamma \)-rays is \( E_{\text{iso}} = 1.27_{-0.15}^{+0.20} \times 10^{55} \text{ erg} \), which is within the highest \( 7\% \) for the KW sample of 338 GRBs with known redshifts (Tsvetkova et al. 2017, 2021). Since \( E_{\text{peak}} \) obtained from our fits differs between the individual emission episodes, we used the spectral peak energy value weighted by the episode fluence, \( E_{\text{peak, \text{avg}}} \geq 145 \text{ keV} \), to estimate the burst time-averaged \( E_{\text{peak, \text{avg}}} \) to \( \sim 1060 \text{ keV} \). This intrinsic peak energy is among the highest \( \lesssim 15\% \) of long KW GRBs. Derived from the peak energy flux, the peak \( \gamma \)-ray luminosity of the burst is \( L_{\gamma} = 1.87_{-0.62}^{+0.56} \times 10^{53} \text{ erg s}^{-1} \).

The rest-frame \( E_{\text{peak}} \) corresponding to the time interval around the peak luminosity is \( \sim 1050 \text{ keV} \). The reported values of \( E_{\text{iso}} \) and \( L_{\text{iso}} \) were calculated in the rest frame 1 keV–10 MeV range. All the quoted errors are at the 1σ confidence level.

With these estimates, GRB 210905A as well as its individual episodes lie inside the 68% prediction interval (PI) of the \( E_{\text{peak, \text{avg}}} - E_{\text{iso}} \) (‘Amati’ relation; Figure 8) for 315 long KW GRBs with known redshifts (Tsvetkova et al. 2021). Likewise, the burst peak luminosity and the corresponding \( E_{\text{peak, \text{avg}}} \) perfectly fit the ‘Yonetoku’ relation for the sample.

Figure 7 shows the GRB 210905A prompt emission in the context of eight GRBs at \( z \geq 6 \). With the rest-frame duration \( T_{90}/(1 + z) \sim 119 \text{ s} \) GRB 210905A is the intrinsically longest high-\( z \) GRB detected to date and is also among the longest \( \lesssim 3\% \) of bursts as compared to the whole KW catalogue\footnote{This sample does not include six KW ultra-long (\( T_{100} > 1000 \text{ s} \)) bursts, all at low-to-moderate redshifts \( z \lesssim 2 \).} (Tsvetkova et al. 2017, 2021), which covers the range \( 0.04 \leq z \leq 9.4 \). In this high-redshift sample, GRBs 210905A and 130606A are the only bursts with well-separated emission episodes. Except for this feature, they are similar to all other bursts which show short spikes with only moderate energy release. The exception is GRB 050904, which is similar to GRB 210905A in terms of energy release (\( E_{\text{iso}} = (1.33 \pm 0.14) \times 10^{54} \text{ erg} \)) but shows a \( \sim 30 \text{ s} \) long emission episode with two extended peaks (and at least a third episode observed in X-rays).

The most powerful burst at low redshift is GRB 130427A at \( z = 0.3399 \) (Selsing et al. 2019). This GRB can be considered as a good analogue of the energetic high-\( z \) population because of its high energy release (Perley et al. 2014; De Pasquale et al. 2016). Its prompt emission parameters are similar to those of Pulse 3 of GRB 210905A (see Table 5): \( E_{\text{peak}} \sim 1415 \text{ keV} \) and \( E_{\text{iso}} \sim 9.4 \times 10^{53} \text{ erg} \) (Tsvetkova et al. 2017). Accordingly, GRB 130427A and Pulse 3 lie very close in the \( E_{\text{peak}} - E_{\text{iso}} \) plane. We should note, however, that the intrinsic durations of GRB...
than the main pulse. So, GRB 130427A is not a ‘genuine’ multi-

Fig. 8. Rest-frame energetics of GRB 210905A in the $E_{p, i, z}$–$E_{\text{iso}}$ and $E_{p, i, z}$–$L_{\text{iso}}$ planes (red stars). Brown stars in the left panel show the values derived for the individual Pulses 1, 2, and 3. The rest-frame parameters of 315 long KW GRBs with known redshifts (Tsvetkova et al. 2021) are shown with circles; the colour of each data point represents the burst’s redshift. In the left plot, rest-frame peak energy values are derived from time-averaged spectral fits ($E_{p, i, z}$). In the right plot, they are derived from spectra, corresponding to the burst’s peak count rate ($E_{p, i}$). The ‘Amati’ and ‘Yonetoku’ relations for this sample are plotted with dashed lines and the dark- and light-grey shaded areas show their 68% and 90% prediction intervals, respectively. The error bars are not shown here for reasons of clarity.

