Searching for obscured AGN in z ∼ 2 submillimetre galaxies

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ABSTRACT

Aims. Submillimetre-selected galaxies (SMGs) at high redshift (z ∼ 2) are potential host galaxies of active galactic nuclei (AGN). If the local Universe is a good guide, >50% of the obscured AGN amongst the SMG population could be missed even in the deepest X-ray surveys. Radio observations are insensitive to obscuration; therefore, very long baseline interferometry (VLBI) can be used as a tool to identify AGN in obscured systems. A well-established upper limit to the brightness temperature of 10 K exists in star-forming systems, thus VLBI observations can distinguish AGN from star-forming systems via brightness temperature measurements.

Methods. We present 1.6 GHz European VLBI Network (EVN) observations of four SMGs (with measured redshifts) to search for evidence of compact radio components associated with AGN cores. For two of the sources, e-MERLIN images are also presented.

Results. Out of the four SMGs observed, we detect one source, J123555.14, that has an integrated EVN flux density of 201 ± 15.2 µJy, corresponding to a brightness temperature of 5.2 ± 0.7 x 10 K. We therefore identify that the radio emission from J123555.14 is associated with an AGN. We do not detect compact radio emission from a possible AGN in the remaining sources (J123325.73, J131225.73, and J163650.43). In the case of J131225.73, this is particularly surprising, and the data suggest that this may be an extended, jet-dominated AGN that is resolved by VLBI. Since the morphology of the faint radio source population is still largely unknown at these scales, it is possible that with a ≈10 mas resolution, VLBI misses (or resolves) many radio AGN extended on kiloparsec scales.

Key words. instrumentation: high angular resolution – techniques: interferometric – galaxies: active – galaxies: nuclei – galaxies: starburst

1. Introduction

Submillimetre galaxies (SMGs) are the most bolometrically luminous sources in the Universe (Swinbank et al. 2014). They are responsible for up to half of the total star formation in the Universe and are likely the massive, dusty progenitors of the largest elliptical galaxies that we see in the local Universe (e.g. Simpson et al. 2014; Swinbank et al. 2006; Casey et al. 2014). While it is clear that the processes of star formation are important in SMGs, it is still unclear as to what fraction also hosts active galactic nuclei (AGN) and whether those AGNs are energetically significant in the bolometric output of those luminous galaxies. Some SMG systems probably host components associated with both AGN activity and star-formation processes (for a review, see Biggs et al. 2010). Previous observations of extragalactic objects have illustrated that the upper limit to the brightness temperature of z > 0.1 star-forming galaxies is expected not to exceed T_e < 10^8 K, and distant SMGs with T_e above this value are most likely powered by AGN (e.g. Condon 1992; Middelberg et al. 2010, 2013).

To date, large surveys focused on detecting AGN systems have been conducted at a range of different wavelengths and, in particular, for the IR and X-ray domains (Alexander et al. 2005, 2008; Ivison et al. 2004; Lutz et al. 2005; Menéndez-Delmestre et al. 2007, 2009; Valiante et al. 2007; Pope et al. 2008; Bonzini et al. 2013; Smolčić et al. 2017; Wang et al. 2013; Stach et al. 2019). Due to the dust extinction in the near-IR, optical, and UV, as well as gas absorption in X-ray bands, these surveys are often incomplete. In the mid- and far-IR, incompleteness of AGN surveys may arise from the fact that not all AGN have significant IR emission from a dusty torus and therefore may not be detected. It is postulated that up to a third of AGN are undetected in these
surveys (Mateos et al. 2017). Moreover, other studies have compared AGN selected from various wavebands and find that their host galaxies tend to have different properties in terms of colour (Hickox et al. 2009) and star-formation rates (SFR; Juneau et al. 2013; Ellison et al. 2016). In particular, Hickox et al. (2009) illustrated that there is only very little overlap between their 122 radio-selected AGN and those selected by X-ray or IR. Therefore, dust-free radio surveys are needed to provide a more complete census of the AGN population. Traditional radio surveys are only sensitive to radio-loud (RL) AGN, which only represent a tiny fraction (10−20%) of the whole AGN population; however, modern radio surveys can achieve a flux depth where radio-quiet AGN can be detected (see Prandoni et al. 2018 for a review). Recent work has focused on the radio as it is sensitive to AGN and star formation, and thus provides a method of surveying AGN and star-formation activity across cosmic time (see e.g. Smolčić et al. 2017; Padovani et al. 2015). A lot of work has also been done to look for AGN-driven radio emission, which has been identified by an excess of radio emission compared to what is expected based on the radio-FIR correlation, holding for star-forming galaxies (e.g. Ivison et al. 2010; Condon et al. 2002; Thomson et al. 2014; Magnelli et al. 2015).

