The ATLASGAL survey: The sample of young massive cluster progenitors

T. Csengeri¹, S. Bontemps², F. Wyrowski¹, S. T. Megeath³, F. Motte⁴,⁵, A. Sanna¹, M. Wienen¹, and K. M. Menten¹

1 Max Planck Institute for Radioastronomy, Auf dem Hügel 69, 53121 Bonn, Germany
e-mail: ctmea@mpifr-bonn.mpg.de
2 Ritter Astrophysical Research Center, Department of Physics and Astronomy, University of Toledo, 2801 W. Bancroft St., Toledo, OH 43606, USA
3 Institut de Planétologie et d’Astrophysique de Grenoble, Univ. Grenoble Alpes – CNRS-INSU, BP 53, 38041 Grenoble Cedex 9, France
4 Laboratoire AIM Paris Saclay, CEA-INSU/CNRS-Université Paris Diderot, IRFU/SAp CEA-Saclay, 91191 Gif-sur-Yvette, France

Received 4 February 2016 / Accepted 16 December 2016

ABSTRACT

Context. The progenitors of high-mass stars and clusters are still challenging to recognise. Only unbiased surveys, sensitive to compact regions of high dust column density, can unambiguously reveal such a small population of particularly massive and cold clumps.

Aims. Here we use the ATLASGAL survey to identify a sample of candidate progenitors of massive clusters in the inner Galaxy.

Methods. We characterise a flux limited sample of compact sources selected from the ATLASGAL survey. Sensitive mid-infrared data at 21–24 μm from the WISE and MIPS GAL surveys were explored to search for embedded objects, and complementary spectroscopic data were used to investigate their stability and their star formation activity.

Results. We identify an unbiased sample of infrared-quiet massive clumps in the Galaxy that potentially represent the earliest stages of massive cluster formation. An important fraction of this sample consists of sources that have not been studied in detail before. We first find that clumps hosting more evolved embedded objects and infrared-quiet clumps exhibit similar physical properties in terms of mass and size, suggesting that the sources are not only capable of forming high-mass stars, but likely also follow a single evolutionary track leading to the formation of massive clusters. The majority of the clumps are likely not in virial-equilibrium, suggesting collapse on the clump scale.

Conclusions. We identify the precursors of the most massive clusters in the Galaxy within our completeness limit, and argue that these objects are undergoing large-scale collapse. This is in line with the low number of infrared-quiet massive clumps and earlier findings that star formation, in particular for high-mass objects is a fast, dynamic process. We propose a scenario in which massive clumps start to fragment and collapse before their final mass is accumulated indicating that strong self-gravity and global collapse is needed to build up rich clusters and the most massive stars.

Key words. surveys – stars: massive – stars: formation

1. Introduction

The formation of massive stars and clusters remains an enigma in modern astrophysics. To explain the origin of high-mass stars, two theoretical scenarios are commonly invoked: the turbulent-core model (McKee & Tan 2003) and a more dynamic view with cluster formation (Bonnell et al. 2001; Bonnell & Bate 2006). The observational challenge to these theories is a systematic, cluster formation (Bonnell et al. 2001; Bonnell & Bate 2006).

With “earliest phases” we refer to the starless/pre-stellar, and the subsequent phase where the submillimetre emission is still dominated by the massive core. This implies that non-luminous massive young stellar object (MYSO) has yet emerged, and no ionising emission from embedded UC-H II regions. This phase is difficult to observe, however, mainly because of the short duration of less than 7.5 × 10³ yr at the scale of ~0.3 pc (Csengeri et al. 2014), and also because of the high complexity of active sites of high-mass star formation. Only a few candidates of massive starless cores have been claimed so far in the literature with masses up to 60 M⊙ in a size of ≥0.1 pc (Tan et al. 2013; Cyganowski et al. 2014), although towards the former object outflow activity has recently been discovered (Tan et al. 2016), and the absence of any molecular emission towards the latter object casts doubts on its Galactic origin. Confirmed examples of massive pre-stellar cores, as well as exploration of the even higher mass regime, are of extreme importance in order to constrain star formation models.

Early searches for the birthplaces of high-mass stars focused on IRAS selected infrared-bright sources, where low- and high-luminosity sources were defined (Molinari et al. 1996). High-luminosity sources were associated with UC-H II regions, i.e. the final stage of high-mass star formation, while the lower luminosity sources were associated with MYSOs. UC-H II regions have typical bolometric luminosities of >10³ L⊙, while (M)YSOs show a broad distribution peaking between 3 × 10⁻²–2 × 10⁴ L⊙.
In Sect. 2.7 we present molecular line data to characterise their evolutionary stage based on the luminosity of the embedded protostars as well as their IR emission and dust temperatures. We study here a flux limited sample of the ATLASGAL compact sources within the highest sensitivity part of the survey, corresponding to a Galactic longitude range of $|l| \leq 60\degree$. The initial sample was selected using a 5 Jy beam-averaged flux density threshold, which translates to $\sim 650 M_\odot$ (see Csengeri et al. 2014) at an average distance of ATLASGAL sources of $d \approx 4.5$ kpc (Wienen et al. 2015). We estimate the corresponding column density by

$$N(H_2) = \frac{F_\nu}{B_\nu(T_\text{d})\Omega \kappa H_2 m_{H_2}} \text{[cm}^{-2}],$$

where $F_\nu$ is the beam-averaged peak flux density; $\Omega$ is the solid angle of the beam calculated by $\Omega = 1.13 \times \Theta^2$, where $\Theta = 19\arcs;\ k_B$ is the mean molecular weight per hydrogen molecule and is equal to 2.8 (Kaufmann et al. 2008); $m_{H_2}$ is the mass of a hydrogen atom; and $R$ corresponds to the gas-to-dust ratio of 100. Our flux density threshold thus corresponds to a column density ($N(H_2)$) of $4.9-40.8 \times 10^{23} \text{cm}^{-2}$ using a centre frequency of 345 GHz at our resolution of $19\arcs$; and dust temperatures of 10–40 K. For the highest column density sources this translates to an extremely high visual extinction of $A_V > 521$ mag$^1$ (see details in Sect. 2.3).

From the ATLASGAL compact source catalogue we find a total of 219 sources above the 5 Jy beam-averaged flux density threshold (Csengeri et al. 2014). In the subsequent steps we first provide distance estimates (Sect. 2.1), and then determine their physical properties (Sects. 2.2 and 2.3). Using mid-infrared data between 21–24 $\mu$m from the WISE point source catalogue (Wright et al. 2010) and photometry based on the MIPS$^2$ GALEX GALEX GALEX catalogues (Carey et al. 2009) and complemented with classification from the MYSO and UC-H ii regions from the RMS (Urquhart et al. 2007) and CORNISH surveys (Hoare et al. 2012; Purcell et al. 2013), we then estimate the evolutionary stage based on the luminosity of the embedded protostars associated with the dust peaks (Sects. 2.4 and 2.5). This leads to the identification of infrared-quiet massive clumps in Sect. 2.6. In Sect. 2.7 we present molecular line data to characterise their associations with shocks from outflow activity to probe ongoing star formation.

### 2.1. Distances and association with massive molecular complexes

The large statistics of the ATLASGAL survey revealed that most of the compact sources, and correspondingly most of the star formation activity is associated with giant molecular complexes (GMCs; e.g. Csengeri et al. 2014). As a first step, therefore, we

\footnote{For the conversion we use $A_V = N_{H_2}(1.04 \times 10^{21} \text{cm}^{-2})$.}
associated the sample with known GMCs (see Table 1). In fact, our selection corresponds to the brightest submillimetre sources in the inner Galactic plane, and it covers well-characterised massive star-forming complexes such as Sgr B2, W51 Main, W49A, W33, and W43. Other examples are found at the edge of expanding bubbles like G348.1825+0.4829, which is the brightest compact source in the vicinity of RCW 120 (Zavagno et al. 2010). We find a few examples of relatively isolated sources, which are not associated with any known GMCs, but smaller molecular clouds, such as G0.5464+0.8521, G333.0164+0.7654, and G43.7950-0.1270, towards which our distance estimates may be less reliable than for the majority of the sample.

We collected distance estimates from the literature using primarily maser parallax measurements (see Reid et al. 2014, for an overview) or spectrophotometric distance estimates (e.g. Russeil 2003; Moisés et al. 2011) for the known GMCs. For the remaining unassociated sources we used the high-density tracer NH$_3$ to determine kinematic distances from Wienen et al. (2012, 2015).

Table 1. Association of the brightest ATLASGAL sources with GMCs.

