

LETTER TO THE EDITOR

Detection of SiC₆ in IRC +10216 through Q-band lines

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ABSTRACT

Ongoing Q-band (31–50 GHz) integrations on IRC +10216 with the Yebes 40m telescope have now reached the level of sub-mK noise, with hundreds of new lines arising above the 3 σ detection limit. Among them is a new, harmonically related series with integer quantum numbers ranging from $J_{\text{up}} = 26$ to 40, which has allowed it to be assigned to SiC₆. We clearly detected a total of nine isolated (or otherwise not critically blended) spectral features, with peak intensities around or slightly below 1 mK in the T_{A}^* scale, thanks to the extremely low noise in the spectra. Five additional lines were seen to be partially blended with lines from other species, but could still be relatively well modeled due to the uniformity of the characteristic U-shaped line profiles emerge from the envelope of IRC+10216. As a result, we built a rotation diagram for SiC₆, providing a column density of $\sim 8 \times 10^{11} \text{ cm}^{-2}$ and a rotational temperature around 19 K. A similar analysis was done for SiC₄, ²⁹SiC₄, and ³⁰SiC₄. Together, these insights offer the first view at the chemistry of linear silicon carbides in IRC+10216.

Key words. astrochemistry – line: identification – molecular data – stars: carbon – stars: individual: IRC +10216

1. Introduction

The search for new molecules in astronomical sources requires extremely deep integrations as the list of already known species continues growing. For this task, it is important to count on constantly updated line catalogs that help us to spot unidentified features among an increasing forest of lines from identified species. These new findings, however, are important not just for extending the list but for the many questions opened up by their presence in the astronomical sources, such as formation paths, physical conditions, and growth of chemical complexity. Long molecules, large polycyclic aromatic hydrocarbons (PAHs), cations and anions, rare isomers, and high vibrational states, are all being found as the sensitivity of the surveys reaches sub-mK levels in TMC-1 and IRC+10216 within the ERC synergy project Nanocosmos¹.

Many new species have been found directly in TMC-1 and IRC+10216 without any previous laboratory information on their frequencies. This is due to the difficulty to obtain their spectra in the laboratory. It is the case for cations, radicals, and metal-bearing species involving Mg, Ca, Al, and Na (Cernicharo et al. 2019; Pardo et al. 2021; Cabezas et al. 2023; Cernicharo et al. 2023a). Nevertheless, in some cases, the main driver for molecular discoveries of metal-bearing species has been new laboratory work (Changala et al. 2022; Gupta et al. 2024). For IRC+10216, additional Q-band observations have been added until September 15, 2024, improving the sensitiv-

ity by a factor of more than 2. This widens the opportunities for identifying new molecular species and new isotopic substitutions, isomers, and vibrationally excited states of already known species, helping to complete the chemical picture of this source (see, e.g., Cernicharo et al. 2023b,a; Cabezas et al. 2023).

In this Letter, we report the detection of SiC₆ through 14 rotational lines, of which 9 are unblended or weakly blended at the observed sub-mK peak intensities, but well above the 3 σ limit. Rotation diagrams, from which a column density and a rotational temperature can be determined, are established for SiC₆, and for SiC₄, ²⁹SiC₄, and ³⁰SiC₄ that also appear in our data. All this information has been analyzed as in our previous studies (Pardo et al. 2022; Cernicharo et al. 2023b,a; Cabezas et al. 2023).

Section 2 gives an update on the status of this large observational program and describes the way we accumulated the data through such a long period of time. Section 3 is aimed at describing the procedure for fitting the data using a double-horned line-shaped function, typical of expanding envelopes. This section also deals with the fact that lines can sometimes be heavily blended. The spectra of SiC₆, SiC₄, ²⁹SiC₄, and ³⁰SiC₄ are presented in Section 4 and their derived physical parameters are discussed in Section 5 on the ground of chemical models, followed by the final conclusions of this work.