Table 5. Parameters of the individual prompt emission pulses.

<table>
<thead>
<tr>
<th>Episode</th>
<th>$E_{\text{peak}, z}$ (keV)</th>
<th>Fluence ($15 - 1500$ keV)$^a$</th>
<th>$E_{\text{iso}}$ (10$^{53}$ erg)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Pulse 1</td>
<td>930$^{+230}_{-140}$</td>
<td>0.471$^{+0.052}_{-0.046}$</td>
<td>3.40$^{+0.41}_{-0.33}$</td>
</tr>
<tr>
<td>Pulse 2</td>
<td>510$^{+45}_{-60}$</td>
<td>0.245$^{+0.050}_{-0.038}$</td>
<td>1.73$^{+0.33}_{-0.32}$</td>
</tr>
<tr>
<td>Pulse 3</td>
<td>1220$^{+60}_{-50}$</td>
<td>1.11$^{+0.38}_{-0.26}$</td>
<td>7.62$^{+1.81}_{-1.59}$</td>
</tr>
<tr>
<td>Total$^c$</td>
<td>1060$^{+470}_{-320}$</td>
<td>1.82$^{+0.29}_{-0.28}$</td>
<td>12.7$^{+2.0}_{-1.9}$</td>
</tr>
</tbody>
</table>

$^a$ Fluences were calculated using the fits with the CPL function from Table 2.

$^b$ Only the KW 3-channel spectrum is used.

$^c$ This fluence is integrated over all three emission episodes.

130427A and pulse 3 of GRB 210905A differ by factor of two ($T_{90, i} \sim 10$ s for GRB 130427A versus $\sim 19$ s for Pulse 3). The initial light curve of GRB 130427A is somewhat similar to that of GRB 210905A since it starts with a large structured peak $\sim 20$ s long in the rest-frame, followed by a third peak starting at $\sim 100$ s. This second pulse is, however, orders of magnitude weaker than the main pulse. So, GRB 130427A is not a ‘genuine’ multi-

4.3. Collimation–corrected energy and central engine

Knowing the value of the jet opening angle is crucially important because it enables us to estimate the ‘true’, collimation-corrected, energetics of the outflow (Frail et al. 2001; Ghirlanda et al. 2007). Numerical and analytical calculations (e.g. Sari et al. 1999) have shown that the half-opening angle of the jet is related to the jet-break time. Following Zhang & MacFadyen (2009) we calculate this angle $\theta_{\text{jet}}$ using the following equation for a uniform jet expanding in a constant-density medium:

$$\frac{\theta_{\text{jet}}}{\text{rad}} = 0.12 \left( \frac{E_{\text{kin, iso}}}{10^{53} \text{erg}} \right)^{-1/8} \left( \frac{n}{\text{cm}^{-3}} \right)^{1/8} \left( \frac{t_{\text{jet}}}{\text{day}} \right)^{3/8} (1 + \frac{z}{6.312})^{-3/8},$$

where $E_{\text{kin, iso}}$ is the kinetic energy of the outflow assuming isotropy; $n = 1$ cm$^{-3}$ is the number density of the medium, assumed to be constant; $t_{\text{jet}}$ is the jet-break time (observer frame, see §3.5), while $z = 6.312$ is the redshift of the event. The kinetic energy is the one left after the prompt phase, and which later dissipates in the afterglow. Together with the energy released as gamma-rays in the prompt phase$^{20}$ $E_{\gamma, \text{iso}}$, it represents part of the total GRB fireball energy $E_{\text{total, iso}} = E_{\text{kin, iso}} + E_{\gamma, \text{iso}}$ (e.g. Zhang & Meszaros 2004; De Pasquale et al. 2016). Assuming an efficiency$^{21}$ $\eta = E_{\gamma, \text{iso}}/E_{\text{total, iso}} = 0.2$ we derive

$^{20}$ Here, $E_{\gamma, \text{iso}}$ is the same as $E_{\text{iso}} = 12.7 \times 10^{53}$ erg of §4.2.

$^{21}$ And thus $E_{\text{kin, iso}} = (1/\eta - 1) E_{\gamma, \text{iso}}$. 