<table>
<thead>
<tr>
<th>Associated GMC or molecular cloud</th>
<th>$d$ [kpc]</th>
<th>Number of sources</th>
</tr>
</thead>
<tbody>
<tr>
<td>W51</td>
<td>5.41$^b$</td>
<td>8</td>
</tr>
<tr>
<td>W-49</td>
<td>11.11$^f$</td>
<td>3</td>
</tr>
<tr>
<td>IRDC G34.4+0.2</td>
<td>1.56$^p$</td>
<td>3</td>
</tr>
<tr>
<td>W43</td>
<td>5.49$^d$</td>
<td>4</td>
</tr>
<tr>
<td>W31</td>
<td>4.95$^f$</td>
<td>3</td>
</tr>
<tr>
<td>W33</td>
<td>2.4$^c$</td>
<td>3</td>
</tr>
<tr>
<td>M17</td>
<td>1.98$m$</td>
<td>4</td>
</tr>
<tr>
<td>G5.89</td>
<td>1.28$^r$</td>
<td>2</td>
</tr>
<tr>
<td>Galactic centre</td>
<td>8.5$^j$</td>
<td>9</td>
</tr>
<tr>
<td>NGC 6334</td>
<td>1.35$^i$</td>
<td>9</td>
</tr>
<tr>
<td>RCW 122</td>
<td>1.8$^{i,j}$</td>
<td>4</td>
</tr>
<tr>
<td>RCW 108-109</td>
<td>1.8$^{i,j}$</td>
<td>3</td>
</tr>
<tr>
<td>G333.1-0.4</td>
<td>4.2$^{a,g}$</td>
<td>6</td>
</tr>
<tr>
<td>RCW106</td>
<td>4.2$^d$</td>
<td>5</td>
</tr>
<tr>
<td>G327</td>
<td>2.5$^h$</td>
<td>3</td>
</tr>
<tr>
<td>RCW 95</td>
<td>2.5$^{a,q}$</td>
<td>5</td>
</tr>
<tr>
<td>RCW 92</td>
<td>3.2$^e$</td>
<td>2</td>
</tr>
<tr>
<td>G305</td>
<td>3.5$^g$</td>
<td>2</td>
</tr>
<tr>
<td>G10.47</td>
<td>8.47$^f$</td>
<td>1</td>
</tr>
<tr>
<td>RCW 120</td>
<td>1.3$^e$</td>
<td>1</td>
</tr>
<tr>
<td>RCW 131/NGC 6537</td>
<td>1.7$^h$</td>
<td>1</td>
</tr>
<tr>
<td>G345.5+1.0</td>
<td>1.8$^{i,a}$</td>
<td>1</td>
</tr>
<tr>
<td>RCW 116b</td>
<td>1.8$^{g,q}$</td>
<td>1</td>
</tr>
<tr>
<td>RCW 117</td>
<td>1.8$^{e,k}$</td>
<td>1</td>
</tr>
<tr>
<td>RCW 99</td>
<td>2.5$^h$</td>
<td>1</td>
</tr>
<tr>
<td>G9.62+0.19</td>
<td>5.15$n$</td>
<td>1</td>
</tr>
<tr>
<td>G351.77</td>
<td>1.0$^o$</td>
<td>1</td>
</tr>
<tr>
<td>SDC335</td>
<td>3.8$^{s,q}$</td>
<td>1</td>
</tr>
</tbody>
</table>

Notes. 
(a) This work; (b) Sato et al. (2010); (c) Zhang et al. (2013); (d) Zhang et al. (2014); (e) Immer et al. (2013); (f) Sanna et al. (2014); (g) Moisés et al. (2011); (h) we revised the spectrophotometric distance of 1.81 kpc by Moisés et al. (2011); (i) Wu et al. (2014); (j) Wu et al. (2012); (k) Reid et al. (2014); (l) Dutra et al. (2003); (m) Straw et al. (1987); (n) Xu et al. (2011); (o) Sanna et al. (2009); (p) Leumini et al. (2011); (q) Kurayama et al. (2011); (r) Russeil (2003); (s) Motogi et al. (2011); (t) López et al. (2011); (u) Peretto et al. (2013) adopts 3.25 kpc.

In short, the rest velocity of the source is compared to the Galactic rotation curve using the parameters derived by Reid et al. (2009) for the Ist and Brand & Blitz (1993) for the IVth Galactic quadrants, respectively. We attempt to solve the distance ambiguity by making use of HI emission and 21 cm continuum absorption against bright HII regions (see Wienen et al. 2015, for details). We then visually inspected near-IR extinction maps (Schneider et al. 2010) and data from the Galactic Ring Survey (GRS, Jackson et al. 2006) to look for spatial and kinematic associations with nearby molecular clouds also considering the larger scale Galactic structure such as spiral arms, and associations with extinction features as well. We also inspected sensitive mid-infrared data from GLIMPSE and MIPS to look for possible associations with infrared-dark dust lanes appearing in high contrast absorption at mid-infrared wavelengths.

In total we associate 78 (36%) sources with available maser parallax measurements. In the Ist Galactic quadrant, such measurements are available for 59 out of 85 sources (70%), providing a robust estimate of their physical parameters. Sources in the IVth Galactic quadrant have less reliable distance measurements owing to the lack of currently available maser parallax distances, except for G339.8841-1.2568 (Krishnan et al. 2015). We were able to provide distance estimates to all but one source in the sample (see Table 4). For simplicity, in the following analysis we exclude the 9 sources associated with the Galactic centre, leaving a total of 210 sources in our sample. An overview of their Galactic distribution is shown in Fig. 1.

2.2. Temperature estimates

Our sample includes clumps in various evolutionary stages such as active star-forming sites with clusters of known (UC-)HII regions (e.g. W51 Main, W49 N, G5.89) and typical hot molecular cores (HMCs; e.g. G10.47+0.3, G34.26+0.15) associated with infrared-bright MYsOs. Over such a large star formation activity, the estimated physical parameters from the 870µm dust emission can be highly sensitive to the used dust temperature. As a simple first-order approximation, a single dust temperature is commonly used for which a $T_d$ of 18 K is a reasonable choice (see e.g. Csengeri et al. 2014) based on the Herschel measured dust temperatures dominated by the interstellar radiation field (Bernard et al. 2010). However, the brightest submillimetre sources are more likely to host embedded heating source(s), which could imply higher dust temperatures. This is supported by statistical studies; for instance, Battersby et al. (2011) show that tracers of star formation are more likely associated with higher dust temperatures, while König et al. (2017) show an increase in dust temperature with evolutionary stage by fitting the spectral energy distribution of ATLASGAL clumps over a range of evolutionary stages. Hence, a temperature correction is necessary for clumps hosting UC-HII regions and infrared-bright protostars, and we discuss below our choice of a higher dust temperature.

At the high densities probed by the ATLASGAL survey it is reasonable to assume that the gas and dust are thermalised. Therefore, to investigate which dust temperature is appropriate for the warmer clumps, we can also rely on complementary molecular line data available in the literature. The NH$_3$ measurements by Wienen et al. (2012) give a gas kinetic temperature between 12−35 K for a large sample of ATLASGAL sources within a beam of ∼40′. Similarly, Dunham et al. (2011) derive a mean gas kinetic temperature of 17 K for a sample of 199 sources from the Bolocam Galactic Plane Survey (BGPS) with a beam of 31′. This corresponds well to the average dust
temperature also used in various studies of ATLASGAL selected sources (e.g. Csengeri et al. 2016a).

A total of 127 sources (corresponding to 58% of the sample) have kinetic temperature estimates from Wienen et al. (2012) and Wienen et al. (2017), and the bulk of these temperature measurements range between 17 and 35 K. Since these measurements do not cover the full sample, they only allow us to estimate a reasonable temperature correction. Furthermore, as these studies mostly rely on the (1, 1) to (2, 2) inversion transitions of NH$_3$, they are more sensitive to the cold gas rather than the hot gas component. Therefore, for the more evolved objects these temperatures should be considered as lower limits on the dust temperature, and thus an upper limit on the mass. When discussing the statistical properties of the sample (Sect. 3), as a second-order approximation, we account for the increasing temperature towards sources with embedded massive stars using a higher $T_d$ of 30 K. This is consistent with the typical values used so far of 15–20 K for cold sources and up to 30–40 K for warm sources (e.g. Motte et al. 2007; Russeil et al. 2010).

2.3. Estimating physical parameters: mass, size, density, and surface density

We use the same formula as described in Schuller et al. (2009) and Csengeri et al. (2014) to estimate the gas mass from the optically thin 870 µm dust emission

$$M_{\text{gas}} = \frac{S_\nu R d^2}{B_\nu(T_d) \kappa_\nu} \approx \frac{M_{\text{gas}}}{M_\odot} = C \times \frac{S_\nu}{[\text{Jy}]} \times \left(\frac{d}{[\text{kpc}]}ight)^2,$$

(2)

where $S_\nu$ is the integrated flux and $d$ corresponds to the distance of the sources. The numerical constant, $C_{18 K} = 6.33$ for $\nu = 345$ GHz, $T_d = 18$ K, and $\kappa_\nu = 0.0185 \text{ cm}^2 \text{ g}^{-1}$ from Ossenkopf & Henning (1994) accounting for a gas-to-dust ratio ($R$) of 100. We note that recent studies suggest a higher gas-to-dust ratio of 150 (Draine 2011), which would directly increase our mass estimates by 50%. Peretto et al. (2013) adopt a 50% lower $\kappa_\nu$, which would also lead to an increase in the estimated mass. Following the analysis in Sect. 2.2, we introduce a temperature corrected mass estimate ($M_{\text{corr}}$) for the more evolved and thus likely warmer objects adopting $T_d = 30$ K, corresponding to $C_{30 K} = 3.08$ from Sect. 3.

