

The ALMA-PILS survey: first detection of methyl isocyanide (CH₃NC) in a solar-type protostar

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ABSTRACT

Context. Methyl isocyanide (CH₃NC) is the isocyanide with the largest number of atoms confirmed in the interstellar medium (ISM), but it is not an abundant molecule, having only been detected towards a handful of objects. Conversely, its isomer, methyl cyanide (CH₃CN), is one of the most abundant complex organic molecules detected in the ISM, with detections in a variety of low- and high-mass sources.

Aims. The aims of this work are to determine the abundances of methyl isocyanide in the solar-type protostellar binary IRAS 16293–2422 and to understand the stark abundance differences observed between methyl isocyanide and methyl cyanide in the ISM.

Methods. We use Atacama Large Millimeter/submillimeter Array (ALMA) observations from the Protostellar Interferometric Line Survey (PILS) to search for methyl isocyanide and compare its abundance with that of its isomer methyl cyanide. We use a new line catalogue from the Cologne Database for Molecular Spectroscopy (CDMS) to identify methyl isocyanide lines. We also model the chemistry with an updated version of the three-phase chemical kinetics model MAGICKAL, presenting the first chemical modelling of methyl isocyanide to date.

Results. We detect methyl isocyanide for the first time in a solar-type protostar, IRAS 16293–2422 B, and present upper limits for its companion protostar, IRAS 16293–2422 A. Methyl isocyanide is found to be at least 20 times more abundant in source B compared to source A, with a CH₃CN/CH₃NC abundance ratio of 200 in IRAS 16293–2422 B and >5517 in IRAS 16293–2422 A. We also present the results of a chemical model of methyl isocyanide chemistry in both sources, and discuss the implications for methyl isocyanide formation mechanisms and the relative evolutionary stages of both sources. The chemical modelling is unable to match the observed CH₃CN/CH₃NC abundance ratio towards the B source at densities representative of that source. The modelling, however, is consistent with the upper limits for the A source. There are many uncertainties in the formation and destruction pathways of methyl isocyanide, and it is therefore not surprising that the initial modelling attempts do not reproduce observations. In particular, it is clear that some destruction mechanism of methyl isocyanide that does not destroy methyl cyanide is needed. Furthermore, these initial model results suggest that the final density plays a key role in setting the abundance ratio. The next steps are therefore to obtain further detections of methyl isocyanide in more objects, as well as undertaking more detailed physico-chemical modelling of sources such as IRAS 16293.

Key words. astrochemistry – stars: formation – stars: protostars – ISM: molecules – ISM: individual objects: IRAS 16293-2422

1. Introduction

Methyl isocyanide (CH₃NC) is the isocyanide with the largest number of atoms confirmed in the interstellar medium (ISM), however, its detections are few and far between. It was first tentatively detected by Cernicharo et al. (1988) in the Sgr B2 cloud, and confirmed with additional transitions by Remijan et al. (2005). Since then, it has only been confirmed in a handful of sources, including the Horsehead photo-dissociation region (PDR) by Gratier et al. (2013), and Orion KL by López et al.

(2014). Conversely, its isomer methyl cyanide (CH₃CN) is one of the most abundant complex organic molecules detected in the ISM, with detections in a variety of low- and high-mass sources (e.g. Beltrán et al. 2005; Bisschop et al. 2008; Zapata et al. 2010).

Such stark differences in the abundances of two isomers raise interesting questions about their chemistry in the ISM, and in particular, whether their abundance differences are a consequence of their formation pathways, a result of differential destruction with chemical evolution, or even other processes. A number of authors have explored this with both theoretical

and laboratory work. Theoretical calculations by DeFrees et al. (1985) of the $\text{CH}_3\text{NC}/\text{CH}_3\text{CN}$ ratio in dense interstellar clouds explored the ratio that is inherited from their protonated precursor ions, protonated methyl cyanide (CH_3CNH^+), and protonated methyl isocyanide (CH_3NCH^+). These ions are formed when CH_3^+ reacts with HCN, forming a complex that can rapidly equilibrate between the two isomeric forms. As the complex relaxes to form a stable molecule, isomerisation occurs until the internal energy of the ion is no longer sufficient to overcome the isomerisation barrier. Ion-electron recombination reactions then lead to the formation of methyl isocyanide and methyl cyanide with methyl cyanide being formed in higher amounts than methyl isocyanide. Methyl cyanide is preferred because it is the more stable of the two isomers.

Laboratory experiments by Hudson & Moore (2004) looked at the effect of non-thermal processes on the formation of isocyanides. They found that bombarding pure ice samples of methyl cyanide with protons and ultraviolet (UV) photons produced methyl isocyanide, however, the introduction of water-rich ices inhibits the process, with OCN^- being formed instead. Mencos & Krim (2016) also used laboratory experiments to show that methyl isocyanide could be formed through the reaction between methyl cyanide and N atoms in ice, triggering isomerisation that produces methyl isocyanide and ketenimine (CH_2CNH) as principal products, with methyl isocyanide being formed in higher amounts.

As well as formation and destruction mechanisms, differences in adsorption properties could be crucial to understanding the observed differences in abundance between methyl cyanide and its isocyanide. Bertin et al. (2017a,b) used temperature programmed desorption (TPD) experiments to study the adsorption of methyl cyanide and methyl isocyanide on interstellar grain surfaces. They found that for amorphous water ice a small amount of the adsorbed molecules desorb through a volcano effect, meaning that some diffusion takes place during the warming up and some of the methyl cyanide and methyl isocyanide molecules get trapped in the corrugation of the amorphous water bulk. The adsorption energies of methyl cyanide were found to be slightly higher than those of methyl isocyanide by 20–40 meV, both experimentally and theoretically. They attribute this to the slightly stronger ability of methyl cyanide to form *H*-bonds because of a local excess of positive charge on the hydrogen atoms of the methyl group. On graphite surfaces, the adsorption energy of each molecule was found to be very sensitive to the structural defects of aromatic carbonaceous surfaces, increasing the average adsorption energy by more than 50% compared to a perfect graphene plane. They again found that the adsorption energy was higher for methyl cyanide than its isomer, which they predict would lead to a gas-phase enrichment of the isonitrile at a given grain temperature compared with the abundance ratio of the two isomers in the condensed phase.

