








LETTER TO THE EDITOR

Perihelion observations of interstellar comet 3I/ATLAS with the IRAM 30-m telescope[★]

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ABSTRACT

Comet 3I/ATLAS is the third interstellar comet identified as passing through the Solar System. Its high outgassing activity and favourable perihelion passage on October 29, 2025 UT provided an excellent opportunity to investigate the composition of its coma gases through millimeter spectroscopy. We present observations undertaken with the IRAM 30-m telescope on November 1–3, 2025 at an heliocentric distance of 1.36–1.37 au. Lines of HCN, CH₃OH, CO, and H₂CO are well detected, and $\sim 4\sigma$ detections are obtained for CS and CH₃CN. The search for H₂S was unsuccessful. Abundances of CO, H₂CO, CH₃OH, and CH₃CN relative to HCN are in the upper ranges of values measured in Solar System comets. The sulfur-to-carbon abundance ratio in 3I/ATLAS's coma is at most the minimum value observed in comets. The unusually low expansion velocity of coma gases suggests a near-nucleus gas flow driven by heavy molecules such as CO₂, and/or a large fraction of the gaseous production coming from sublimating icy grains.

Key words. comets: general – comets: individual: 3I/ATLAS

1. Introduction

Comets are fingerprints of the formation of planetary systems. Water and the numerous complex molecules composing the ices of their nucleus were likely synthesized in the molecular cloud precursor of the Solar System (Altwegg et al. 2019; Ceccarelli et al. 2023; Biver et al. 2024b). The passing of a bright cometary-like interstellar object through the Solar System was long awaited to investigate the physical and chemical properties of icy planetesimals in other planetary systems. The first interstellar object, 1I/'Oumuamua, was discovered in October 2017 a few days after its closest approach to Earth. There was no detection of coma gases and dust particles ('Oumuamua ISSI Team 2019, and references therein), but astrometric measurements suggested the presence of non-gravitational forces related to cometary activity (Micheli et al. 2018). The second interstellar object 2I/Borisov, discovered in 2019, showed cometary activity and was extensively observed. Several simple volatile species commonly observed in Solar System comets (OH, C₂, CN, HCN, NH, NH₂, CO, Ni, Fe) were detected (e.g., Opitom et al. 2019, 2021; Cordiner et al. 2020; Deam et al. 2025). Spectroscopic properties and mixing ratios were found to be remarkably similar to Solar System comets.

Comet 2I/Borisov was found to be carbon-chain depleted, similar to the class of carbon-chain depleted comets of Solar System origin (Opitom et al. 2019; Kareta et al. 2020; Lin et al. 2020). However, a high enrichment in CO relative to H₂O was revealed, suggesting specific formation conditions in its natal protoplanetary disk (Bodewits et al. 2020; Cordiner et al. 2020).

In this letter, we report on spectroscopic observations with the 30-m telescope of the Institut de radioastronomie millimétrique (IRAM) of the third identified interstellar object, 3I/ATLAS (also named 3I/2025 N1 (ATLAS)), which were performed near its perihelion on October 29.48, 2025 UT, at $r_h = 1.356$ au from Sun. This interstellar object was discovered with the robotic telescope ATLAS in Chile on July 1, 2025, as it was at 5 au from the Sun, and was confirmed to be a comet the day after (CBET 5578, Seligman et al. 2025). Subsequent studies have revealed a fast increase of its gaseous activity as it approached the Sun, making the detection of molecular species at millimeter wavelengths possible. Detections of HCN and CH₃OH were reported from observations with the James Clerk Maxwell Telescope at $r_h = 2.1$ au (Coulson et al. 2026) and $r_h = 1.39$ au (Kuan et al. 2025), respectively. The HCN and CH₃OH molecular lines were mapped with the ALMA Compact Array (ACA) in a time range spanning 2.6–1.7 au pre-perihelion, from which it was inferred that the CH₃OH/HCN ratio in 3I/ATLAS is among the largest values measured in any comet (Roth et al. 2026). In this letter, we present the detection of HCN, CH₃OH, CO, H₂CO, CH₃CN and CS at

[★] Based on observations carried out with the IRAM 30-m telescope. IRAM is supported by INSU/CNRS (France), MPG (Germany) and IGN (Spain).