An additional factor of uncertainty in the mass determination originates from optically thick free-free emission which may contribute to the 870 µm flux density towards the clumps hosting UC-H$\alpha$ regions. For a uniform density H$\alpha$ region this contribution can be negligible because the thermal dust emission dominates the spectral energy distribution (Hunter et al. 2000). However, for a non-uniform density source with high temperature and density, a more significant contribution from free-free emission is possible (e.g. Beuther et al. 2004; Keto 2002, 2003). This may lead to an overestimate of the mass. Schuller et al. (2009), however, estimate this contribution to be considerably lower than the thermal emission from dust at 345 GHz, not more than 20% for the most extreme cases. In addition, the already formed ionising sources are still compact and therefore are expected to have a minor influence on the mass estimates. To conclude, for sources associated with known embedded UC-H$\alpha$ regions we conservatively consider the derived masses as upper limits (Table 4).

In Csengeri et al. (2014) the ATLASGAL sources have been fitted with a 2D Gaussian, from which we estimate the source size following the method of Bontemps et al. (2010). We define the source radius ($R_{90}$) which contains 90% of the total...
mass, assuming a 2D Gaussian profile, by $R_{90} = 1.95 \times FWHM_{\text{deconvolved}} / \sqrt{2 \ln 2}$, where $FWHM_{\text{deconvolved}}$ corresponds to the deconvolved source size. From the mass and size, we use the following equations to calculate the mean volume density and surface density:

\[ \bar{n} = \frac{M_{\text{gas}}}{\pi R_{90}^2} \]

\[ \Sigma = \frac{M_{\text{gas}}}{\pi R_{90}^2} \]

The estimated physical parameters are listed in Table 4.

### 2.4. Mid-infrared tracers of embedded protostars

Here we investigate the properties of the embedded YSOs closely associated with the dust peaks in order to reveal their star formation activity and, in particular, the youngest clumps within the sample. To classify an embedded source as a high-mass young stellar object, we rely on the IRAS-based colour properties for embedded high-mass stars from Wood & Churchwell (1989) following the method of Motte et al. (2007) and Russeil et al. (2010). Based on stellar models we adopt a luminosity threshold corresponding to objects with $L_{\text{bol}} > 10^4 L_{\odot}$, which potentially host a $> 15 M_{\odot}$ zero age main sequence (ZAMS) star with a spectral type earlier than B0.5 (see also Mottram et al. 2011). In practice, using the expression from Motte et al. (2007), we translate this luminosity to a flux density threshold of $\sim 289$ Jy at 21 \(\mu\)m for an object at 1 kpc. This limit is a factor of ten higher than was used for Cygnus-X (Motte et al. 2007) and NGC 6334 (Russeil et al. 2010), and we evaluate the implication of the higher luminosity limit in Sect. 2.5. We extrapolate between the 22 \(\mu\)m and 24 \(\mu\)m data assuming blackbody radiation for the embedded protostar following Russeil et al. (2010).

Below this luminosity threshold we find high-mass protostars or their potential precursors, which are still deeply embedded in their mass reservoir, while their energetics is dominated by accretion processes. Although theoretical models are strongly dependent on the accretion rates, they predict Class 0-like high-mass protostars up to $L_{\text{bol}} < 10^3 L_{\odot}$ with an envelope mass of $> 50 M_{\odot}$. Our threshold is also consistent with evolutionary tracks of high-mass stars assuming constant or varying accretion with rates of $10^{-4} - 10^{-3} M_{\odot}$ yr\(^{-1}\) (Duarte-Cabral et al. 2013; Kuiper & Yorke 2013). In fact, the potentially highest mass protostar in Cygnus-X, N63, has an envelope mass of $\sim 44 M_{\odot}$ and $L_{\text{bol}} \sim 339 L_{\odot}$ estimated from SED models using Herschel data (Duarte-Cabral et al. 2013).

To extract the mid-infrared flux density of the sample, we first use the most sensitive and highest angular-resolution data from the MIPS/GAL survey (Carey et al. 2009). The MIPS detector saturates at 2 Jy for point sources, and confusion due to nearby point sources or bright extended emission may hinder the photometry at the position of the dust peaks. Considering only the point source saturation limit, the MIPS/GAL survey reveals B3 type (or earlier) stars only beyond 3.8 kpc. Less luminous but closer objects are likely to appear saturated in the MIPS images. Therefore, in order to cover a continuous sensitivity range it is necessary to complement the MIPS/GAL data with data from the WISE satellite. The PSF photometry of the WISE point source catalogue is sensitive up to a $\sim 330$ Jy flux density threshold (Cutri et al. 2012).

The MIPS/GAL survey at 24 \(\mu\)m has a 3\(\sigma\) point source sensitivity of 2 mJy at a 6\arcmin\ angular resolution, while the WISE survey provides a 5\(\sigma\) point source sensitivity of 6 mJy at a 12\arcmin\ resolution at 22 \(\mu\)m in unconfused regions of the Galaxy. Owing to extended emission from nebulosity and polycyclic aromatic hydrocarbons (PAHs), these sensitivity levels may be higher towards confused regions. The factor of 4 difference in beam areas may also result in different flux densities determined for the same source from the MIPS/GAL and WISE data. In Appendix A we discuss the cross-calibration between the MIPS and WISE photometry.

Although ultimately the sensitivity threshold is only a factor of few better for the MIPS/GAL survey, its angular-resolution (6\arcmin\), which is more than three times higher, is a key benefit in our study. This angular resolution corresponds to a physical scale of $\sim 0.14$ pc at a 4.5 kpc distance, which is the size scale of massive dense cores embedded in massive clumps (Motte et al. 2007; Zhang et al. 2009; Csengeri et al. 2011b). This angular resolution allows us to resolve and distinguish between infrared-bright and -quiet massive dense cores for objects located at distances closer than $\sim 4.5$ kpc.

All sources were therefore first visually inspected using the publicly available calibrated data tiles from the MIPS/GAL survey, and where no point source was found within 6\arcmin\ of the ATLASGAL peak position, we extracted an upper limit on the flux density (for more details see Appendix A.1). Aperture photometry was performed on point sources where no saturated pixels are present (Appendix A.2). Where the MIPS/GAL data was not useable, we used the flux density from the WISE point source catalogue Wright et al. (2010, Appendix A.3)\(^3\).

In Fig. 2 we show two examples of massive clumps with (no or weak) embedded mid-infrared sources, the upper panel showing the WISE 22\(\mu\)m image of G340.9698-1.0212 and the lower panel showing the MIPS/GAL 24\(\mu\)m image towards G34.43+00.24 (Garay et al. 2004; Rathborne et al. 2005); the rest of the sample is shown in Appendix B, Fig. B.1. This particular example shows the importance of angular resolution because this region hosts an already formed UC-H II region in the vicinity of G34.4+0.2-MM1 and MM2, while the brightest dust peak itself is clearly shifted from the position of the UC-H II region (Rathborne et al. 2005). The other embedded dust peaks in this clump are active sites of high-mass star formation in their earliest phase (Rathborne et al. 2008).

In Fig. 3, we show the resulting photometry measurements versus clump mass using $T_d = 18$ K (see Sect. 2.3). The photometry is scaled to a common wavelength of 22 \(\mu\)m and to a distance of 1 kpc. Using a homogeneous dust temperature facilitates the comparison with the larger sample of sources discussed in Sect. 2.5. While the brightest mid-infrared sources correspond to active high-mass star-forming sites in the inner Galaxy, we reveal a population of massive clumps which are clearly fainter at mid-infrared wavelengths (see also Sect. 2.6).

### 2.5. Tracers of embedded high-mass stars

According to stellar evolution models, stars of types B3–B0.5 may also have luminosities between $L_{\text{bol}} = 10^3 - 10^4 L_{\odot}$, hence they fall below our threshold introduced in Sect. 2.4. We assess their potential contamination of the sample of the youngest massive clumps by using sensitive radio-surveys such as CORNISH (Hoare et al. 2012; Purcell et al. 2013) and

\[ F_{22,\mu m} = 8.2839 \times 10^{\text{flux density}}^{-2.5}. \]

\(2\) We use 19\arcsec\ for the resolution of the ATLASGAL maps.

\(3\) The WISE 22 \(\mu\)m flux density is calculated from the W4 band magnitudes by using the formula $F_{22,\mu m} = 8.2839 \times 10^{\text{mag.W4}}^{-2.5}$. 