What is lacking in the study of methyl isocyanide in the ISM is its detection in more objects, to provide statistically significant abundance ratios and chemical modelling of its formation and destruction pathways in different environments. To achieve the first point, high sensitivity observations are needed as it has shown very weak emission in the objects it has been detected in so far. In this work, we therefore utilise high sensitivity Atacama Large Millimeter/submillimeter Array (ALMA) observations from the protostellar interferometric line survey (PILS) to present the first detection of methyl isocyanide in a solar-type protostar, IRAS 16293–2422 (hereafter IRAS 16293). IRAS 16293 is a Class 0 solar-type protostellar binary, located only 141 pc away (Dzib et al. 2018), in the Ophiuchi star-forming

complex. It consists of at least two sources, A and B, which are separated by $5''$ (705 AU), and each shows a rich spectrum of complex organic molecules (Bisschop et al. 2008; Jørgensen et al. 2016; Lykke et al. 2017). We present details of the observations in Sect. 2, the results of an abundance analysis and chemical modelling in Sect. 3, and a discussion of the astronomical implications in Sect. 4. Finally, the conclusions of this work are presented in Sect. 5.

2. Observations

The observations presented here were taken as part of the PILS programme, an ALMA observing programme to study IRAS 16293 in Band 7, between 329.147 GHz and 362.896 GHz. The phase centre of the observations is located between the two components of the binary system at $\alpha_{J2000} = 16^{\text{h}}32^{\text{m}}22^{\text{s}}.72$, $\delta_{J2000} = -24^{\circ}28'34''.3$. The data are a combination of the main 12 m array and Atacama compact array (ACA) observations, which have been combined to have a restoring beam of $0''.5$ at a spectral resolution of 0.2 km s^{-1} . They reach a sensitivity of about $7\text{--}10 \text{ mJy beam}^{-1} \text{ channel}^{-1}$, that is, approximately $\sim 5 \text{ mJy beam}^{-1} \text{ km s}^{-1}$ across the entire frequency range. Further details of the data reduction and continuum subtraction procedure can be found in Jørgensen et al. (2016).

In this work we analyse spectra extracted from two positions, where a large number of complex organic molecules have been previously detected and the line emission is not obscured by dust. The first position at $\alpha_{J2000} = 16^{\text{h}}32^{\text{m}}22^{\text{s}}.58$, $\delta_{J2000} = -24^{\circ}28'32''.8$, is offset from the continuum peak position of IRAS 16293B in the southwest direction by one beam diameter ($0''.5$, 70 AU). This position is used as the lines are particularly bright, do not have strong absorption features, and do not suffer from high continuum optical depth (Coutens et al. 2016, Lykke et al. 2017). The full width half maximum (FWHM) of lines around source A are much broader ($\sim 2\text{--}3 \text{ km s}^{-1}$) than around source B ($\sim 1 \text{ km s}^{-1}$) making line blending a significant problem for identification and analysis of molecular emission. We, therefore, choose a position at $\alpha_{J2000} = 16^{\text{h}}32^{\text{m}}22^{\text{s}}.90$, $\delta_{J2000} = -24^{\circ}28'36''.2$, that is $0''.6$ (85 AU) northeast from the peak continuum position of IRAS 16293A. This position has been used previously to study the nitrile chemistry in IRAS 16293 (Calcutt et al. 2018), the oxygen-bearing molecules (Manigand et al. 2018), and methyl isocyanate emission (Ligterink et al. 2017).

3. Results and analysis

Both offset positions towards A and B were searched for methyl isocyanide. Isotopologues were not searched for as the line lists are not available. Methyl isocyanide is detected towards IRAS 16293B, with ten unblended lines and six blended lines above 3σ present in the spectra. All of the lines that are predicted by a local thermodynamic equilibrium (LTE) spectral model are detected in the spectra. A list of the detected transitions, frequencies, upper energy levels, Einstein coefficients, and whether the transitions are blended are given in Table 1. This is the first detection of this molecule in a solar-type protostar. No lines are detected towards the offset position of IRAS 16293A. We additionally searched positions closer to the A peak continuum position, however, no lines could be identified as blending became too severe.

3.1. Spatial extent

Figure 1 shows emission maps for methyl cyanide and methyl isocyanide in IRAS 16293. In the left panel, a velocity-corrected

Table 1. Spectroscopic information for lines of CH₃NC detected in IRAS 16293B.

Transition	Frequency (MHz)	E_u (K)	A_{ij} (s ⁻¹)	Blended ^a
17 7–16 7	341 329.6	493	2.83×10^{-3}	U
17 6–16 6	341 429.4	402	2.98×10^{-3}	B
17 5–16 5	341 513.8	324	3.11×10^{-3}	B
17 4–16 4	341 582.9	261	3.22×10^{-3}	U
17 3–16 3	341 636.7	211	3.31×10^{-3}	B
17 2–16 2	341 675.1	176	3.37×10^{-3}	B
17 1–16 1	341 698.2	155	3.40×10^{-3}	U
17 0–16 0	341 705.9	148	3.41×10^{-3}	U
18 7–17 7	361 396.3	511	3.43×10^{-3}	U
18 6–17 6	361 501.8	419	3.60×10^{-3}	B
18 5–17 5	361 591.2	341	3.74×10^{-3}	U
18 4–17 4	361 664.3	278	3.85×10^{-3}	B
18 3–17 3	361 721.2	229	3.94×10^{-3}	U
18 2–17 2	361 761.9	193	4.01×10^{-3}	U
18 1–17 1	361 786.3	172	4.04×10^{-3}	U
18 0–17 0	361 794.4	165	4.06×10^{-3}	U

Notes. ^aU denotes an unblended line and B denotes a blended line.

integrated emission (VINE) map shows the emission from methyl cyanide (the $18_{\pm 4} - 17_{\pm 4} v_8 = 1 I = +1$ transitions, which are unresolved). VINE maps are a method for determining the spatial scale of molecules in spectrally dense regions, with large velocity gradients. They remove the problem of contamination by nearby lines in the emission map. More details of how a VINE map is created and the caveats with this method can be found in Calcutt et al. (2018). The right panel of Fig. 1 shows a zoom in of the methyl cyanide emission and an integrated emission map of methyl isocyanide (the $18_1 - 17_1 v = 0$ line). The emission traces the disk component of IRAS 16293B and shows a similar morphology and extent to methyl cyanide emission. The other lines of methyl isocyanide also show the same extent and morphology.