Very long baseline interferometry (VLBI) provides an alternative method of identifying AGN that is not affected by star-formation-related radio emission via the detection of high brightness temperature compact radio components as can be seen in Garrett et al. (2001, 2005), Chi et al. (2013), Middelberg et al. (2011, 2013), Herrera Ruiz et al. (2017), for example. In particular, Chi et al. (2013) identified 12 AGNs in the Hubble Deep Field-North (HDF-N) using a global array of VLBI telescopes. As suggested in Chi et al. (2009), the radio-enhanced AGNs are probably obscured AGNs which remain undetected even in the deepest X-ray surveys. Indeed, if the local Universe is a good guide, roughly half of the AGNs in SMGs are Compton thick (e.g. Risaliti et al. 1999).

Chi et al. (2009) argued that deep, high-resolution radio observations are required in order to generate complete samples of obscured AGNs at high redshifts. However, until recently, deep, wide-field VLBI surveys were difficult to realise and the field of view (FoV) is still relatively limited. The employment of new analysis techniques, as seen in Radcliffe et al. (2016) and Deller et al. (2011) for example, has permitted much deeper and wider VLBI surveys of AGN to be conducted (e.g. Middelberg et al. 2013; Herrera Ruiz et al. 2017; Radcliffe et al. 2018). VLBI has become a sensitive tool in distinguishing between AGN and star-formation processes, including the ability to detect Compton-thick AGNs in dusty systems that would otherwise go undetected. These advances have proven that the radio morphological information and brightness temperature measurements provided by VLBI are sensitive enough to isolate the pure star-forming regions from AGN within individual SMGs (e.g. Muxlow et al. 2005; Chi et al. 2013).

In this paper, we present 1.6 GHz VLBI observations of four SMGs at \( z \sim 2 \) using the European VLBI Network (EVN). For a subset of sources we also use e-MERLIN observations from the e-MERLIN Galaxy Evolution (eMERGE) Great Observatories Origins Deep-North (GOODS-N) survey (Muxlow et al. 2005, 2020). The sources are SMG J123555.14+620901.7, SMG J122600.10+620253.5, SMG J131225.73+423941.4, and SMG J163650.43+405734.5. These are located in the HDF-N, SSA-13, and ELAIS-N2 fields and were selected from Chapman et al. (2005). The four sources that were chosen have \( z \sim 2 \), which is the mean redshift of the Chapman et al. (2005) sample, an epoch that may represent the peak of quasar activity (Wolf et al. 2003; Shankar et al. 2009). On the basis of the results of Chi et al. (2013) who detected 12 radio sources brighter than 150 \( \mu Jy \) in the HDF-N, the targets were also chosen to have a total VLA 1.4 GHz flux density >200 \( \mu Jy \), which is much higher than the average value of ~110 \( \mu Jy \) of the whole sample. The VLA 1.4 GHz luminosities of the four sources are between 10 and 30 times more luminous than the local ultra-luminous starburst Arp 220.

This paper is organised as follows. In Sect. 2 we describe the VLBI observations and the data reduction methodology; in Sect. 3 we present and discuss our results and the derived source properties compared with other data at different frequencies; and finally, we note the main conclusions of the paper in Sect. 4. For this paper, we assume a flat \( \Lambda \)CDM universe with \( H_0 = 67.8 \pm 0.9 \; \text{km} \; \text{s}^{-1} \; \text{Mpc}^{-1}, \Omega_m = 0.308 \pm 0.012 \) and \( \Omega_\Lambda = 0.692 \pm 0.012 \) (Planck Collaboration XIII 2016).

### 2. Observations and data analysis

#### 2.1. EVN observations

The source properties, including multi-wavelength flux densities in the literature, redshift information, as well as derived 1.4 GHz luminosities and \( q \) values (see Sect. 3 for details) are listed in Table 1. The flux densities at 1.4 GHz (VLA) and at 850 \( \mu m \), as well as the redshift information are taken from Chapman et al. (2005). The 350 \( \mu m \) fluxes are from the Herschel point-source catalogue\(^1\) except for the upper limit of J131225.73, which was obtained by Dowell et al. (2003) with second-generation Submillimetre High Angular Resolution Camera (SHARC-2) observations. The VLA 1.4 GHz luminosities were derived from their observed 1.4 GHz fluxes without applying any \( k \)-correction. The SFR, assuming that the radio emission arises purely from star-formation processes, was calculated following Kennicutt (1998):

\[
\text{SFR} \left( M_\odot \text{yr}^{-1} \right) = 1.4 \times 10^{-28} \ell_\nu (\text{ergs s}^{-1} \text{Hz}^{-1}),
\]

where \( \ell_\nu \) represents the luminosity at a certain frequency \( \nu \); we used the VLA 1.4 GHz radio luminosities to derive the corresponding SFRs. For the VLBI detected source, J123555.14, the SFR values in Table 1 were derived with and without the contribution from the compact VLBI-detected AGN component. No \( k \)-correction was applied to the far-IR-radio correlation parameter \( q_{450 \mu m} \) and \( q_{180 \mu m} \). The values reported here (see Sect. 3.2 for details).