A60, page 5 of 25
Urquhart et al. (2007) to search for ionising emission from stars in our sample of types later than B0.5. Since there is no homogeneous high angular-resolution radio dataset available covering the entire sample towards the IVth Galactic quadrant, we estimate a contamination level by performing the same analysis as above for a larger sample of CORNISH sources.

In Fig. 4 we show the same plot of distance scaled photometry at 22 µm versus mass as in Fig. 3 (see Sect. 2.4), but for a larger sample of ATLASGAL sources hosting embedded UC-H II regions from Urquhart et al. (2013b), as well as MYSOs from the RMS survey (Lumsden et al. 2013). This sample has been studied in multwavelength follow-up observations (for a summary, see Lumsden et al. 2013), and the sources have been classified according whether they host (M)YSOs, H II regions, or a mixture of both. We extracted the 22 µm fluxes for these samples from the WISE catalogue considering matches between the ATLASGAL dust peaks and mid-infrared point sources with a small angular separation (<6′′). We considered only good quality flag photometry and a high signal-to-noise ratio of >10 in the 22 µm band. More importantly, for the RMS and CORNISH sources we only show those with the most reliable distance estimates based on maser parallax measurements (Reid et al. 2014).

We find that massive clumps ($M_{\text{clump}} > 650 M_\odot$) hosting embedded UC-H II regions occupy a well-defined region, and mostly lie above our threshold of $S_{22 \mu m} > 289$ Jy at 1 kpc. Out of the 78 sources of the shown ATLASGAL-CORNISH matches, 5 lie below this threshold, corresponding to 6.4% of the entire sample. Such weak, compact radio continuum emission has been found towards low bolometric luminosity ($L_{\text{bol}} \leq 5 \times 10^4 L_\odot$) MYSOs (see e.g. Guzmán et al. 2010; Lumsden et al. 2013; Avison et al. 2015). They are likely associated with shock-ionisation from an early outflow activity (e.g. Hoare & Franco 2007).

Our flux limited selection of ATLASGAL sources with embedded UC-H II regions are also shown as a comparison, and find that 3 sources out of the 31 fall below our mid-infrared threshold, which is 9.7%. We inspected these sources visually, and find that one source lies in a confused region where either the mid-infrared photometry or the radio continuum may be affected by extended emission, or the radio continuum emission may be associated with the extended H II region rather than the dust peak. The rest could correspond to the above discussed genuine low bolometric luminosity radio sources associated with MYSOs. To summarise, the
lowest luminosity massive clumps likely lack embedded UC-H\textsc{ii} regions and bright MYSOs, and their contamination below our luminosity threshold is estimated to be less than ten per cent.

2.6. Classification of the sample

Using the information described in the previous sections, we classify our sample primarily based on their mid-infrared flux density and indicate the presence of UC-H\textsc{ii} regions if radio continuum data is available. We do not rely on radio continuum data, however, because sensitive unbiased surveys are not yet available in the IVth quadrant, and thus the sample of UC-H\textsc{ii} regions associated with our clumps is incomplete. We introduced the threshold of 289 Jy at 22\,$\mu$m at 1 kpc as discussed in Sect. 2.4, and as empirically justified in Sect. 2.5. Sources above and below this threshold are termed “infrared bright” and “infrared quiet”, respectively. This classification is consistent, but simplified compared to previous works (e.g. Csengeri et al. 2014; Giannetti et al. 2014; Csengeri et al. 2016a). In addition, it quantitatively selects the young clumps based on the luminosity of the embedded protostar(s). Contamination of more evolved objects is estimated to be less than 10\% (Sect. 2.5). Nevertheless, in the following we indicate the subsample of massive clumps hosting genuine UC-H\textsc{ii} regions for the infrared-bright clumps.

From the sample of 210 sources we were able to classify 209 sources: 139 (66\%) are infrared-bright including 51 (24\%) unambiguously associated with UC-H\textsc{ii} regions, and 80 (34\%) are infrared quiet. Considering associations with infrared sources only from the WISE catalogue, we find 56 sources (26\%) with no counterpart; the remaining 74\% are associated with a mid-infrared source, which is consistent with earlier findings from Csengeri et al. (2014). Some objects are located in regions with complex extended emission at mid-infrared wavelengths, which may affect the photometry, hence their nature being quiet or bright may remain somewhat uncertain. As shown before, this definition of mid-infrared-quiet clumps picks up massive clumps hosting high-mass protostars or their precursors; where no ancillary radio data is available, the contamination from 8\,GHz submillimetre sources and bright MYSOs, and their contamination below our luminosity threshold is estimated to be less than ten per cent.

Fig. 5. SiO (8–7) line observed towards the infrared-quiet massive clump G340.9698–1.0212. The dashed line shows the rest velocity of the source from the C\textsuperscript{18}O (3–2) observations at 337.06110 GHz from the same dataset.

The data have been reduced using the GILDAS software following standard procedures. The spectra at the centre, i.e. the dust peak position, have been extracted from the maps. Baselines of order zero have been subtracted and the data have been scaled to the $T_{\text{mb}}$ temperature using 73\% efficiency\textsuperscript{4}. The average noise level is 48 mK in 1.3 km s\textsuperscript{-1} resolution.

A total of 27 out of the 46 objects (59\%) show emission of this line above 4$\sigma$ as measured on the line area with Gaussian fitting. We show an example in Fig. 5, while all the detections are shown in Appendix C, Fig. C.1. The observed detection rate is similar to that of Klaassen et al. (2012) despite the lower sensitivity, and is somewhat lower than reported by Leurini et al. (2014), who had, on the other hand, a factor of ~2 higher sensitivity for the same transition. These studies focused on more evolved, mostly infrared-bright clumps, many of them hosting UC-H\textsc{ii} regions. The fact that the high detection rates are broadly similar towards the less luminous thus likely younger clumps is consistent with the findings of Csengeri et al. (2016a) using lower-J SiO lines.

3. Statistical properties of the brightest submillimetre sources

In the following we investigate the distribution and the statistics of the estimated physical properties of the entire sample. As defined in Sect. 2.6 we distinguish between infrared-quiet massive clumps and those classified as infrared-bright, and also show the subsample hosting MYSOs and/or UC-H\textsc{ii} regions. Their average physical properties are summarised in Table 2.

3.1. Mass–distance distribution and completeness

In Fig. 6 we show the distribution of the sample. As discussed in Sect. 2.1, a substantial fraction of the sources are part of active star-forming complexes. Their typical distance is around 3–4 kpc, close to the average distance of 4.5 kpc for ATLASGAL clumps (Wienen et al. 2012). Overall 60\% of the sample is located within 4.5 kpc; a total of 86 sources are located at greater distances, of which 29 are infrared quiet.

\footnote{\url{http://www.apex-telescope.org/telescope/efficiency/}}

http://www.apex-telescope.org/telescope/efficiency/
more distant sources are therefore dominated by infrared-bright clumps. We also show a threshold of 650 $M_\odot$, introduced in Csengeri et al. (2014), which is an extrapolation of the massive dense core definition by Motte et al. (2007) to the $0.3$ pc scale of ATLASGAL clumps. Empirically this selection has proved to select cores hosting high-mass protostars by Bontemps et al. (2010).

The mass estimation corresponding to the threshold of 5 Jy used for our flux limit is indicated in Fig. 6 as green and red dotted lines for infrared-bright ($T_d = 30$ K) and infrared-quiet clumps ($T_d = 18$ K), respectively. It shows that our selection is complete for infrared-quiet clumps above 650 $M_\odot$ up to a distance of 4.5 kpc, while for infrared-bright clumps above 310 $M_\odot$, since we expect the more evolved clumps to have lost a part of their mass in the central region through cluster formation and dissipation, we can argue that this sample is roughly complete for all clumps having 650 $M_\odot$ at birth at all their stages of evolution from infrared quiet to infrared bright and possibly also UC-H II regions. A total of 28 infrared-quiet and 55 infrared-bright clumps are found above 650 $M_\odot$ and 310 $M_\odot$ at a distance closer than 4.5 kpc. The area covered by the sector of our Galaxy from $\ell = -60^\circ$ to $60^\circ$ and for a distance lower than 4.5 kpc is found to represent $\sim 10\%$ of the total area of the galactic disk within the solar circle after removing the inner 3 kpc. In a rough approximation, we can consider this 10% as the fraction of star formation in the Galactic disk which is occurring in our distance limited sample.

