

New ultra metal-poor stars from SDSS: follow-up GTC medium-resolution spectroscopy[★]

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ABSTRACT

Context. The first generation of stars formed in the Galaxy left behind the chemical signatures of their nucleosynthesis in the interstellar medium, visible today in the atmospheres of low-mass stars that formed afterwards. Sampling the chemistry of those low-mass provides insight into the first stars.

Aims. We aim to increase the samples of stars with extremely low metal abundances, identifying ultra metal-poor stars from spectra with modest spectral resolution and signal-to-noise ratio (S/N). Achieving this goal involves deriving reliable metallicities and carbon abundances from such spectra.

Methods. We carry out follow-up observations of faint, $V > 19$, metal-poor candidates selected from SDSS spectroscopy and observed with the Optical System for Imaging and low-Intermediate-Resolution Integrated Spectroscopy (OSIRIS) at GTC. The SDSS and follow-up OSIRIS spectra were analyzed using the FERRE code to derive effective temperatures, surface gravities, metallicities and carbon abundances. In addition, a well-known extremely metal-poor star has been included in our sample to calibrate the analysis methodology.

Results. We observed and analyzed five metal-poor candidates from modest-quality SDSS spectra. All stars in our sample have been confirmed as extremely metal-poor stars, in the $[\text{Fe}/\text{H}] < -3.3$ regime. We report the recognition of J173403+644632, a carbon-enhanced ultra metal-poor dwarf star with $[\text{Fe}/\text{H}] = -4.3$ and $[\text{C}/\text{Fe}] = +3.1$.

Key words. stars: abundances – stars: Population II – Galaxy: abundances – galaxies: formation – galaxies: halos

1. Introduction

Metal-poor stars are extremely rare objects. Their study is key to understanding the early Galaxy. The fraction of stars with low metal content increases with magnitude (Robin et al. 2003) and increases at high Galactic latitudes. In addition, the number of halo stars at $[\text{Fe}/\text{H}] \sim -4$ is only a small percent of those at $[\text{Fe}/\text{H}] \sim -3$, and those are only a small percent of those at $[\text{Fe}/\text{H}] \sim -2$ (Allende Prieto et al. 2016). About two dozen stars are known at $[\text{Fe}/\text{H}] < -4$ (Bonifacio et al. 2015; Placco et al. 2015; Aguado et al. 2016) even though it is the most interesting regime on the yields from the very first generation of stars. Developing new procedures to identify faint metal-poor stars is an important task.

We have been trying to identify these rare objects from large ongoing spectroscopy surveys (SDSS, LAMOST, etc.), to re-observe them at medium-resolution, and to derive reliable stellar atmospheric parameters, including metallicity and whenever possible carbon abundances (Aguado et al. 2016, 2017). This method can be used to identify extremely metal-poor stars as faint as $V \sim 19$. Before performing high-resolution spectroscopy to derive multiple chemical abundances, having an intermediate medium-resolution step provides a powerful time-saving tool

with regard to large telescopes (6–10 m class) since it dramatically removes false positives in the analysis of large-scale spectroscopic surveys. With this methodology, a large sample of ultra-metal poor stars can be assembled in preparation for the arrival of the next 30 m class generation of telescopes.

This is the third paper in a series devoted to follow-up spectroscopy of extremely and ultra metal-poor stars identified from SDSS and other large-scale spectroscopic surveys. In this work we report on faint candidates observed with the Optical System for Imaging and low-Intermediate-Resolution Integrated Spectroscopy (OSIRIS) at Gran Telescopio de Canarias (GTC). We also report on the recognition of J173403+644632, which at $g_{\text{mag}} = 19.6$ is the faintest carbon-enhanced ultra metal-poor dwarf star known. The target selection is explained in Sect. 2, which provides a full description of the procedure originally described in Aguado et al. (2016). The observations and data reduction are detailed in Sect. 3. The derived atmospheric parameters are analyzed and discussed in Sect. 4. In addition, a study of the most metal-poor star known, J1029+1729, is developed to test our methodology in Sect. 4.1. Conclusions are given in Sect. 5.

2. SDSS analysis and target selection

We have re-analyzed more than a million spectra ($\sim 90\%$ stars) from the original Sloan Digital Sky Survey (York et al. 2000), the Sloan Extension for Galactic Understanding and Exploration

[★] Based on observations made with the Gran Telescopio Canarias (GTC), installed in the Spanish Observatorio del Roque de los Muchachos of the Instituto de Astrofísica de Canarias, on the island of La Palma. Programme ID GTC2E-16A and ID GTC65-16B.

