Progress on the calibration of surface brightness–color relations for early- and late-type stars[★]

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ABSTRACT

Context. Surface brightness–color relations (SBCRs) are widely used for estimating angular diameters and deriving stellar properties. They are critical to derive extragalactic distances of early-type and late-type eclipsing binaries or, potentially, for extracting planetary parameters of late-type stars hosting planets. Various SBCRs have been implemented so far, but strong discrepancies in terms of precision and accuracy still exist in the literature.

Aims. We aim to develop a precise SBCR for early-type B and A stars using selection criteria, based on stellar characteristics, and combined with homogeneous interferometric angular diameter measurements. We also improve SBCRs for late-type stars, in particular in the *Gaia* photometric band.

Methods. We observed 18 early-type stars with the VEGA interferometric instrument, installed on the CHARA array. We then applied additional criteria on the photometric measurements, together with stellar characteristics diagnostics in order to build the SBCRs.

Results. We calibrated a SBCR for subgiant and dwarf early-type stars. The RMS of the relation is $\sigma_{F_{V_0}} = 0.0051$ mag, leading to an average precision of 2.3% on the estimation of angular diameters, with 3.1% for $V-K < -0.2$ mag and 1.8% for $V-K > -0.2$ mag. We found that the conversion between Johnson- K and 2MASS- K_s photometries is a key issue for early-type stars. Following this result, we have revisited our previous SBCRs for late-type stars by calibrating them with either converted Johnson- K or 2MASS- K_s photometries. We also improve the calibration of these SBCRs based on the *Gaia* photometry. The expected precision on the angular diameter using our SBCRs for late-type stars ranges from 1.0 to 2.7%.

Conclusions. By reaching a precision of 2.3% on the estimation of angular diameters for early-type stars, significant progress has been made to determine extragalactic distances, such as M31 and M33 galaxies, using early-type eclipsing binaries.

Key words. stars: fundamental parameters – distance scale – techniques: interferometric

1. Introduction

Determining the expansion of the Universe, that is the Hubble constant (H_0) to better than 2%, is required in order to understand the nature of dark energy. However, the two most accurate methods for that, the cosmic microwave background (Planck Collaboration VI 2020; Riess et al. 2019) and the distance ladder, are inconsistent today, which is referred to as the “tension” (Wong et al. 2020). One of the keys to resolve this tension is based on the calibration of the Leavitt period–luminosity (PL) law of Cepheids (Leavitt & Pickering 1912). Cepheids are indeed the backbone of the extragalactic distance ladder because their pulsation periods, which are easily determined observationally, directly correlate with their luminosities. Another method consists in using eclipsing binaries to constrain extragalactic distances.

Recently, a new estimate of the distance to the Large Magellanic Cloud (LMC), based on 20 late-type eclipsing binaries,

has been obtained by the Araucaria team¹ (Pietrzyński et al. 2019). Their precision of 1% is mostly due to the precision of the surface brightness–color relation (SBCR), calibrated on 41 nearby red clump giant stars using infrared interferometry (Gallenne et al. 2018). The same was done to derive the distance of the Small Magellanic Cloud (SMC) with a precision of better than 2% (Graczyk et al. 2020). Deriving the distance from eclipsing binaries is simple: the radius of both components is estimated from the combination of photometry and spectroscopy, and angular diameters are estimated from the magnitude and color of stars through a SBCR. The combination of radii and angular diameters provides the distance. The influence of interstellar attenuation in neighboring galaxies has been studied using several techniques so far (Bonanos et al. 2006; Pietrzyński et al. 2019; Graczyk et al. 2020), and it is still under investigation.

The situation concerning the early-type eclipsing binaries is more complex because the calibration of the SBCR requires high angular resolution measurements, and early-type stars are

[★] Based on CHARA/VEGA observations.

¹ <https://aracaria.camk.edu.pl/>

particularly active (Martins et al. 2015; Gordon et al. 2018). Challouf et al. (2014) improved, by a factor of 2, the precision on SBCR of early-type stars from about 15 to 7%, corresponding to the most accurate SBCR developed so far for early-type stars. However, this precision is still not sufficient to derive the distance of extragalactic early-type eclipsing binaries with a precision of a few percent. It is also worthwhile to mention that the distances to M31 and M33 are currently based on models of early-type eclipsing binaries, but not SBCRs (Bonanos et al. 2006; Vilardell et al. 2010). Another interesting approach is to analyze O- and B-type detached eclipsing binaries in the LMC, for which the distance is known, and to derive the surface brightness (Taormina et al. 2019). In this paper, we aim to calibrate the SBCR for early-type stars precisely, following the strategy of Paper I.

Moreover, with the work done on the infrared photometry of early-type stars, it is now possible to improve the SBCRs for late-type stars shown in (Salsi et al. 2020, Paper I hereafter). It turns out that some refinements are necessary concerning the infrared photometric systems that are used. Finally we have also reconsidered the calculation of the extinction in the *Gaia* band for proposing a new improvement of the SBCRs in this work.

Sections 2–4 are devoted to the strategy and data selection, the VEGA observations, and the SBCR calibration of early-type stars, respectively. A subsequent discussion is provided in Sect. 5. The revision of SBCR for late-type stars is presented in Sect. 6, while some general conclusions are given in Sect. 7.

2. Strategy and data selection

2.1. Criteria on stellar characteristics

In Paper I, we have shown that any stellar activity or characteristics (multiplicity, binarity, variability, etc.) may impact the calibration of SBCRs. We therefore implemented a set of criteria to select a correct sample of early-type stars.

We restricted the calibration of our SBCR to V and K photometries, as it is the set of color which provides the lowest dispersion (Kervella et al. 2004). We started the selection from the SIMBAD Astronomical Database². We first selected early-type stars, thus O, B, and A stars with $V-K < 1$ mag. We then considered only subgiants and dwarfs. From Challouf et al. (2015), we know that the projected rotational velocity affects the surface brightness of the star. They show that a rotational velocity lower than 85% of the critical velocity has an impact of at most $\sigma_{F_V} = 0.003$ mag on the RMS of the SBCR. Alternatively, if one cannot access the critical velocity of the star, they demonstrate that considering stars with projected rotational velocity $V_{\text{rot}} \sin i$ lower than 100 km s^{-1} also results in a dispersion of 0.003 mag. Therefore, to reach a precision of $\sim 2\%$ on the angular diameter estimate, we made the choice to consider only stars with $V_{\text{rot}} \sin i$ lower than 75 km s^{-1} .

We excluded all known binary stars in our sample. The SEDs of all the stars were checked with the VO Sed Analyzer (VOSA) software³.

We have shown in Paper I that a variability above 0.1 mag could significantly affect the SBCR. Following this strategy, we searched for information about the variability of the stars in Samus' et al. (2017). We then rejected variable stars with a variability above 0.1 mag. We quantitatively study this point later in Sect. 5.1.

We finally searched for stars with expected angular diameters between 0.3 and 0.8 milli-second of arc (mas). This is optimal for the VEGA instrument (Mourard et al. 2009, 2011), installed at the CHARA array in Mount Wilson, USA (ten Brummelaar et al. 2005). We finally end up with a total of 18 stars to be observed in the northern hemisphere. No O-type stars were selected due to a roughly equal combination of variability and multiplicity criteria.

2.2. Photometric selection and interstellar attenuation

As demonstrated in Paper I, precise photometries are of course mandatory for the calibration of SBCR. The Kharchenko & Roeser (2009) catalog offers precise visible magnitudes and gathers measurements from several other catalogs (HIPPARCOS-Tycho catalogs, Carlsberg Meridian Catalog, and the Positions and Proper Motions catalog). All the visible magnitudes are given in the Johnson- V filter. With such a catalog, the precision on the V magnitude of our stars ranges from 0.002 mag to 0.008 mag.

Finding precise infrared K photometry is more complex. Indeed, we faced some issues with the 2MASS catalog (Cutri et al. 2003), where seven out of the 18 stars in our sample are affected by imprecise infrared measurements (i.e., an uncertainty higher than 0.1 mag). This problem has already been identified in Paper I and is due to saturation issues. We identified accurate measurements in Ducati (2002) for these seven stars.

The precision on the K/K_s photometry in our sample ranges from 0.1 to 2.7%. However, the drawback of this approach is that the selected photometric values are inhomogeneous in terms of the filter pass band. We could convert Johnson- K photometries into 2MASS using transformation equations, as we propose later in Sect. 6 for late-type stars, but this leads to other issues, in particular for early-type stars, that are discussed in Sect. 5.2.

For the reddening correction, we used the *Stilism*⁴ online tool (Lallement et al. 2014; Capitanio et al. 2017) to compute the color excess $E(B-V)$, considering early *Gaia* DR3 parallaxes (Gaia Collaboration 2021, 2018). The interest of this tool is the tridimensional maps of the local interstellar matter (ISM) it offers, based on measurements of starlight absorption by dust or gaseous species. The interstellar attenuation A_V in the visible band is defined as follows

$$A_V = A_V \times E(B-V), \quad (1)$$

where R_V is the total-to-selective extinction ratio in the visible band, for which we set $R_V = 3.1$, and we used $A_K = 0.114 \times A_V$ (Cardelli et al. 1989).