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\( \theta_{\text{jet}} = 0.147 \pm 0.017 \) rad, or 8.41 \pm 0.97 degrees. If we consider that the outflow is collimated, the ‘true’ gamma-ray energy of the jet is \( E_\gamma = E_{\text{iso}} (1 - \cos(\theta_{\text{jet}})) \approx 1 \times 10^{52} \) erg. The assumed efficiency is justified theoretically (e.g. Guetta et al. 2001) and by recent studies of GRB afterglows in the optical, X-rays and GeV gamma-rays (e.g. Beniamini et al. 2015). However, higher values are also possible, as suggested by some observations (Zhang et al. 2007a; Lü et al. 2018) and theoretical models (e.g. Kobayashi & Sari 2001; Zhang & Yan 2011). As shown in De Pasquale et al. (2016), the minimum \( E_{\text{total}} \) is obtained for \( \eta = 3/4 \). Lower efficiencies correspond to higher total energies. Therefore, with \( \eta = 0.2 - 0.75 \) we can estimate the ‘total collimated energy’ of the jet to be \( E_{\text{total}} = E_{\gamma}/\eta \approx 3 - 8 \times 10^{52} \) erg. We note that the dependence of \( \theta_{\text{jet}} \) on \( n \) and the kinetic energy is rather weak (Eq. 1). Thus, the total energy is not sizably affected by the exact values of \( n \) and \( E_{\text{kin}} \).

The most widely discussed models of central engines of GRBs are accreting magnetars or accreting black holes. We can assume for a standard neutron star with mass \( M = 1.4 M_\odot \) the maximum rotation energy to be in the range \( 3 \times 10^{52} \) erg (Lattimer & Prakash 2016) – \( 7 \times 10^{52} \) erg (Haensel et al. 2009). Therefore, our analysis allows us to disfavour a standard magnetar as central engine of this GRB. Only the most extreme magnetar models with \( M \approx 2.1 M_\odot \) and rotation energy \( \sim 10^{53} \) erg are not excluded (see Metzger et al. 2015; Dall’Osto et al. 2018; Stratta et al. 2018). On the other hand, according to the Kerr metric (Kerr 1963) the rotational energy \( E_{\text{rot}} \) of a black hole can reach up to 29% of its total mass, which exceeds that of neutron stars by a full order of magnitude. Indeed, rotating black holes of mass \( M \sim 3 M_\odot \) possess rotational energies up to \( E_{\text{rot}} \sim 10^{54} \) erg (e.g. van Putten & Della Valle 2017). Therefore, an energy budget of \( \sim 10^{53} \) erg can be conveniently extracted via the Blandford-Znajek mechanism (Blandford & Znajek 1977), thereby suggesting that the central engine of GRB 210905A may well be a rotating black hole.

4.4. The early X-ray and optical/NIR afterglow

As shown in §3.3, although the optical-to-X-ray SED at 0.1 d is in agreement with the cooling break lying within the X-ray band, both their light curves are not well explained by the standard fireball scenario before the common shallow break at \( \lesssim 0.7 \) d. Here, we can investigate whether our data can justify the early decay and the shallow break.

First, the early break at \( \lesssim 0.7 \) d is not well constrained but we can exclude that it is due to a wind-to-constant-density transition as the light-curve decline, in such a scenario, would become shallower and not steeper (e.g. Panei et al. 2007; Schulze et al. 2011). The times and the slopes instead make it an example of a ‘canonical’ GRB X-ray afterglow light curve (Nousek et al. 2006; Zhang et al. 2007b). Studying the canonical light curve, Zhang et al. (2007b) interpreted the break between the shallow segment with \( \alpha = 0.7 \) to the more ‘normal’ segment with \( \alpha = 1 \) as the end of an ‘energy injection’ phase. During energy injection, the ejecta is still receiving energy, either from a long-lived central engine, or by slower ejecta shells that catch up with the leading shell. In other words, the mild break should be interpreted as cessation of energy injection. Following the relations in Zhang et al. (2006), where \( q \) the energy injection index, we have (for ISM and \( p = 2.2 \)): \( \alpha_{\text{eq}} = (2(p - 4) + (p + 3)q)/4 \), from which follows \( q = 0.84 \) and \( \alpha = 0.69 \) for \( \nu \gg \nu_c \). For X-rays we obtain \( \alpha_X = (2(p - 4) + (p + 2)q)/4 \), so \( q = 0.84 \) and \( \alpha = 0.98 \) for \( \nu \gg \nu_c \). Using a stratified shell model with ejected mass \( M(\nu) \propto \nu^{-4} \), where \( \nu \) is the Lorentz factor of the shell (Rees & Mészáros 1998) and the relation between \( s \) and \( q \) parameters (\( s = (10 - 7q)/(2 + q) \), Zhang et al. 2006), we find that a value of \( s \approx 1.45 \) fits the pre-break behaviour. Equally, a magnetar central engine model that continuously injects energy as \( L(t) \propto t^{-\mu} \) (e.g. Dai & Lu 1998), can model the early decay with \( q = 0.84 \). Therefore, we cannot discard one model over the other, specifically stratified shell versus magnetar. However, as discussed in §3.5, the energy constraints likely limit the viability of a new-born magnetar as the power source of the energy injection.