Our EVN survey was split into two 12h observing sessions with project codes EC029A and EC029B, respectively (PI S. Chi). EC029A was observed on 6 November 2009 targeting J131225.73 (SSA-13) and J163650.43 (ELAIS-N2), and EC029B was observed on 7 November 2009 targeting J123555.14 and J122600.10 (both are located in GOODS-N). The eight telescopes listed in Table 2 were used in both sessions, including the 100 m Effelsberg telescope, the 76 m Lovell telescope, and the 25 m Urgunji telescope in China. The array provided an angular resolution of \( \sim 10 \text{mas} \) and a 1-\( \sigma \) sensitivity of \( \sim 16 \mu Jy \text{beam}^{-1} \).

The sources were observed in the standard phase referencing mode at 1.6 GHz (\( \lambda \approx 18 \text{cm} \)). For EC029A, we used the bright quasar 3C345 as a “fringe finder”. Two well-established VLBI calibrators, J1317+4115 (~200 mJy) and J1640+3946 (~890 mJy), lying 1.7 and 1.4” away from the target sources were used as phase calibrators. Each target, along with its phase

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\(^1\) [http://archives.esac.esa.int/hsa/whsa/]
Table 1. Source properties, including multi-wavelength flux densities in the literature, redshift information, as well as derived 1.4 GHz luminosities, SFR, and \( q \) values.

<table>
<thead>
<tr>
<th>Source name</th>
<th>( S_{1.4,\text{GHz}} ) [mJy]</th>
<th>( S_{850\mu m} ) [mJy]</th>
<th>( S_{350\mu m} ) [mJy]</th>
<th>( z )</th>
<th>( L_{1.4,\text{GHz}} ) [10^{25} , \text{W Hz}^{-1}]</th>
<th>SFR ( q_{1.4,\text{GHz}} ) [10^{-2} , \text{M}_\odot , \text{yr}^{-1}]</th>
<th>( q_{850\mu m} )</th>
<th>( q_{350\mu m} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>J123555.14</td>
<td>212.0 ± 13.7</td>
<td>5.4 ± 1.9</td>
<td>23.1 ± 0.6</td>
<td>1.875</td>
<td>0.55 ± 0.04</td>
<td>0.77 ± 0.06</td>
<td>1.41 ± 0.16</td>
<td>2.04 ± 0.03</td>
</tr>
<tr>
<td>J123555.14 ( ^{(+)} )</td>
<td>– – – –</td>
<td>– – – –</td>
<td>– – – –</td>
<td>– – – –</td>
<td>0.028 ± 0.002</td>
<td>0.039 ± 0.003</td>
<td>2.42 ± 0.15</td>
<td>– – – –</td>
</tr>
<tr>
<td>J123600.10</td>
<td>262.0 ± 17.1</td>
<td>6.9 ± 2.0</td>
<td>57.2 ± 0.6</td>
<td>2.710</td>
<td>1.66 ± 0.11</td>
<td>2.32 ± 0.15</td>
<td>1.42 ± 0.13</td>
<td>2.34 ± 0.03</td>
</tr>
<tr>
<td>J131225.73</td>
<td>752.5 ± 4.2</td>
<td>4.1 ± 1.3</td>
<td>&lt;14.7</td>
<td>1.554</td>
<td>1.23 ± 0.01</td>
<td>1.72 ± 0.01</td>
<td>0.74 ± 0.14</td>
<td>&lt;1.29</td>
</tr>
<tr>
<td>J163650.43</td>
<td>221.0 ± 16.0</td>
<td>8.2 ± 1.7</td>
<td>45.9 ± 2.9</td>
<td>2.378</td>
<td>1.02 ± 0.07</td>
<td>1.42 ± 0.10</td>
<td>1.57 ± 0.10</td>
<td>2.32 ± 0.04</td>
</tr>
</tbody>
</table>

Notes. The SFR and \( q_{1.4\,\text{GHz}} \) value for J123555.14 derived from both the total and AGN-subtracted fractional \( L_{1.4\,\text{GHz}} \), excluding the contribution from the compact AGN core measured by VLBI, are reported in this table. \(^{(+)AGN}\)-subtracted components.

calibrator, was observed with a cycle time of 10 min (8 min the target and 2 min on the phase calibrator). The fringe finder was observed for 4 min in the middle of the session. For EC029B, we had an \( \sim 190 \, \text{mJy} \) primary phase calibrator J1241+6020, which is located \( \sim 1.5^\circ \) from the centre of the HDF-N, and an \( \sim 7 \, \text{mJy} \) secondary phase calibrator J1234+619 lying \( \sim 25^\circ \) away from the HDF-N centre. In this case, the observations cycled between the primary calibrator, the secondary calibrator, J123600.10 (the target), the secondary calibrator, and J123555.14 (the target). Details about the calibrators are listed in Table 3. The secondary calibrator of EC029B, J1234+619, was found to be displaced with respect to the correlated position by \( \sim 20 \, \text{mas} \). The improved position measured by these observations via the AIPS task \textit{JMFIT} is given in Table 3.