2. Observations and data accumulation

The data presented in this Letter come from an ongoing 31–50 GHz spectral survey of IRC+10216 carried out with the 40 meter antenna of the Yebes Observatory (IGN, Spain; hereafter Yebes-40m), after several improvements in its equipment

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¹ <https://nanocosmos.iff.csic.es/>

funded by Nanocosmos. The data was obtained after accumulation of 1360 hours of on-source telescope time in 480 observing sessions from April 14, 2019 to September 15, 2024. The experimental setup has been described in detail by [Tercero et al. \(2021\)](#). The antenna provides a main beam efficiency from 0.6 at 31 GHz to 0.43 at 50 GHz and a beam size in the range 36–56'' for those frequencies (45'' at 39.2 GHz). The Nanocosmos *Q*-band receiver uses two HEMT cold amplifiers covering the 31.0–50.3 GHz frequency range with horizontal and vertical polarizations. The backends are 16×2.5 GHz fast Fourier transform spectrometers (FFTS), with a primary spectral resolution of 38.1 kHz. The observing mode was position switching with an off position at 300'' away in azimuth from the reference $\alpha_{J2000} = 09^{\text{h}} 47^{\text{m}} 57.36^{\text{s}}$ and $\delta_{J2000} = +13^{\circ} 16' 44.4''$ position used for IRC+10216. Pointing corrections were obtained by observing the SiO masers of R Leo, and errors were always within 2–3''. The intensity scale of the final calibrated data is antenna temperature (T_A^*) corrected for atmospheric absorption using the ATM package ([Cernicharo 1985](#); [Pardo et al. 2001, 2025](#)). The calibration uncertainties are estimated to be within 10%.

The receiver temperature improvements during this long project, along with the changing atmospheric conditions and different integration times in each observing session means that the addition of the data has to be done in an especially careful way to get the most accurately calibrated final spectrum. Nevertheless, we would like to note that the presented data correspond to time averaged intensities as the lines in this source show a dependency of their intensity with stellar phase ([Menten et al. 2012](#); [Cernicharo et al. 2014](#); [Pardo et al. 2018](#)). We used reference strong lines to discard whole (or part of) observing sessions with technical or pointing problems. Then, we fitted the baselines to each FFTS section for each session after accumulating all valid scans in it. We avoid observations at very low elevations, less than ~ 23 degrees, which would result in high noise. Then, we accumulated all sessions for each FFTS section, using the rms from the previous step as weighting factor. We fitted a new baseline to this result in smaller frequency ranges (200 MHz) than the FFTS sections themselves (2.5 GHz) to do a much finer job without degrading the original spectral resolution (38 kHz) at first. However, in the final product we applied a six-channel smoothing. This resulted in a final frequency resolution of $\Delta\nu = 229$ kHz and a frequency dependent σ level, which is indicated in the upper panel of [Fig. 1](#) and in each panel of [Fig. 2](#).

3. Data analysis

For the final accumulated spectral results, we applied a global fit strategy, described in this section. First, we use as input the central line frequencies from the MADEX catalog ([Cernicharo 2012](#)), corresponding to molecular species that we already know to be present or expected in the data. To these we added all central frequencies associated to unidentified spectral features in the data. We developed a simple extended routine for multi-line calculations using the same basic shape function from the SHELL method of the GILDAS² software package, expressed as

$$\phi(\nu) = \frac{A}{\Delta\nu} \frac{1 + 4H[(\nu - \nu_0)/\Delta\nu]^2}{1 + H/3}, \quad (1)$$

with A being the integrated intensity, ν_0 the central frequency, $\Delta\nu$ the full width at zero level (obtained from twice the termi-

² <http://www.iram.fr/IRAMFR/GILDAS>

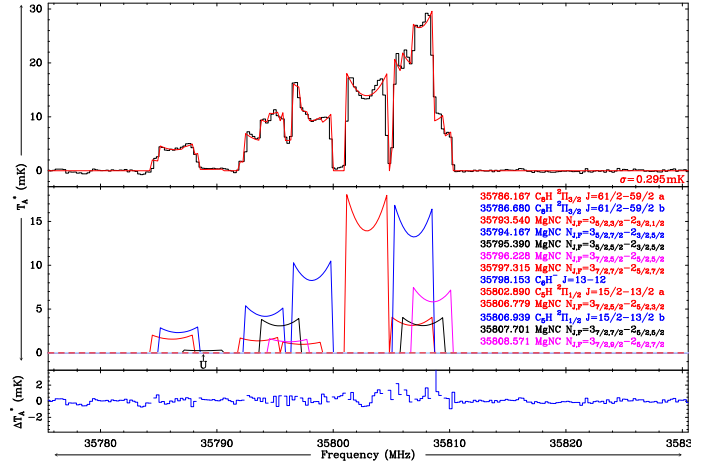


Fig. 1. Extreme example of global fit of IRC+10216 *Q*-band data to a series of individual lines (for $H = 0.3$) with integrated areas that can be obtained in a quite confident way, described in Section 3, despite the blending. The residuals from the global fit are plotted in the bottom panel. All spectral features present in the data, except a very small one, can be identified at present.

nal expansion velocity of the envelope, see below), and H the horn to center intensity ratio. The profile varies from a parabola ($H = -1$) for optically thick lines, a flat topped shape ($H = 0$), which corresponds to spatially unresolved optically thin lines, and a double-peaked profile ($H > 0$), which is produced by spatially resolved optically thin lines.