We note also that the theoretical energy-injected \( \alpha_X \sim 0.98 \) is larger than the value observed (0.74). However, the theoretical value assumes \( \nu > \nu_c \) but in Table 3 we see that \( \nu_c \) is well within the X-ray band in the first day after the burst trigger. The energy injection changes the way the cooling frequency evolves, which is \( \nu_c \propto \nu^{-0.58} \) for \( q = 0.84 \), and thus the cooling frequency evolves slightly faster whilst energy injection is happening: if the cooling frequency is at \( \sim 2 \) keV at 0.1 d, then at 1 d it would have been at 0.5 keV, and consistent with what we observe. Moreover, one should also consider that the \( \nu_c \) break is likely smooth and covers a relatively large interval (e.g. Granot & Sari 2002). Therefore, the observed temporal decay index may well be somewhere between the values predicted for the \( \nu < \nu_c \) and the \( \nu > \nu_c \) cases, i.e. \( 0.69 < \alpha_X < 0.98 \), in agreement with the observed value before \( \sim 70 \) ks.

4.5. The nature of the H-band flattening at late times

The likely discovery of the host of GRB 210905A is a rare discovery, given that up to July 2022 only three hosts (those of GRBs 050904, 130606A, and 140515A) had been confirmed at \( z > 6 \) (McGuire et al. 2016), and four if we consider the possible detection of the GRB 060522 host (Tanvir et al. 2012).

The observed brightness of the source detected with \( HST \) in the \( F140W \) band corresponds to a rest-frame \( m_{1400} \sim 21 \) mag, which is consistent with the characteristic magnitude at 1600 \( \AA \) of \( z = 6 - 7 \) galaxies (e.g. Bouwens et al. 2021). Therefore, such a galaxy is not unusual, although it is more luminous in the UV than galaxies that contribute the most to the star formation at these redshifts. In the following, we make use of the brightness of \( H_{AB} = 25.8 \) mag resulting from the light curve fitting. A host galaxy at \( z = 6.3 \) with such a brightness and thus a rest-frame UV luminosity of \( L_\nu = 1.47 \times 10^{50} \) erg s\(^{-1}\) Hz\(^{-1}\) would have a SFR \( \sim 16 M_\odot \) yr\(^{-1}\) using equation 1 in Kennicutt (1998). This is certainly an acceptable value (see also the discussion in Saccardi et al. 2022), and in fact McGuire et al. (2016) find that the \( z > 6 \) GRB hosts known to date likely have similar SFR, assuming a short-lived burst of star formation (see also Tanvir et al. 2012). If its brightness is confirmed with further observations, the host of 210905A would also be the brightest. We caution however, that at this stage it is not possible to separate some contamination from a possible foreground source discussed in Saccardi et al. (2022), and thus the host can be fainter and the inferred SFR lower.

One could also speculate whether a SN can contribute to the final observation. However, note that a SN should reach an absolute magnitude of \( m_{2200} \sim 21 \) mag in the far UV (\( H \)-band in the observer frame). This is four times more than the most luminous GRB-SN confirmed spectroscopically, SN2011kl associated with GRB 111209A (\( m_{2353} \sim 19.6 \) mag at peak, Greiner et al. 2015; Kann et al. 2019), although Kann et al. (2021) have
recently claimed the existence of an even more luminous SN associated with GRB 140506A with $M_Z \approx -20.5$ mag. As we find no evidence that GRB 210905A is more than just a very energetic but otherwise typical long GRB, there is no reason to claim the GRB would be accompanied by an extremely UV-luminous SN of a type not seen associated with GRBs before.