Our observations were recorded at 1024 Mbits s\(^{-1}\) (Nyquist sampled with two-bit encoding, dual-polarisation, 8 \times 16 MHz IF channels) for a total observing time of \( \sim 24 \, \text{h} \). Two-second integrations and 16 spectral channels per 16 MHz baseband were adopted for the correlation parameters. Our data were correlated at the Joint Institute for VLBI ERIC\(^{2}\) (JIVE).

2 www.jive.eu/

Table 2. Telescopes used in the observations (ordered alphabetically).

<table>
<thead>
<tr>
<th>Name</th>
<th>Location</th>
<th>Diameter (m)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Effelsberg</td>
<td>Germany</td>
<td>100</td>
</tr>
<tr>
<td>Lovell</td>
<td>UK</td>
<td>76</td>
</tr>
<tr>
<td>Medicina</td>
<td>Italy</td>
<td>32</td>
</tr>
<tr>
<td>Noto</td>
<td>Italy</td>
<td>32</td>
</tr>
<tr>
<td>Ossala</td>
<td>Sweden</td>
<td>25</td>
</tr>
<tr>
<td>Torun</td>
<td>Poland</td>
<td>32</td>
</tr>
<tr>
<td>Urumqi</td>
<td>China</td>
<td>25</td>
</tr>
<tr>
<td>Westerbork</td>
<td>The Netherlands</td>
<td>25</td>
</tr>
</tbody>
</table>

22. Data analysis

The observed data were analysed using the Astronomical Image Processing System (AIPS)\(^{3}\). All of the sources and the calibrators were calibrated with the following strategy: an initial amplitude calibration that was derived from the system temperature and the gain curves of the telescopes (available from the JIVE archive as a calibration table generated by the pipeline processes) was applied first. Thereafter, bad data with abnormally high or corrupted amplitude or phase information were removed using the AIPS tasks SPFLG and CLIP. The dispersive delays were then corrected for using an ionospheric map that was implemented within VLBATECR. The instrumental delays (fixed delay offsets between the IF channels were calibrated by running FRING on data from a single scan on a strong source 3C345 and J1241+6020 for EC029A and EC029B, respectively). We performed fringe-fitting on the fringe finder and primary phase calibrators to calibrate the phase and phase-rates by also using FRING. Finally, we performed a bandpass calibration using BPASS by again employing 3C345 and J1241+6020 for EC029A and EC029B, respectively.

We employed different strategies for EC029A and EC029B for self-calibration. For EC029A, we used SPLIT to separate the phase-calibrators (J1317+4115 and J1640+3946) and generated the best possible self-calibrated maps of these sources. The clean-components of these maps were then used to update the source model used by FRING, and the phase and phase-rates were re-determined. After applying the new corrections, we used SPLIT again to separate the calibrators, then successive loops of \textit{IMAGR} and \textit{CALIB} resulted in the final maps for these sources. The amplitude and phase corrections derived from \textit{CALIB} were then applied to the target sources using \textit{CLCAL}. The target data were then separated using \textit{SPLIT} and dirty images generated with \textit{IMAGR}.

For EC029B, a similar approach was implemented with corrections from the primary calibrator, which were determined by FRING and later refined by CALIB, including amplitude corrections. These were applied to the secondary calibrator and the targets. The initial images made of the secondary calibrator show it to be shifted about 20 mas from the phase centre (correlated position) of the map. This position was known to be an error but it remained uncorrected at the time the data were correlated. FRING was re-run on the secondary calibrator using the new map with the correct position, and the phase and phase-rate corrections were also applied to the targets.

Of the four targeted sources, only one was unambiguously detected by VLBI, J123555.14, which is located in the HDF-N. The other three sources were not detected. We generated maps using highly tapered uv-data (excluding the Urumqi telescope), but no further detections were found. We used the Common Astronomy Software Applications (CASAl McMullin et al. 2007) task \textit{viewer} to generate the radio contour maps of the sources.

Table 4 summarises the VLBI results. In this table, the angular size of J123555.14 was measured by fitting a single 2D Gaussian to the source using the AIPS task \textit{JMFIT}. The linear size at the distance of this source was calculated following Wright (2006)\(^{4}\).

4 http://www.astro.ucla.edu/~wright/CosmoCalc.html

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It reaches an rms noise level of \( \sim \) J123555.14 with a resolution that matched the published eMERGE DR1 maximum extent of the Lovell Telescope primary beam response. We also created an additional re-weighted eMERGE 1.5 GHz map to derive the spectral index (see Sect. 3.3), has an rms noise level of \( \sim \) J123555.14 \( \pm \) 110.2 \( \pm \) 15.2 \( \times \) 201.1 \( \pm \) 40.2 \( \times \) 0.95 \( \times \) \( 10^5 \) \( K \) being an AGN, it is not certain that AGN in sources can reach, this value can typically only be attained by a Gaussian fitting using 3MFIT. \(^{16}\) Fomalont et al. (2006). \(^{17}\)VLBA calibrator catalogue (http://www.vlba.nrao.edu/astro/calib/). \(^{18}\)Ivison et al. (2002).
Fig. 1. Contoured maps for the VLBI-detected source J123555.14+620901.7. 