In our analysis, the *LSR* velocity of IRC+10216 has been fixed to -26.5 km s⁻¹ and the terminal expansion velocity, v_{end} , of the envelope has been fixed to 14.5 km s⁻¹ ([Cernicharo et al. 2000](#)). The H and A parameters depend for each line on the spatial distribution of the emission and on the line opacity. For most lines, we know that the emission appears in a shell with inner and external radii between 10 and 20'' and a thickness of 5–10'' ([Agúndez et al. 2017](#)). With very few exceptions, we found that $H = 0.3$ is a good value for our global fit calculations. This choice does not substantially affect the derived areas of the lines used for the rotation diagrams presented in this work.

The overall fit to the line profiles is built by creating a first guess table with values of ν_0 , A , and H for each expected (isolated or blended) spectral line. A synthetic spectrum is created from a line-by-line fit of A , with ν_0 and H kept fixed, starting from the lowest frequency upwards. An iterative process is achieved to adjust the observed global line profile to the contribution of the individual lines. This iterative process runs nearly automatically, except for a few cases where the blending is so heavy that no realistic parameters can be obtained for some lines. In these cases, intervention into the analysis is necessary. After just a handful of steps, it is possible to have a first guess file from which the fitting of A for all lines produces variations of less than 10%. The adopted strategy provides excellent results: although it is common in our data to have several lines blended, in almost all cases it becomes possible to separate each one due to the similarity of the profiles and the use of information of other lines in each species' rotational ladder. [Figure 1](#) gives an example of a very satisfactory line decomposition in IRC+10216 *Q*-band data. This has also been applied in previous IRC+10216 publications (see, e.g., [Changala et al. 2022](#)).