In the following we explore possible alternatives to the above interpretation of a jet break well visible in X-rays but hidden in the NIR by a constant source that becomes dominant. Thus, we consider the possibility that the afterglow still contributes substantially and the $H$ band can be modelled with a single power-law, with a chromatic break in the X-rays. It is also possible to speculate that the late light curve is the consequence of a spectral break moving between optical and X-rays. However, this not only contradicted by the elongated, extended, and offset nature of the $HST$ detection but, as we have shown in Section §3.3, our SED analysis which shows $\nu_0$ being already between the optical and X-ray bands after 0.7 d. Therefore, it is not possible to invoke the presence of an additional break in the slow-cooling regime moving into the band after this time. Moreover, the change in the temporal index is inconsistent with the passage of the cooling break, which should be $\Delta n = 0.25$ in the slow-cooling and uniform-medium environment, and additionally incompatible with other regimes, for example fast cooling or a wind-blown environment.

Another possibility is to consider a bright reverse shock, but this requires either a large energy gradient or a big difference in shell velocity, both of which are inconsistent with the gradual energy injection scenario which explains the early light curve until 0.7 d. We could invoke a second, discrete shell with energy that is less than or comparable to the first (post initial energy injection) shell but much faster i.e. a delayed launch. However, the shell would have to conveniently collide with the leading shell at about the jet-break time and the optical excess would be the contribution from the reverse shock. This is not only an incredible coincidence, but would also approximately double the total energy requirements to explain a second shell, making this event even more extreme. We cannot exclude that a more mild shock could however explain the X-ray data at 10–20 d, which lies just above the analytical modelling of the light curve. We also cannot, confirm this possibility with the few data points available, which are anyway within 2$\sigma$ of the analytical model.

In conclusion, we consider the detection of the host and/or an intervening galaxy (or a mix of the two) as the strongest and most plausible explanation for the flattening of the $H$-band light curve.

4.6. The X-ray afterglow in context

To put the X-ray emission in the context of other GRB afterglows, in particular high-redshift GRBs, we retrieved from the Swift Burst Analyst (Evans et al. 2010) the X-ray light curves of 421 long-duration Swift GRBs with detected X-ray afterglows (detected in at least two epochs) and known spectroscopic redshifts, which were discovered before the end of July 2022. We processed the data and moved them to their rest-frames following Schulze et al. (2014). Figure 9 shows the parameter space occupied by long-duration GRBs as a density plot and the X-ray light curve of GRB 210905A in blue. We have also included the X-ray light curves of the high-redshift GRBs 090423, 090429B and 100905A that have only a photometric redshift. The uncertainty in luminosity for these three bursts is indicated by red-shaded regions around the light curves at their redshifts.

GRB 210905A's X-ray afterglow is among the most luminous at all times. Even compared to other GRBs at $5 < z < 6$ (green; GRBs 060522, 060927, 130606A, 131227A, 140304A, 201221A, 220521A) and $z > 6$ (red; 050904, 090813, 090423, 090429B, 100905A, 120521C, 120923A, 140515A), GRB 210905A has an exceptionally high luminosity. Furthermore, its X-ray afterglow is fading slower than those of most GRBs, at least until the jet break at $\sim 5 \times 10^5$ s in the rest frame (§3.5). Here we note that some of the other bursts at high-z do not show a clear light-curve break in X-rays (GRBs 050904, 090813, 090423, 130606A), although some of them show a break in the optical (GRBs 050904, 090423, 090429B, 120521C), and GRB 140515A has just one single detection that suggests a possible break similar to GRB 210905A. This is because of the low observed flux of these very high-z afterglows, as only the most luminous events are bright enough for Swift/XRT.

4.7. The optical/NIR afterglow in context

Following the method devised by Kann et al. (2006), we are able to put the NIR afterglow into the context of the (optical/NIR) total afterglow sample. We derive the observer-frame $R_C$ magnitude by shifting all data to the $H$ band, then extrapolating the spectral slope into the observer-frame $R_C$ band, which is completely suppressed at the redshift of the GRB (assuming that there would be no Lyman absorption). The spectral slope, redshift, and the lack of extinction are then used to derive the magnitude shift $dR_C = -5.12_{-0.21}^{+0.20}$ mag to $z = 1$. The derived $R_C$-band light curve still represents an observed magnitude, it is as if the GRB were at $z = 1$ in a completely transparent universe.