(a): EVN 1.6 GHz image plotted with contour levels (3, 4, 5, 6, and 7) \times \text{rms noise level}; the white cross presents the VLA coordinates measured by Richards (2000). The FoV of this image is shown by a red box in each of the other three sub-figures. 

(b): Hubble CANDELS F814W ACS image in which the source shows a face-on disc-like morphology with a close companion for clarity; the image was smoothed with a 2D Gaussian convolution with FWHP of \(\sim 0.1''\). The overlaid contour was plotted at 3, 6, 12, 24, and 48 \times \text{rms noise of the eMERGE-JVLA maximum sensitivity image at 1.5 GHz.}

(c): eMERGE re-weighted 1.5 GHz map presented with contour levels 3, 6, 12, and 24 \times \text{rms noise level.}

(d): eMERGE 5.5 GHz map plotted with contour levels (3, 4, and 5) \times \text{rms noise level.}

The detected compact AGN core is clearly shown in the images. In each panel, the beam pattern of the contours is illustrated with a black ellipse.

fraction of their flux is from the compact central core (Jarvis et al. 2019; Muxlow et al. 2005).

The detected source J123555.14 has a brightness temperature of \(5.2 \pm 0.7 \times 10^5\ K\), which is approximately five times higher than the star-forming envelope. This supports the interpretation that J123555.14 contains an AGN core. The undetected sources, however, have upper limits on their brightness temperatures exceeding \(10^5\ K\), which cannot completely rule out AGN activity in these objects.

3.2. Infrared-radio correlation

Our four sources are part of a larger sample of SMGs from Chapman et al. (2005) that have good positions and measured
redshifts. Table 1 presents the source flux densities at 1.4 GHz and 850 \( \mu \)m provided in Chapman et al. (2005) and the derived \( q_{1.4 \text{GHz}} \) values. We define the latter as:

\[
q_{1.4 \text{GHz}} = \log_{10} \left( \frac{L_{850 \mu m}}{L_{1.4 \text{GHz}}} \right),
\]

where \( L_{850 \mu m} \) and \( L_{1.4 \text{GHz}} \) are the luminosities at 850 \( \mu \)m and 1.4 GHz. We also calculated the \( q_{350 \mu m} \) values for the sources using the same strategy (Table 1).

A fairly tight far-infrared-radio correlation applies to a wide-range of galaxy types in the local Universe (see Solarz et al. 2019 and references therein). The relation also applies to radio-loud AGNs that are observed to have much lower values of \( q \) on average, compared to radio quiet or star-forming systems. The value of \( q \) can therefore help to distinguish between AGN activity and star-forming processes in extragalactic systems (e.g. Condon et al. 2002; Sargent et al. 2012; Delhaize et al. 2017).

Figure 2 presents a plot of \( q_{1.4 \text{GHz}} \) versus the redshift for the 76 SMG sources in Chapman et al. (2005), including the four sources observed by VLBI. Since we were mostly interested in seeing whether our sources deviated from the rest of the sample, a simple linear fit was applied to the observed band flux ratios without the application of any \( k \)-correction. The Chapman et al. (2005) sample has a mean \( q_{1.4 \text{GHz}} \) of \( \sim 1.85 \) with a standard deviation from the linear fit of 0.35. Three of our four sources appear to follow the FIR-radio correlation, which is represented by \( q_{1.4 \text{GHz}} \) with offsets that are smaller than 2\( \sigma \) from the linear fit.

It is interesting, however, that all three objects are below the correlation, thus, on average, they have larger radio luminosities than expected. However, for the detected source J123555.14, the \( q_{1.4 \text{GHz}} \) value increases to 2.42 \( \pm \) 0.15 when the contribution of the AGN core is subtracted from the 1.4 GHz radio luminosity. This value is actually higher than most of the sources in Fig. 2. Considering that the linear fit was performed on a sample of SMG, which is possibly contaminated by AGN, the correlation shown in Fig. 2 might be displaced towards lower \( q \) values with respect to that of a sample consisting purely of star-forming galaxies.

One of the sources, J131225.73, has a value of \( q_{1.4 \text{GHz}} \sim \) 0.74, which departs from the mean linear fit by \( \sim 3\sigma \). A further discussion on this point is given in Sect. 3.3.