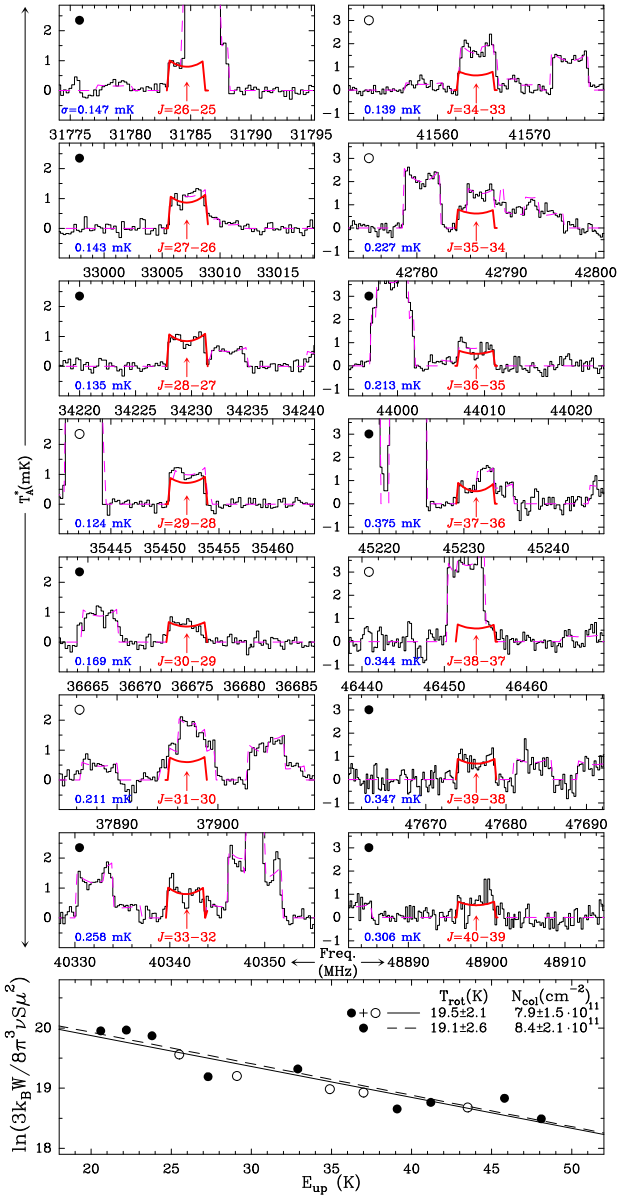


Fig. 2. SiC₆ lines observed in this work, with the fit for them in red. Black dots correspond to lines with no critical blending from other species while open dots identify lines heavily blended. The dashed magenta curves represent the overall fit for all lines according to the parameters given in Table 1. The bold red curves correspond to the contribution of the SiC₆ lines alone. The rms or σ values given for each spectrum in dark blue are calculated from the channels, with no signal in the fitted spectrum. The best-fit rotational diagrams using only the not heavily blended SiC₆ rotational transitions are shown separately as solid and dashed lines, respectively, in the bottom panel.

4. Results

Laboratory spectroscopy for SiC₆ is available from McCarthy et al. (2000). The molecule was implemented in MADEx adopting a dipole moment of 8.25 D as derived from ab initio calculations by the same authors. Although laboratory data only reach up to 20.8 GHz ($J = 17$), predictions for the rotational lines in the Q-band range ($J = 26$ – 25 up to $J = 40$ – 39) have uncertainties between 2 and 10 kHz. Hence, the estimated frequencies are well suited for searching this molecule in our data.

Table 1. SiC₆ lines and the spectral features near them, as modeled in Figure 2 with H from Equation (1) fixed to 0.3 and $v_{\text{end}} = 14.5 \text{ km s}^{-1}$.

Molecule	μ (Debye)	Transition ID	ν_{ul} (MHz)	E_{up} (K)	Line Str.(S)	W (mK- km s^{-1})
U31779			31778.709			4.573
SiC ₆	8.220	J = 26 → 25	31784.664	20.60	26.00	24.500
C ₁₀ H ⁻	15.200	J = 53 → 52	31785.863	41.20	53.00	63.157
C ₄ H $v_7=1$	2.100	see [1]	31786.650	247.10	2.39	41.758
SiC ₆	8.220	J = 27 → 26	33007.120	22.20	27.00	27.754
U33008			33008.406			5.471
U33010			33010.335			3.649
SiC ₆	8.220	J = 28 → 27	34229.572	23.80	28.00	28.047
H ¹³ CCC ¹⁵ N	3.732	J = 4 → 3	34233.231	4.10	4.00	8.319
CH ₃ C ₅ N	5.400	J = 22 ₀ → 21 ₀	34233.376	18.90	22.00	7.577
U35441			35440.971			9.568
HC ₉ N	5.200	J = 61 → 60	35442.454	52.70	61.00	143.684
SiC ₆	8.220	J = 29 → 28	35452.020	25.50	29.00	22.721
U35453			35452.518			8.466
l-HC ₄ N	4.330	$N_J = 8_7 \rightarrow 7_6$	36666.112	8.40	6.91	26.789
SiC ₆	8.220	J = 30 → 29	36674.464	27.30	30.00	17.271