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Table 1. Production rates and abundances in comet 3I/ATLAS on November 1–3, 2025.

Molecule	Prod. Rate ^a (10 ²⁵ molec.s ⁻¹)	Abundance relative to water		Abundance relative to HCN	
		3I/ATLAS ^b	Comets	3I/ATLAS	Comets
HCN	4.1 ± 0.2	0.062 ± 0.003%	0.08–0.25%	1	1
HNC	0.5 ± 0.2 ^c	<0.009%	0.0015–0.035% ^d	<0.13	0.01–0.29 ^d
CH ₃ CN	1.5 ± 0.4	0.023 ± 0.006%	0.008–0.054%	0.36 ± 0.09	0.07–0.52
CH ₃ OH	338 ± 5	5.1 ± 0.1%	0.4–6.1%	83 ± 4	4–64 ^e
H ₂ CO	35 ± 4 ^c	0.53 ± 0.06%	0.13–1.4%	8.6 ± 1.0	1.3–13
CO	680 ± 11	10.3 ± 1.7%	0.4–35%	167 ± 8	8–270 ^e
H ₂ S	<39	<0.59%	0.09–1.5%	<9	1.5–9
CS	1.7 ± 0.4 ^c	0.024 ± 0.006%	0.05–0.20% ^d	0.34 ± 0.11	0.2–3.2 ^d
HC ₃ N	<2.1	<0.031%	0.002–0.068%	<0.52	0.016–0.42
OCS	<50	<0.75%	0.05–0.40%	<12	0.43–3.5
CH ₃ CHO	14 ± 7	<0.32%	0.04–0.08%	<5.1	0.1–0.8
(CH ₂ OH) ₂	33 ± 13	<0.59%	0.07–0.35%	<9.5	0.8–2.9
SO	<27 ^c	<0.41%	<0.03–0.30%	<7	0.3–4.6
HNCO	25 ± 11	<0.52%	<0.01–0.62%	<8	0.2–4.8
HCOOH	<109	<1.7%	<0.04–0.68%	<27	0.6–7.2
NH ₂ CHO	<4.8	<0.08%	0.015–0.022%	<1.2	0.09–0.3
HDO	<121	<1.83%	0.028–0.12%	–	–
CH ₃ OD	<17.9	<0.27%	<0.02%	–	–

Notes. ^(a)Uncertainties correspond to propagated 1 σ uncertainties on the line areas, and do not include an estimated additional uncertainty of 7% related to model parameters. ^(b)Assuming the water production rate $Q_{\text{H}_2\text{O}} = 6.58 \times 10^{28}$ molec.s⁻¹ measured from MAJIS/JUICE (Bockelée-Morvan et al. 2026). Uncertainties in $Q_{\text{H}_2\text{O}}$ are not taken into account. ^(c)Daughter distributions with $L_p = 0$ –1600, 1700, 2000, and 3000 km are assumed for HNC, CS, H₂CO and SO, respectively. Due the large beam size (15 000–26 000 km), distributed sources with $L_p \leq 3000$ km provide similar results. ^(d)Decreasing abundance with increasing r_h . ^(e)Excluding values in C/2016 R2 (PanSTARRS): 280, 2.6×10^4 for CH₃OH/HCN and CO/HCN, respectively.

2000 km (i.e., 0.5 times the H₂CO photodissociative scalelength) based on measurements in Solar System comets (Roth et al. 2021; Biver et al. 2024a). The derived H₂CO production rate is not significantly dependent on the assumed L_p since the beam size is large (~ 17 000 km diameter at the frequency of the H₂CO line)¹. For SO and CS, we assumed $L_p = 3000$ km and 1700 km, respectively (Biver et al. 2024a). For other molecules (except HCO⁺, which we do not consider in the analysis), we assumed direct release from the nucleus. Derived production rates (or their upper limits) are listed in Table 1.