While \( q_{1.4 \text{GHz}} \) and \( q_{850 \mu m} \) are more sensitive to the dust emission measurement, other \( q \)-values with wider IR waveband coverage, such as \( q_{IR} \) and \( q_{1} \), are more sensitive to the actual IR emission from the source. We note that \( q_{IR} \) is defined as the logarithmic ratio of the rest-frame 8–1000 \( \mu \)m flux and 1.4 GHz radio flux. Bell (2003) measured a median \( q_{IR} \) value of 2.64 \( \pm \) 0.02 over 164 SMGs, showing no signs of AGN. This value is similar to the medium \( q_{IR} \) value of 2.59 \( \pm \) 0.05 in Thomson et al. (2014) involving 76 SMGs. Del Moro et al. (2013) classified sources with \( q_{IR} < 1.68 \) as radio excess AGNs. Additionally, \( q_{1} \) is the logarithmic ratio of the far-IR and 1.4 GHz radio luminosity, where the far-IR luminosity is estimated by the flux in a wide band centred at 80 \( \mu \)m. Kovács et al. (2006) derived an average \( q_{1} \) value of 2.14 with an intrinsic spread of 0.12. Sources with a \( q_{1} \) value that is more than 2\( \sigma \) lower than the mean value are likely to be hosts of a radio-loud AGN. In Sect. 3.3, we use the \( q \) values of the sources from the literature to support our arguments. Both the \( q_{IR} \) value of J123555.14 and the \( q_{1} \) values of J131225.73 and J163650.43 are below the mean values mentioned above, which is in agreement with the \( q_{1.4 \text{GHz}} \) distribution characteristics of Fig. 2.

### 3.3. Source description

The various contour maps of our four targets are shown in Figs. 1–4. The crosses on the figures are centred on their expected a priori positions and scaled with respect to the associated error. In Fig. 1a, the white cross presents the VLA-measured coordinates of J123555.14 given in Richards (2000). The cross in Fig. 3a represents the expected position of J123600.10 measured in the re-processed eMERGE image. For the other two undetected sources in Fig. 4, the expected positions and errors...
**Fig. 3.** (a): EVN 1.6 GHz VLBI image of the field associated with J123600.10 (undetected). The white cross on the map presents the position of the source as measured in the re-processed eMERGE image. The map size corresponds to the 2σ uncertainties in the VLA position. This image was plotted with contour levels $-3$ and $3\times$ the rms noise level. (b): Hubble NICMOS NIC2 $F_{160}$ image of J123600 in which the source shows a disc-like structure with close projected companions. For clarity, the image was smoothed with a 2D Gaussian convolution with a FWHP of $\sim0.15''$. The overlaid contour was plotted at 3, 6, 12, and 24$\times$ the rms noise level of the eMERGE-JVLA moderate resolution image at 1.5 GHz. The red box shows the FoV of (a). As shown in (b), the radio and optical images have similar morphologies with reasonable position offsets of $\sim0.7''$. In each panel, the beam pattern of the contours is illustrated with a black ellipse.

**Fig. 4.** EVN 1.6 GHz images of (a) J131225.73 (undetected) and (b) J163650.43 (also undetected). The expected positions and errors in the maps represented by the crosses are provided in Fomalont et al. (2006) and Ivison et al. (2002) for (a) and (b), respectively. The maps are presented with contour levels $-3$ and $3\times$ the rms noise. The white crosses on the maps are centred on the expected positions of the sources. For clarity, the map sizes correspond to the 2σ and 1.5σ uncertainties in the VLA positions for (a) and (b), respectively. In each panel the VLBI beam pattern is illustrated with a black ellipse.
northern component was detected at the southern component, but we note that this is only detectable in our image at the 4-σ level, which is consistent with the source detected by Read et al. (2018). They did not detect the fainter southern component, but we note that this is only detectable in our image at the 4-σ level, which has a weak reliability. Therefore, we do not discuss this component in this paper. The brighter northern component was detected at the ~7-σ level with a brightness temperature of 5.2 ± 0.7 × 10^8 K.

The source shows a face-on disc-like morphology with a close companion (~10–20 pc in projection) in the rest-frame near-UV image (Fig. 1b). This image was taken by the Advanced Camera for Surveys (ACS) of the Hubble Space Telescope (HST) with the F814W filter in the Cosmic Assembly Near-infrared Deep Extragalactic Legacy Survey (CANDELS; Koekemoer et al. 2011) and was retrieved from the Hubble Legacy Archive. The HST image was processed by SAO Image DS9. For clarity, the image was smoothed with a 2D Gaussian convolution with a FWHM of ~0.1″.

The HST CANDELS image was contoured with the 1.5 GHz eMERGE DR1 maximum sensitivity map, where the source is detected with a peak and integrated flux density of 132.3 ± 2.1 μJy beam^-1 and 185.2 ± 4.6 μJy, respectively. These values are consistent with the VLBI flux within the errors. Moreover, this source has a total 1.4 GHz VLA flux density of 212.0 ± 13.7 μJy, which corresponds to a radio luminosity of 5.50 ± 0.4 × 10^{24} W Hz^-1. As shown in Fig. 5, this source has a prominent VLBI-to-VLA flux density ratio of 0.95, which is much higher than the mean value of ~0.6 measured by Read et al. (2017) over a larger sample of Very Long Baseline Array (VLBA) detected sources, which were observed with a similar, approximate milli-arcsecond resolution and an ~10 μJy sensitivity. This suggests that most of the radio emission is associated with a compact milli-arcsecond central core. Another possible explanation for a large value of S_{VLBI}/S_{VLA} is source variability. In particular, the VLBI and VLA data in this paper were taken at different epochs. Nevertheless, only a few percent of faint radio sources are expected to be variable and this particular source was not observed as being variable in the recent study of Radcliffe et al. (2019). The source has a SFR of 0.77 ± 0.06 × 10^8 M_⊙ yr^-1, which was not calculated with the total L_1.4 GHz. In removing the 95% VLBI recovered luminosity which is considered to be associated with an AGN, a corresponding SFR of 385 ± 50 M_⊙ yr^-1 is indicated by the 1.4 GHz radio luminosity contributed by star-forming processes (0.28 ± 0.02 × 10^{24} W Hz^-1).