3. VEGA/CHARA interferometric measurements

We observed the sample of 18 early-type stars from 23 February 2019 to 16 December 2020. The calibrated oifits files are available on OidB⁵. The data were processed using the standard VEGA pipeline (Mourard et al. 2009) and the squared visibilities were calibrated using reference stars selected with the SearchCal tool⁶ (Bonneau et al. 2006). The list of the calibrators is included in Table 1. The uniform-disk angular diameter in the R band $\theta_{\text{UD}}[R]$ is taken from the JMMC Stellar Diameters

⁴ The online tool is available at <http://stilism.obspm.fr>

⁵ <http://oidb.jmmc.fr/index.html>

⁶ The tool is available at <https://www.jmmc.fr/english/tools/proposal-preparation/search-cal/>

² Available at <http://simbad.u-strasbg.fr/simbad/>

³ <http://svo2.cab.inta-csic.es/theory/vosa/>

Table 1. Reference stars used for VEGA observations listed with their spectral type, their V magnitude, and their uniform-disk angular diameter in the R band.

Target	Reference stars	Sp. Type	V [mag]	$\theta_{UD}[R]$ [mas]
HD 11415	HD 10221	A0V	5.59	$0.224_{\pm 0.013}$
	HD 12301	A0I	5.61	$0.403_{\pm 0.013}$
	HD 6210	F6V	5.83	$0.489_{\pm 0.039}$
HD 114330	HD 107070	A5IV/V	5.90	$0.289_{\pm 0.021}$
	HD 112846	A3III	5.79	$0.294_{\pm 0.009}$
	HD 116831	A8V	5.97	$0.279_{\pm 0.019}$
HD 145389	HD 140728	A0V	5.48	$0.232_{\pm 0.015}$
	HD 143584	F0IV	6.03	$0.327_{\pm 0.023}$
	HD 144206	B9III	4.71	$0.314_{\pm 0.019}$
HD 145570	HD 143459	A0V	5.53	$0.272_{\pm 0.019}$
	HD 145607	A2IV	5.42	$0.307_{\pm 0.023}$
HD 148112	HD 144874	A7V	5.64	$0.340_{\pm 0.024}$
	HD 152614	B8V	4.38	$0.333_{\pm 0.024}$
HD 149438	HD 146624	A1V	4.78	$0.339_{\pm 0.023}$
	HD 148605	B3V	4.79	$0.213_{\pm 0.015}$
HD 152107	HD 143584	F0IV	6.03	$0.327_{\pm 0.023}$
	HD 144206	B9III	4.71	$0.314_{\pm 0.019}$
	HD 149303	A2V	5.68	$0.286_{\pm 0.020}$
	HD 155860	A5III	6.13	$0.231_{\pm 0.017}$
HD 192640	HD 191610	B2.5V	4.93	$0.210_{\pm 0.018}$
	HD 193369	A2V	5.58	$0.263_{\pm 0.016}$
HD 195810	HD 193472	A5	5.94	$0.319_{\pm 0.023}$
	HD 196544	A1IV	5.42	$0.270_{\pm 0.008}$
	HD 196740	B5IV	5.05	$0.210_{\pm 0.012}$
HD 27819	HD 25202	F4V	5.87	$0.376_{\pm 0.026}$
	HD 28226	Am	5.71	$0.345_{\pm 0.025}$
HD 27962	HD 27459	F0IV/V	5.24	$0.409_{\pm 0.029}$
	HD 28226	Am	5.71	$0.345_{\pm 0.025}$
HD 3360	HD 1976	B5IV	5.58	$0.200_{\pm 0.006}$
	HD 2054	B9IV	5.72	$0.207_{\pm 0.006}$
	HD 3240	B7III	5.08	$0.238_{\pm 0.015}$
	HD 6676	B8V	5.77	$0.207_{\pm 0.014}$
HD 33959	HD 34452	A0	5.37	$0.214_{\pm 0.006}$
	HD 34578	A5II	5.03	$0.597_{\pm 0.055}$
	HD 35239	B9III	5.93	$0.227_{\pm 0.016}$
	HD 35520	A1p	5.91	$0.301_{\pm 0.008}$
HD 35468	HD 34203	A0V	5.52	$0.251_{\pm 0.016}$
	HD 34658	F3III/IV	5.32	$0.498_{\pm 0.038}$
	HD 37490	B3V	4.59	$0.229_{\pm 0.023}$
	HD 38899	B9IV	4.88	$0.298_{\pm 0.019}$
HD 58142	HD 47100	B8III	5.33	$0.236_{\pm 0.014}$
	HD 56963	F2V	5.74	$0.405_{\pm 0.030}$
	HD 60652	A5m	5.91	$0.283_{\pm 0.019}$
	HD 70313	A3V	5.54	$0.300_{\pm 0.020}$
HD 886	HD 1439	A0IV	5.88	$0.207_{\pm 0.013}$
	HD 560	B9V	5.53	$0.219_{\pm 0.013}$
HD 89021	HD 85795	A3III	5.27	$0.297_{\pm 0.021}$
	HD 90470	A3V	6.01	$0.264_{\pm 0.007}$
	HD 90840	A4V	5.78	$0.276_{\pm 0.018}$
	HD 91312	A7IV	4.72	$0.554_{\pm 0.036}$
	HD 94334	A1V	4.66	$0.368_{\pm 0.011}$
HD 97633	HD 92825	A3V	5.07	$0.346_{\pm 0.022}$
	HD 93702	A2V	5.31	$0.301_{\pm 0.022}$
	HD 95608	A1V	4.40	$0.453_{\pm 0.036}$

Notes. Section 3 describes the parameters of the calibrators.

Catalogue version 2 (Bourges et al. 2017, JSDC2), but we consider the uncertainties from the JDSC version 1 (Lafrasse et al. 2010), which is more conservative. The observing log is given in Table B.1. The systematic uncertainties stem from the uncertainties on the calibrator diameters, given in Table 1, and they are negligible with respect to the statistical ones.

We used the JMMC LitPro⁷ tool (Tallon-Bosc et al. 2008) to fit a model of a linear limb-darkened disk on the calibrated squared visibilities. This model has the following two parameters: the limb-darkened angular diameter and the limb-darkening coefficient, u_R . Considering the precision of the VEGA measurements and the range of spatial frequencies that have been covered, we cannot adjust the coefficient of the limb darkening. The u_R coefficient for each star is fixed and taken from the Claret & Bloemen (2011) catalog. We searched for the effective temperature T_{eff} , the gravity $\log g$, and the metallicity Z of the star and we took the closest value of each parameter available in the catalog in order to determine u_R . Claret’s grids have a step of 250 K in temperature, thus the largest error we can make on the temperature is 125 K without any interpolation. This error on the temperature leads to an average error of 0.05% on the angular diameter, which is well below our typical errors.

This parameter is fixed in the fitting process of the LITpro tool. The coefficient u_R is given in Table 2 together with the derived limb-darkened angular diameter. The corresponding visibility curves are presented in Fig. A.1. The precision on the angular diameter that we obtain ranges from 0.78 to 5.10%, with a median value of 1.8%.

4. Calibration of the SBCR for early-type stars

The surface brightness S_λ of a star is correlated to its limb-darkened angular diameter θ_{LD} and its apparent magnitude corrected from the extinction m_{λ_0} by the following formula (Wesselink 1969)

$$S_\lambda = m_{\lambda_0} + 5 \log \theta_{LD}. \quad (2)$$

Wesselink (1969) used this definition to highlight the correlation between the surface brightness and the color of the star by the relation below:

$$S_{\lambda_1} = \sum_{n=0}^N C_n (m_{\lambda_1} - m_{\lambda_2})_0^n, \quad (3)$$

which defines the so-called SBCR. Later, Barnes & Evans (1976) developed another definition of the SBCR, which we denote as F_λ . To be consistent with the strategy of Paper I, we consider the Barnes & Evans (1976) definition in the rest of our study, namely

$$F_{\lambda_1} = 4.2196 - 0.1 \times \left\{ \sum_{n=0}^N C_n (m_{\lambda_1} - m_{\lambda_2})_0^n \right\}. \quad (4)$$

The 18-star sample covers a range of $V-K$ color from -1 to $+0.6$ mag. We computed the surface brightness of the 18 stars following Eq. (4). Our fitting strategy uses the orthogonal distance regression (ODR), which considers both F_V and $V-K$ errors. A more detailed description can be found in Appendix A of Paper I. The final SBCR for early-type stars combining Johnson

⁷ The tool is available at <https://www.jmmc.fr/english/tools/data-analysis/litpro/>

Table 2. VEGA angular diameter measurements for the 18 early-type stars (see Sect. 3 for a detailed description of the method used to derive the angular diameter of stars).