3.2. Temperature distribution

From the NH$_3$ data discussed in Sect. 2.2 we determine a median kinetic temperature of 25.5 K for 60% of the sample. Figure 7 shows the temperature distribution of these sources. In this sample 28 have a firm association with UC-HII-regions, and the majority of the sources are classified as infrared bright, in total 71. These types of sources have a median kinetic temperature of 29.6 and 27.2 K, with a standard deviation of 8–11 K. Towards infrared-quiet sources we find a median of 23.8 K with a standard deviation of 4.7 K. This suggests a trend of lower temperatures towards younger sources; however, given their large standard deviation, this difference is not statistically significant. Following our argumentation in Sect. 2.4, and considering the similar temperatures, here we do not make a distinction between massive clumps with UC-HII regions and those hosting infrared-bright sources. A more accurate temperature assignment to the group of infrared-quiet massive clumps still remains uncertain since the sources with available NH$_3$ measurements are dominated by the more evolved sources, and the targeted infrared-quiet clumps may not be representative of the entire subsample.

3.3. Distribution of physical properties

In Fig. 8 we show the mass distribution of the sample, where we adopted $T_d = 30$ K for the infrared-bright sources and kept $T_d = 18$ K for the infrared-quiet clumps (see Sect. 2.3). We give an overview of the statistical properties of the physical parameters of all sources and the two source types in Table 2. The global statistical properties of the infrared-bright and -quiet samples are similar, and we find a remarkably flat mass distribution for all source types, the majority of the sources lying between 650 $M_\odot$ and 8000 $M_\odot$. Lower mass objects correspond to nearby massive dense cores with masses between $\sim 86$–$120 M_\odot$. Towards sources at greater distances we find masses of $\sim 10^4 M_\odot$ corresponding to cloud fragments rather than clumps. The mean clump mass is $\sim 4.5 \times 10^3 M_\odot$, which is about a factor of two higher compared to the star formation selected larger sample of ATLASGAL clumps studied by Urquhart et al. (2014).

The size distribution of the sample is shown in Fig. 9, and as above, there is no significant difference between the different source types. Adopting the nomenclature of Williams et al. (2000, see also Bergin & Tafalla 2007), we find a small fraction of sources that correspond to cores or entire cloud fragments,
Table 2. Statistics of the physical properties of all sources, and the infrared-bright and infrared-quiet samples.

<table>
<thead>
<tr>
<th></th>
<th>All sources</th>
<th>Infrared-bright massive clumps</th>
<th>Infrared-quiet massive clumps</th>
</tr>
</thead>
<tbody>
<tr>
<td>Distance [kpc]</td>
<td>Min 1.00</td>
<td>Max 15.6</td>
<td>Min 1.0</td>
</tr>
<tr>
<td></td>
<td>Max 14.20</td>
<td>Mean 4.9</td>
<td>Max 14.6</td>
</tr>
<tr>
<td></td>
<td>Median 3.9</td>
<td>Stdev 3.30</td>
<td>Median 3.70</td>
</tr>
<tr>
<td></td>
<td></td>
<td>R90 [pc]</td>
<td>R90 [pc]</td>
</tr>
<tr>
<td>M_{18K} [M_\odot]</td>
<td>Min 1.20</td>
<td>Max 5.83</td>
<td>Min 1.27</td>
</tr>
<tr>
<td></td>
<td>Median 5.0</td>
<td>Stdev 5.9</td>
<td>Median 5.87</td>
</tr>
<tr>
<td></td>
<td></td>
<td>σ_{18K} [g cm^{-2}]</td>
<td>σ_{18K} [g cm^{-2}]</td>
</tr>
<tr>
<td></td>
<td>Min 0.27</td>
<td>Max 5.98</td>
<td>Min 0.27</td>
</tr>
<tr>
<td></td>
<td>Median 1.11</td>
<td>Stdev 0.75</td>
<td>Median 0.99</td>
</tr>
<tr>
<td>σ_{corr} [×10^3 cm^{-3}]</td>
<td>Min 0.14</td>
<td>Max 44.15</td>
<td>Min 0.14</td>
</tr>
<tr>
<td></td>
<td>Median 3.17</td>
<td>Stdev 1.56</td>
<td>Median 1.56</td>
</tr>
</tbody>
</table>

Fig. 8. Temperature corrected mass distribution (using T_d = 30 K for the infrared-bright clumps). Colours are the same as in Fig. 7. The clumps hosting genuine UC-H II regions are a subsample of the infrared-bright clumps.

Fig. 9. Distribution of the deconvolved physical size (R_90) of the sample. Colours are the same as in Fig. 7. The dashed lines indicate an approximate distinction between cores, clumps, and cloud fragments (Bergin & Tafalla 2007).

while the majority of the sources correspond to clumps. They have typical sizes of 0.3–0.4 pc, which is comparable to the mean size of ~0.36 pc found by Russeil et al. (2010) in NGC 6334/NGC 6357. Compared to the population of massive dense cores in Cygnus-X, the sizes are ~3 times larger, while the masses are on average ~5 times larger, suggesting their potential to form high-mass stars.

In Fig. 10 we show the mass surface density distribution of the sample. Lower surface densities are found towards the infrared-bright sources, a considerable fraction of which host UC-H II regions. This is likely due to an observational bias, driven both by a larger fraction of such sources at a larger distance, which then correspond to longer structures with lower average surface density, and the lower mass selection threshold for the infrared-bright sources. The mean surface density is above the theoretical value of \( \Sigma_{\text{cl}} \geq 1 \text{ g cm}^{-2} \) by Krumholz & McKee (2008), suggesting that a substantial fraction of our sources have the potential to form high-mass stars. Correspondingly, the median volume-averaged density of 3.13 × 10^5 cm^{-3} is higher than that found in other samples in the literature (e.g. Russeil et al. 2010), and suggests that we probe here high density, star-forming gas.

4. Nature of the brightest submillimetre sources in the inner Galaxy

The brightest submillimetre sources in the inner Galaxy are generally associated with active sites of star formation. Here we
first investigate the star formation potential and evolutionary aspect of the entire sample (Sect. 4.1), and argue that the majority of the infrared-quiet massive clumps are potential precursors of massive stars and/or clusters. We then show that our sample of infrared-quiet clumps contains sources previously not recognised in other surveys using infrared emission, or absorption (Sect. 4.2). Finally, we search for signposts of embedded protostars in the infrared-quiet subsample to assess the stage of ongoing star formation (Sect. 4.3). We then discuss the stability of the entire sample (Sect. 4.4), and provide predictions for their small-scale fragmentation (Sect. 4.5).

### 4.1. Precursors of massive clusters

Comparing the physical properties in terms of mass, size, and density, in Table 2 we show that our sample exhibits rather homogeneous properties, and no statistically significant differences are found between the two source types. This includes well-studied Galactic mini-starburst regions, for instance the W43 ridge (e.g. Motte et al. 2003; Bally et al. 2010; Nguyen Luong et al. 2011), W51 Main (e.g. Mufson & Liszt 1979; Zapata et al. 2008; Ginsburg et al. 2015), and W49N (e.g. Peng et al. 2010; Galván-Madrid et al. 2013). While indications of ongoing star formation is prevalent in the sample, are all clumps prone to form mini-starbursts and rich clusters? We compare our results with the empirical relation of Kauffmann & Pillai (2010) in Fig. 11, which suggests that all of our sources have the potential to form high-mass stars. While the comparative statistical analysis of infrared-quiet and -bright sources reveals similar physical characteristics in terms of mass and size, they seem to differ only in their embedded (proto-)stellar population. Low bolometric luminosity sources with similar physical characteristics to their more evolved counterparts have been frequently interpreted as being in an earlier evolutionary stage (e.g. Motte et al. 2007; Molinari et al. 2008). This suggests that infrared-quiet clumps are the precursors of the infrared-bright clumps and UC-H II regions, and the sample can be arranged along a similar evolutionary path.

The large mass reservoir of infrared-bright sources suggests that they likely continue producing high-mass stars. Unless specific conditions are required for the formation of rich clusters, we shall have selected here their true precursors.

---

**Fig. 10.** Mass surface density (Σ) distribution of the sample using $T_d = 30$ K for the infrared-bright clumps. Colours and symbols are the same as in Fig. 7.

**Fig. 11.** Temperature corrected mass estimates (using $T_d = 30$ K for the infrared-bright clumps) versus deconvolved physical size ($R_{90}$) from Sect. 2.3. Dashed lines show the relation of Kauffmann et al. (2010) scaled to the Ossenkopf & Henning (1994) dust opacity following Dunham et al. (2011). In the upper left corner we show the factor of two uncertainty in the mass estimate. Colours are the same as in Fig. 7.

### 4.2. Massive cluster progenitors in isolation?

Because the clustered nature of high-mass star formation was recognised early on, several studies searched for younger objects in the vicinity of bright, luminous sources (e.g. Garay et al. 2004; Thompson et al. 2005). To look for neighbouring, more evolved objects with UC-H II regions in the close vicinity of the infrared-quiet sample, we used the CS (2–1) survey of IRAS sources (Bronfman et al. 1996), which is complete towards IRAS selected UC-H II regions in the Galactic plane. In total we found 92 associations (42%) within 0.8 arcmin offset, which corresponds to 0.26 pc for the nearest source and 2.9 pc for the most distant source in our sample. However, in the infrared-quiet sample only 29 (36%) are found in the vicinity of an IRAS source, and their offsets correspond to a physical separation between 0.3–16.3 pc. Compared with the more sensitive and higher angular resolution Red MSX survey (RMS, Lumsden et al. 2013), we find six infrared-quiet clumps in the vicinity of a bright MYSO with $L_{bol} > 2 \times 10^7 L_{\odot}$.