The determination of mixing ratios with respect to water relies on contemporaneous observations of H₂O or of its photodissociation products (OH, H). Bockelée-Morvan et al. (2026) report $Q_{\text{H}_2\text{O}} = (6.58 \pm 0.13) \times 10^{28}$ molec.s⁻¹ on November 2 from infrared H₂O data with a $3 \times 3'$ field of view obtained with the Moons And Jupiter Imaging Spectrometer (MAJIS) on board the Jupiter Icy Moons Explorer (JUICE). On the other hand, Combi et al. (2026), from Lyman- α data obtained with SOHO/SWAN, report a large water production rate of $Q_{\text{H}_2\text{O}} = 3.2 \times 10^{29}$ molec.s⁻¹ at 1.4 au post-perihelion (November 6)². The large field of view of SOHO/SWAN was possibly sensitive to large-scale water production from subliming icy grains. Extended water production from icy grains (i.e., hyperactivity) at perihelion is suggested from the icy area fraction (see Lis et al. 2019) of $\sim 900\%$ at 1.4 au that we derived from the $Q_{\text{H}_2\text{O}}$ mea-

sured by SOHO/SWAN and setting a nucleus radius of 1.3 km as inferred from HST observations (Hui et al. 2026). To interpret the IRAM data ($\sim 11''$ field of view), we used the $Q_{\text{H}_2\text{O}}$ value of 6.6×10^{28} molec.s⁻¹ deduced from the MAJIS data. Observations of OH 18-cm lines at the Nançay radio telescope ($3.5' \times 19'$ beam) yielded $Q_{\text{OH}} = (7.0 \pm 0.6) \times 10^{28}$ molec.s⁻¹ for October 13–19 ($r_h = 1.43$ au) (Crovisier et al. 2025)³ and $Q_{\text{OH}} = (6.1 \pm 1.4) \times 10^{28}$ molec.s⁻¹ for November 18–22 ($r_h = 1.57$ au).

Derived production rates and mixing ratios can be compared to pre-perihelion measurements from when the comet was farther from the Sun. The CO/H₂O mixing ratio at $r_h = 1.36$ au ($\sim 10\%$) is much lower than the value at 3.32 au ($\sim 165\%$) derived from JWST observations (Cordiner et al. 2025b). This is not surprising, as water sublimation is not efficient at $r_h \geq 3$ au. The CH₃OH production rate at $r_h = 1.36$ au is a factor of two lower than the trend with r_h established from ACA observations in the range 2.6–1.7 au (Roth et al. 2026), suggesting that the activity of 3I/ATLAS leveled off when approaching perihelion. The CH₃OH/H₂O ratio ($5.1 \pm 0.1\%$) is, within 1.5 σ , consistent with the pre-perihelion ($r_h = 1.43$ au) value of $8 \pm 2\%$ derived by Roth et al. (2026).

4. Discussion

The molecular species detected in this millimeter spectral survey of 3I/ATLAS are those which are commonly observed in Solar System comets (e.g., Biver et al. 2024a,b), indicating that the parent protoplanetary disk (and natal molecular cloud) of this icy object share the same chemical properties, at least at a first approximation. However, the derived abundances in 3I/ATLAS's coma are somewhat different from those of Solar System comets

¹ The ACA observations suggest that H₂CO is more extended in 3I/ATLAS (M. Cordiner, personal communication). A production rate 9% and 18% higher than given in Table 1 is obtained for $L_p = 6000$ and 8000 km, respectively.

² Tan et al. (2026) report $Q_{\text{PH}_2\text{O}}$ values five to six times smaller using the same SOHO/SWAN data set but also conclude that 3I/ATLAS is an hyperactive comet. Still unpublished MAVEN data are consistent with results from Combi et al. (2026) (J. Deighan, personal communication).

³ Updated value from a reevaluation of the maser inversion.

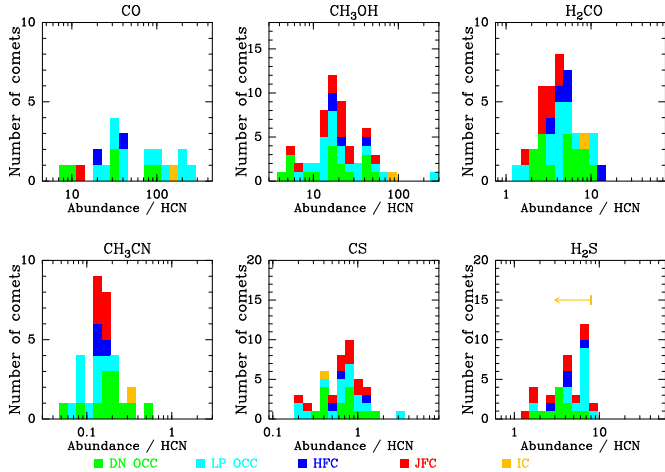


Fig. 3. Histograms of the abundances relative to HCN measured from observations at radio wavelengths. Dynamically new (DN) and long-period (LP) comets from the Oort cloud (OCC), Halley-family comets (HFC) and Jupiter-family comets (JFC) are indicated in green, cyan, blue, and red colors, respectively. The orange color is for 3I/ATLAS. See Biver et al. (2024b) and references therein.