3.2. Column densities and excitation temperatures

To determine abundances and excitation temperatures, the spectral modelling software CASSIS¹ was used to calculate synthetic spectra and determine best-fit spectral models to the observations assuming LTE. A large grid of models was run with column densities between 1×10^{13} and 9×10^{14} cm⁻² and excitation temperatures between 80 and 300 K. Line widths were varied between 0.8 and 1.2 km s⁻¹, which are the typical line widths found in other PILS studies (e.g. Jørgensen et al. 2016; Ligterink et al. 2017). The peak velocities were varied between 2.5 and 2.9 km s⁻¹, corresponding to one channel on either side of the V_{LSR} (2.7 km s⁻¹). The results of these models were compared to the spectra of the most optically thin lines to determine the best-fit model by computing the minimum χ^2 .

The brightness temperature of the line emission is calculated by CASSIS according to

$$T_B(\nu) = T_0 \left(\frac{1}{e^{T_0/T_{ex}} - 1} - \frac{1}{e^{T_0/T_{bg}} - 1} \right) (1 - e^{-\tau_\nu}), \quad (1)$$

¹ CASSIS has been developed by IRAP-UPS/CNRS: <http://cassis.irap.omp.eu/>

where $T_0 = h\nu/k$. A derivation of this formula using the equation of radiative transfer can be found in Vastel (2016). Typically, the background temperature (T_{bg}) is the cosmic microwave background (CMB) of 2.73 K. Both sources in IRAS 16293, however, exhibit bright continuum emission from the dense dust in each source. This means a significant background radiation temperature should be considered when calculating the brightness temperature of the line emission. This can be applied to the column densities as a correction factor. The correction factor is 1.1 at 150 K towards IRAS 16293B.

The initial modelling of the methyl isocyanide emission lines was performed using the jet propulsion laboratory (JPL) catalog (Pickett et al. 1998) entry. Lines were assigned tentatively, but the predictions had large uncertainties (more than 0.5 MHz or 0.4 km s⁻¹), deviated to lower frequencies by almost one line width at $K = 0$, rapidly increasing to higher K . The JPL entry was based on Bauer & Bogey (1970) with an estimate of the rotational parameter $A - B$. These early laboratory data, however, extend only to 141 GHz with $J = 7-6$ and $K \leq 6$. A recent study provides additional rotational data up to 643 GHz supplemented by ground-state combination differences from the ν_4 infrared (IR) spectrum (Pracna et al. 2011). The rotational parameter $A - B$ and the centrifugal distortion parameter D_K , necessary for describing the K level spacing, were taken from an earlier IR study (Pliva et al. 1995). The resulting predictions will be made available in the Cologne Database for Molecular Spectroscopy² (CDMS; Endres et al. 2016). Vibrational contributions to the partition function at 150 K were estimated from ν_8 and its overtones from that earlier IR study as 1.17.

An error analysis was also carried out for the column densities, excitation temperatures, and abundances determined in this work. The flux calibration accuracy of these data are better than 5% (Jørgensen et al. 2016). A larger source of error comes from the quality of the fit of the spectral model to the data. We, therefore, determine the errors by determining the minimum change to each quantity, which produces a significant difference to the goodness of fit.

Table 2 shows the best-fit excitation temperature and column density for methyl isocyanide in IRAS 16293B and the upper limits for IRAS 16293A, as well as the CH₃CN/CH₃NC abundance ratio using the methyl cyanide abundances determined in Calcutt et al. (2018). Figure 2 shows the LTE model overlaid on the data for all detected transitions in IRAS 16293B overlaid with the best-fit LTE model. The upper limits are determined for the brightest methyl isocyanide line in the frequency range of the observations (the $17_3 - 16_3$ line at 341.637 GHz). We use $1.05 \times 3 \times RMS \times \sqrt{\Delta V \times FWHM}$ to compute the 3σ limit, where 1.05 is a factor to account for a 5% flux calibration uncertainty. We assume a line width of 2.2 km s⁻¹ based on the line width found for nitriles towards this source by Calcutt et al. (2018), and an excitation temperature of 150 K. The CH₃CN/CH₃NC abundance ratio in IRAS 16293B is 200, which is much lower than that found in IRAS 16293A (>5517).