HC ₉ N $v_{10}=3$	5.200	$J_L = 65_{-1} \rightarrow 64_{-1}$	37888.364	276.60	65.00	14.666
C ¹³ CCCH	2.100	see [2]	37896.002	4.50	2.29	9.480
SiC ₆	8.220	J = 31 → 30	37896.904	29.10	31.00	19.114
HC ¹³ CC ₅ N	4.820	J = 34 → 33	37897.938	31.80	34.00	28.044
CH ₂ CHCNa	3.821	$J_{K,k} = 4_{0,4} \rightarrow 3_{0,3}$	37904.851	4.50	4.00	29.085
U37907			37906.800			11.841
U40332			40332.097			37.223
U40335			40335.000			10.764
SiC ₆	8.220	J = 33 → 32	40341.773	32.90	33.00	25.301
MgC ₅ N	7.300	$N_J = 35_{69/2} \rightarrow 34_{67/2}$	40348.015	34.90	34.50	53.699
MgC ₅ N	7.300	$N_J = 35_{71/2} \rightarrow 34_{69/2}$	40349.822	34.90	35.50	43.479
U41559			41558.500			8.209
SiC ₆	8.220	J = 34 → 33	41564.200	34.90	34.00	20.054
HCC ¹³ CC ₄ N	4.820	J = 37 → 36	41564.473	37.90	37.00	30.450
U41567			41567.000			7.111
HC ₄ C ¹³ CCN	4.820	J = 37 → 36	41574.391	37.90	37.00	41.439
HC ₅ N $v_{11}=2$	4.330	J = 16 → 15	42780.600	325.70	48.00	63.374
SiC ₆	8.220	J = 35 → 34	42786.623	37.00	35.00	19.810
U42788			42787.655			26.468
H ¹³ CCC ¹⁵ N	3.732	J = 5 → 4	42791.453	6.20	5.00	20.019
U42794			42794.200			18.733
C ₈ H ² $\Pi_{3/2}$	6.500	$J = \frac{75}{2} \rightarrow \frac{73}{2}$ a 1	43999.034	40.50	74.90	58.626
C ₈ H ² $\Pi_{3/2}$	6.500	$J = \frac{75}{2} \rightarrow \frac{73}{2}$ b	43999.793	40.50	74.90	57.243
HCC ¹³ C ¹⁵ N	3.732	J = 5 → 4	44007.159	6.30	5.00	6.288
SiC ₆	8.220	J = 36 → 35	44009.042	39.10	36.00	16.118
HCC ¹³ CCCN	4.330	J = 17 → 16	45217.843	19.50	17.00	252.488
HCC ¹³ CCN	4.330	J = 17 → 16	45223.345	19.50	17.00	244.703
SiC ₆	8.220	J = 37 → 36	45231.455	41.20	37.00	19.149
H ¹³ CC ₆ N	4.820	J = 41 → 40	45233.690	45.60	41.00	21.366
C ₇ H ² $\Pi_{3/2}$	5.945	$J = \frac{83}{2} \rightarrow \frac{81}{2}$	46452.599	68.20	106.00	86.785
SiC ₆	8.220	J = 38 → 37	46453.864	43.50	38.00	18.668
HC ₇ N $v_{15}=3$	4.820	$J_L = 40_{-3} \rightarrow 39_{-3}$	46466.395	315.50	40.80	3.494
HC ₇ N $v_{15}=3$	4.820	$J_L = 40_3 \rightarrow 39_3$	46466.591	315.50	40.80	3.494
SiC ₆	8.220	J = 39 → 38	47676.268	45.80	39.00	23.017
S3		$J = \frac{81}{2} \rightarrow \frac{79}{2}$	47683.300			24.534
U47691			47691.450			23.674
S4		$J = \frac{83}{2} \rightarrow \frac{81}{2}$	48883.000			14.165
SiC ₆	8.220	J = 40 → 39	48898.666	48.10	40.00	17.244