Name	Sp. Type	V [mag]	A_V [mag]	$(V-K)_0$ [mag]	K -ref	u_R	$\theta_{LD \pm \sigma_{\theta_{LD}}}$ [mas]	χ_r^2
HD 11415	B2V	3.35	0.050	$-0.484_{\pm 0.012}$	Ducati (2002)	0.281	$0.471_{\pm 0.007}$	1.097
HD 114330	A1IV	4.38	0.012	$-0.022_{\pm 0.04}$	Cutri et al. (2003)	0.413	$0.443_{\pm 0.023}$	1.107
HD 145389	B9V	$4.23^{(*)}$	0.028	$-0.115_{\pm 0.017}$	Cutri et al. (2003)	0.358	$0.452_{\pm 0.007}$	0.089
HD 145570	A1V	4.93	0.003	$0.297_{\pm 0.02}$	Cutri et al. (2003)	0.512	$0.431_{\pm 0.008}$	0.114
HD 148112	A2V	$4.57^{(*)}$	0.016	$0.036_{\pm 0.022}$	Ducati (2002)	0.410	$0.423_{\pm 0.010}$	0.089
HD 149438	B0V	2.82	0.214	$-0.939_{\pm 0.032}$	Ducati (2002)	0.235	$0.338_{\pm 0.011}$	0.528
HD 152107	A1V	$4.82^{(*)}$	0.022	$0.229_{\pm 0.022}$	Cutri et al. (2003)	0.434	$0.432_{\pm 0.009}$	0.134
HD 192640	A2V	$4.95^{(*)}$	0.009	$0.520_{\pm 0.02}$	Cutri et al. (2003)	0.447	$0.489_{\pm 0.011}$	0.244
HD 195810	B6IV	$4.03^{(*)}$	0.053	$-0.398_{\pm 0.038}$	Cutri et al. (2003)	0.319	$0.394_{\pm 0.003}$	0.038
HD 27819	A2V	4.80	0.003	$0.347_{\pm 0.014}$	Ducati (2002)	0.468	$0.489_{\pm 0.007}$	1.592
HD 27962	A2IV	$4.30^{(*)}$	0.003	$0.199_{\pm 0.033}$	Cutri et al. (2003)	0.350	$0.538_{\pm 0.013}$	0.587
HD 3360	B2IV	$3.67^{(*)}$	0.090	$-0.618_{\pm 0.038}$	Cutri et al. (2003)	0.278	$0.350_{\pm 0.004}$	0.299
HD 33959	A9V	$5.00^{(*)}$	0.006	$0.610_{\pm 0.026}$	Cutri et al. (2003)	0.476	$0.515_{\pm 0.004}$	0.472
HD 35468	B2V	$1.64^{(*)}$	0.003	$-0.686_{\pm 0.038}$	Ducati (2002)	0.267	$0.785_{\pm 0.007}$	1.162
HD 58142	A0.5V	4.61	0.006	$0.033_{\pm 0.018}$	Cutri et al. (2003)	0.410	$0.419_{\pm 0.008}$	0.303
HD 886	B2IV	$2.83^{(*)}$	0.034	$-0.756_{\pm 0.032}$	Ducati (2002)	0.277	$0.435_{\pm 0.004}$	0.478
HD 89021	A1IV	3.43	0.006	$0.045_{\pm 0.094}$	Ducati (2002)	0.453	$0.757_{\pm 0.007}$	0.389
HD 97633	A2IV	3.32	0.003	$0.017_{\pm 0.094}$	Cutri et al. (2003)	0.431	$0.769_{\pm 0.010}$	1.443

Notes. From left to right: the name of the star, its spectral type taken from the SIMBAD Astronomical Database, the V magnitude (Kharchenko & Roeser 2009), the visual interstellar extinction (Lallement et al. 2014; Capitanio et al. 2017), the $(V-K)_0$ color corrected from the extinction, the reference used for the infrared K photometry, the limb-darkening coefficient, the limb-darkened angular diameter, and the reduced chi-squared. $(*)$ Variable stars.**Table 3.** Parameters of K/K_s and converted SBCRs for early-type stars (see Sect. 5.2 for a description of the conversions).

N_c / N	C_0	C_1	C_2	C_3	σ_{RMS} [mag]	Expected $\frac{\sigma_{\theta_{LD}}}{\theta_{LD}}$ [%]
K/K_s	$2.6675_{\pm 0.0149}$	$1.6556_{\pm 0.0571}$	$-0.6084_{\pm 0.0667}$	$0.4350_{\pm 0.1530}$	0.00508	2.34
K 11/18	$2.7162_{\pm 0.0196}$	$1.6267_{\pm 0.0773}$	$-0.6191_{\pm 0.1011}$	$0.5733_{\pm 0.2276}$	0.00770	3.55
K_s 7/18	$2.6434_{\pm 0.0199}$	$1.6845_{\pm 0.0792}$	$-0.6947_{\pm 0.0924}$	$0.5734_{\pm 0.2276}$	0.00770	3.55

Notes. The coefficients C_n follow the definition of Eq. (4). The N_c/N column stands for the number of converted photometries as a fraction of the total number of stars in the sample. The RMS and the corresponding expected precision on the angular diameter, computed from Eq. (6), are shown in the last columns.

and 2MASS photometries (K/K_s SBCR hereafter) is then shown in Fig. 1. The coefficients of our SBCR are shown in the first row of Table 3, together with their uncertainties. We discuss the possibility of converting all the photometric measurements into the same system in Sect. 5.2. We restricted our SBCR to the third order.

Using Eq. (2) with visible V magnitudes, we have:

$$\theta_{LD} = 10^{8.4392 - 0.2V_0 - 2F_{V_0}}. \quad (5)$$

Applying the partial derivative method on Eq. (5) gives

$$\frac{\sigma_{\theta_{LD, rms}}}{\theta_{LD}} = 2 \ln(10) \sigma_{F_{V_0}}. \quad (6)$$

The average RMS of the relation is found to be $\sigma_{F_{V_0}} = 0.00508$ mag. This corresponds to a relative precision on the angular diameter of 2.3%, according to Eq. (6). However, the lack of measurements in the blue part (i.e., $V-K < -0.2$ mag) has to be taken into account. We split the $V-K$ validity domain into two ranges, namely $-0.8 \text{ mag} < V-K < -0.2 \text{ mag}$ and

$-0.2 \text{ mag} < V-K < 0.6 \text{ mag}$, and we computed the RMS of the SBCR on both ranges. For $V-K < -0.2$ mag, we expect a precision of 3.1% on the angular diameter. On the other hand, the expected precision is 1.8% for $V-K > -0.2$ mag.

5. Discussion

5.1. Impact of the variability on the SBCR

As in Paper I, we aimed to quantify the impact of the variability on the SBCR. In our sample of early-type stars, ten out of 18 stars are flagged as variables. Their variability ranges from ± 0.01 mag to ± 0.06 mag according to Samus' et al. (2017), with a median value at 0.04 mag. We computed a SBCR considering an offset on the V magnitude of the ten variable stars corresponding to their maximum amplitude. The resulting SBCR is consistent at a level of less than 1σ with the current SBCR all over the $V-K$ validity domain. We conclude that a variability under 0.06 mag does not have any consequence on our result. Keeping these variable stars in the sample is therefore acceptable.

Fig. 1. F_{V_0} derived in this work as a function of $(V-K)_0$. Empty squares are data with a K Johnson photometry, while filled squares are those with a K_s photometry. The third-order SBCR for early-type dwarfs is shown by the red solid line. The gray area denotes the RMS of the relation. The bottom panel presents the residual in units of the RMS of the relation. See Sect. 4 for a description of the fitting strategy.

5.2. The K -band photometry of early-type stars

The Johnson- K and 2MASS- K_s photometries of our sample are not equally distributed in terms of $V-K$ colors (see Fig. 1), which prevents one from calibrating a purely homogeneous SBCR, which is either based on K or K_s , respectively. We therefore made the choice of combining K and K_s data to calibrate the SBCR. The precision we expect on the angular diameter using this relation is 2.3%, but we do not exclude a bias due to the fact that we mixed different infrared photometric bands. However, using the conversion relations K to K_s or the reverse also implies some difficulties. Indeed, the typical transformation equations (Bessell & Brett 1988; Carpenter 2001) are indirectly deduced from the 2MASS-CIT and 2MASS-SAO equations of Bessell & Brett (1988) and Carpenter (2001). Moreover, the CIT equation of Bessell & Brett (1988) is based on the observation of only 21 stars, and the bluer spectral type of the sample is B7. Regarding our sample, this corresponds to a $V-K$ color of -0.3 mag. Nevertheless, the resulting converted SBCRs are shown as blue and green solid lines in Fig. 2 for Johnson- K and 2MASS- K_s photometries respectively. The coefficients are shown in Table 3. The converted Johnson- K SBCR is consistent with the inhomogeneous SBCR at less than 1σ over all the $V-K$ validity domain. Concerning the uniform 2MASS relation, such a conversion does not influence the calibration of the SBCR at more than 1σ for $V-K > -0.4$ mag. The inconsistency, however, reaches more than 4σ for the bluest part of the relation. The expected precision on the angular diameter using a SBCR based on a uniform set of photometry is of 3.6%.