The source appears to be slightly resolved in the eMERGE 1.5 and 5.5 GHz images, the latter has a peak brightness of ~5σ. The total flux densities of the eMERGE 1.5 and 5.5 GHz images yield a spectral index of ~−0.49. This relatively flat spectral index is indicative of an AGN component being the dominant source of radio emission. This source was also confirmed to be an AGN in Radcliffe et al. (2018). Its mid-IR spectrum also shows evidence of an AGN in this source (Pope et al. 2008). Hainline et al. (2011) suggest that an estimated fraction of 0.7–0.8 of its Spitzer-Infrared Array Camera (IRAC) 8 μm emission was contributed by a power-law component (f ∝ λ^α where α = 3 was the best fit), which is considered to likely originate from an AGN. By fitting local spectral energy distributions (SEDs) with 24 μm Spitzer photometry, Murphy et al. (2009) suggested that more than 50% of its 8–1000 μm total IR energy budget could be contributed by AGN activity. This source has a 580 μm and 350 μm of 5.4 ± 1.9 mJy and 23.1 ± 0.6 mJy; this yields a q_1.4GHz to q_0.4GHz of 1.41 ± 0.16 and 2.04 ± 0.03, respectively (see Table 1). Although the q_1.4GHz and q_0.4GHz of these source seem to follow the FIR-radio correlation, it has a qIR of 1.51 (Murphy et al. 2009). This value is ~4σ lower than the median qIR value of ~2.60 (Bell 2003; Thomson et al. 2014). According to the argument that Del Moro et al. (2013) raise that sources with qIR < 1.68 are likely radio excess AGNs, the low qIR value of this source is a strong indicator of an AGN. After subtracting the AGN-contributed radio luminosity, the source has a q_1.4GHz and q_0.4GHz values of 2.42 ± 0.15 and 3.05 ± 0.02,
respectively. These values may reflect the \( q \) values for purely star-forming systems. As shown in Fig. 2, the star-formation-associated value of \( q_{1.4\, \text{GHz}} \) is higher than the majority of the Chapman et al. (2005) SMG sample. The source also shows evidence of an AGN at X-ray wavelengths, as suggested by Alexander et al. (2005) and Del Moro et al. (2013).

\( J123600.10 \) was not detected in our VLBI observations with a 3-\( \sigma \) brightness temperature limit of \( <5.6 \times 10^5 \) \( K \) and it is located at a redshift of 2.710 (Chapman et al. 2005). This source is detectable in the eMERGE re-weighted image with an angular size of 859\( \times \)253 mas\(^2\), and the measured peak and integrated flux densities are 52.4 \pm 1.7 \mu Jy and 83.5 \pm 4.1 \mu Jy, respectively. This yields a brightness temperature of \( \sim 523 \) \( K \). Since the peak brightness measured in the eMERGE re-weighted image is only approximately five times the VLBI detecting threshold, it is not surprising that we did not detect this source with VLBI.

The contours of the eMERGE re-weighted image for this source are shown in Fig. 3b on top of the HST image taken with the Near Infrared Camera and Multi-Object Spectrometer (NICMOS) using the NIC2 camera, and the F160W filter in Swinbank et al. (2010) (project ID: 9506). The source shows a disc-like morphology with close companions. The Hubble NICMOS image was retrieved from the Hubble Legacy Archive and was processed by SAO Image DS9. For clarity, the image was smoothed with a 2D Gaussian convolution with a FWHP of \( \sim 0.15'' \).

This source has a \( S_{850\, \mu m} \) and \( S_{350\, \mu m} \) of 6.9 \pm 2.0 mJy and 57.2 \pm 0.6 mJy; this yields a \( q_{1.4\, \text{GHz}} \) and \( q_{350\, \mu m} \) of 1.42 \pm 0.13 and 2.34 \pm 0.03, respectively (Table 1). Additionally, both follow the FIR-radio correlation. The source has a low VLBI-to-VLA flux ratio upper limit of 0.16, suggesting that a significant fraction of the radio emission is extended and probably associated with star-formation processes. This is supported by other measurements – Chapman et al. (2003) suggest that this galaxy is possibly an edge-on merger from its optical morphology. However, this source has a total 1.4 GHz VLA flux density of 262.0 \pm 17.1 \mu Jy, which corresponds to a radio luminosity of \( 1.66 \pm 0.11 \times 10^{25} \) \( W\, Hz^{-1} \). This value can only be reached by an AGN or extremely strong star-forming galaxies with SFR of \( 2.32 \pm 0.15 \times 10^3 \, M_\odot \, yr^{-1} \). The high luminosity indicates that extended jet emission from a jet associated with quasar activity may exist in this source (Jarvis et al. 2019; Muxlow et al. 2005), which was undetected in these VLBI observations.