Surveys sensitive to cold dust are better suited to directly probe the youngest clumps. To evaluate the fraction of previously identified infrared-quiet clumps, we first looked for associations with cloud structures identified as infrared-dark clouds (IRDCs) by correlating the peak dust positions with the catalogues of IRDCs identified in the literature. In total we find 7 associations with MSX (Jackson et al. 2008; Simon et al. 2006) and Spitzer (Peretto & Fuller 2009). In total we find 20 associations with an MSX Dark Cloud within 120 $′′$, 15 of which are infrared-quiet sources. For the association with Spitzer Dark Clouds we used the Herschel-confirmed clouds from Peretto et al. (2016) with a smaller positional offset of 50 $′′$ corresponding to their average equivalent size of $25′′ + 1\sigma$, where $\sigma$ is their standard deviation. In total, 65 of our clumps have associations, 31 of which are infrared-quiet clumps. This means that ~61% of the infrared-quiet clumps have not been recognised before in these surveys.

Thus, we find that the infrared-quiet sample is found in relative isolation compared to the bright clumps which already undergo active star formation. Less than half of the ATLASGAL selected massive clumps have IRDC counterparts, and comparison with the star formation selected ATLASGAL clumps of
Urquhart et al. (2014) shows that 25 of the infrared-quiet sources are newly identified. In summary, we find that ~23% of our infrared-quiet sample has not been recognised by these previous studies.

### 4.3. Signposts of already formed embedded high-mass protostars

To tackle the stage of ongoing star formation, here we compare our sample with tracers of high-mass star formation. We found a total of 23 clumps coinciding with extended green objects (EGOs); (Cyganowski et al. 2008; and Chen et al. 2013) most of them are infrared-quiet clumps, 2 sources are infrared-bright, and 3 are UC-H II regions. Since EGOs are believed to trace powerful outflows mostly from (presumably massive) Class 0 YSOs, it confirms that the infrared-quiet clumps host the youngest protostellar phases. However, only 25% of the infrared-quiet clumps are associated with EGOs. This may indicate that the EGO phase is shorter than the infrared-quiet clump phase, and therefore does not represent well the entire early evolutionary phase of massive clumps. The few EGOs associated with infrared-bright clumps may indicate that young massive protostars can still form at the time of the infrared-bright phase, and even perhaps at the UC-H II phase if star formation is a continuous process in clumps.

On the other hand, the lack of EGOs does not necessarily imply a lack of outflow activity, since high extinction, even at 4.5 μm, may prevent their detection. In our SiO (2–1) spectroscopic study, ~60% of infrared-quiet clumps show emission, which is a much larger fraction than that of EGOs. Using the lower energy SiO (2–1) line, Csengeri et al. (2016a) find that an even higher fraction of the infrared-quiet massive clumps exhibit emission in this shocked gas tracer. This is consistent with a picture that infall and accretion processes may already be occurring in at least some of the infrared-quiet clumps, indirectly pointing at an underlying population of deeply embedded protostars (see the example of NGC 6334 I(N) by Megeath & Tieftrunk 1999).

Another tracer of the nature of embedded protostars is Class II methanol maser emission. These masers are excited by radiative pumping requiring the intense mid-infrared emission from the warm, dense material of deeply embedded hot-cores or UC-H II regions (Menten 1991; Sobolev et al. 1997). We therefore cross-correlated our sample with the Methanol MultiBeam survey (Caswell et al. 2010, 2011; Green et al. 2010; Breen et al. 2015), and found 138 (66%) associations within 20″, of which 57 are infrared-quiet massive clumps corresponding to ~71% of the subsample. Methanol masers thus seem to trace a larger fraction of the infrared-quiet phase, although less selectively since they are also found towards the more evolved stages of clump evolution. Altogether we find that a considerable fraction of the infrared-quiet sample show signs of embedded (high-mass) protostars and ongoing star formation process.

### 4.4. Stability estimates and timescales

Using high-density tracers we estimate here the virial mass and compare it to the mass estimates based on the dust emission. As a first step we used the NH$_3$ (1, 1) line width measurements of Wienen et al. (2012) and Wienen et al. (2017) where available, and then complemented these data with measurements of the CS (2–1) line from Bronfman et al. (1996), and finally of the H$^{13}$CO$^+$ (1–0) line from recent surveys using the Mopra (Wyrowski et al., in prep.) and the IRAM 30 m telescopes (Csengeri et al. 2016a). We have line width estimates for a total of 173 sources, 80% of the whole sample. While this is a smaller sample than studied in Wienen et al. (2012) and Urquhart et al. (2014), it is more focused on the brightest ATLASGAL clumps.

We estimate the virial mass assuming a density profile of $n(r) \sim r^{-2}$, and using $M_{\text{vir}} = 3 \times \frac{R_90 \times \sigma_{\text{tot}}^2}{G}$ (Bertoldi & McKee 1992), where $R_90$ is the radius of a sphere, $\sigma_{\text{tot}}$ is the velocity dispersion of the gas, and $G$ is the gravitational constant,

$$M_{\text{vir}} = \frac{697 \times R_90}{[\text{pc}]} \times \frac{\sigma_{\text{tot}}^2}{[\text{km s}^{-1}]}.$$  \hspace{1cm} (5)

As shown in Fig. 12, we find surprisingly small $M_{\text{vir}}/M$ ratios towards the majority of the sources, on average 0.47 with a median of 0.29, strongly suggesting that they are gravitationally unstable. While the used molecular tracers selectively probe the high-density gas, uncertainties due to the different tracers and beam sizes could affect our virial mass estimates. However, despite the various datasets used, we see no systematics in the estimated virial masses. As Fig. 12 shows, 155 sources (89%) have a smaller virial mass than their estimated gas mass. We find the same trend for the more distant sources of the sample, which are beyond our completeness limit. Uncertainties in the mass estimates due to the dust opacity and gas-to-dust ratio would lower the $M_{\text{vir}}/M$ ratio by a factor of two. Even considering the case of a uniform $n(r) \sim r^{-1}$ density profile, we would still find a large fraction (67%) of the clumps having $M_{\text{vir}}/M < 1$, which could thus still be considered as gravitationally unstable. Magnetic fields could, however, slow down the rate of collapse. For a large number of clouds Kauffmann et al. (2013) provides estimates of the virial parameter defined by $M_{\text{vir}}/M$, and find similarly low values. They suggest that magnetic fields as strong as 1 mG would be required to slow down the collapse, otherwise they undergo rapid collapse.

The size-scale of clumps, 0.2–1 pc, is large compared to the scale of MDCs (<0.15 pc) and cores (<0.05 pc) for which collapse signatures are frequently observed. Our results therefore suggest that these massive cluster progenitors are already collapsing on a global scale, i.e. a clump scale. This is particularly intriguing because it suggests that gravity could be the...
Table 3. Comparison of the properties of the infrared-quiet sources with the literature.

<table>
<thead>
<tr>
<th>ATLASGAL infrared-quiet(^d) clumps</th>
<th>Herschel-IRDC(^a) clumps</th>
<th>IRDC clumps</th>
<th>CygX clumps</th>
</tr>
</thead>
<tbody>
<tr>
<td>Distance [kpc]</td>
<td>1.15–15.60 (4.47)</td>
<td>0.12–16.45 (4.18)</td>
<td>1.8–7.1 (3.7)(^b)</td>
</tr>
<tr>
<td>Size [pc]</td>
<td>0.08–16.3 (0.36)</td>
<td>0.02–2.6 (0.6)</td>
<td>1–3(^d)</td>
</tr>
<tr>
<td>(M_{\text{C}}) [M(_{\odot})]</td>
<td>200–21 000 (1400)</td>
<td>0.11–16 000 (700)</td>
<td>120–16 000 (950)(^b)</td>
</tr>
<tr>
<td>(\rho_{\text{crit}} \times 10^5 \text{ cm}^{-3})</td>
<td>0.20–18.46 (2.39)</td>
<td>0–0.58 (0.01)</td>
<td>0.001–0.1(^b)</td>
</tr>
</tbody>
</table>

Notes. (\(^a\)) Range of minimum to maximum values, the median is given between parentheses. (\(^b\)) For the adopted dust temperature see their original work. (\(^c\)) Adopted from the values listed in Table 2 from Traficante et al. (2015). (\(^d\)) Adopted from the values listed in Table 3 from Rathborne et al. (2006). (\(^e\)) Adopted from the values listed in Table 1 from Rathborne et al. (2006). (\(^f\)) Distance from Rygl et al. (2012). (\(^g\)) Scaled to 1.4 kpc from Motte et al. (2007).

dominant force in the mass assembly process over size-scales of entire clumps, which are larger than previously considered (e.g. McKee & Ostriker 2007).