3.3. Chemical modelling

In order to investigate the chemistry of CH₃NC, we use an updated version of the three-phase chemical kinetics model MAGICKAL (Garrod 2013). This model has been updated with the back-diffusion correction of Willis & Garrod (2017). The chemical network is based on that of Belloche et al. (2017), using gas-phase, grain-surface, and bulk ice chemistry, with formation

² <http://www.astro.uni-koeln.de/cdms>

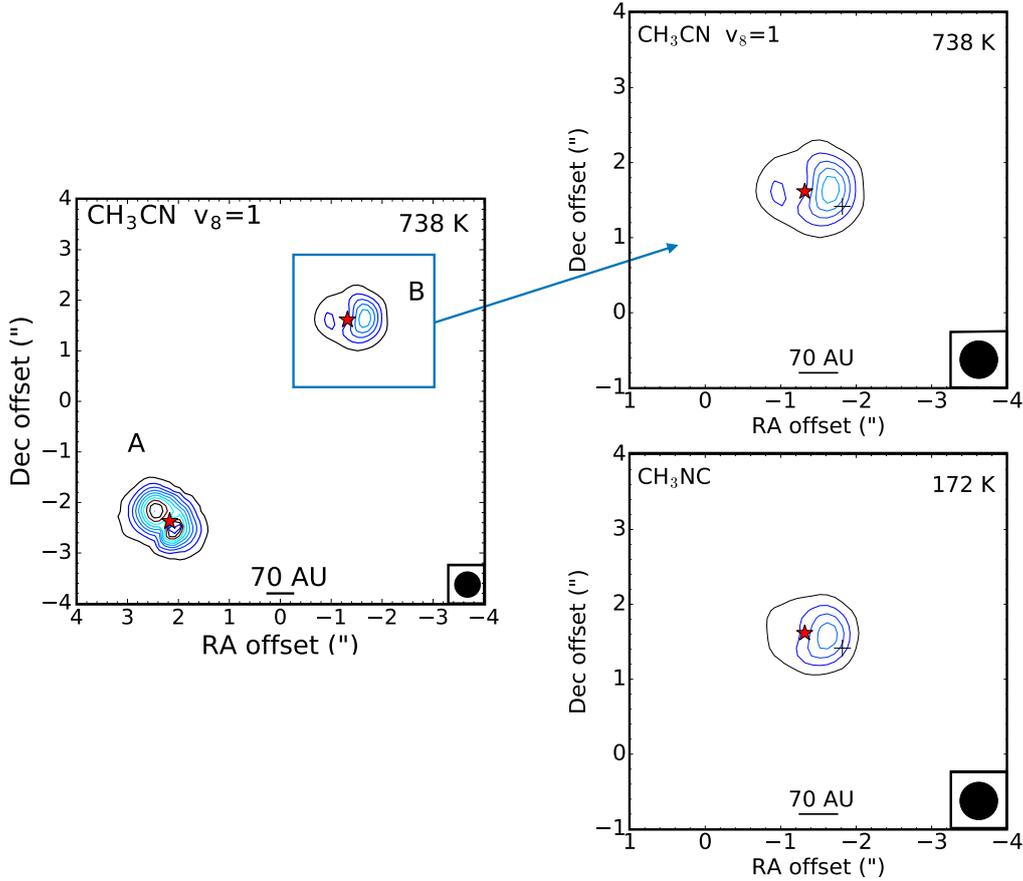


Fig. 1. *Left panel:* velocity-corrected integrated emission (VINE) map of the $18_{\pm 4} - 17_{\mp 4} v_8 = 1 I = +1$ transitions (332.029 GHz) of methyl cyanide, which are unresolved. *Right panel:* zoom in of the methyl cyanide emission towards IRAS 16293B and integrated emission map of the $18_1 - 17_1 v = 0$ line (361.786 GHz) of methyl isocyanide. For IRAS 16293B, six channels (1.2 km s^{-1}) centred on the V_{LSR} are integrated over. For IRAS 16293A, 16 channels (3.2 km s^{-1}) centred on the V_{peak} of each pixel are integrated over. The axes show the position offset from the phase centre of the observations. Contour levels start at 30 mJy km s^{-1} and increase in steps of 50 mJy km s^{-1} . The red star marks the peak continuum position in each source and the black cross marks the offset position where the spectra analysed in this work are extracted from. The RA and Dec offsets are relative to the phase centre of the observations.

Table 2. Excitation temperature (T_{ex}), column density (N_{tot}), and the abundance ratio, $N(\text{CH}_3\text{CN})/N(\text{CH}_3\text{NC})$, in IRAS 16293B and IRAS 16293A.

Source	T_{ex} (K)	N_{tot} (cm^{-2})	$\frac{N(\text{CH}_3\text{CN})^c}{N(\text{CH}_3\text{NC})}$
IRAS 16293B	150 ± 20	$2.0 \pm 0.2 \times 10^{14}$	200
IRAS 16293A	150^a	$<1.45 \times 10^{13b}$	>5517

Notes. All models assume LTE and a source size of $0''.5$. The FWHM of IRAS 16293B lines is $1 \pm 0.2 \text{ km s}^{-1}$ and the peak velocity is $2.7 \pm 0.2 \text{ km s}^{-1}$, based on a fit of the data. ^(a)The excitation for IRAS 16293A is adopted to be the same as IRAS 16293B and similar to other molecules detected towards this source. ^(b)Determined for the 3σ upper limit assuming an FWHM of 2.2 km s^{-1} , based on the FWHM of other molecules detected in this source in Calcutt et al. (2018). ^(c)Determined using CH_3CN column densities from Calcutt et al. (2018).

and destruction mechanisms for CH_3NC added. This is, to the authors' knowledge, the first time that CH_3NC has been incorporated into an astrochemical kinetics model. This network is a preface to a more complete chemical treatment for isocyanides (Willis et al., in prep.). The two-stage physical model used is also similar to that of Belloche et al. (2017), in which a cold collapse is followed by a static warm-up to 400 K. The cold collapse in

the model has an isothermal gas temperature (10 K), and the dust temperature cools from an initial value of 16 K to a final value of 8 K. The chemical model is a single-point model, thus it has a uniform density. Following Coutens et al. (2018), we run two models, using two different final densities for the collapse phase: $n_{\text{H}} = 6 \times 10^{10} \text{ cm}^{-3}$, corresponding to the continuum peak of IRAS 16293B and $n_{\text{H}} = 1.6 \times 10^7 \text{ cm}^{-3}$, corresponding to the density of the filament between IRAS 16293A and IRAS 16293B (Jacobsen et al. 2018). The subsequent warm-up phase starts at a dust temperature of 8 K, and reaches a final temperature of 400 K at $2.8 \times 10^5 \text{ yr}$. This timescale is used to represent an intermediate-timescale warm-up, where $2 \times 10^5 \text{ yr}$ is the time spent to reach a dust temperature of 200 K. This is taken from Garrod & Herbst (2006). The model presented here follows Garrod (2013) in using a warm-up to 400 K instead of 200 K, by extending the “intermediate” temperature function beyond 200 K.