(Table 1). The first striking difference is the low ($\sim 0.06\%$) HCN/H₂O mixing ratio. For comparison, the HCN/H₂O ratio measured in the radio is $\sim 0.1\%$ for most Solar System comets (Biver et al. 2024b). Also, striking is the low CS abundance, which is in the low end of the range of values in Solar System comets. On the other hand, CO/H₂O and CH₃OH/H₂O are enhanced with respect to the mean abundances in Solar System comets, while H₂CO/H₂O and CH₃CN/H₂O are typical (Fig. E.1). From histograms of the abundance ratios with respect to HCN (Fig. 3), 3I/ATLAS displays high CO/HCN, H₂CO/HCN, CH₃OH/HCN, and CH₃CN/HCN abundance ratios and a rather low CS/HCN ratio. As concluded by Roth et al. (2026), 3I/ATLAS’s coma is especially highly enriched in its CH₃OH/HCN abundance, being only surpassed by the peculiar CO-rich and N₂-rich comet C/2016 R2 (PanSTARRS). It would be tempting to conclude to a high elemental C/N ratio in 3I/ATLAS nuclear ices, but HCN is not the dominant carrier of nitrogen in cometary ices (Biver et al. 2024b). Figure F.1 shows a histogram of the ratio of the total abundance of sulfur-bearing species to the total abundance of carbon-bearing molecules in comets, considering only the CO, CH₃OH, H₂CO, CS and H₂S molecules observed in 3I/ATLAS. The value for 3I/ATLAS is at most the minimum value observed in comets, illustrating the peculiar chemistry of this interstellar object.

From the observation of the HDO line at 241.56 GHz (Table B.1), we derived an upper limit for the D/H in H₂O of $<0.91\%$ (at 3σ). This value is a factor of 10 or more higher than the value measured in comets (e.g., Müller et al. 2022; Biver et al. 2024b; Cordiner et al. 2025a). For CH₃OH, we obtained an upper limit on the D/H value of 5.3% by stacking 17 CH₃OD lines between 226 and 272 GHz. This is a factor of 4 larger than the value measured in comet 67P (Drozdovskaya et al. 2021), thereby excluding extreme deuteration levels for this molecule in 3I/ATLAS.

Perhaps, the most surprising property of 3I/ATLAS’s coma is the very low expansion velocity, which is a factor of two below the value of $\sim 0.7\text{--}0.8\text{ km s}^{-1}$ measured in comets of similar activity levels and at comparable heliocentric distances. This

low velocity might be the signature of a gas flow dominated by heavy molecules such as CO₂ (Fig. G.1). This molecule was found to be overabundant in 3I/ATLAS at 3.3 au (Cordiner et al. 2025b). Comet 3I/ATLAS would then share similarities with the hyperactive comet 103P/Hartley 2, whose nucleus released large amounts of CO₂ at its perihelion, dragging water ice grains and chunks that subsequently sublimated in the coma (A’Hearn et al. 2011). The amounts of water and CO₂ molecules released from the nucleus of 103P/Hartley 2 are comparable (Fougere et al. 2013), which might explain why the gas expansion velocity in this comet (0.6 km s^{-1} , Hartogh et al. 2011) was not abnormally low (Appendix G, Fig. G.1). Alternatively, the low velocity measured in 3I/ATLAS could be directly related to a large amount of coma gases released from icy grains. This is predicted by gas flow kinetic models including release both from the nucleus and from icy grains using a direct simulation Monte Carlo approach, and the effect is all the more significant since the amount of icy grains is high (Fougere et al. 2012).

In conclusion, the observations of comet 3I/ATLAS with the IRAM 30-m telescope allowed us to identify six molecules in its coma with relative abundances that are at the extreme ends of the range of values measured in Solar System comets. Advanced modelling is needed to understand the unique dynamic properties of the atmosphere of this comet.