In conclusion, if one wants to derive the angular diameter of a star with K photometry, we suggest using the K/K_s or K SBCRs

Fig. 2. Comparison between our SBCR with mixed Johnson-2MASS photometries (red solid line), SBCR with converted 2MASS K_s photometry (green solid line), and Johnson- K photometry (blue solid line) for early-type stars. Light-gray dots are eclipsing binaries from Taormina et al. (2019). The orange dot shows the data for HD 149438 using the HIPPARCOS parallax. See Sect. 5.2 for detailed information on the converted SBCRs.

in Table 3, with a good level of confidence, the K/K_s relation being the most precise. If instead one has K_s photometry for his or her star, using the SBCR based on K_s is more consistent, but a bias due to the conversion from K to K_s in the calibration process of the SBCR is not excluded, as already discussed. Future investigations are needed to evaluate the consistency of such photometric conversion relations on the early-type stars' color range.

5.3. Comparison with the literature

HD 35468 was already observed by Challouf et al. (2014) and they obtained an angular diameter of 0.715 ± 0.005 mas, whereas our measurement is 0.786 ± 0.007 mas. Challouf et al. (2014) took the uniform-disk angular diameter of their calibrators in the JSDC (Lafosse et al. 2010), while we used the second version of the catalog (Bourges et al. 2017). We processed the measurements of Challouf et al. (2014) using the JSDC2 angular diameters measurements for the three calibrators of Challouf et al. (2014), and with the same strategy described in this work. We found an angular diameter of 0.807 ± 0.026 mas, which is consistent with our value at a level of 1σ . This clearly shows some bias in the initial data reduction of Challouf et al. (2014) due to differences between JSDC and JSDC2 angular diameters.

In addition to HD 35468, two other stars in Table 2 have been observed previously at CHARA. For HD 97633, Maestro et al. (2013) measured an angular diameter of 0.740 ± 0.024 mas using PAVO, which is consistent with our result. The PAVO measurements of both Maestro et al. (2013) and Gordon et al. (2019) led to a smaller value for the angular diameter for HD 3360

Fig. 3. *Top panel:* comparison between our SBCR for early-type stars and relations of Challouf et al. (2014), Boyajian et al. (2014), Kervella et al. (2004), and Adams et al. (2018). *Bottom panel:* differences in angular diameter estimates between our SBCR and the other relations. See Sect. 5.3 for more details on the comparison.

(0.311 ± 0.010 mas and 0.280 ± 0.018 mas, respectively). The diameter measurements in both papers are based on only two observations, so they could be more susceptible to systematic errors. Also, Casagrande et al. (2014) and White et al. (2018) found increased discrepancies as the angular diameters approached the resolution limits, which is the case here with such a value for the angular diameter for HD 3360.

We also compared our extinction values with those of Challouf et al. (2014). We computed the extinction of our 18 stars by combining the Q -method from Johnson & Morgan (1953) and the intrinsic colors method by Wegner (2014). We found consistent results with *Stilism*, except for one star, namely HD 149438. This star has a smaller *Gaia* parallax value than HIPPARCOS, but the values of the extinction are consistent ($A_V = 0.21$ mag for *Gaia* versus $A_V = 0.15$ mag for HIPPARCOS). The uncertainty of the *Gaia* parallax is larger than the one of HIPPARCOS. We suggest this star could be oversaturated in *Gaia* photometry broadbands. In Fig. 2, we included the data for HD 149438 using the HIPPARCOS parallax (orange dot). The choice of the parallax does not affect the SBCR at more than 1σ . For consistency, we decided to keep the *Gaia* parallax value for HD 149438 in this work.

In Fig. 3, we included a comparison between our SBCR for early-type stars and the relations of Boyajian et al. (2014), Kervella et al. (2004), Challouf et al. (2014), and Adams et al. (2018). The bottom panel shows the normalized difference we expect in terms of the angular diameter between our SBCR and these relations on the $-0.6 < V-K < 0.6$ mag color range. The first two references have been largely used so far, and their $V-K$ color domain of validity cover a large part of the early-type range. Conversely to our SBCR, these relations are fully linear. We expect a difference on the angular diameter of more

than 5% for $V-K > 0$ mag using the relations of Kervella et al. (2004) and Adams et al. (2018). We find an agreement of less than 4% on the angular diameter with the linear relation of Boyajian et al. (2014) for $V-K > -0.2$ mag. This comparison demonstrates that a linear SBCR for early-type stars is no longer valid for $V-K < -0.2$ mag. Comparing our new relation with the dwarfs relation of Challouf et al. (2014) leads to a good agreement of less than 5% on $-0.6 < V-K < 0.6$ mag. The agreement is even more evident for $V-K > 0$ mag, with a difference of at most 2% on the angular diameter estimate.

Among the eight stars measured by Challouf et al. (2014), six of them do not fulfill the stellar characteristics criteria, which can also explain this gap. The inconsistency we see for $V-K < -0.2$ mag can also be explained by the consideration of photometric uncertainties. Indeed, Challouf et al. (2014) have assumed respective errors of 0.015 mag and 0.03 mag on the V and K photometries of their sample, while, via this work, we see that the infrared K photometry plays a major role in the calibration of a SBCR. The uncertainty on the K photometry often exceeds 0.03 mag in our sample, considering an arbitrary error could therefore induce an underestimation of the photometric uncertainty and finally a bias in the calibration of the relation.

Taormina et al. (2019, 2020) analyzed two early-type eclipsing binaries in the LMC with the aim to obtain precise and accurate stellar parameters of each early-type stars, which were then used to derive the surface brightness. They compare the measurements of their first binary system BLMC-01 of class IV or V, with the all-classes relation of Challouf et al. (2014) and they find good agreement. However, their measurements are inconsistent with our inhomogeneous K/K_s SBCR (see light-gray dots on Fig. 2) at more than 5.5σ . The disagreement is lower but still significant when considering the SBCR converted into K_s system. Such inconsistencies require more investigation.

6. Updating late-type stars SBCRs and linking to early-type stars

6.1. Update on late-type SBCRs

In Paper I, we implemented SBCRs using *Gaia* photometry. We noticed an error in the extinction calculation for the G -band. Indeed, Eq. (11) from Paper 1 should have been rewritten in the following way (Danielski et al. 2018):

$$A_G = A_V \times (a_1 + a_2(G-K)_0 + a_3(G-K)_0^2 + a_4(G-K)_0^3 + a_5A_V + a_6A_V^2 + a_7(G-K)_0A_V), \quad (7)$$

with $a_1 = 0.935556283$, $a_2 = -0.090722012$, $a_3 = 0.014422056$, $a_4 = -0.002659072$, $a_5 = -0.030029634$, $a_6 = 0.000607315$, and $a_7 = 0.002713748$. The converted SBCRs using the *Gaia* photometry are shown in Fig. 4, together with their parameters in Table 4. The expected precision on the angular diameter ranges from 1.1 to 2.5%.

We also propose to improve the calibration of our SBCRs for late-type stars obtained in Paper I by using transformation equations to build a uniform infrared photometry set. We used Eq. (A.1) from Carpenter (2001), namely $K_s = K - 0.044$ mag, to transform our photometric sets. The resulting SBCRs are shown in Fig. 5 for each sample (i.e., F5/K7-II/III, F5/K7-IV/V, M-II/III, and M-V stars). Their parameters are included in Table 5. The angular diameter precision ranges from 1.0 to 2.7%.

Fig. 4. Converted SBCRs for late-type stars using *Gaia* photometry, plotted above the data (black squares). Relations with converted Johnson K -photometries are shown as a blue solid line, while 2MASS K_s ones are in a green solid line. See Sect. 6.1 for more information.

6.2. Connecting late-type to early-type SBCRs

Our F5/K7-IV/V SBCR covers a color range from $V-K = 1$ mag to $V-K = 2.7$ mag, while our SBCR for early-type dwarf stars is valid up to $V-K = 0.6$ mag. There is therefore a gap in the validity domain of our relations between $V-K = 0.6$ mag and $V-K = 1$ mag. One way to solve this problem would be to merge

both samples and deduce a unique SBCR. The methodology used to develop both SBCRs differs regarding the interferometric measurements we have selected to calibrate the relations. The SBCR for early-type stars is a homogeneous relation calibrated using measurements from a single instrument, namely VEGA. On the other hand, SBCRs for late-type stars were implemented with measurements taken from different instruments. Also, we

Fig. 5. SBCRs for late-type stars using a converted photometric set for each sample, plotted above the data (black squares). Relations with converted Johnson K -photometries are shown as a blue solid line, while 2MASS K_s ones are in a green solid line. See Sect. 6.1 for more information.

have shown in Paper I that SBCRs depend on the spectral type of stars, where we raised strong discrepancies between M stars and F5/K7 stars. We thus made the choice not to calibrate a single SBCR for early- and late-type stars to avoid any systematics due to a mix of spectral types, and also due to a mix of uniform and inhomogeneous sets of interferometric measurements. By extending our F5/K7-IV/V SBCR as calibrated in

Paper I (based on K and K_s photometry) until $V-K = 0.6$ mag, we find a difference of 2σ between both early-type (using K/K_s) and late-type relations. They are however consistent in the error bars at 1σ . The same result is found when comparing consistently uniform K and K_s SBCRs, respectively. We expect the CHARA/SPICA instrument to solve this gap by providing more F0 to F5 measurements.