The primary formation mechanism of CH_3NC in our model follows DeFrees et al. (1985). First, the radiative association of CH_3^+ and HCN forms CH_3CNH^+ . This intermediate is formed with enough energy that a certain percentage of these molecules can isomerise to CH_3NCH^+ ; we use the suggestion by DeFrees et al. (1985) of 15%. The CH_3NCH^+ then recombines with free electrons, forming CH_3NC and H as the primary channel ($\sim 65\%$ of recombinations), with the remainder of CH_3NCH^+ molecules recombining to HCN and CH_3 . The binding energies on water

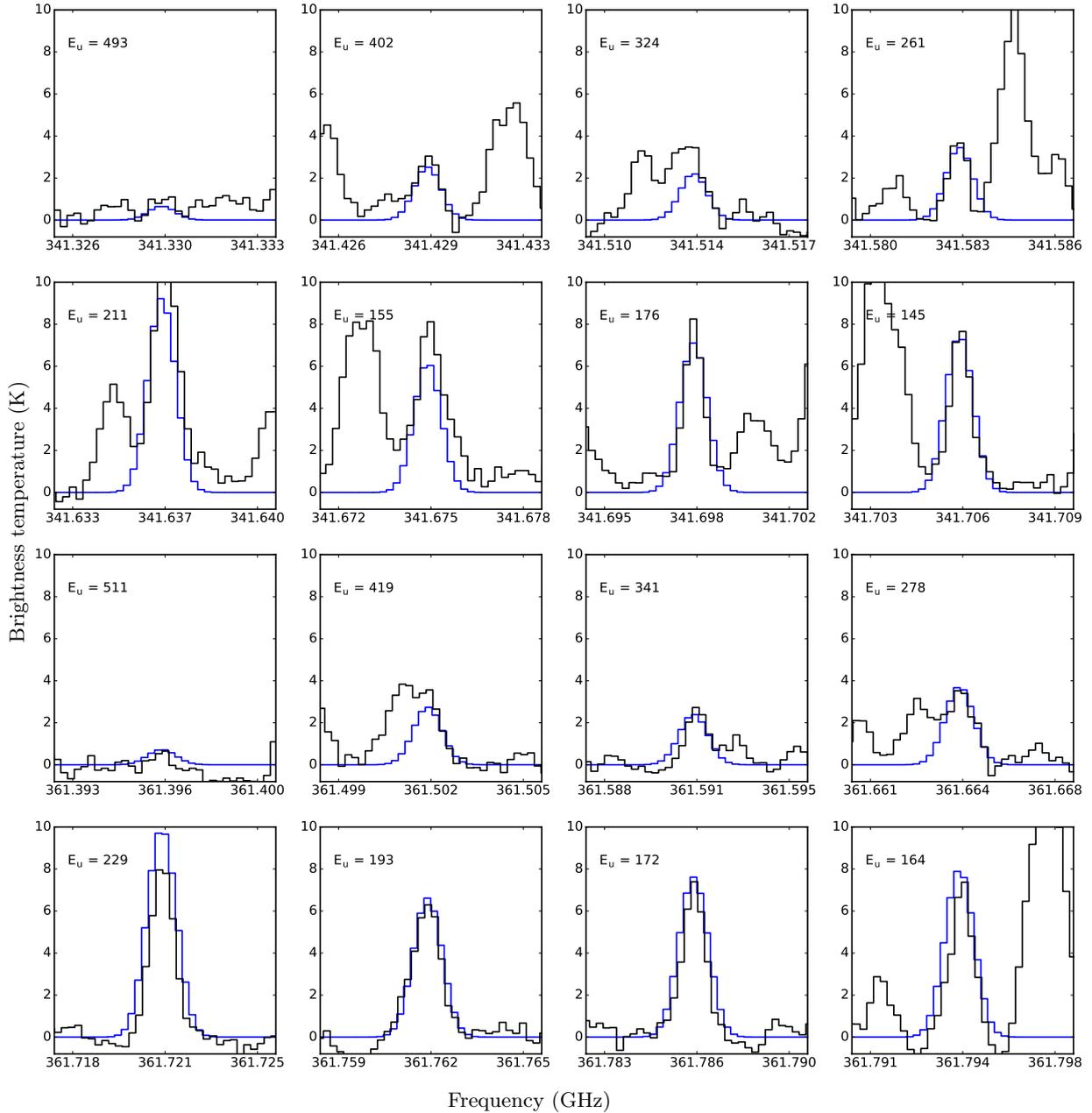


Fig. 2. Lines of CH₃NC detected in IRAS 16293B (black) overlaid with an LTE spectral model (blue). The upper energy level of each line is given in the top left corner of each plot in K.

ice for methyl cyanide and methyl isocyanide were updated in the network using new values from [Bertin et al. \(2017a\)](#); 6150 K and 5686 K, respectively). The primary destruction mechanism for CH₃NC on the grain surface is reaction with H, forming HCN and the radical CH₃. This reaction is assumed to have a barrier of 1200 K, based on recent calculations of the barrier for H + HNC ([Graninger et al. 2014](#)). The same reaction is also allowed to occur in the gas phase, assuming the same barrier.

Since we are comparing CH₃NC abundances to values for CH₃CN, it is important to review how CH₃CN is formed in our network as well. CH₃CN is formed primarily via the hydrogenation of the CH₂CN radical on the grain surface and in the ice mantle, with a small contribution from the reaction of CH₂CN with HCO. There is also a viable, less efficient gas-phase pathway to production of CH₃CN, previously discussed in regards to CH₃NC. The CH₃CNH⁺ molecules that do not

isomerise to CH₃NCH⁺ can recombine with electrons to form CH₃CN and H as the primary products, with some instances of this channel leading to HNC and CH₃.