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Appendix A: Log of the observations

Table A.1. Log of observations.

UT date (yyyy/mm/dd.d–dd.d)	$\langle r_h \rangle$ (au)	$\langle \Delta \rangle$ (au)	pwv	Integ. time (min)	Freq. range ^a (GHz)
2025/11/01.33–01.48	1.361	2.276	1.2 mm	37	248.3–256.4, 264.3–272.4
01.51–01.54	1.361	2.274	0.8 mm	33	224.6–232.7, 240.6–248.7
01.57–01.58	1.361	2.274	2 mm	16	146.8–154.9, 162.8–170.9
2025/11/02.34–02.45	1.364	2.265	12–5 mm	94	146.8–154.9, 162.8–170.9
02.47–02.51	1.364	2.264	3 mm	32	248.3–256.4, 264.3–272.4
02.53–02.56	1.365	2.263	1.9 mm	37	224.6–232.7, 240.6–248.7
2025/11/03.32–03.36	1.368	2.254	2.7 mm	37	146.8–154.9, 162.8–170.9
03.38–03.46	1.368	2.253	2.6 mm	65	248.3–256.4, 264.3–272.4
03.48–03.58	1.369	2.252	3–4 mm	82	224.6–232.7, 240.6–248.7

Notes. ^(a) Frequency ranges covered in USB and LSB.

Appendix B: Detected lines and their characteristics

Table B.1. Line intensities from IRAM observations

Date (yyyy/mm/dd.dd)	Molecule	Transition	Frequency (MHz)	Offset ^a ($''$)	$\int T_{mB} dv$ (K km s ⁻¹)	Doppler shift (km s ⁻¹)
2025/11/02.36	HCN	3 – 2	265886.434	3.0	0.381 ± 0.020^b	-0.04 ± 0.03
2025/11/03.15				5.7	0.333 ± 0.073^b	$+0.17 \pm 0.13$
2025/11/02.36	HNC	3 – 2	271981.142	3.0	0.047 ± 0.018	-0.32 ± 0.26
2025/11/02.57	CH ₃ CN	8 – 7 ^c	147149–147175	2.9	0.039 ± 0.013	
		9 – 8 ^c	165540–165570	2.9	0.052 ± 0.021	
2025/11/02.57	CH ₃ OH	$J_1 - J_0 E (J=1-5)$	165050–165369	2.9	$0.058 - 0.107 (\pm 0.009)$	-0.07 ± 0.01
2025/11/02.57	CH ₃ OH	$J_1 - J_0 E (J=6-9)$	165678–167931	2.9	$0.078 - 0.041 (\pm 0.011)$	-0.05 ± 0.05
2025/11/02.57	CH ₃ OH	$3_2 - 2_1 E$	170060.581	2.9	0.114 ± 0.013	-0.01 ± 0.06
2025/11/02.40	CH ₃ OH	$8_{-1} - 7_0 E$	229758.811	2.9	0.109 ± 0.012	-0.01 ± 0.06
2025/11/02.36	CH ₃ OH	$5_K - 4_K A^+, A^-, E^d$	241700–243916	2.9	$0.101 - 0.337 (\pm 0.013)$	-0.03 ± 0.01
2025/11/02.33	CH ₃ OH	$J_3 - J_2 A^+, A^- e$	251164–252253	2.9	$0.027 - 0.068 (\pm 0.012)$	
2025/11/02.33	CH ₃ OH	$11_0 - 10_1 A^+$	250507.016	2.9	0.068 ± 0.012	-0.04 ± 0.09
2025/11/02.36	CH ₃ OH	$5_2 - 4_1 E$	266838.123	2.9	0.199 ± 0.012	-0.04 ± 0.03
2025/11/03.15				5.7	0.176 ± 0.051	$+0.05 \pm 0.17$
2025/11/02.36	CH ₃ OH	$9_0 - 8_1 E$	267403.394	2.9	0.056 ± 0.016	-0.20 ± 0.17
2025/11/02.40	CO	2 – 1	230538.000	2.9	0.069 ± 0.011	-0.11 ± 0.08
2025/11/02.36	CS	5 – 4	244935.557	2.9	0.040 ± 0.012	-0.36 ± 0.20
2025/11/02.57	CS	3 – 2	146969.049	2.9	0.014 ± 0.006	$+0.04 \pm 0.20$
2025/11/02.40	H ₂ CO	$3_{12} - 2_{11}$	225697.772	2.9	0.087 ± 0.010	$+0.07 \pm 0.06$
2025/11/02.57	H ₂ CO	$2_{11} - 1_{10}$	150498.334	2.9	0.032 ± 0.007	$+0.11 \pm 0.11$
2025/11/02.36	HCO ⁺	3 – 2	267557.625	3.0	0.099 ± 0.025	$+0.64 \pm 0.38$
2025/11/02.57	H ₂ S	$1_{10} - 1_{01}$	168762.762	2.9	0.015 ± 0.011	
2025/11/02.33	SO	6, 5 – 5, 4	251825.770	3.0	< 0.036	
2025/11/02.4	(CH ₂ OH) ₂	$J = 20 - 29^f$	225689–270647	3.0	0.0039 ± 0.0014	-0.11 ± 0.20
2025/11/02.4	CH ₃ CHO	$J = 12 - 14^g$	226552–271069	3.0	0.0042 ± 0.0011	$+0.19 \pm 0.13$
2025/11/02.36	HNCO	$J = 11 - 12, K = 0, 1^c$	240876–264694	2.9	0.051 ± 0.023	-0.11 ± 0.26
2025/11/02.36	HDO	$2_{11} - 2_{12}$	241561.641	2.9	< 0.033	