Table 1. Coordinates and atmospheric parameters for the program stars based in the analysis of the low-resolution spectra with the FERRE code.

Star	g mag	RA h ' ''	Dec ° ' ''	T_{eff} K	$\log g$ cm s^{-2}	[Fe/H]	[C/Fe]	$\langle S/N \rangle^a$
SDSS J012512+070319	19.0	01:25:12.5	+07:03:19.8	6454	4.9	-3.2	0.8	22
SDSS J015131+163944	18.9	01:51:31.2	+16:39:44.9	5917	3.9	-3.6	1.3	19
SDSS J041800+062308	19.1	04:18:00.7	+06:23:08.1	5579	0.5	-4.3	-	21
SDSS J094708+461010	19.1	09:47:08.2	+46:10:10.1	5916	4.9	-4.7	0.2	14
SDSS J173403+644632	19.6	17:34:03.9	+64:46:32.9	6470	4.9	-3.8	-0.2	15
Star	Plate	mjd	Fiberid	T_{SSPP} K	$\log g_{\text{SSPP}}$ cm s^{-2}	[Fe/H] _{SSPP}		
SDSS J012512+070319	4556	55912	180	-	-	-		
SDSS J015131+163944	5118	55830	336	-	-	-		
SDSS J041800+062308	2826	54389	171	6387	2.70	-3.26		
SDSS J094708+461010	4695	55 957	944	-	-	-		
SDSS J173403+644632	2561	54 597	501	6229	3.11	-3.02		

Notes. The SDSS identification and parameters derived by the automatic SSPP are listed below (when available). ^(a) Signal-no-noise ratios have been calculated as the average value for the entire SDSS spectrum.

(SEGUE, [Yanny et al. 2009](#)) and the Baryonic Oscillations Spectroscopic Survey (BOSS, [Eisenstein et al. 2011](#); [Dawson et al. 2013](#)). The main stellar parameters (effective temperature T_{eff} , surface gravity $\log g$, and metallicity $[\text{Fe}/\text{H}]^1$) and the carbon abundance have been derived.

The SDSS optical spectra are from SDSS Data Release 9 (DR9, [Ahn et al. 2012](#)) for observations with the original SDSS spectrograph, and DR12 ([Alam et al. 2015](#)) for data obtained with the upgraded BOSS spectrographs ([Smee et al. 2013](#)). The resolving power of SDSS spectra is about 2000, covering the range $\sim 3800\text{--}9100 \text{ \AA}$. Since the integration time is the same for all the targets, the signal-to-noise ratio (S/N) is widely variable. We use the fortran FERRE code ([Allende Prieto et al. 2006](#)) to perform a massive analysis of the spectra, as explained in [Aguado et al. \(2016 and 2017\)](#). For this work we focus on promising candidates with low S/N spectra and, consequently, higher error bars in the parameter determination (see Table 1). These targets are too faint for follow-up with the 4.2 m WHT, and so a parallel program was started on the 10.4 m GTC.

3. Observations and data reduction

We performed medium-resolution, long-slit spectroscopy with OSIRIS at the 10.4 m telescope GTC. The program was granted 28 h under proposals ID GTC2E-16A and ID GTC65-16B in service mode to carry out long-slit spectroscopy with the R2500U grism and a 1.0 arcsec slit, providing a spectral range 3600–4500 Å with a resolving power ($R \equiv \lambda/\delta\lambda \sim 2300$). The slit was in all cases oriented at the parallactic angle. The individual exposures were never longer than 30 min to minimize the impact of cosmic rays (see Table 2).

Data reduction consisted in bias subtraction, flat-fielding, wavelength calibration using HgAr + Xe calibration lamps, and a combination of individual spectra. The processing was performed with the *twodspec* and *onespec* packages in IRAF²

¹ We use the bracket notation to report chemical abundances: $[a/b] = \log\left(\frac{N(a)}{N(b)}\right) - \log\left(\frac{N(a)}{N(b)}\right)_\odot$, where $N(x)$ represents number density of nuclei of element x .

² IRAF is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy (AURA) under cooperative agreement with the National Science Foundation.

Table 2. Observing log.