We also calculate the free-fall times of these clumps by \(t_{\text{ff}} = \sqrt{32 \pi G n}, \) where \(G\) is the gravitational constant, and \(n\) is the volume density. We find on average \(5.64 \times 10^4\) yr, which is somewhat shorter than the lifetime estimate of \(7.5 \times 10^4\) yr by Csengeri et al. (2014). Such relative collapse timescales (i.e. a few times the free-fall timescale) are typically observed towards low-mass cores (see Ward-Thompson et al. 2007, and references therein), and have recently been suggested for the formation of high-mass cores by Duarte-Cabral et al. (2013).

Collapse timescales of a few times \(10^4\) yr suggests that all massive clumps, precursors of massive clusters, are transient entities and should disappear on short timescales, typically less than \(10^5\) yr (see also Zinnecker & Yorke 2007; Motte et al. 2007). Embedded MYSOs and UC-H\(\Pi\) evolve, however, on longer timescales, of the order of \(10^6–10^7\) yr according to stellar evolution models with accretion rates of \(M_{\text{acc}} \sim 10^{-3} M_{\odot}\) yr\(^{-1}\) (e.g. Churchwell 2002; Hoare et al. 2007; Kuiper & Yorke 2013). This is at least up to an order of magnitude longer than estimated for the collapse of their parent clump. This means that the clump material should have already collapsed leaving a considerably less dense and less massive clump towards the more evolved objects. This apparent contradiction to the observations can only be explained if there is mass replenishment onto the clump on large scales fuelling a continuous star formation, or if we selectively pick here the youngest of the galactic UC-H\(\Pi\) regions still surrounded by their natal clump. Case studies of massive cluster progenitors show evidence both for global collapse and for flows of dense gas (e.g. Schneider et al. 2010; Csengeri et al. 2011a,b; Peretto et al. 2013; Tackenberg et al. 2014). This points to a scenario where massive clumps undergo star formation, while gas replenishment arrives on larger scales, and maintains the observed similar mass range for these clumps.

4.5. Initial stages of massive proto-clusters: prediction for fragmentation properties

We overview the physical properties of our infrared-quiet sample in comparison with typical IRDCs (Rathborne et al. 2006), Herschel clumps embedded in IRDCs (Traficante et al. 2015), and clumps in the Cygnus-X high-mass star-forming region (Motte et al. 2007) in Table 3. While IRDCs are typically more extended structures, they have lower masses and exhibit smaller volume densities. Herschel selected clumps span a much broader range of physical conditions and cover similarly high-mass clumps as our sample. They are found to be, however, more extended and are therefore on average lower density structures. The ATLASGAL counterparts of IRDCs selected by Kainulainen & Tan (2013) correspond to sources below the flux density threshold used for this study. The ATLASGAL sources used here are the brightest and therefore the most massive among the compact millimetre sources in the inner Galaxy. Their typical size of \(\sim 0.4\) pc corresponds to protoclusters. We can therefore discuss our sample in terms of the precursors of the richest clusters in the Galaxy.

In Fig. 13 we show the temperature corrected mass versus the surface density distribution of the entire sample following Tan et al. (2013), with different colours corresponding to the infrared-bright and infrared-quiet subsamples. It is clear that the sample discussed here lies, on average, close to the theoretical limit for high-mass star formation around \(\Sigma \sim 1 \text{ g cm}^{-2}\) (Krumholz & McKee 2008) supporting the scenario that it is capable of forming high-mass stars and clusters. Taking the two extreme profiles of constant density \(n(r) \sim r^0\) and power law with \(n(r) \sim r^{-2}\), we mark with stars the predicted fragmentation on scales smaller by a factor of ten, approaching the size-scales of a few thousand AU corresponding to individual collapsing objects (e.g. Bontemps et al. 2010). The uniform density profile, where the surface density stays constant, predicts fragments below solar mass, which would correspond to clumps forming a rich cluster of low-mass stars. The case of the steeper density profile predicts more massive structures on smaller scales, reaching a mass range of several tens of solar masses. Envelope masses of this order are expected to form high-mass stars.

Some of the well-studied sources of our selection have been subject to high angular resolution follow-ups. The extreme clumps of the W43 complex, MM1 to MM3 (Motte et al. 2003; Sridharan et al. 2014) hosts one of the most extreme sites of active high-mass star formation in terms of mass and youth (Louvet et al. 2014). Massive envelopes have been identified towards other clumps as well, for example G35.03+0.35 (G35.0250+0.3501) (Beltrán et al. 2014), G35.20–0.74 (G35.1976–0.7427) (Sánchez-Monge et al. 2013), and SDC335-MM1(G335.5857–0.2906) (Peretto et al. 2013). High angular resolution observations with ALMA towards a large fraction of the infrared-quiet sources of the presented sample indeed reveal a significant population of massive dense cores (Csengeri et al. 2017).

5. Summary and conclusions

We investigate the properties of a flux limited sample of massive clumps from the ATLASGAL survey. Based on mid-infrared photometry, we divide the sample into infrared-quiet and infrared-bright clumps. Infrared-bright clumps peak on
Table 4. Flux limited sample of ATLASGAL sources with distance estimate and the corresponding dust properties.

<table>
<thead>
<tr>
<th>ATLASGAL*</th>
<th>Gal. long.</th>
<th>Gal. lat.</th>
<th>$F_{\nu}$</th>
<th>$S_\nu$</th>
<th>$\theta_{\text{maj}} \times \theta_{\text{min}}$</th>
<th>$P_{\text{PA}}$</th>
<th>$d_{\text{maj}}$</th>
<th>$d_{\text{min}}$</th>
<th>Ref.</th>
<th>$M_{\text{dust}}$</th>
<th>$R_{\text{dust}}$</th>
<th>$\zeta$</th>
<th>$n$</th>
<th>$\Delta R_{\text{dust}}$</th>
<th>$\sigma$ (cm$^{-3}$)</th>
<th>$# \mathrm{IRQ}$</th>
<th>$# \mathrm{SRC/DRC}$</th>
<th>EGO</th>
</tr>
</thead>
<tbody>
<tr>
<td>24.0 15.0</td>
<td>138.6 69.8</td>
<td>138.6 69.8</td>
<td>5000 5000</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td></td>
</tr>
<tr>
<td>30.4 60.9</td>
<td>30.4 60.9</td>
<td>30.4 60.9</td>
<td>5000 5000</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td></td>
</tr>
<tr>
<td>36.8 66.8</td>
<td>36.8 66.8</td>
<td>36.8 66.8</td>
<td>5000 5000</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td></td>
</tr>
<tr>
<td>43.2 62.2</td>
<td>43.2 62.2</td>
<td>43.2 62.2</td>
<td>5000 5000</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td></td>
</tr>
<tr>
<td>49.6 58.3</td>
<td>49.6 58.3</td>
<td>49.6 58.3</td>
<td>5000 5000</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td></td>
</tr>
<tr>
<td>56.0 54.4</td>
<td>56.0 54.4</td>
<td>56.0 54.4</td>
<td>5000 5000</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td></td>
</tr>
<tr>
<td>62.4 50.5</td>
<td>62.4 50.5</td>
<td>62.4 50.5</td>
<td>5000 5000</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td></td>
</tr>
<tr>
<td>68.8 46.6</td>
<td>68.8 46.6</td>
<td>68.8 46.6</td>
<td>5000 5000</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td></td>
</tr>
<tr>
<td>75.2 42.7</td>
<td>75.2 42.7</td>
<td>75.2 42.7</td>
<td>5000 5000</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td></td>
</tr>
<tr>
<td>81.6 38.8</td>
<td>81.6 38.8</td>
<td>81.6 38.8</td>
<td>5000 5000</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td>0.3 0.3</td>
<td>100 100</td>
<td></td>
</tr>
</tbody>
</table>

Notes: The full table is available only in electronic form at the CDS. Column 1 gives the corresponding ATLASGAL source name from Csengeri et al. (2014) including the Galactic coordinates, whileCols. 2 and 3 list the Galactic coordinates. Columns 4, 5 give the peak flux density, and the integrated flux density assuming a Gaussian flux distribution from Csengeri et al. (2014), respectively. Columns 6–7 list the major and minor axis of the fitted Gaussian FWHM to the source, and the position angle. Column 8 gives the used distance estimate. Columns 9, 10 give the minimum and the maximum distances. Column 11 gives the reference for the distance. Columns 12–15 give the physical parameters of the sources (see text for details.) (1) We revised the spectrophotometric distance of 1.81 kpc by (Moisés et al. 2011). (1) Sources also included in the ATLASGAL top 100 selection Gianetti et al. (2014).