\( J131225.73 \) was not detected in our VLBI observations with a 3-\( \sigma \) brightness temperature limit of \( <3.5 \times 10^5 \) \( K \) and it is located at a redshift of 2.378 (Chapman et al. 2005). This source has a \( S_{850\, \mu m} \) and \( S_{350\, \mu m} \) of 8.2 \pm 1.7 mJy and 45.9 \pm 2.9 mJy; this yields a \( q_{1.4\, \text{GHz}} \) and \( q_{350\, \mu m} \) of 1.57 \pm 0.10 and 2.32 \pm 0.04, respectively (see Table 1). Additionally, both seem to follow the FIR-radio correlation. The source recovers \( <22\% \) of its 1.4 GHz flux; the non-detection on VLBI scales may suggest that a significant fraction of the radio emission is associated with star-formation processes. This is supported by other measurements – Engel et al. (2010) classify this source as a close binary galaxy merger because the CO(3–2) and CO(7–6) data show two peaks separated by \( \sim 3\, kpc \). A \( q_1 \) value of 1.75 \pm 0.17 for this source was measured by Kovács et al. (2006), this value is slightly lower than the 2-\( \sigma \) constraint from the mean value of 2.14, and it was considered to be consistent with the far-infrared to radio correlation.

However, \( J163650.43 \) was not detected in our VLBI observations with an implied 3-\( \sigma \) brightness temperature limit of \( <3.5 \times 10^5 \) \( K \) and it is located at a redshift of 2.378 (Chapman et al. 2005). This source has a \( S_{850\, \mu m} \) and \( S_{350\, \mu m} \) of 6.5 \pm 2.5 mJy and 22% of its 1.4 GHz flux; the non-detection on VLBI scales may suggest that a significant fraction of the radio emission is associated with star-formation processes. This is supported by other measurements – Engel et al. (2010) classify this source as a close binary galaxy merger because the CO(3–2) and CO(7–6) data show two peaks separated by \( \sim 3\, kpc \). A \( q_1 \) value of 1.75 \pm 0.17 for this source was measured by Kovács et al. (2006), this value is slightly lower than the 2-\( \sigma \) constraint from the mean value of 2.14, and it was considered to be consistent with the far-infrared to radio correlation.

4. Summary

We have conducted EVN 1.6 GHz observations of four SMGs (J123555.14, J123600.10, J131225.73, and J163650.43). Out of the four targets, we detected J123555.14 once, which is a source located in the GOODS-N field with a brightness temperature of 5.2 \pm 0.7 \times 10^5 \, K. This value exceeds the maximum brightness temperature of \( \sim 10^5 \, K \) that a star-forming galaxy is expected to reach. We therefore suggest that the radio emission associated with J123555.14 largely arises from AGN processes. We also derived brightness temperature upper limits of the three
undetected sources. The non-detections may suggest that most of their radio emission is powered by star-forming processes or extended radio jets. Notably, the three undetected sources all have relatively high 1.4 GHz radio luminosities and/or show evidence of an AGN in other wavebands. This fact may suggest that while SGMs are potential hosts of AGN, star-forming processes and AGN activity probably exist in such systems concordantly.

In particular, at least one of the sources, J131225.73, shows multi-wavelength properties that would lead us to expect it to be detected by VLBI. It is highly luminous in the \( 752.5 \pm 2.4 \mu Jy \) or \( 1.23 \pm 0.01 \times 10^{25} \, \text{W} \, \text{Hz}^{-1} \) at 1.4 GHz, and has an exceptionally low value of \( q \), and is unresolved in both the 1.4 and 8.4 GHz VLA observations. We suggest that this source might be associated with an AGN that is dominated by extended jet emission. It would be interesting to observe the structure of this radio source with intermediate resolution, such as e-MERLIN at 5 GHz, in order to better understand the nature of the source. Since the morphology of the faint radio source population is still largely unknown on these scales, it is possible that VLBI misses (or resolves) many extended radio AGN of this type. As illustrated by the upper limits of the VLBI-to-VLA flux ratio of the undetected sources, a bright radio source (~200 \( \mu Jy \) at 1.4 GHz) with less than ~20% of its radio flux contributed by an AGN would probably be missed by VLBI observations at a 1-\( \sigma \) sensitivity of ~10\( \mu Jy \). Surveys with multiple resolution (e.g. VLA, e-MERLIN, and VLBI) are therefore needed to determine what fraction of the extragalactic objects have similar properties and thus would probably be missed by VLBI.

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