Table 4. Parameters of the new converted SBCRs for late-type stars calibrated with *Gaia* photometry.

		N_c / N	SBCR F_{G_0} vs. $(G-K)_0$	$(G-K)$ range [mag]	σ_{RMS} [mag]	Expected $\frac{\sigma_{\theta_{\text{LD}}}}{\theta_{\text{LD}}}$ [%]
F5/K7-II/III	Paper I ^(*)		$F_{G_0} = -0.1289_{\pm 0.0010}(G-K)_0 + 3.9341_{\pm 0.0022}$	[1.6; 3.1]	0.00244	1.12
	K	0/69	$F_{G_0} = -0.1289_{\pm 0.0010}(G-K)_0 + 3.9341_{\pm 0.0022}$	[1.6; 3.1]	0.00244	1.12
	K_s	69/69	$F_{G_0} = -0.1289_{\pm 0.0010}(G-K_s)_0 + 3.9398_{\pm 0.0022}$	[1.7; 3.2]	0.00244	1.12
F5/K7-IV/V	Paper I ^(*)		$F_{G_0} = -0.1436_{\pm 0.0017}(G-K)_0 + 3.9587_{\pm 0.0027}$	[0.8; 2.7]	0.00494	2.27
	K	16/38	$F_{G_0} = -0.1472_{\pm 0.0015}(G-K)_0 + 3.9604_{\pm 0.0025}$	[0.8; 2.7]	0.00494	2.27
	K_s	22/38	$F_{G_0} = -0.1472_{\pm 0.0015}(G-K_s)_0 + 3.9669_{\pm 0.0024}$	[0.8; 2.8]	0.00494	2.27
M-II/III	Paper I ^(*)		$F_{G_0} = -0.1295_{\pm 0.0025}(G-K)_0 + 3.9324_{\pm 0.0088}$	[2.9; 4.9]	0.00543	2.50
	K	0/24	$F_{G_0} = -0.1295_{\pm 0.0025}(G-K)_0 + 3.9324_{\pm 0.0088}$	[2.9; 4.9]	0.00543	2.50
	K_s	24/24	$F_{G_0} = -0.1295_{\pm 0.0025}(G-K_s)_0 + 3.9381_{\pm 0.0089}$	[3.0; 4.9]	0.00543	2.50
M-IV/V	Paper I ^(*)		$F_{G_0} = -0.1388_{\pm 0.0028}(G-K)_0 + 3.9393_{\pm 0.0093}$	[3.0; 4.6]	0.00382	1.76
	K	13/15	$F_{G_0} = -0.1416_{\pm 0.0021}(G-K)_0 + 3.9444_{\pm 0.0069}$	[3.0; 4.6]	0.00407	1.87
	K_s	2/15	$F_{G_0} = -0.1416_{\pm 0.0021}(G-K_s)_0 + 3.9506_{\pm 0.0069}$	[3.0; 4.6]	0.00407	1.87

Notes. The corrected relations from Paper I are shown in the first row for each sample of stars. We describe the correction applied here in Sect. 6.1. The N_c/N column stands for the number of converted photometries as a fraction of the total number of stars in the sample. The RMS and the corresponding expected precision on the angular diameter are shown in the last columns. ^(*)SBCR of Paper I corrected from the extinction error (see Sect. 6.1).

Table 5. Parameters of the SBCRs for late-type stars described in Paper I, but converted into the same photometric system (see Sect. 6.1).

		N_c / N	SBCR F_{V_0} vs. $(V-K)_0$	$(V-K)$ range [mag]	σ_{RMS} [mag]	Expected $\frac{\sigma_{\theta_{\text{LD}}}}{\theta_{\text{LD}}}$ [%]
F5/K7-II/III	K	0/70	$F_{V_0} = -0.1220_{\pm 0.0006}(V-K)_0 + 3.9278_{\pm 0.0016}$	[1.8; 3.8]	0.00223	1.03
	K_s	70/70	$F_{V_0} = -0.1220_{\pm 0.0006}(V-K_s)_0 + 3.9332_{\pm 0.0016}$	[1.8; 3.9]	0.00223	1.03
F5/K7-IV/V	K	16/38	$F_{V_0} = -0.1404_{\pm 0.0014}(V-K)_0 + 3.9603_{\pm 0.0025}$	[1.0; 3.4]	0.00575	2.65
	K_s	22/38	$F_{V_0} = -0.1404_{\pm 0.0014}(V-K_s)_0 + 3.9665_{\pm 0.0025}$	[1.0; 3.4]	0.00575	2.65
M-II/III	K	0/29	$F_{V_0} = -0.1165_{\pm 0.0012}(V-K)_0 + 3.9051_{\pm 0.0055}$	[3.7; 7.0]	0.00461	2.12
	K_s	29/29	$F_{V_0} = -0.1165_{\pm 0.0012}(V-K_s)_0 + 3.9102_{\pm 0.0055}$	[3.7; 7.0]	0.00461	2.12
M-IV/V	K	14/16	$F_{V_0} = -0.1247_{\pm 0.0009}(V-K)_0 + 3.9133_{\pm 0.0044}$	[3.8; 7.4]	0.00404	1.86
	K_s	2/16	$F_{V_0} = -0.1247_{\pm 0.0009}(V-K_s)_0 + 3.9188_{\pm 0.0044}$	[3.8; 7.4]	0.00404	1.86

Notes. The N_c/N column stands for the number of converted photometries as a fraction of the total number of stars in the sample. The RMS and the corresponding expected precision on the angular diameter are shown in the last columns.

7. Conclusions and perspectives

We carefully selected 18 early-type stars according to selection criteria in order to measure their angular diameter with the VEGA combiner at the CHARA Array interferometer. The mean precision we obtain on our angular diameter measurements is 2%. Using our SBCR leads to an expected statistical precision of 2.3% on the derived angular diameter, but the user should take care of photometric conversion issues discussed in Sect 5.2.

This work supports the results shown by Paper I. Indeed, to reach such statistical precision, we demonstrated the necessity of implementing selection criteria in order to calibrate the SBCRs. We also showed the need of including the $V-K$ uncertainties in the fitting process. In Paper I, we showed that SBCRs for late-type stars depend on the class of stars. In this work, and as a first step, we considered only dwarfs and subgiants stars. We also improved SBCRs for late-type stars presented in Paper I by converting the photometric data into the same system, and also by correcting the calculation of the extinction in the *Gaia* band.

The lack of a large set of homogeneous angular diameter interferometric measurements remains a major problem in the implementation of SBCRs. With the future SPICA instrument (Mourard et al. 2018) that will be installed at the CHARA array, we expect to derive the angular diameter of 800 stars all over the HR diagram, with a 1% precision level. CHARA/SPICA will also study rotation, multiplicity, wind, and environment by performing images in the visible domain. That should therefore reinforce our knowledge of SBCRs for early-type stars and in particular help to better understand the impact of stellar activity.

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⁸ Available at <http://cdsweb.u-strasbg.fr/>

NASA/ADS system. This work has made use of data from the European Space Agency (ESA) mission *Gaia* (<https://www.cosmos.esa.int/gaia>). This research has made use of the Jean-Marie Mariotti Center OIFits Explorer service⁹. This research also made use of Astropy, a community-developed core Python package for Astronomy (Astropy Collaboration 2018). MT acknowledges financial support from the Polish National Science Center grant PRELUDIUM 2016/21/N/ST9/03310. The research leading to these results has received funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation program (grant agreement No 695 099) and from the National Science Center, Poland grant BEETHOVEN UMO-2018/31/G/ST9/03050. We acknowledge support from the DIR/WK/2018/09 grant of the Polish Ministry of Science and Higher Education.

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⁹ Available at <http://www.jmmc.fr/oifitseexplorer>

Appendix A: VEGA visibility curves

Fig. A.1. Squared visibility versus spatial frequency for all stars in our sample with their corresponding statistical uncertainties (red dots). The solid blue lines indicate the best uniform-disk model obtained from the LITpro fitting software. See Sect. 3 for a detailed description of the fitting strategy.

Fig. A.1. continued.

Fig. A.1. continued.

Appendix B: Observing log

Table B.1. VEGA observing log of the 18 early-type stars.