Notes. The area of the line W in mK-kms⁻¹ is derived from the only free parameter for each line in the fit which is the integrated intensity A in Equation (1). S3 and S4 are the unknown carriers of two series of lines identified in the data to which even the quantum numbers can be assigned. Other unidentified lines are labeled as UXXXXXX with the rounded value in MHz from their central frequencies. [1] (N,P,J,F) = (3,1,7/2,3) → (2,-1,5/2,2) [2] (N,J)_{f,F} = (4,7/2)_{3,5/2} → (3,5/2)_{2,3/2}.

Figure 2 shows the accumulated results of our Yebes-40m IRC+10216 Q-band integration as of September 15, 2024, around the expected positions of SiC₆ rotational transitions, with the global fit superimposed. From this figure, we get the line parameters listed in Table 1 that have been used for deriving the rotational temperature and column density of SiC₆ from a rotational diagram linear regression (also plotted on Figure 2). One of the lines ($J = 32$ → 31) has been eliminated from the analysis because it is strongly affected by RFI signals around 39.0–39.25

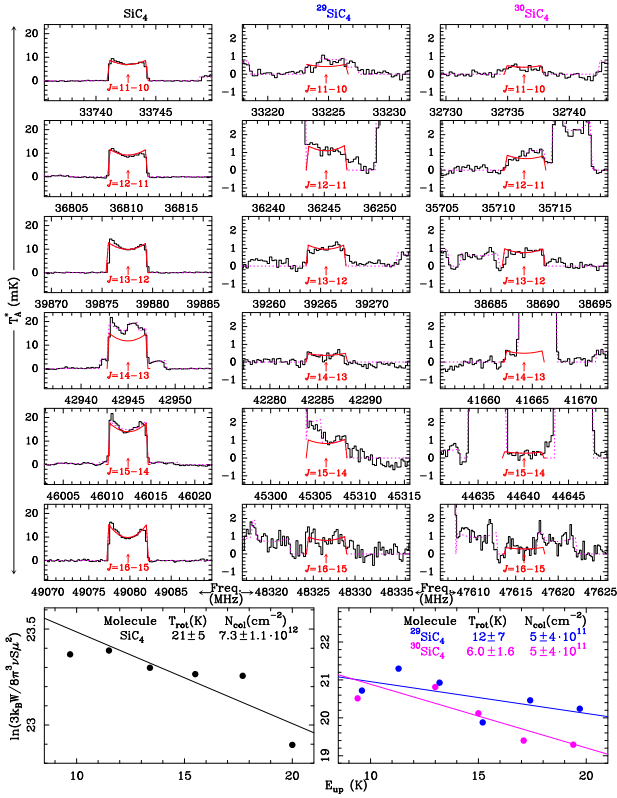


Fig. 3. SiC_4 , $^{29}\text{SiC}_4$, and $^{30}\text{SiC}_4$ lines found in this work, with their fitted line profile shown in red. The dotted magenta curves represent the overall fit that takes into account all lines in the data.

GHz. The $J = 30 \rightarrow 29$ line, kept in the figure and in the analysis, may also suffer from the same problem but to a much lower extent. Otherwise, the rotation diagram appears to be very consistent and the results of the linear regression do not show any significant changes when the most heavily blended SiC_6 lines are included and when they are not.

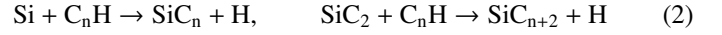
A similar analysis was performed on SiC_4 , with the results shown on Figure 3. From the intensities found for the $J = 11 \rightarrow 10$ to $J = 16 \rightarrow 15$ transitions of this molecule and the $^{28}\text{Si}/^{29}\text{Si}$ and $^{28}\text{Si}/^{30}\text{Si}$ isotopic ratios in this source (15.4 and 20.3, respectively; Cernicharo et al. 2000), along with the current sensitivity of the survey, we expect to detect $^{29}\text{SiC}_4$ (comfortably) as well as $^{30}\text{SiC}_4$ (more marginally). Both species are also included Figure 3. The four ^{13}C substitutions were searched for, but they do not appear in the data above the current noise level, which was anticipated. Some of the lines in Figure 3 are blended with other species but, due to space limitation, we omit these details.

We could think about adding SiC_2 to this analysis. However, this species is a ring with a C_{2v} symmetry, it is light and has only three detected transitions in our Q -band survey with line strengths spanning three orders of magnitude, making the simple rotation diagram analysis difficult and uncertain with only those lines. We can say, however, that the column densities of SiC_4 and SiC_6 are more than two orders of magnitude below those derived for SiC_2 and Si_2C by Cernicharo et al. (2015, 2018) and Massalkhi et al. (2018) using an extended range of frequencies.

5. Discussion and conclusions

The chemistry of SiC_6 in IRC +10216 has not been studied to our knowledge, but that of SiC_4 has been discussed to some extent. Ohishi et al. (1989) proposed that SiC_4 could be formed in

neutral-neutral reactions of SiC_2 with C_2H_2 and C_2H or through ion schemes starting with the reaction of Si^+ and C_2H_2 . In the chemical model of MacKay & Charnley (1999), SiC_4 is mainly formed in the external circumstellar layers through the reaction between SiC_2 and C_2H . Schemes based on neutral-neutral reactions or ion-neutral reactions could be behind the synthesis of SiC_4 and SiC_6 in IRC +10216. Within the first category, reactions of the type



arise as very promising routes to the synthesis of SiC_4 , SiC_6 , and even longer Si-bearing carbon chains, although detailed theoretical calculations are needed to shed light on whether these reactions are barrierless and proceed through H atom elimination. In the second category, experimental data support the reactivity of the cation Si^+ with unsaturated carbon chains. However, the products formed, namely, $\text{SiC}_2\text{H}^+ + \text{H}$ in the case of acetylene and $\text{C}_4\text{H}^+ + \text{SiH}$ in the case of diacetylene (Wlodek et al. 1991), indicate that the synthesis is probably not straightforward and may involve several steps of carbon chain growth. This could include reactions between cations SiC_nH^+ and unsaturated hydrocarbons such as C_2H_2 and C_2H . A mechanism similar to that proposed by Petrie (1996) for the formation of metal-bearing carbon chains would be based on the radiative association of Si^+ with a neutral carbon chain, followed by the dissociative recombination of the ionic complex formed. This approach could also contribute to the formation of silicon-bearing carbon chains in IRC +10216, similarly to the scenario reported here.

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