Stars	Date	N_{exp}	t_{exp} s	Seeing
J012512+070319	27-Sep.-2016	6	1600	0.8
J015131+163944	22/28-Sep.-2016	7	1600	1.1
J041800+062308	26/28-Sep.-2016	6	1600	1.0
J094708+461010	28/31-Mar.-2016	10	1827	0.9
J173403+644632	29/30-Apr.-2016	10	1827	1.0
G64-12	1-Jan.-2017	3	30	1.2

Notes. A slit width of 1 arcsec. The moon was gray during all the observations and clouds were not noticeable.

([Tody 1993](#)). Cross-correlation between individual spectra and correction of radial velocity offsets was made with the *rv* package.

4. Analysis and discussion

Following the same procedure described by [Aguado et al. \(2017\)](#), we employed a grid of synthetic spectra computed with the ASSεT code ([Koesterke et al. 2008](#)). The model atmospheres were computed as in [Mészáros et al. \(2012\)](#). The spectral synthesis was performed as described by [Allende Prieto et al. \(2014\)](#), but the grid we used differs in the parameter coverage: $-6 \leq [\text{Fe}/\text{H}] \leq -2$; $+1 \leq [\text{C}/\text{Fe}] \leq +5$; $4750 \text{ K} \leq T_{\text{eff}} \leq 7000 \text{ K}$, and $1.0 \leq \log g \leq 5.0$. The α -element abundance and microturbulence were fixed at $[\alpha/\text{Fe}] = 0.4$ and 2 km s^{-1} , respectively. Lacking a reliable flux calibration for our observations, a running-mean filter with a window of 30 nearby pixels was used for continuum normalization. The same algorithm was applied to both observed and synthetic spectra after resampling the latter to the wavelength OSIRIS spectra of the former. Further details are given in [Aguado et al. \(2017\)](#). FERRE is used to simultaneously derive the three main stellar parameters – effective temperature, surface gravity, metallicity – as well as the carbon abundance.

Table 3 summarizes the results of the FERRE analysis for the sample stars. A high-quality OSIRIS spectrum of the well-known metal-poor star G64-12 was analyzed as a reference. The good agreement between the optimal set of parameters

Table 3. The stellar parameters and main results obtained from OSIRIS spectra.

Star	$\langle S/N \rangle$	T_{eff} [K]	ΔT_{eff} [K]	$\log g$ [cm s ⁻²]	$\Delta \log g$ [cm s ⁻²]	[Fe/H]	Δ [Fe/H]	[C/Fe]	Δ [C/Fe]
SDSS J012512+070319	77	6312	34	5.0	0.07	-3.3	0.1	1.2	0.6
SDSS J015131+163944 ¹	62	6036	26	5.0	0.07	-3.8	0.08	1.1	0.5
SDSS J041800+062308	73	6247	35	5.0	0.08	-3.4	0.1	0.7	0.7
SDSS J094708+461010	80	5858	19	5.0	0.07	-4.1	0.06	1.0	0.3
SDSS J173403+644632	60	6183	35	5.0	0.09	-4.3	0.11	3.1	0.1
G64-12	132	6393	29	4.8	0.07	-3.2	0.1	1.0	0.4

Notes. Δ are the internal uncertainties of parameters derived with FERRE.

References. (1) [Aguado et al. \(2017\)](#).

($T_{\text{eff}} = 6393$ K, $\log g = 4.8$, $[\text{Fe}/\text{H}] = -3.2$) and the literature ([Aoki et al. 2006](#); [Placco et al. 2016](#)) ($T_{\text{eff}} = 6390$ K, $\log g = 4.38$, $[\text{Fe}/\text{H}] = -3.2$), together with the quality of the fit, suggest that this is a highly reliable method of confirming extremely metal-poor stars. Further details on FERRE tests are given by [Allende Prieto et al. \(2014\)](#), [Aguado et al. \(2017\)](#).

FERRE offers multiple options for deriving internal uncertainties. For our program we chose the estimates obtained from Monte Carlo experiments after injecting noise into the observations and the error treatment is the same as that included in [Aguado et al. \(2017\)](#). We derived the error bars by adding in quadrature the internal uncertainties showed in Table 3 with other sources of error: $\Delta T_{\text{eff}} = 70$ K, $\Delta \log g = 0.5$, $\Delta[\text{Fe}/\text{H}] = 0.1$ dex and $\Delta[\text{C}/\text{Fe}] = 0.2$ dex. The spectral range in common between the ISIS and the OSIRIS spectra is about 820 \AA wider, with the H_{β} line missing in the latter case. We examine the results for each object in Sect. 4.2.