The basic chemical network we construct here for CH₃NC chemistry is intended to incorporate the most obvious formation and destruction mechanisms, in keeping with the larger network. It is nevertheless possible that alternative mechanisms exist, while the efficiencies of those included here cannot be reliably known without careful experimental studies. In particular, the reaction of H+CH₃NC, which is included in both the gas-phase and the grain-surface networks, is quite speculative. The mechanics of the reaction with H and HNC are fundamentally different from this process, but the lack of experimental work makes it difficult to produce a better prediction. In particular, there is a large uncertainty on the barrier of the CH₃NC+H reaction as it likely does not equal the barrier on

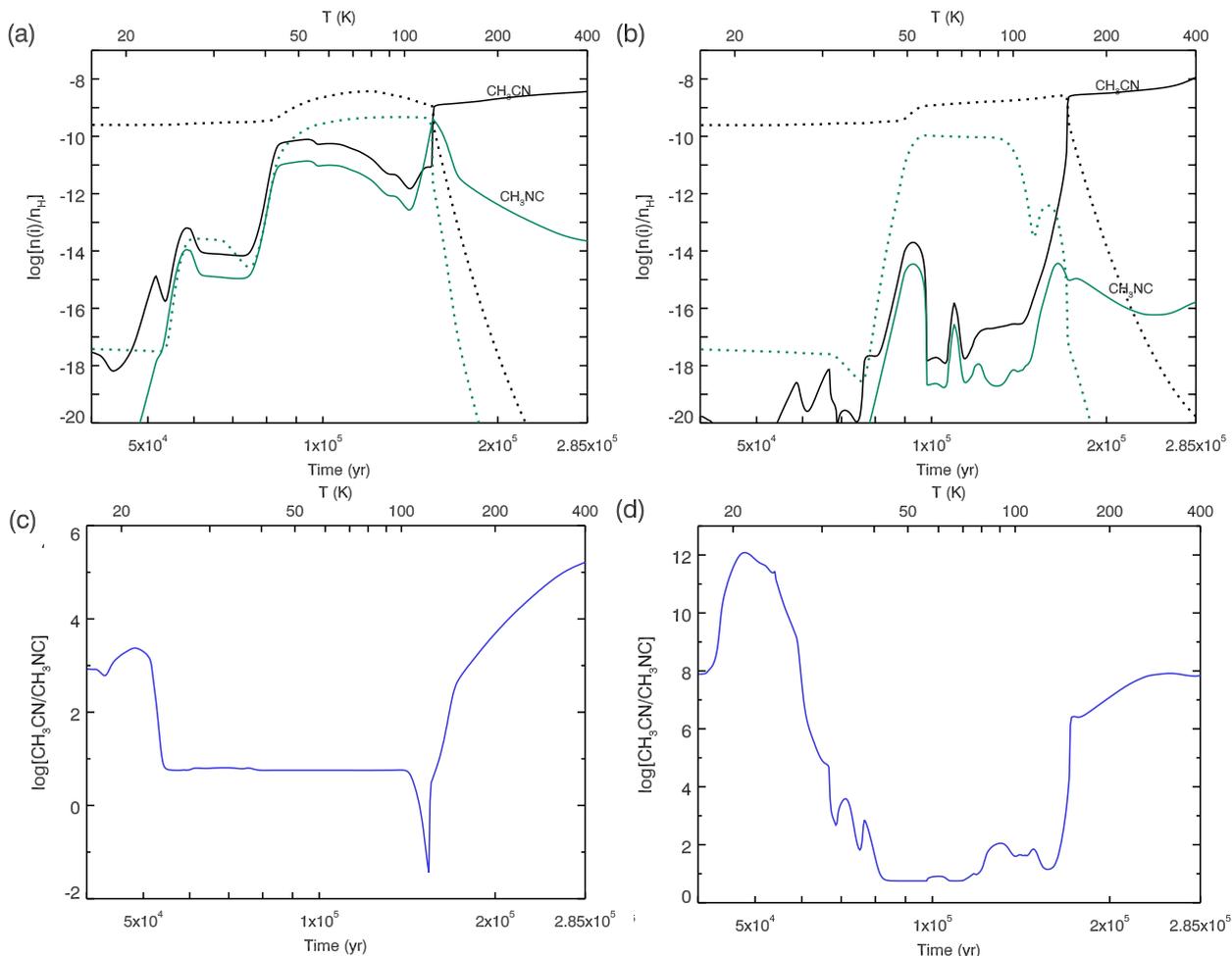


Fig. 3. Chemical model abundances of CH_3CN (black) and CH_3NC (green) for the warm-up stage of a hot core type model. *Panel a:* a model with a final collapse density of $n_{\text{H}} = 1.6 \times 10^7 \text{ cm}^{-3}$; *panel b:* a model with a final collapse density of $n_{\text{H}} = 6 \times 10^{10} \text{ cm}^{-3}$. Solid lines denote gas-phase abundances, while dashed lines indicate grain-surface abundances. *Panel c:* $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ ratio (blue) for the $n_{\text{H}} = 1.6 \times 10^7 \text{ cm}^{-3}$ model; *panel d:* $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ ratio (blue) for the $n_{\text{H}} = 6 \times 10^{10} \text{ cm}^{-3}$ model.

the $\text{HNC}+\text{H}$ reaction. We hope to continue to refine this network for CH_3NC and other isocyanides (Willis et al., in prep.), but this result highlights the need for more experimental work on these molecules. The results of the models are shown in Fig. 3, with panel (a) corresponding to $n_{\text{H}} = 1.6 \times 10^7 \text{ cm}^{-3}$, panel (b) to $n_{\text{H}} = 6 \times 10^{10} \text{ cm}^{-3}$, and panels (c) and (d) showing the $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ ratio for each of the models. The different densities exhibit different behaviour, particularly with regards to CH_3NC . The peak abundance of CH_3NC in the low-density model is $\sim 4 \times 10^{-10}$, while the high-density model produces a much lower peak abundance, $\sim 4 \times 10^{-15}$. This is due to the increased efficiency of grain-surface destruction of CH_3NC via reaction with H in the high-density model. As a result of this, as well as a higher peak abundance of CH_3CN in the high-density model, the $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ ratios at 150 K (the approximate excitation temperature of CH_3NC) vary greatly between the two models. In the low-density case, the ratio at 150 K is ~ 450 , which is the same order as the observed ratio in IRAS 16293B. The ratio at the same point in the high-density model is $\sim 2.5 \times 10^6$, which is consistent with the upper limits for IRAS 16293A.