Notes. ^(a) Pointing offset. ^(b) Taking into account all hyperfine components. ^(c) Sum of four lines. ^(d) 7 lines. ^(e) 16 lines. ^(f) 45 lines. ^(g) 60 lines.

Appendix C: Synthetic line profiles

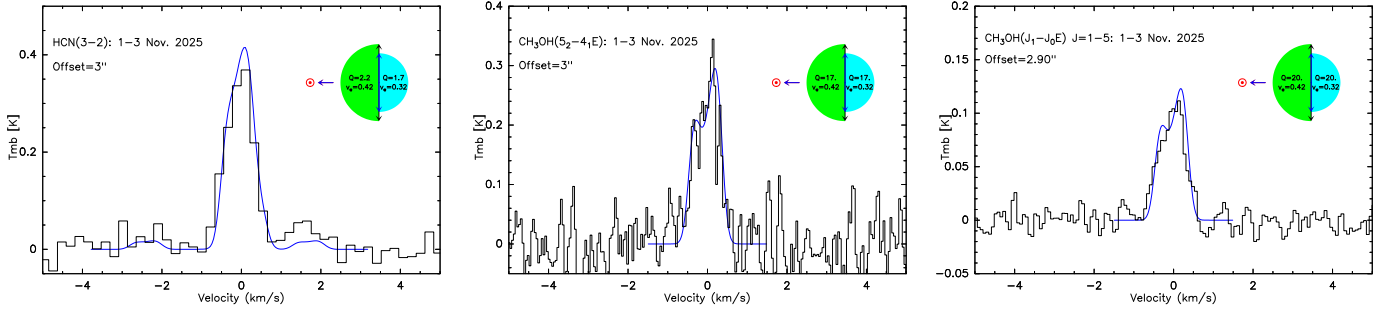


Fig. C.1. Synthetic line profiles (blue) superimposed on observed line profiles (black) of HCN (left panel) and CH₃OH (middle and right panels). The inserts in the upper right of the plots show the assumed geometry of the outgassing with respect to the Sun direction, with the production rates given in units of 10²⁵ molec. s⁻¹ for HCN, and in units of 10²⁶ molec. s⁻¹ for CH₃OH. The assumed expansion velocities (km s⁻¹) in the sunward and antisunward hemispheres are also indicated. Thermal broadening assuming a kinetic temperature of 60 K (see main text) is considered.