4.1. Testing the methodology with J1019+1729

[Caffau et al. \(2011, 2012\)](#) recognized J1019+1729, a dwarf star at a metallicity of about $[\text{Fe}/\text{H}] \sim -5$, from SDSS/SEGUE spectroscopic data. The low carbon and oxygen abundances of this star challenge the picture of star formation that suggests no low-mass stars were formed ([Bromm & Loeb 2003](#)). To test our methodology, we analyzed the X-shooter spectrum of this star ($R = 7900$ in the UVB arm). FERRE derives $T_{\text{eff}} = 5834$ K, $\log g = 5.0$, $[\text{C}/\text{Fe}] = -0.47$, and $[\text{Fe}/\text{H}] = -4.55$. The effective temperature derived by the authors (5811 K) is perfectly recovered by the code and is in agreement with our analysis of the ISIS spectrum of this target presented in [Aguado et al. \(2017\)](#). In addition, FERRE is able to identify J1029+1729 as a dwarf star. Since it is not visible, we are able to place an upper limit $[\text{C}/\text{Fe}] < +0.7$. However, our derived metallicity is higher than that determined by [Caffau et al. \(2012\)](#), due to the ISM contribution to the Ca II H and K lines. The metallicity is compatible with our analysis in [Aguado et al. \(2017\)](#) with ISIS and UVES-smoothed data. Furthermore, we ran the FERRE code only in the spectral region where the strongest Fe I lines are (3810–3865 Å) and assuming the previously derived T_{eff} and $\log g = 5.0$ over the entire X-shooter spectrum, and we were able to recover $[\text{Fe}/\text{H}] = -4.75$ (see Fig. 1), which is in agreement with the 1D-LTE Fe I abundance derived by the authors in the original paper from higher resolution UVES spectroscopy. Finally, we smoothed the X-shooter spectrum of J1019+1729 to the OSIRIS resolution ($R = 2500$) and, again, we ran FERRE and arrived at the values $[\text{Fe}/\text{H}] = -4.5$ and $T_{\text{eff}} = 5826$ K. This exercise

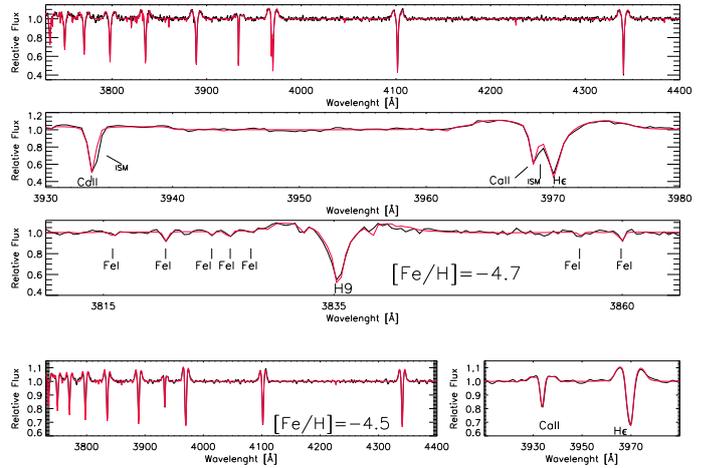


Fig. 1. VLT/X-shooter spectrum ($R = 7900$) of J1019+1729 (black line) and the best fit derived by FERRE (red line). *From top to bottom:* the entire spectrum (3700–4400 Å). The range of the Ca II lines. The H9 spectral region with several Fe I lines assuming $T_{\text{eff}} = 5834$ K, $\log g = 5.0$ from our previous analysis over the entire spectra. The derived iron abundance is displayed. Finally the X-shooter spectrum smoothed to the OSIRIS resolution ($R = 2500$) and re-analyzed by FERRE: *left panel* shows the entire spectrum together with the derived metallicity while in the *right panel* a detail of the Ca II lines is depicted.

demonstrates that our methodology is robust and works even at higher resolution ($R = 7900$). We are able to recover the same metallicity and effective temperature from the analysis of SDSS, ISIS, and UVES-smoothed data ([Aguado et al. 2017](#)), and from X-shooter and X-shooter-smoothed spectra.

4.2. Discussion of the new set of faint metal-poor stars

Figure 2 (upper panel) shows the spectrum of G64-12 (black solid line) and its best fit derived with FERRE. All spectra analyzed in this work are plotted below.