We also ran several models incorporating the radiative association of CH_3^+ and HNC to test the effect that this proposed reaction has on the abundances of CH_3CN and CH_3NC . This process was given the same rate as the equivalent association of HCN with CH_3^+ , and results in CH_3NCH^+ . This ion is

then assumed to have enough internal energy to isomerise as in the HCN reaction. Several different branching ratios were tested, including those that lead predominantly to CH_3CN and CH_3NC . However, it was observed that the inclusion of this process did not affect the abundance of CH_3CN in any significant way, and had only a small effect on the abundance of CH_3NC in the low-density model, increasing the peak abundance to a maximum of $\sim 9 \times 10^{-10}$ when the reaction favours CH_3NC production.

Additionally, we have explored disabling the $\text{H}+\text{CH}_3\text{NC}$ grain surface reaction and gas-phase reaction to understand how important this speculative reaction is on the $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ ratio. Disabling the grain surface reaction had practically no effect, however, disabling the gas-phase reaction had a significant impact at higher temperatures (Fig. 4). CH_3NC is not destroyed efficiently in the gas-phase and this keeps the $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ ratio much closer to unity. It is in fact quite close to the 85%/15% split seen in the formation pathways. This result shows that there needs to be some efficient mechanism for the destruction of CH_3NC in the gas-phase that does not operate on CH_3CN . The $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ ratio varies greatly in the temperature range of 130–170 K. This is due to multiple factors. First, both species come off of the grain surface in this window, with CH_3CN coming off slightly later due to its somewhat higher binding energy. Once both species are off

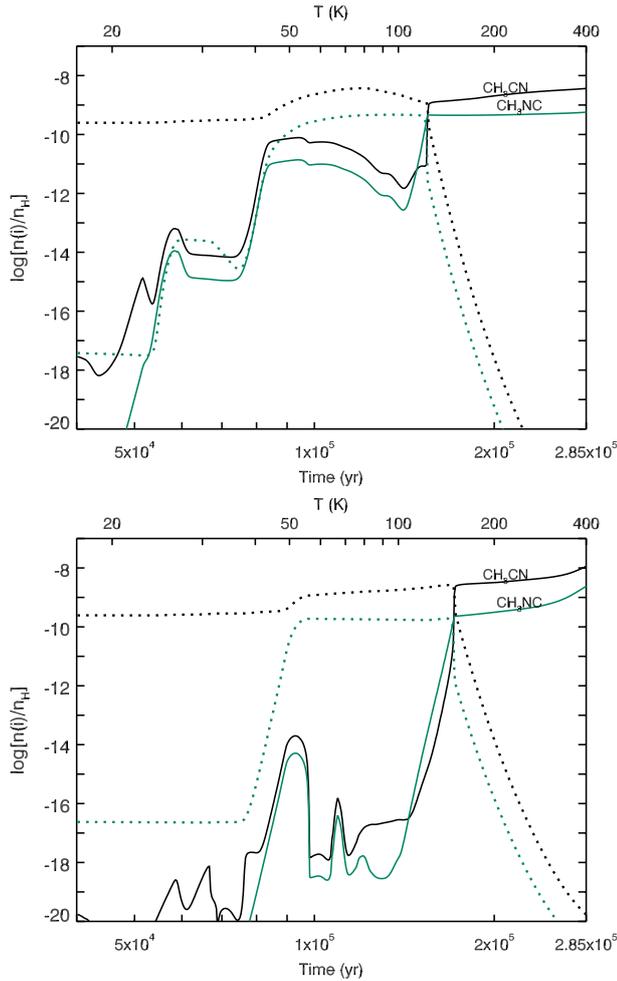


Fig. 4. Chemical model abundances of CH_3CN (black) and CH_3NC (green) for the warm-up stage of a hot core type model where the main destruction pathway for CH_3NC ($\text{H}+\text{CH}_3\text{NC}$) has been disabled. Solid lines denote gas-phase abundances, while dashed lines indicate grain-surface abundances. The *top panel* shows the model with a final collapse density of $n_{\text{H}} = 1.6 \times 10^7 \text{ cm}^{-3}$ and the *bottom panel* shows the model with a final collapse density of $n_{\text{H}} = 6 \times 10^{10} \text{ cm}^{-3}$.

the grain, CH_3NC can be destroyed in the gas phase by reaction with H, which is not the case for CH_3CN . This leads to a dip in CH_3NC abundance relative to CH_3CN and a change in the abundance ratio of these molecules. This may be the reason for the difference in $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ between sources A and B. Also, although the low-density model fits the observations better than the high-density model, it may not be a representative density for the regions being probed by these observations. This and recent results from our other models (Coutens et al. 2018) highlight the need for chemical models in which the collapse and warm-up phases occur simultaneously, and which are tailored for specific objects. Models in which the maximum density is reached prior to any warm-up seem to provide a poor representation of the chemistry in cases where densities reach extreme values. This density threshold appears to be around 10^9 cm^{-3} , based on results from Coutens et al. (2018) for NH_2CN .

4. Discussion

The lack of detection of methyl isocyanide in the second component in the IRAS 16293 binary, IRAS 16293A, shows that there

Table 3. $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ ratios in different objects.