Star	Date [yyyy.mm.dd]	Peak	HA [h]	λ [nm]	$\lambda_{\min}$ [nm]	$\lambda_{\max}$ [nm]	B_p [m]	Arg [deg]	S/N	$V_{\text{cal}}^2 \pm \text{stat} \pm \text{syst}$
HD 11415	2019.10.07	1	3.39	710	700	720	65.87	-165.17	15.33	0.922 $\pm 0.114 \pm 0.003$
	2019.10.07	1	3.40	730	720	740	65.87	-165.27	20.17	0.898 $\pm 0.045 \pm 0.004$
	2019.10.07	1	3.79	730	720	740	65.87	-170.48	12.92	0.925 $\pm 0.072 \pm 0.005$
	2019.10.07	1	4.14	730	720	740	65.86	-175.28	18.26	0.932 $\pm 0.051 \pm 0.004$
	2019.10.07	1	4.53	730	720	740	65.86	179.49	11.13	0.939 $\pm 0.084 \pm 0.004$
	2020.07.18	1	-3.41	730	720	740	219.83	-51.60	5.60	0.353 $\pm 0.063 \pm 0.007$
	2020.07.18	1	-2.50	730	720	740	229.86	-65.39	5.26	0.286 $\pm 0.054 \pm 0.005$
	2020.07.19	1	-6.57	730	720	740	241.34	34.36	7.29	0.197 $\pm 0.027 \pm 0.017$
	2020.07.19	1	-5.50	730	720	740	244.59	21.34	5.29	0.180 $\pm 0.034 \pm 0.015$
	2020.07.19	1	-4.98	730	720	740	245.45	14.99	3.33	0.225 $\pm 0.067 \pm 0.019$
	2020.07.20	1	-1.90	730	720	740	142.75	-89.44	11.57	0.613 $\pm 0.053 \pm 0.015$
HD 114330	2019.02.23	1	0.94	730	720	740	182.53	12.90	4.48	0.442 $\pm 0.099 \pm 0.025$
	2019.02.23	1	2.18	710	700	720	177.66	-2.94	6.54	0.535 $\pm 0.082 \pm 0.032$
	2019.02.23	1	2.21	730	720	740	177.73	-3.34	7.09	0.566 $\pm 0.080 \pm 0.024$
	2019.02.23	1	2.68	710	700	720	180.08	-9.39	7.49	0.545 $\pm 0.073 \pm 0.033$
	2019.02.23	1	2.70	730	720	740	180.21	-9.61	6.96	0.522 $\pm 0.075 \pm 0.022$
	2020.03.04	1	-0.41	705	695	715	62.90	-114.85	17.61	0.853 $\pm 0.048 \pm 0.005$
	2020.03.04	1	-0.41	725	715	735	62.90	-114.85	14.76	0.809 $\pm 0.055 \pm 0.002$
	2020.03.04	1	-0.08	705	695	715	61.16	-115.11	11.59	0.829 $\pm 0.061 \pm 0.002$
	2020.03.04	1	-0.08	725	715	735	61.16	-115.11	9.20	0.863 $\pm 0.077 \pm 0.002$
	2020.03.07	1	-1.26	725	715	735	65.54	-114.99	10.20	0.858 $\pm 0.084 \pm 0.005$
	2020.03.07	1	-0.42	705	695	715	62.96	-114.84	17.71	0.903 $\pm 0.051 \pm 0.011$
	2020.03.07	1	-0.42	725	715	735	62.96	-114.84	9.01	0.848 $\pm 0.094 \pm 0.002$
HD 145389	2019.05.02	1	2.22	730	720	740	151.21	-141.88	6.65	0.586 $\pm 0.088 \pm 0.008$
	2019.05.02	1	2.69	710	700	720	149.07	-147.98	9.56	0.614 $\pm 0.094 \pm 0.008$
	2019.05.02	1	2.68	730	720	740	149.11	-147.84	11.63	0.642 $\pm 0.055 \pm 0.010$
	2019.05.02	1	2.98	710	700	720	147.74	-151.96	12.80	0.590 $\pm 0.046 \pm 0.009$
	2019.05.02	1	3.00	730	720	740	147.66	-152.22	7.14	0.605 $\pm 0.085 \pm 0.009$
	2020.03.04	1	-1.86	725	715	735	61.42	-105.17	11.66	0.914 $\pm 0.078 \pm 0.002$
HD 145570	2019.06.14	1	-0.78	710	700	720	154.56	-110.12	9.34	0.631 $\pm 0.069 \pm 0.013$
	2019.06.17	1	-1.96	710	700	720	154.79	-113.22	11.50	0.619 $\pm 0.054 \pm 0.010$
	2019.06.17	2	-1.94	710	700	720	83.50	87.05	5.14	0.807 $\pm 0.157 \pm 0.006$
	2019.06.17	1	-1.62	710	700	720	156.03	-112.12	6.60	0.598 $\pm 0.091 \pm 0.010$
	2019.06.17	2	-1.53	710	700	720	90.32	88.32	13.09	0.871 $\pm 0.067 \pm 0.004$
	2019.06.17	1	-1.22	710	700	720	156.14	-111.08	9.13	0.592 $\pm 0.065 \pm 0.009$
	2019.06.17	2	-1.21	710	700	720	94.85	89.21	8.81	0.784 $\pm 0.089 \pm 0.004$
	2019.06.17	2	-0.86	710	700	720	99.11	90.15	5.86	0.858 $\pm 0.147 \pm 0.006$
HD 148112	2019.06.15	1	-1.77	730	720	740	151.68	-108.73	8.61	0.662 $\pm 0.105 \pm 0.033$
	2019.06.15	1	-1.20	730	720	740	155.38	-110.31	9.89	0.612 $\pm 0.062 \pm 0.026$
	2019.06.15	1	-0.68	730	720	740	156.24	-112.09	9.80	0.647 $\pm 0.066 \pm 0.032$
	2019.08.16	1	4.37	710	700	720	33.77	145.58	12.78	1.002 $\pm 0.078 \pm 0.002$
	2019.08.16	1	4.39	730	720	740	33.76	145.54	14.82	0.995 $\pm 0.067 \pm 0.002$
HD 149438	2019.06.17	1	-0.12	710	700	720	144.32	-102.73	13.50	0.758 $\pm 0.047 \pm 0.010$
	2019.06.17	1	0.27	710	700	720	138.51	-100.46	16.78	0.791 $\pm 0.047 \pm 0.010$
	2019.06.17	1	0.28	730	720	740	138.35	-100.40	10.72	0.799 $\pm 0.075 \pm 0.013$
	2019.06.17	1	1.52	710	700	720	112.40	-92.91	15.05	0.875 $\pm 0.049 \pm 0.011$
	2019.06.17	1	1.52	730	720	740	112.27	-92.88	8.79	0.816 $\pm 0.093 \pm 0.014$
	2020.03.05	1	-0.04	705	695	715	131.76	36.02	12.03	0.737 $\pm 0.082 \pm 0.011$

Notes. The columns list the date, the hour angle HA, the mean, the minimum, and the maximum wavelengths over which the squared visibility is calculated, the projected baseline length B_p , its orientation Arg, and the signal-to-noise ratio (S/N). The last column provides the calibrated squared visibility V_{cal}^2 together with the statistic and the systematic errors. The data are available on the Jean-Marie Mariotti Center OiDB service (available at <http://oidb.jmmc.fr>). The interferometric observations are described in Sect. 3.

Table B.1. continued.