4.2.1. J012512+070319

The quality of the BOSS spectrum of this target is the highest in the sample. The initially derived set of parameters appears to be reliable; the much higher S/N OSIRIS spectra lead to an effective temperature that differs by only 140 K. The slightly lower temperature we find decreases the derived metallicity by 0.1 dex. The final values for the parameters and metallicities of this object are $T_{\text{eff}} = 6312 \pm 78$ K, $\log g = 5.0 \pm 0.5$,

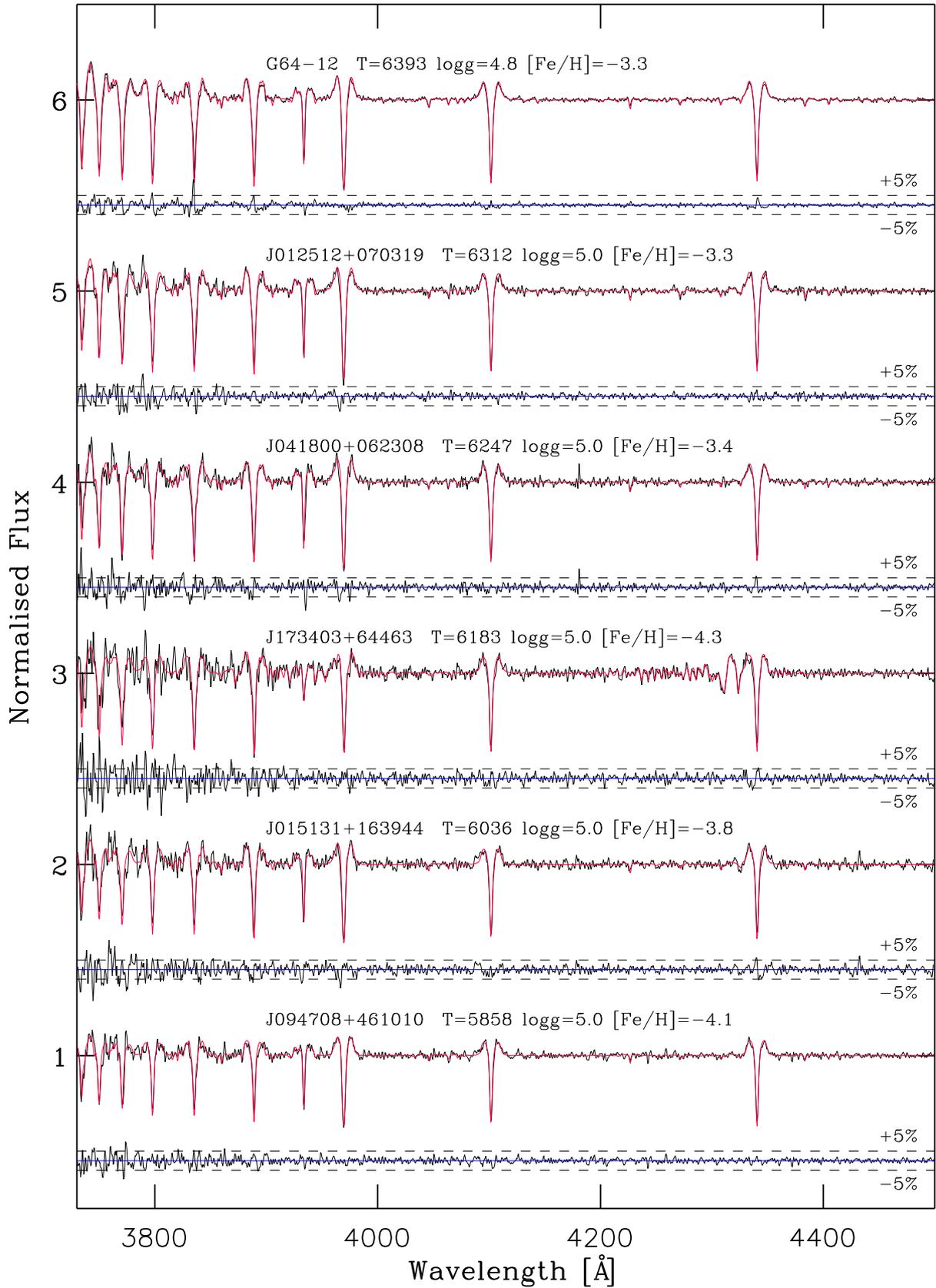


Fig. 2. GTC/OSIRIS spectra (3700–4500 Å) of our stellar sample (black line) and the best fit calculated with FERRE (red line). Under each of the spectra we also depict the residuals (difference between the observed spectrum and the best fit) together with the +5% and -5% reference lines. On top of each spectrum the main stellar parameters are displayed.