We then compare the afterglow with the GRB afterglow light curve samples of Kann et al. (2006, 2010, 2011) as well as samples from upcoming publications (Kann et al., 2022a,b,c, in prep.). The result is shown in Figure 10, with GRB 210905A highlighted in red. The sample of Kann et al. (2022a), in prep. focuses on $z \geq 6$ GRBs, and these light curves are highlighted.
4.8. Dust absorption and equivalent hydrogen column densities

As in other high-\(z\) bursts (see e.g. Zafar et al. 2010, 2011, 2018; Melandri et al. 2015), GRB 210905A is characterised by negligible absorption in the optical/NIR, in agreement with those expected for high-\(z\) galaxies populating the faint end of the luminosity function (e.g. Salvaterra et al. 2011). In particular, McGuire et al. (2016) studied three \(z > 5.9\) GRB hosts and noted that afterglow analyses in each case pointed to low-line-of-sight dust extinction.

Although a low \(A_V\) is expected to correlate with a low \(N_{H, X}\), the high \(N_{H, X}\) value of \(7.7^{+1.5}_{-1.2} \times 10^{22} \text{ cm}^{-2}\) is also not exceptional. It is also observed in other environments, for example in AGNs, and can be naturally explained by the absorption of intervening metals along the line-of-sight (Starling et al. 2013; Campana et al. 2015), which reside almost entirely in the neutral gas at \(z > 4.5\) (e.g. Péroux & Howk 2020), although one cannot exclude the contribution of increasing gas density in the vicinity of the GRB (Heintz et al. 2018). The high \(N_{H, X}\) is also in contrast with the \(N_{HI} \approx 1.35 \times 10^{21} \text{ cm}^{-2}\) measured via the Lyman-\(\alpha\) absorption-line by Fausey et al. (in prep.).

The difference can be explained by the very high number of ionising photons produced by the GRB that could ionise the IGM along the line-of-sight up to several hundreds of pc (Saccardi et al. 2022). We discuss the IGM contribution in more detail in Fausey et al. (in prep.).

4.9. X-ray afterglow luminosity versus prompt energy

The X-ray luminosity and the isotropic gamma-ray energy release seem to broadly follow a linear relation as already shown by De Pasquale et al. (2006) (see also Nysewander et al. 2009), suggesting a roughly universal efficiency for converting a fraction of the initial kinetic energy\(^{23}\) into gamma-ray photons. This was later further confirmed by D’Avanzo et al. (2012). GRBs at \(z > 6\) also follow this relation. We test here whether GRB 210905A follows this relation despite its luminosity. We estimate the afterglow X-ray integral flux in the 2–10 keV rest-frame common energy band and compute the corresponding rest-frame X-ray luminosity at different rest-frame times. The 2–10 keV rest-frame flux was computed from the observed integral power law.

\[
f_{X, ff}(2–10 \text{ keV}) = f_X(0.3–10 \text{ keV}) \frac{(10^{20} \text{ erg s}^{-1})^{-2} \Gamma}{10^{20} \text{ erg s}^{-1} - 0.3^{2} \Gamma}. \tag{2}
\]

The obtained X-ray light curve was then fitted with a multiply broken power-law, after removing the time intervals showing significant flaring, and then the fits were interpolated or extrapolated to the rest-frame times \(t_{ff} = 5 \text{ min}, t_{ff} = 1 \text{ hr}, t_{ff} = 11 \text{ hr}, t_{ff} = 24 \text{ hr}\). As shown in Figure 11 the properties of GRB 210905A are fully consistent with the \(E_{iso} - L_X\) correlations found for long GRBs by D’Avanzo et al. (2012).

4.10. Long-GRB progenitors at high redshift

At high redshift the Universe is expected to be populated by pop-III stars, the first stars that formed out of gas clouds of as thick black curves. At early times, the afterglow of GRB 210905A is seen to be among the most luminous known, albeit still fainter than the early afterglows of high-\(z\) GRBs 130606A and especially 050904 (Kann et al. 2007). Interestingly, the early flash of GRB 210905A aligns well in rest-frame time (between 70 and 110 s) with those seen in GRB 050904 (Boer et al. 2006). GRB 160625B (Troja et al. 2017, an extremely energetic lower-redshift GRB, highlighted in blue), and, with less contrast, in GRB 130606A (Castro-Tirado et al. 2013). On the other hand, several bright prompt-associated flashes happen significantly earlier, such as the cases of GRB 080319B (Racusin et al. 2008; Bloom et al. 2009) and GRB 120711A (Martin-Carrillo et al. 2014, Kann et al. 2022a, in prep.). Therefore, this similarity in time is likely just a chance coincidence.

An interesting result is found towards the end of the light curve. After removing the potential constant component, the combination of a late break and an early shallow decay makes the afterglow of this burst the most luminous ever detected for a certain time span, before the shallower post-break decay of the afterglow of GRB 160625B (which itself had a very late jet break. Kangas et al. 2020) makes the latter the most luminous known at very late times again (Kann et al. 2022c, in prep.). This provides further evidence for the extremely energetic nature of GRB 210905A.