References. (1) This work; (2) Sato et al. (2010); (3) Zhang et al. (2013); (4) Zhang et al. (2014); (5) Immer et al. (2013); (6) Sanna et al. (2014); (7) Moisés et al. (2011); (8) Wu et al. (2014); (9) Wu et al. (2012); (10) Reid et al. (2014); (11) Dutra et al. (2003); (12) Stray et al. (1987); (13) Xu et al. (2011); (14) Sanna et al. (2009); (15) Leinumi et al. (2011); (16) Kurayama et al. (2011); (17) Russell et al. (2003); (18) Reid et al. (2009); (19) Brunthaler et al. (2009); (20) Urquhart et al. (2014b); (21) Krishnan et al. (2015); (22) Motogli et al. (2011).
embedded objects with $L_{bol} > 10^4 L_\odot$ corresponding to MYOSs or high-mass stars with UC-H II regions, while infrared-quiet clumps lack such bright embedded objects. From the initial sample of 210 sources above a 5 Jy beam-averaged peak flux density, we identify 80 infrared-quiet massive clumps, of which 50 are located within $d < 4.5\,\text{kpc}$. About one-fourth of them have not been recognised by other surveys.

We estimate their physical properties based on the dust emission, and find that statistically both samples exhibit similar properties in terms of mass and size. Only their embedded (proto)stellar content is different, the infrared-quiet sample lacks luminous embedded objects. This suggests that infrared-quiet clumps are the precursors of the bright clumps, and thus they likely follow the same evolutionary path. A substantial fraction of infrared-quiet clumps exhibit signatures of ongoing star formation activity such as emission from shocked gas associated with material ejection. In particular, selective star formation tracers, such as EGO activity and Class II methanol maser emission, suggest that at least some infrared-quiet massive clumps host high-mass protostars. This makes them excellent candidates to be the potential precursors of the most massive stars forming in our Galaxy.

The overall low ratio of virial mass to mass of the sample suggests a collapse on the clump scale, which is consistent with a fast, dynamic star formation process. In this scenario massive clumps start to fragment and collapse before their final mass is accumulated indicating that strong self-gravity and global collapse is needed to build up rich clusters and the most massive stars.

High angular-resolution observations of infrared-quiet massive clumps are the next step to reveal their protostellar content and uncover the initial conditions for high-mass star formation.

Acknowledgements. We thank the referee for the careful reading of the manuscript. This work was partially funded by the ERC Advanced Investigator Grant GLOSTAR (247078). T.C.s acknowledges support from the Deutsche Forschungsgemeinschaft. DFG via the SPP (priority programme) 1573 “Physics of the ISM”. T.C.s thanks R. Gutermuth for making available the PhotVis tool. This paper is based on data acquired with the Atacama Pathfinder EXperiment (APEX). APEX is a collaboration between the Max Planck Institute for Radioastronomy, the European Southern Observatory, and the Onsala Space Observatory. This research made use of data products from the Midcourse Space Experiment (MSX) Processing of the data and the Ballistic Missile Defense Organization with additional support from NASA Office of Space Science. This research has also made use of the NASA/IPAC Infrared Science Archive, which is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration. This publication makes use of data products from the Wide-field Infrared Survey Explorer, which is a joint project of the University of California, Los Angeles, and the Jet Propulsion Laboratory/California Institute of Technology, funded by the National Aeronautics and Space Administration. This research made use of Montage, funded by the National Aeronautics and Space Administration’s Earth Science Technology Office, Computation Technologies Project, under Cooperative Agreement Number NCCS-626 between NASA and the California Institute of Technology. Montage is maintained by the NASA/IPAC Infrared Science Archive.

References

Appendix A: Mid-infrared photometry

A.1. MIPSGAL upper limits

Using the publicly available calibrated MIPSGAL tiles, we searched for embedded compact objects spatially coinciding with the ATLASGAL dust peaks. From the 219 brightest submm sources we find that 27 do not show any clear point source within 5′′ from the dust peaks. After a visual inspection, we extracted an upper limit for these sources by taking the median of the pixel values within a radius of 3′′, i.e. half of the FWHM of the PSF at 24 μm. We then multiply this value by the beam area assuming a Gaussian beam with a FWHM of 6′′ to obtain an upper limit on the source flux. We estimate the error by calculating the REDMSQ\(^5\) value in the same area, in the footprint of the PSF, as was defined in Megeath et al. (2012). We extract upper limits between 0.041–1.368 Jy/beam with an average of 0.4 Jy/beam.

A.2. MIPSGAL aperture photometry at 24 μm

We performed aperture photometry towards the dust peaks that are clearly associated with a 24 μm point source and without saturated pixels. In greater detail, we extracted the flux density of point sources from the MIPSGAL data using the PhotVis v1.10 tool (Gutermuth et al. 2008) positionally coinciding with the ATLASGAL source. We adopted a zero point magnitude of 12.448, which was calculated assuming a pixel size of 1′′/25, an aperture correction of 2.05, which is used for apertures of 7′′ with a sky annulus between 20′′ to 32′′ (Engelbracht et al. 2007)\(^6\). We chose these radii because the background emission varies substantially and a larger sky annulus helps to avoid contamination by nearby bright sources. However, our background values still show variations and so may still be contaminated by extended emission. Therefore, we caution that our values are subject to uncertainties.

We then converted the measured magnitudes to flux density using \(F_\nu \equiv c_{\text{corr}} \times F_{\nu 0} \times 10^{-\text{mag}/2.5}\), where \(F_{\nu 0}\) corresponds to the zero magnitude flux level, which is 7.17 Jy and mag is the magnitude given by PhotVis. The targeted sample consists of deeply embedded objects, where \(\beta \sim 1–2\) corresponds to rising SEDs towards longer wavelengths. Therefore, we applied a colour correction \((c_{\text{corr}})\) of 0.981 for objects with \(\beta \sim 1\) (Engelbracht et al. 2007).

A.3. Comparing MIPSGAL and WISE flux densities

We measure 24 μm flux densities (and upper limits) between 0.041 and 6.88 Jy for these sources, with a mean of 1.84 and a median of 0.81 Jy. For 75 sources in the sample we also have photometry from the WISE point source catalogue (Wright et al. 2010). Since there is an overlap in the wavelength ranges of the WISE band at 22 μm and the MIPS detector at 24 μm, we attempted to cross-calibrate the photometry of the two measurements. We found that numerous sources appearing in the WISE catalogue and not saturated with MIPS are flagged as likely spurious detections. This is likely due to source confusion and extended nebulosity surrounding the majority of our weak sources, which results in bad quality flags for these sources in the WISE point source catalogue. Therefore, we only have 21 sources with reliable flux measurement from WISE, of which only 19 have smaller than 5′′ angular offset from the ATLASGAL peak position.

Our attempt to cross-calibrate the two instruments is shown in Fig. A.1, where we also correct for a spectral index of 1.5 that has been found between MSX and MIPSGAL according to the MIPSGAL data delivery document\(^7\). With some scatter, we find similar flux density measurements below ~4 Jy with the two instruments. Since the MIPS detector starts to saturate at 2 Jy, the flattening of the MIPSGAL fluxes between ~2–4 Jy is likely due to the saturation or non-linearity of the detector. In addition, since WISE has a beam that is two times larger (the MIPS PSF is 6′′ at 24 μm, while WISE has a 12′′ angular resolution), it may include emission from the surrounding diffuse emission, not only the protostar. Since saturation and non-linearity of the detector may significantly influence the determined values, in our analysis for sources above 2 Jy we rely on the WISE point source catalogue values.

\(^5\) REDMSQ = \(\sqrt{\text{median}(S[i,j]) - \text{median}(S[i,j])^2}\).

\(^6\) We used the following formula to calculate the zero point magnitude, mag0 = \(2.5 \times \log_{10}\left(\frac{F_{\nu}}{\Omega \times \text{pixel size}}\right)\), where \(\Omega\) corresponds to the pixel size in steradian and \(F_{\nu}\) is 7.17 Jy.

\(^7\) http://irsa.ipac.caltech.edu/data/SPITZER/MIPSGAL/images/mipsgal_delivery_guide_v3_29aug08.pdf
Appendix B: Infrared-quiet massive clumps

Fig. B.1. Left: mid-infrared emission from MIPSGAL at 24 µm or WISE at 22 µm in colour scale with the 870 µm contours from ATLASGAL. The arrows mark the position of infrared-quiet clumps. Right: 870 µm emission with contours from ATLASGAL. The blue crosses mark the positions of UC-H II regions from CORNISH. The figure label shows the corresponding ATLASGAL source name.
Fig. B.1. continued.
Fig. B.1. continued.
Fig. B.1. continued.
Fig. B.1. continued.
Fig. B.1. continued.
Fig. B.1. continued.
Fig. C.1. SiO (8–7) detections with the APEX/FLASH+ receiver of a selection of infrared-quiet massive clumps. The red dashed line shows the fitted Gaussians, and the grey histogram shows the line area. The blue dashed line indicates the $v_{lsr}$ of the sources.
Fig. C.1. continued.