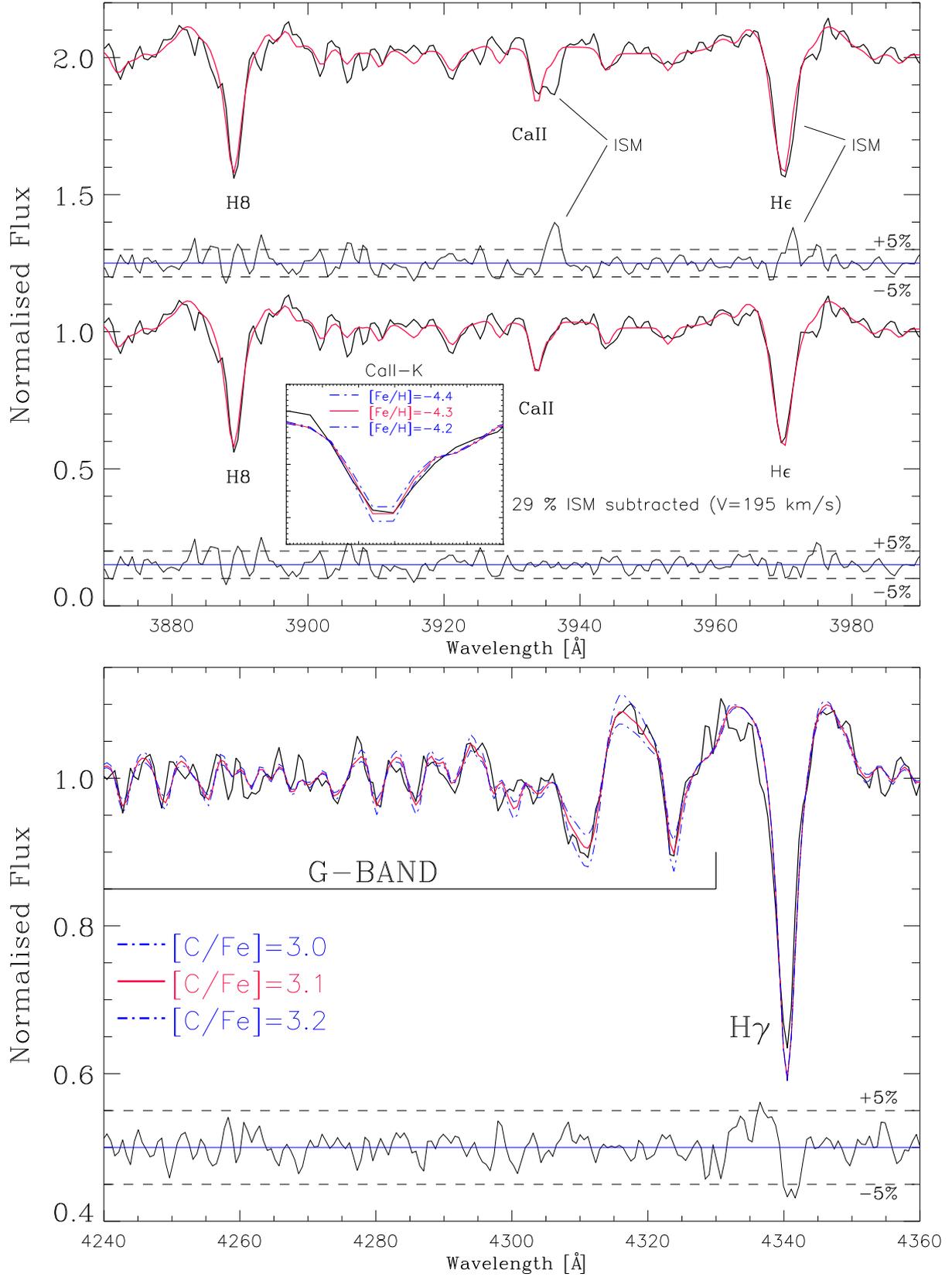


Fig. 3. *Upper panel:* spectral range in the vicinity of the Ca II H-K lines for the OSIRIS spectrum of the star J134157+513534 (black solid line) with the best fit (red line) before ISM subtraction (*top panel*) and after ISM subtraction (*bottom panel*). Below the spectra we display the residuals together with the +5% and -5% reference lines. A detail of the Ca II K line after ISM subtraction with an upper and lower metallicity limit (dashed blue line) is plotted. *Lower panel:* spectral range of the G-band of the OSIRIS spectrum of J134157+513534 (black solid line) with the synthetic spectra for the best fit (red line), an upper and lower limit of carbon abundances (dashed blue line).

and $[\text{Fe}/\text{H}] = -3.3 \pm 0.2$. The derived carbon abundance is $[\text{C}/\text{Fe}] = 1.2 \pm 0.6$.