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\(^{23}\) Not to be confused with \(E_{iso}\), which is the energy left after the prompt phase.
pristine composition. Chemical feedback from the supernova explosions of these very massive stars produces metal enrichment within star-forming clouds, raising the metallicity above a critical threshold above which we expect a slow transition of the SFR from massive pop-III to solar-size pop-II and pop-I stars (e.g. Schneider et al. 2006; Maio et al. 2010). Determining how this transition takes place is one of the main missing ingredients to understand galaxy formation in the early Universe. All models (e.g. Mészáros & Rees 2010; Toma et al. 2011; Piro et al. 2014) predict pop-III GRBs to be very energetic events, and with very long intrinsic durations of $10^4$ s, making their detection possible even at the highest redshifts. In particular, Toma et al. (2011) suggested that they can release an equivalent isotropic energy up to $\sim 10^{56-57}$ erg.

In Figure 12 we compare $\theta_{\text{jet}}$ and the collimated energy $E_{\gamma}$ of GRB 210905A with the KW sample of 43 long GRBs with reliable jet-break time estimates (Tsvetkova et al. 2017, 2021). Considering the uncertainty on the collimation-corrected energy, GRB 210905A lies just outside the 1σ confidence level of the $E_{\text{peak},z} \sim E_{\gamma}$ (‘Ghirlanda’ relation, see Ghirlanda et al. 2004, 2007) and thus well compatible with this relation. The energy values involved in GRB 210905A, both isotropic and collimated, are large but do not significantly differ from those at low redshift (see Figs. 8 and 12). At lower $z$, other events have produced $E_{\gamma} \gtrsim 10^{54}$ erg isotropically and $E_{\gamma} \approx 10^{52}$ erg collimation-corrected (see also Cenko et al. 2011). The most outstanding example is GRB 130427A at $z = 0.3399$ (see §4.1), the most powerful GRB at $z < 0.9$ (e.g. Maselli et al. 2014; De Pasquale et al. 2016).

GRB 210905A has the highest $E_{\gamma}$ in the Konus-Wind catalogue. This and the large $E_{\text{peak},z}$ suggest a large bulk Lorentz factor $\Gamma_0$ of the jet. The afterglow light curve, as reported in Figure 3, decays as a power-law in both the optical and X-ray band from $\gtrsim 5000$ s onwards (observer frame). This suggests that the afterglow deceleration time happened before this epoch. Following the method in Molinari et al. (2007), an upper limit on this peak time provides a lower limit to the maximum bulk Lorentz factor of the jet, namely $\Gamma_0 \gtrsim 200$, assuming a constant density medium $n_0 = 1$ cm$^{-3}$ and the isotropic energy of GRB 210905A ($\S4.2$). With this estimate, and the inferred half-opening angle, the burst is consistent with the $\theta_{\text{jet}} \sim \Gamma_0$ broad anti-correlation reported in Ghirlanda et al. (2012, see their Figure 4).

Therefore, GRB 210905A, although extremely bright, is not separated markedly from other classical GRBs at low redshift. In summary, no features of this event point to a pop-III origin.

### 4.11. Star-formation rate at very high redshift

The rate of GRBs can be used to estimate the SFR in the remote Universe (see §1). Recently, Lloyd-Ronning et al. (2019, 2020a,b, and references therein) have argued that at high redshift, the GRB jets were, on average, narrower than those of closer GRBs (see also Laskar et al. 2014, 2018). This would imply that more stars formed at high redshift than previously estimated, unless the GRB properties, and thus their rate, are extremely environment-sensitive (Kistler et al. 2008, 2009; Robertson & Ellis 2012; Jakobsson et al. 2012; Tanvir et al. 2012; Japelj et al. 2016; Palmerio et al. 2019). In the left panel of Figure 12, we report the relation found by the above authors in the $\theta_{\text{jet}} = (1 + z)$ plane. GRB 210905A is an outlier event located at $2 - 3\sigma$ above this relation. We observe that the half-opening angle of this GRB at $z = 6.312$ ($\theta_{\text{jet}} \approx 8$ deg) is consistent with the median value of $\theta_{\text{jet}} = 7.4_6^{+11}_{-6}$ deg for GRBs at $z \sim 1$ but larger than the mean of $\theta = 3.6 \pm 0.7$ deg found using three $z > 6$ bursts (GRBs 050904, 090423, 120521C, Laskar et al. 2014, 2018). When we include GRB 210905A, the mean for the $z \geq 6$ bursts is $\theta = 4.8 \pm 0.6$ deg, closer to the best value for $z \sim 1$ events. These findings would argue against a putative inverse correlation between $z$ and $\theta_{\text{jet}}$.