Appendix D: CH₃OH rotation diagram

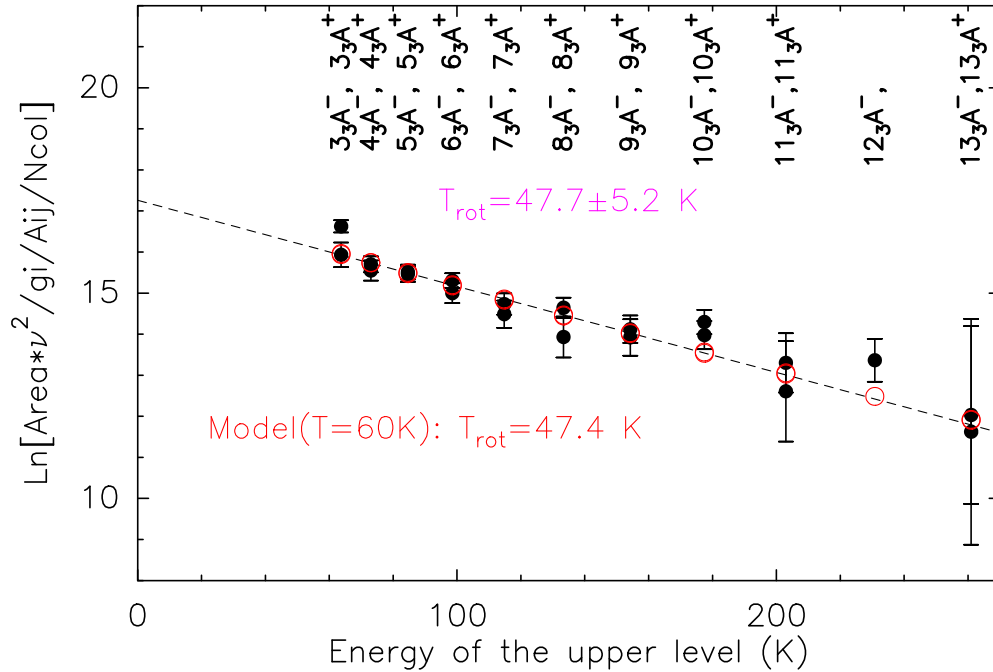


Fig. D.1. Rotation diagram of CH₃OH, using $J_3 - J_{3A^+}$ and A^- lines near 252 GHz. Red circles show the output of the excitation model assuming a kinetic temperature of 60 K. It reproduces nicely the measured rotational temperature of 47.7±5.2 K. The labels of the upper levels of the transitions are given at the top of the plot.

Appendix E: Histograms of abundances relative to water

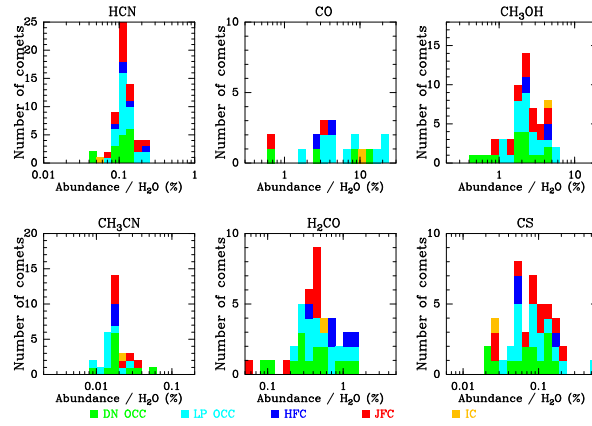


Fig. E.1. Same as Fig. 3 for abundances relative to H₂O.

Appendix F: Histogram of the sulfur-carbon ratio

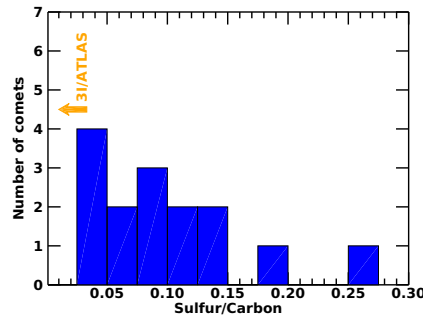


Fig. F.1. Histograms of the sulfur to carbon abundance ratio measured from observations at radio wavelengths. Considered sulfur and carbon-bearing molecules are CO, CH₃OH, H₂CO, H₂S and CS. Only comets for which these five species are detected are considered. Note however that H₂CO has a minor abundance, compared to CH₃OH and CO, and CS is much less abundant than H₂S in most comets. The upper limit for 3I/ATLAS is shown by an orange arrow. See Biver et al. (2024b) and references therein.