4.2.2. J015131+163944

This object was observed with ISIS on the WHT and analyzed by [Aguado et al. \(2017\)](#). The set of derived parameters by these authors is $T_{\text{eff}} = 6025 \pm 77$ K, $\log g = 4.8 \pm 0.6$, $[\text{Fe}/\text{H}] = -3.8 \pm 0.2$ and $[\text{C}/\text{Fe}] = 1.3 \pm 0.4$, from a spectrum with very similar resolving power ($R \sim 2400$) and slightly lower S/N than the OSIRIS spectrum presented here. The values given in [Table 3](#) are essentially the same: $T_{\text{eff}} = 6036 \pm 75$ K, $\log g = 5.0 \pm 0.5$, $[\text{Fe}/\text{H}] = -3.8 \pm 0.2$ and $[\text{C}/\text{Fe}] = 1.1 \pm 0.5$.

4.2.3. J041800+062308

The SEGUE spectrum of this object has a fairly poor quality, and, as a result, the derived effective temperature, $T_{\text{eff}} = 5579$ K, is not very reliable. The set of parameters from the OSIRIS analysis ($T_{\text{eff}} = 6247 \pm 78$ K, $\log g = 5.0 \pm 0.5$, $[\text{Fe}/\text{H}] = -3.4 \pm 0.2$) are more robust, while the *G*-band is close to the detection limit ($[\text{C}/\text{Fe}] = 0.7 \pm 0.7$).

4.2.4. J094708+461010

The S/N of this BOSS spectrum is the lowest in the sample, and for this reason we consider the derived atmospheric parameters in [Table 1](#) unreliable. However, both the effective temperature and the gravity derived from the OSIRIS spectrum are compatible with those from the BOSS observations. Accepting the FERRE derived values, we arrive at $T_{\text{eff}} = 5858 \pm 73$ K, $\log g = 5.0 \pm 0.5$, $[\text{Fe}/\text{H}] = -4.1 \pm 0.2$, and $[\text{C}/\text{Fe}] = 1.0 \pm 0.4$. Interestingly enough, the Yoon-Beers diagram described in [Yoon et al. \(2016\)](#) points out that there are two different “families” of CEMP-no stars, which they designate as Group II and Group III. J094708+461010 is a Group II with $[\text{Fe}/\text{H}] = -4.1$ and $A(\text{C}) = 5.3$.

4.2.5. J173403+644632

The spectrum of this target shows the clearest and unresolved calcium ISM contribution, and the strongest *G*-band absorption, making this faint object very interesting candidate for a detailed study. Since the Ca II H-K spectral region (the one containing more information) is affected by CH lines, we perform the analysis with FERRE as follows:

- We derived simultaneously the effective temperature ($T_{\text{eff}} = 6183$ K (see [Fig. 3](#), upper panel), the carbon abundance ($[\text{C}/\text{Fe}] = 3.1$: see [Fig. 3](#), lower panel), and a surface gravity of $\log g = 5.0$.
- A Gaussian profile for the calcium ISM contribution is adopted. A grid of varying absorption coefficients (ϵ) from 5% to 30% in step sizes of 1% and relative velocities from 180 km s^{-1} to 220 km s^{-1} in steps of 5 km s^{-1} were subtracted, and thus the resulting spectrum was reanalyzed with FERRE.
- The minimum χ^2 is obtained for the point $\epsilon = 0.29$ and $V = 195 \text{ km s}^{-1}$ which reproduces best both the Ca II and the H_γ lines (see [Fig. 3](#), upper panel).