### 5. Conclusions

GRB 210905A was a long burst at redshift $z = 6.312$. Our extensive and prompt follow-up observations from optical/NIR to X-ray and gamma-ray bands, starting in the first seconds, have allowed us to study in detail both the prompt and the afterglow phases. We carried out a joint time-resolved analysis of the last of the three pulses of the prompt emission, which is shown to be in agreement with synchrotron emission, similar to other bursts at lower redshifts. Among the sample of ten $z \geq 6$ GRBs known to date, GRB 210905A stands out (together with GRB 050904), having the highest isotropic energy release and among the highest afterglow luminosity at late times, while still being consistent with the range of values found for other long GRBs.

The temporal evolution of the afterglow can be interpreted as due to energy injection followed by a decay well in agreement with the slow-cooling scenario and a constant-density (‘ISM’) circumburst medium profile within the standard fireball theory. However, the optical and X-ray afterglows are among the most luminous ever detected, in particular in the optical range at $t \gtrsim 0.5$ d in the rest frame, due to very slow fading and a late jet break. In late HST imaging, we find evidence for an underlying host with UV luminosity slightly larger than that of galaxies contributing the most to star formation at $z = 6 - 7$. If confirmed with further observations, the host of GRB 210905A would be the fourth and the brightest GRB host at $z > 6$ detected to date. It would also be bright enough to be characterised via spectroscopy.

\[ \text{To derive the peak time, we assumed smoothness } k = 1 \text{ and decay indices } \alpha_0 = -0.7, \alpha_0 = \alpha_1, \text{opt} = 0.69, \text{ before and after the peak, respectively.} \]
with the JWST (e.g. McGuire et al. 2016), providing one of the first and better estimates on the SFR, metallicity and dust content of a GRB host at very high redshift.

The jet break at ∼ 50 d (observer frame) results in a half-opening angle that is larger than that of other z > 6 bursts, thus putting into question the putative inverse dependence of the half-opening angle on redshift. The large total energy budget of $E_{\text{total}} > 10^{52}$ erg associated with this GRB likely excludes all but the most extreme magnetar models as a central engine of this GRB. Therefore, our analysis leaves the Kerr black hole as the preferred scenario for the central engine of GRB 210905A. Finally, the shallow evolution before 1 day suggests that the black hole injected energy via stratified mass ejecta with different Lorentz factors.

In summary, this burst is consistent with the ‘Amati’, ‘Ghirlanda’, and ‘Yonetoku’ relations. This fact, and the agreement with the $E_{\text{iso}} - L_X$ plane show that GRB 210905A is a very energetic event but still in the upper tail of the prompt energy and X-ray luminosity distributions of long GRBs. It is not unexpected that our view of the high-z GRB Universe is biased towards the most luminous events, simply because our instruments are limited in sensitivity. In other words, despite its outstanding luminosity it is unlikely that the origin of this GRB is different from those of low-redshift GRBs, such as a pop-III progenitor.

Gamma-ray bursts at $z \geq 6$ are rare events from the perspective of today’s follow-up capabilities, but they are just a small part of a larger population that future proposed missions promise to uncover (e.g. THESEUS, Amati et al. 2018; Gamow, White et al. 2021) and, in synergy with the largest ground- and space-based telescopes (such as the James Webb Space Telescope), to answer open questions in modern astrophysics such as the identification of the sources responsible for cosmic reionisation, and the evolution of SFR and metallicity across the transition from pop-III stars to pop-II and pop-I stars.

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References


![Diagram](image-url) Fig. 12. Collimated parameters of GRB 210905A (red symbols) compared to a KW sample of 43 long GRBs from Tsvetkova et al. (2017, 2021). We assumed $n = 0.2$ and $r = 1$ cm$^{-1}$ for all bursts. Left: Half-opening angle $\theta_{0,z}$ versus redshift. The dashed line within the grey area shows the relation found in Lloyd-Ronning et al. (2020a) with its error. Right: $E_{\gamma} - E_{\text{peak},z}$ diagram. As in Figure 8, the colour of each data point represents the burst’s redshift. The ‘Ghirlanda’ relation is plotted together with its 68% and 90% PIs.