Star	Date [yyyy.mm.dd]	Peak	HA [h]	λ [nm]	$\lambda_{\min}$ [nm]	$\lambda_{\max}$ [nm]	B_p [m]	Arg [deg]	S/N	$V^2_{\text{cal+stat+sys}}$
HD 152107	2019.06.15	1	0.91	710	700	720	155.89	−126.31	8.65	0.627 $\pm 0.073 \pm 0.013$
	2019.06.15	1	0.89	730	720	740	155.91	−126.155	8.33	0.614 $\pm 0.074 \pm 0.010$
	2019.06.15	1	1.53	710	700	720	154.36	−133.49	8.40	0.633 $\pm 0.075 \pm 0.010$
	2019.06.15	1	1.52	730	720	740	154.39	−133.41	11.92	0.610 $\pm 0.051 \pm 0.009$
	2020.03.04	1	−1.09	705	695	715	63.97	−112.32	9.26	0.935 $\pm 0.101 \pm 0.002$
	2020.03.04	1	−1.09	725	715	735	63.96	−112.30	4.81	0.873 $\pm 0.182 \pm 0.002$
	2020.03.04	1	−0.56	705	695	715	65.10	−117.60	12.67	0.905 $\pm 0.071 \pm 0.003$
	2020.03.04	1	−0.56	725	715	735	65.09	−117.54	9.07	0.856 $\pm 0.094 \pm 0.002$
HD 192640	2019.07.06	1	0.79	710	700	720	176.39	−27.36	5.96	0.413 $\pm 0.094 \pm 0.005$
	2019.07.06	1	0.79	730	720	740	176.39	−27.38	6.88	0.427 $\pm 0.062 \pm 0.005$
	2019.07.06	1	1.17	730	720	740	175.53	−30.27	8.47	0.438 $\pm 0.052 \pm 0.005$
	2019.08.15	1	2.47	730	720	740	33.15	151.85	13.76	0.933 $\pm 0.052 \pm 0.001$
	2019.08.15	2	2.47	730	720	740	201.98	−37.34	5.97	0.383 $\pm 0.064 \pm 0.006$
	2019.08.16	1	1.75	710	700	720	33.59	156.70	17.96	1.002 $\pm 0.044 \pm 0.001$
	2019.08.16	1	1.74	730	720	740	33.60	156.77	11.38	0.972 $\pm 0.085 \pm 0.001$
	2019.08.16	1	2.36	710	700	720	33.23	152.55	23.78	0.979 $\pm 0.041 \pm 0.001$
	2019.08.16	1	2.33	730	720	740	33.25	152.72	19.09	0.952 $\pm 0.050 \pm 0.001$
HD 195810	2019.07.06	1	1.67	710	700	720	174.90	−33.63	12.65	0.571 $\pm 0.073 \pm 0.010$
	2019.07.06	1	1.66	730	720	740	174.86	−33.57	7.17	0.625 $\pm 0.087 \pm 0.011$
	2019.07.06	1	2.09	710	700	720	176.42	−35.59	11.32	0.566 $\pm 0.050 \pm 0.011$
	2019.07.06	1	2.08	730	720	740	176.40	−35.56	7.55	0.580 $\pm 0.077 \pm 0.010$
	2019.07.06	1	2.43	710	700	720	177.17	−36.93	9.82	0.585 $\pm 0.060 \pm 0.013$
	2019.07.06	1	2.43	730	720	740	177.17	−36.93	8.89	0.589 $\pm 0.066 \pm 0.011$
	2019.08.17	1	2.99	710	700	720	33.92	149.87	23.50	0.986 $\pm 0.042 \pm 0.001$
	2019.08.17	1	3.00	730	720	740	33.92	149.82	19.92	0.990 $\pm 0.050 \pm 0.001$
HD 27819	2019.10.07	1	−2.12	730	720	740	63.50	−112.26	17.10	0.851 $\pm 0.050 \pm 0.008$
	2019.10.07	1	−1.70	730	720	740	64.91	−113.64	10.64	0.871 $\pm 0.082 \pm 0.007$
	2019.10.07	1	−1.29	730	720	740	65.67	−115.17	11.30	0.915 $\pm 0.081 \pm 0.008$
	2019.10.09	1	−0.89	710	700	720	242.78	25.72	8.39	0.212 $\pm 0.025 \pm 0.031$
	2019.10.09	1	−0.88	730	720	740	242.71	25.63	6.94	0.196 $\pm 0.028 \pm 0.028$
	2019.10.09	1	−0.50	710	700	720	240.55	22.80	6.89	0.238 $\pm 0.035 \pm 0.033$
	2019.10.09	1	−0.51	730	720	740	240.58	22.84	12.26	0.206 $\pm 0.017 \pm 0.028$
	2019.10.09	1	−0.08	710	700	720	238.10	19.40	8.07	0.226 $\pm 0.026 \pm 0.028$
	2019.10.09	1	0.39	710	700	720	235.49	15.25	8.39	0.176 $\pm 0.020 \pm 0.022$
	2019.10.09	1	0.38	730	720	740	235.52	15.31	11.87	0.184 $\pm 0.015 \pm 0.024$
HD 27962	2019.10.07	1	−0.90	710	700	720	65.87	−116.84	8.26	0.817 $\pm 0.099 \pm 0.010$
	2019.10.07	1	−0.92	730	720	740	65.87	−116.77	12.64	0.885 $\pm 0.070 \pm 0.009$
	2019.10.07	1	−0.45	710	700	720	65.52	−119.09	14.64	0.864 $\pm 0.059 \pm 0.009$
	2019.10.07	1	−0.45	730	720	740	65.51	−119.11	9.38	0.866 $\pm 0.092 \pm 0.007$
	2019.10.07	1	0.10	710	700	720	64.24	−122.33	16.21	0.962 $\pm 0.059 \pm 0.010$
	2019.10.10	1	−1.67	710	700	720	165.56	−6.20	1.81	0.432 $\pm 0.040 \pm 0.019$
	2019.10.10	1	−1.67	720	700	740	165.56	−6.17	5.53	0.314 $\pm 0.057 \pm 0.019$
	2020.12.16	1	−1.51	710	700	720	93.93	105.62	8.24	0.700 $\pm 0.085 \pm 0.023$
	2020.12.16	1	−0.88	710	700	720	100.92	101.70	9.30	0.747 $\pm 0.080 \pm 0.017$
	2020.12.16	1	−0.89	730	720	740	100.89	101.72	14.54	0.739 $\pm 0.051 \pm 0.021$
	2020.12.16	1	−0.22	710	700	720	105.76	98.17	9.23	0.686 $\pm 0.074 \pm 0.016$
	2020.12.16	1	−0.23	730	720	740	105.74	98.19	18.87	0.690 $\pm 0.037 \pm 0.020$
	2020.12.16	1	0.50	710	700	720	107.88	94.72	14.49	0.716 $\pm 0.049 \pm 0.020$
HD 3360	2019.08.15	1	−0.68	710	700	720	32.08	175.73	11.37	0.939 $\pm 0.083 \pm 0.001$
	2019.08.15	2	−0.71	710	700	720	201.90	−13.57	6.21	0.591 $\pm 0.091 \pm 0.007$
	2019.08.15	1	−0.70	730	720	740	32.09	175.85	19.57	0.952 $\pm 0.049 \pm 0.001$
	2019.08.16	1	−1.26	730	720	740	32.11	−179.09	13.46	1.018 $\pm 0.076 \pm 0.001$
	2019.08.16	2	−1.26	730	720	740	202.87	−8.25	9.36	0.624 $\pm 0.067 \pm 0.011$
	2020.08.26	1	−2.75	705	695	715	133.59	−83.32	15.32	0.806 $\pm 0.053 \pm 0.005$
	2020.08.26	1	−1.96	705	695	715	142.47	−93.06	18.79	0.742 $\pm 0.039 \pm 0.010$

Table B.1. continued.