Thus, the final set of parameters are $T_{\text{eff}} = 6183 \pm 78$ K, $\log g = 5.0 \pm 0.5$, $[\text{Fe}/\text{H}] = -4.3 \pm 0.2$, and $[\text{C}/\text{Fe}] = 3.1 \pm 0.2$. Again, regarding the diagram mentioned in [Sect. 4.2.4](#), J173403+644632 with $[\text{Fe}/\text{H}] = -4.3$ and $A(\text{C}) = 7.2$ would clearly be in the Group III. This target is the only ultra metal-poor star known at $V > 19$. Finally, the fact that the derived

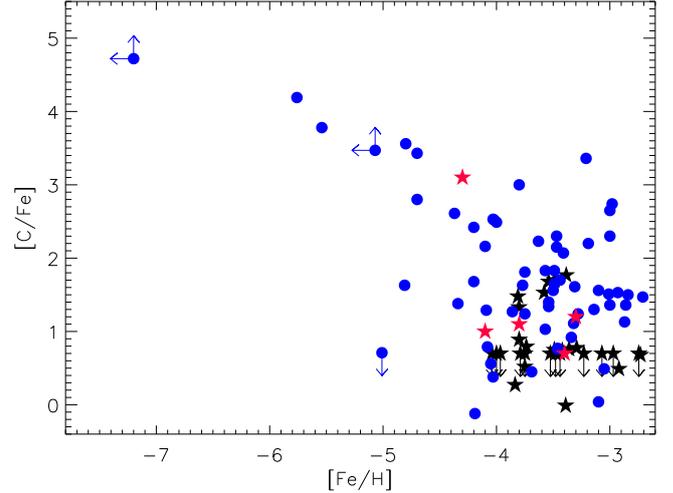


Fig. 4. Carbon over iron versus metallicities of stars presented in this work (red filled stars) and those from [Aguado et al. \(2017\)](#) (black filled stars) analyzed with the improved methodology, together with stars from the literature ([Sivarani et al. 2006](#); [Yong et al. 2013](#); [Frebel et al. 2005, 2006](#); [Caffau et al. 2014](#); [Allende Prieto et al. 2015](#)) represented by blue filled circles.

gravities are always on the limit of the grid is compatible with the relatively high error bars adopted. However, the influence on the derived metallicity and carbon abundance is small and already considered in error bars through other sources of error ([Aguado et al. 2017](#)).

5. Conclusions

The multiple techniques used to identify and confirm metal-poor stars have been summarized by [Beers & Christlieb \(2005\)](#), [Frebel & Norris \(2015\)](#), and references therein, and include high proper-motion surveys, Schmidt objective-prism surveys, photometric surveys, and spectroscopic surveys. The most iron-poor star known, SM 0313–6708 with $[\text{Fe}/\text{H}] < -7.3$, was discovered from a photometric survey, the Skymapper Southern Sky Survey ([Keller et al. 2014](#)). The most metal-poor star known, J1029+1729, was in turn recognized in the SDSS spectroscopic survey by [Caffau et al. \(2012\)](#). Obviously, the observational efforts and resources employed demand less telescope time when using photometry to identify metal-poor stars ([Starkenburg et al. 2017](#)), but the only path to confirm and derive reliable metallicities involves spectroscopic techniques.

Usually a low-resolution spectroscopic analysis is required to securely identify metal-poor candidates, and the subsequent and significantly more expensive high-resolution analysis confirms the metallicity and provides multiple chemical abundances. Avoiding false positives (incorrect identifications) optimizes the invested time in 6–10 m class telescopes. In this paper we use an intermediate step (see [Beers et al. 1985](#)) to increase this ratio. Using medium-resolution spectroscopy we are able to derive reliable metallicities and sometimes carbon abundances, even in the $[\text{Fe}/\text{H}] < -4.0$ regime.

We have identified a number of very faint metal-poor stars, and offer to the community a sizable sample of faint ultra metal-poor stars to observe with the next generation of 30–40 m telescopes. [Figure 4](#) shows the carbon abundances derived from OSIRIS spectra for the targets studied in this work, together with those from [Aguado et al. \(2017\)](#), and other CEMP stars from the literature. The chance to carry out a follow-up program for observing stars with $V > 19$ significantly improves our success ratio of identifying ultra metal-poor stars. The more

carbon-enhanced metal-poor stars are those that are below the $[\text{Fe}/\text{H}] = -4.0$ regime.

In addition, we demonstrate the advantages of our methodology when the calcium ISM contribution is not resolved at medium resolution which is the vast majority. However, the most interesting star of this work, J173403+644632, is an excellent candidate to observe at higher resolution spectroscopy and to derive its chemical composition in detail, giving important constraints to the early Universe. Finally, we have observed and analyzed, with highly reliable results, well-known metal-poor stars (G64-12 and J1029+1729), together with those presented in Aguado et al. (2017), and this allows us to trust our methodology when looking for ultra metal-poor stars.

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