Source	$\frac{N(\text{CH}_3\text{CN})}{N(\text{CH}_3\text{NC})}$	Object type	Ref.
IRAS 16293A	$>5517^a$	LM	-
IRAS 16293B	200	LM	-
Sgr B2(N) molecular cloud	50	HM	[1]
Orion KL	500	HM	[2]
W51 e1/e2	96^b	HM	[3]
NGC 7129 FIRS 2	350^b	IM	[4]
TMC-1	≥ 11	CC	[5]
Horsehead	7	PDR	[6]

Notes. ^(a)Based on upper limits of CH_3NC . ^(b) CH_3NC only tentatively detected. LM stands for low-mass hot core, HM stands for high-mass hot core, IM stands for intermediate-mass, CC stands for cold core, and PDR stands for photon-dominated region.

References. (1) Remijan et al. (2005); (2) López et al. (2014); (3) Kalenskii & Johansson (2010); (4) Fuente et al. (2014); (5) Irvine & Schloerb (1984); (6) Gratier et al. (2013).

is significant chemical differentiation between these sources. It is difficult to determine how common such differences are for different protostars, with only a handful of methyl isocyanide detections in the literature available for comparison. Table 3 shows the $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ abundance ratio from these detections in low-mass hot cores, high-mass hot cores, intermediate-mass hot cores, cold cores, and PDRs. Several of these ratios are based on either tentative detections or on abundance values that are not robust, making a detailed comparison with them problematic, but they can give a general picture of the chemistry. The $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ abundance ratio for IRAS 16293B falls in the middle of the range of values determined in intermediate and high-mass sources, and is an order of magnitude higher than in the TMC-1 cold core and the Horsehead PDR. The value in IRAS 16293A is more than an order of magnitude higher than in the other hot core objects.

Comparisons of the spatial scale of methyl isocyanide to other hot core objects is limited due to a lack of information in this regard. For the most part it is assumed to be co-spatial with methyl cyanide emission, emanating from the hot core component as is found in IRAS 16293B. A different scale is found in Sgr B2(N) by Remijan et al. (2005), where it is deficient in the hot core component, showing only large-scale emission. Such differences between Galactic Centre objects and objects outside the Galactic Centre is unsurprising given the significant differences in the physical environments of such objects. In fact, differences in the spatial scale between Galactic Centre objects and objects outside the Galactic Centre are also seen for other molecules such as glycolaldehyde (Hollis et al. 2004; Beltrán et al. 2009).

The chemical modelling presented in Sect. 3.3 fails to reproduce the observed $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ abundance ratio at the high densities of IRAS 16293 B, but the abundance ratio produced is consistent with the upper limits of IRAS 16293A. It does show that the ratio depends on the temperature, and consequently the evolutionary stage of the source. An earlier stage of evolution would correspond to a lower $\text{CH}_3\text{CN}/\text{CH}_3\text{NC}$ ratio, when CH_3NC has not been significantly destroyed by the gas-phase reaction with H. This combined with additional formation mechanisms that have not been included in the model at present could result in abundance ratios closer to the observational values.

Such a scenario, if applied to IRAS 16293, suggests that the B source is younger than the A source or at least less evolved. This agrees broadly with the outflow picture of the two sources, where outflows have only been detected from IRAS 16293A (Kristensen et al. 2013; Girart et al. 2014), and not IRAS 16293B, where it may not have had time to establish outflows. It is, however, difficult to disentangle the large-scale emission around the B source and therefore determine if there are any indications of previous outflow activity from the source. It could be that B is the less evolved source and therefore its outflows have not yet ‘switched on’ or that it is currently going through a quiescent accretion or ejection phase.

This view of the evolutionary stage of these two sources is contradicted by the vinyl cyanide C_2H_3CN abundances towards both sources. Garrod et al. (2017) find that the vinyl cyanide abundance only becomes significant during the late stages of the warm-up phase of a protostar’s formation. Calcutt et al. (2018) find vinyl cyanide to be nine times more abundant towards IRAS 16293B, which could indicate that the B source is the more evolved protostar. They also suggest that both sources have the same physical age and mass but that the A source has higher accretion rates, leading to a warm-up timescale so short that vinyl cyanide is not efficiently formed. At present it is difficult to determine whether either scenario describes the IRAS 16293 system.

5. Conclusions

In this work we have presented the first detection of methyl isocyanide in a solar-type source, IRAS 16293B, and the first chemical modelling of methyl isocyanide. This represents a significant step forward in understanding methyl isocyanide chemistry in the ISM. Many key questions, however, still remain in order to explain the observed CH_3CN/CH_3NC abundance ratios in different objects. In particular, is the isomerisation of CH_3CNH^+ the dominant reaction pathway to form this molecule? The chemical modelling of methyl isocyanide has also shown that the gas-phase destruction mechanism, $H+CH_3NC$, is necessary to increase the CH_3CN/CH_3NC ratio to that found in hot corino and hot core objects. Whilst this reaction is only speculative, having not yet been studied experimentally, and has a very uncertain barrier, it does highlight the need for an efficient mechanism for the destruction of methyl isocyanide in the gas-phase that does not operate on methyl cyanide.

The chemical modelling has also shown that the CH_3CN/CH_3NC abundance ratio is very sensitive to the final collapse density. At the high densities representative of both A and B, this results in higher CH_3CN/CH_3NC abundance ratios that do not match the observed ratio towards source B. Only the upper limits towards the A source are consistent with the chemical modelling in the high density case.

The next step to understanding methyl isocyanide chemistry is further detections of methyl isocyanide to establish its abundance ratio in a number of objects. Such detections will also need to be compared to further experimental and theoretical work, to explore other formation and destruction mechanisms, and to build a more complete picture of this molecule in the ISM. Crucially, to explore methyl isocyanide chemistry with chemical modelling, a more comprehensive physical model is needed, where warm-up is performed after the final density has been reached. This is especially crucial in objects such as IRAS 16293 where densities reach extreme values.

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