Star	Date [yyyy.mm.dd]	Peak	HA [h]	λ [nm]	$\lambda_{\min}$ [nm]	$\lambda_{\max}$ [nm]	B_p [m]	Arg [deg]	S/N	$V^2_{\text{cal} \pm \text{stat} \pm \text{syst}}$
	2020.08.26	1	-1.96	725	715	735	142.47	-93.06	24.95	0.771 $\pm 0.031 \pm 0.010$
	2020.08.26	1	-1.40	705	695	715	147.43	-99.61	25.44	0.765 $\pm 0.030 \pm 0.015$
	2020.08.26	1	-1.40	725	715	735	147.43	-99.61	24.16	0.741 $\pm 0.031 \pm 0.013$
	2020.08.26	1	-1.00	725	715	735	150.34	-104.32	22.90	0.760 $\pm 0.033 \pm 0.014$
	2020.08.26	1	0.13	705	695	715	155.25	-117.31	12.55	0.687 $\pm 0.055 \pm 0.012$
	2020.08.36	1	0.13	725	715	735	155.25	-117.31	22.27	0.695 $\pm 0.031 \pm 0.009$
	2020.08.29	1	-1.94	710	700	720	171.79	-3.35	7.91	0.629 $\pm 0.080 \pm 0.015$
	2020.08.30	1	-0.87	710	700	720	151.08	-105.72	7.99	0.697 $\pm 0.087 \pm 0.011$
	2020.08.30	1	-0.26	710	700	720	154.04	-112.78	23.36	0.709 $\pm 0.030 \pm 0.015$
	2020.08.30	1	-0.26	730	720	740	154.04	-112.78	20.52	0.718 $\pm 0.035 \pm 0.012$
	2020.08.30	1	0.18	730	720	740	155.36	-117.88	20.73	0.732 $\pm 0.035 \pm 0.010$
HD 33959	2019.08.17	1	-3.64	710	700	720	34.00	-159.96	11.11	0.980 $\pm 0.088 \pm 0.001$
	2019.08.17	1	-3.65	730	720	740	34.00	-159.95	13.75	0.996 $\pm 0.072 \pm 0.001$
	2019.08.17	1	-3.26	710	700	720	34.05	-162.80	14.53	0.989 $\pm 0.068 \pm 0.001$
	2019.08.17	2	-3.24	710	700	720	210.32	10.48	8.18	0.255 $\pm 0.031 \pm 0.004$
	2019.08.17	1	-3.27	730	720	740	34.05	-162.76	14.37	1.007 $\pm 0.070 \pm 0.001$
	2019.08.17	2	-3.24	730	720	740	210.31	10.42	16.78	0.256 $\pm 0.015 \pm 0.004$
	2019.08.17	1	-2.81	710	700	720	34.07	-166.35	10.39	0.963 $\pm 0.093 \pm 0.001$
	2019.08.17	1	-2.83	730	720	740	34.07	-166.15	11.85	0.925 $\pm 0.078 \pm 0.001$
	2019.08.17	2	-2.82	730	720	740	210.06	6.55	3.78	0.297 $\pm 0.079 \pm 0.005$
	2020.08.26	1	-3.41	705	695	715	120.39	-91.44	22.24	0.629 $\pm 0.028 \pm 0.056$
	2020.08.26	1	-2.97	705	695	715	129.56	-94.82	21.21	0.575 $\pm 0.027 \pm 0.084$
	2020.08.26	1	-2.97	725	715	735	129.56	-94.82	18.09	0.630 $\pm 0.035 \pm 0.072$
	2020.08.26	1	-2.61	705	695	715	135.92	-97.39	25.38	0.613 $\pm 0.024 \pm 0.071$
	2020.08.26	1	-2.61	725	715	735	135.92	-97.39	16.74	0.591 $\pm 0.035 \pm 0.005$
	2020.08.26	1	-2.18	705	695	715	142.58	-100.48	21.21	0.568 $\pm 0.027 \pm 0.083$
	2020.08.26	1	-2.18	725	715	735	142.58	-100.48	19.17	0.554 $\pm 0.029 \pm 0.068$
HD 35468	2019.10.05	1	-1.20	710	700	720	65.85	-115.71	18.80	0.786 $\pm 0.042 \pm 0.004$
	2019.10.05	1	-1.19	730	720	740	65.85	-115.72	21.47	0.712 $\pm 0.033 \pm 0.003$
	2019.10.05	1	-0.73	730	720	740	65.27	-116.74	18.46	0.768 $\pm 0.042 \pm 0.003$
	2019.10.05	1	0.08	730	720	740	62.49	-119.41	25.37	0.801 $\pm 0.032 \pm 0.003$
	2020.08.30	1	-3.28	710	700	720	131.90	-111.05	13.14	0.278 $\pm 0.021 \pm 0.017$
	2020.08.30	1	-3.29	730	720	740	131.81	-111.06	13.84	0.272 $\pm 0.020 \pm 0.017$
	2020.08.30	1	-2.77	710	700	720	141.22	-110.42	10.25	0.209 $\pm 0.020 \pm 0.020$
	2020.08.30	1	-2.77	730	720	740	141.28	-110.41	8.12	0.243 $\pm 0.030 \pm 0.004$
	2020.08.30	1	-2.33	710	700	720	147.69	-110.21	10.09	0.189 $\pm 0.019 \pm 0.018$
	2020.08.30	1	-2.32	730	720	740	147.80	-110.21	8.17	0.220 $\pm 0.027 \pm 0.004$
	2020.08.30	1	-1.88	710	700	720	152.37	-110.29	11.20	0.182 $\pm 0.016 \pm 0.017$
	2020.08.30	1	-1.88	730	720	740	152.37	-110.29	15.35	0.154 $\pm 0.010 \pm 0.003$
HD 58142	2019.02.23	1	1.07	710	700	720	243.55	8.38	2.58	0.265 $\pm 0.103 \pm 0.006$
	2019.02.23	1	1.66	710	700	720	244.00	2.83	3.37	0.259 $\pm 0.077 \pm 0.008$
	2019.02.23	1	1.66	730	720	740	244.00	2.77	4.75	0.296 $\pm 0.062 \pm 0.011$
	2019.02.23	1	2.14	710	700	720	244.03	-1.76	3.56	0.266 $\pm 0.075 \pm 0.009$
	2019.02.23	1	2.63	710	700	720	243.76	-6.49	2.12	0.199 $\pm 0.094 \pm 0.007$
	2019.02.23	1	2.62	730	720	740	243.77	-6.34	3.50	0.248 $\pm 0.111 \pm 0.012$
	2020.03.07	1	1.83	705	695	715	65.319	-144.30	12.88	0.961 $\pm 0.075 \pm 0.008$
	2020.03.07	2	1.83	705	695	715	154.43	-137.53	5.42	0.670 $\pm 0.124 \pm 0.037$
	2020.03.07	1	1.83	725	715	735	65.322	-144.25	10.16	0.923 $\pm 0.091 \pm 0.004$
	2020.03.07	2	1.83	725	715	735	154.43	-137.53	11.41	0.650 $\pm 0.057 \pm 0.037$
	2020.03.07	1	2.43	705	695	715	64.76	-151.88	13.43	0.953 $\pm 0.071 \pm 0.008$
	2020.03.07	2	2.43	705	695	715	152.58	-145.25	13.37	0.691 $\pm 0.052 \pm 0.037$
	2020.03.07	2	2.43	725	715	735	152.58	-145.25	10.00	0.654 $\pm 0.065 \pm 0.044$
HD 886	2019.08.20	1	1.42	730	720	740	175.54	-31.95	17.64	0.531 $\pm 0.030 \pm 0.005$
	2019.08.20	1	1.95	710	700	720	176.97	-34.74	8.98	0.518 $\pm 0.058 \pm 0.005$
	2019.08.20	1	1.95	730	720	740	176.97	-34.74	18.83	0.480 $\pm 0.025 \pm 0.004$
	2019.08.20	1	2.49	710	700	720	177.44	-37.09	9.13	0.504 $\pm 0.055 \pm 0.005$

Table B.1. continued.

Star	Date [yyyy.mm.dd]	Peak	HA [h]	λ [nm]	$\lambda_{\min}$ [nm]	$\lambda_{\max}$ [nm]	B_p [m]	Arg [deg]	S/N	$V^2_{\text{cal+stat+syst}}$
	2019.08.20	1	2.49	730	720	740	177.44	−37.09	16.47	0.518 $\pm 0.031 \pm 0.005$
	2019.08.21	1	2.48	710	700	720	177.44	−37.05	11.85	0.521 $\pm 0.044 \pm 0.005$
	2019.08.21	1	2.48	730	720	740	177.44	−37.05	8.78	0.535 $\pm 0.061 \pm 0.005$
	2019.08.22	1	1.97	710	700	720	177.02	−34.86	23.36	0.519 $\pm 0.022 \pm 0.005$
	2019.08.22	1	1.97	730	720	740	177.02	−34.86	17.84	0.508 $\pm 0.028 \pm 0.005$
	2019.08.22	1	2.28	710	700	720	177.39	−36.26	20.02	0.506 $\pm 0.025 \pm 0.005$
	2019.08.22	1	2.28	730	720	740	177.39	−36.24	17.81	0.492 $\pm 0.028 \pm 0.004$
	2019.08.22	1	2.63	710	700	720	177.37	−37.62	16.20	0.528 $\pm 0.033 \pm 0.005$
	2019.08.22	1	2.63	730	720	740	177.37	−37.62	16.68	0.520 $\pm 0.031 \pm 0.005$
HD 89021	2019.02.26	1	1.99	710	700	720	64.05	−145.57	8.43	0.760 $\pm 0.069 \pm 0.007$
	2019.02.26	2	2.00	710	700	720	151.20	−138.74	7.89	0.182 $\pm 0.060 \pm 0.007$
	2019.02.26	1	2.00	730	720	740	64.04	−145.63	5.66	0.763 $\pm 0.135 \pm 0.006$
	2019.02.26	2	2.00	730	720	740	151.20	−138.74	3.17	0.249 $\pm 0.079 \pm 0.010$
	2019.02.26	2	2.50	710	700	720	148.56	−145.21	7.36	0.262 $\pm 0.050 \pm 0.007$
	2020.06.15	1	4.88	710	700	720	140.52	179.97	8.83	0.240 $\pm 0.027 \pm 0.002$
	2020.06.15	1	4.89	730	720	740	140.52	179.75	7.65	0.280 $\pm 0.037 \pm 0.003$
	2020.06.15	1	5.35	710	700	720	140.91	172.67	8.54	0.250 $\pm 0.017 \pm 0.002$
	2020.06.15	1	5.34	730	720	740	140.90	172.78	7.76	0.269 $\pm 0.035 \pm 0.002$
	2020.12.19	1	−0.15	710	700	720	107.28	101.47	7.04	0.588 $\pm 0.083 \pm 0.045$
	2020.12.20	1	−0.14	730	720	740	107.31	101.33	5.28	0.546 $\pm 0.104 \pm 0.020$
	2020.12.21	1	0.80	730	720	740	107.62	91.73	7.59	0.530 $\pm 0.070 \pm 0.021$
HD 97633	2019.05.03	1	1.13	710	700	720	142.96	−122.06	8.38	0.252 $\pm 0.030 \pm 0.007$
	2019.05.03	1	1.59	710	700	720	136.12	−125.81	8.83	0.304 $\pm 0.041 \pm 0.009$
	2019.05.03	1	1.60	730	720	740	136.02	−125.87	14.95	0.257 $\pm 0.017 \pm 0.006$
	2019.05.03	1	2.01	710	700	720	129.31	−129.75	8.94	0.309 $\pm 0.052 \pm 0.013$
	2019.05.03	1	2.01	730	720	740	129.25	−129.78	9.18	0.372 $\pm 0.040 \pm 0.013$
	2020.12.16	1	−2.35	710	700	720	80.64	110.16	14.83	0.688 $\pm 0.046 \pm 0.031$
	2020.12.16	1	−2.35	730	720	740	80.59	110.18	17.82	0.673 $\pm 0.038 \pm 0.029$
	2020.12.16	1	−1.74	710	700	720	90.23	105.65	24.23	0.636 $\pm 0.026 \pm 0.029$
	2020.12.16	1	−1.74	730	720	740	90.23	105.65	26.70	0.646 $\pm 0.024 \pm 0.028$
	2020.12.16	1	0.14	710	700	720	107.21	96.11	12.62	0.477 $\pm 0.038 \pm 0.022$
	2020.12.16	1	0.14	730	720	740	107.21	96.11	23.73	0.463 $\pm 0.020 \pm 0.020$
	2020.12.16	1	0.61	710	700	720	107.92	94.19	13.77	0.432 $\pm 0.031 \pm 0.019$
	2020.12.16	1	0.62	730	720	740	107.91	94.18	21.60	0.454 $\pm 0.021 \pm 0.019$