Appendix G: Expansion velocity

Table G.1. Line widths

Line	Average half-widths ^a (VHM) (m s ⁻¹)	
	blue wing	red wing
HCN 265.9 GHz	463±33	387±33
CH ₃ OH 266.8 GHz	426±30	359±45
CH ₃ OH 165 GHz	672±59	305±35
CO 230.5 GHz	616±153	316±50

Notes. ^(a) Values derived by fitting a half-Gaussian in the blue and red wings of the profiles.

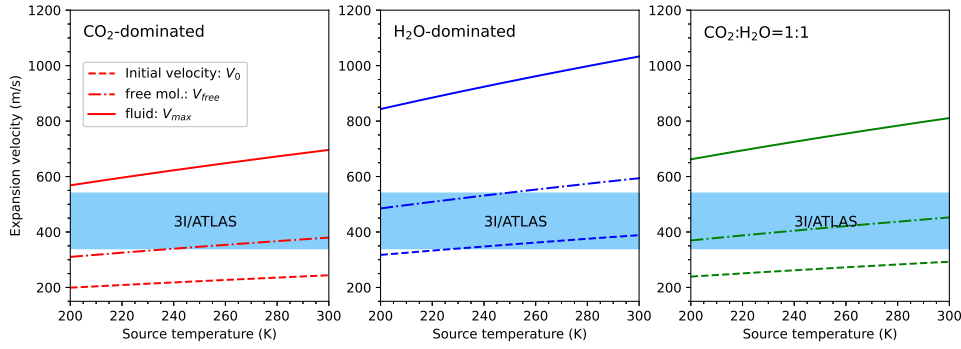


Fig. G.1. Initial velocity V_0 , maximal velocity V_{max} , and terminal velocity for free-molecular expansion V_{free} as a function of source temperature (T_N). From left to right, CO₂-dominated, H₂O-dominated, and H₂O-CO₂ gas mixture. The range for 3I/ATLAS (in blue) delineates the expansion velocities measured in the blue and red wings of the line profiles (i.e. VHM values, see main text).

Figure G.1 compares the expansion velocity measured in 3I/ATLAS from molecular line profiles to the gas velocity expected in three limiting cases.

We define the gas initial velocity V_0 as the sonic velocity:

$$V_0 = \sqrt{\frac{\gamma k_B T_0}{m_g}}, \quad (\text{G.1})$$

where γ is the specific heat ratio (equal to 1.33 for H₂O, 1.28 for CO₂), T_0 is the gas temperature at sonic line (equal to ~ 0.82 times the temperature of the source releasing gases, referred to as T_N), and m_g is the molecular mass of expanding gases.

The maximal terminal gas velocity, which is reached in case of complete transfer of initial thermal energy into kinetic energy of the flow, is (e.g., Zakharov et al. 2021)

$$V_{\text{max}} = V_0 \sqrt{\frac{\gamma + 1}{\gamma - 1}}. \quad (\text{G.2})$$

This terminal velocity is almost reached only for active comets, with large regions of fluid flow. This is shown by Zakharov et al. (2023) who studied flow conditions from fluid to free molecular using a kinetic approach (direct simulation Monte Carlo method, so-called DSMC).

The maximum velocity reached in free-molecular expansion (i.e., in case of very weak gas production) is

$$V_{\text{free}} = \sqrt{\frac{8}{\pi} \frac{k_B T_N}{m_g}}. \quad (\text{G.3})$$

In Figure G.1, we consider comas dominated by H₂O, by CO₂, and composed of both gases in equal quantities ($m_g = 31$ g/mole, $\gamma = 1.3$). Calculations are made for source temperatures from 200 K (temperature of exposed ice at typically $r_h=1$ au) to 300 K (appropriate if one considers gas diffusion through a layer of warm non-ice material). For an H₂O-dominated coma, V_{max} is close to the typical expansion velocity of 0.7–0.8 km s⁻¹ measured in moderately active comets ($Q_{\text{H}_2\text{O}} = 10^{28}$ – 10^{29} molec. s⁻¹). This outgassing regime is excluded for 3I/ATLAS. The measured velocities are consistent with an H₂O-dominated outgassing only in the case of very weak water production (with free molecular expansion)(Fig. G.1, central panel). A CO₂-rich coma is favored as in this case the expected gas expansion velocity is significantly lower (Fig. G.1, left and right panels).