

LETTER TO THE EDITOR

# Quasi-periodic pulsations with periods that change depending on whether the pulsations have thermal or nonthermal components

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## ABSTRACT

*Context.* Quasi-periodic pulsations (QPPs) typically display periodic and regular peaks in the light curves during the flare emissions. Sometimes, QPPs show multiple periods at the same wavelength. However, changing periods in various channels are rare.

*Aims.* We report QPPs in a solar flare on 2014 October 27. They showed a period change that depended on whether thermal or nonthermal components were included. The flare was simultaneously observed by many instruments.

*Methods.* Using the fast Fourier transform (FFT), we decomposed the light curves at multiple wavelengths into slowly varying and rapidly varying signals. Then we identified the QPPs as the regular and periodic peaks from the rapidly varying signals. The periods are derived with the wavelet method and confirmed based on the FFT spectra of the rapidly varying signals.

*Results.* We find a period of  $\sim 50$  s from the thermal emissions during the impulsive phase of the flare, that is, in the soft X-ray bands. At the same time, a period of about 100 s is detected from the nonthermal emissions, such as hard X-ray and microwave channels. The period ratio is exactly 2.0, which might be due to the modulations of the magnetic reconnection rate by the fundamental and harmonic modes of magnetohydrodynamic waves. Our results further show that the  $\sim 100$  s period is present over a broad wavelength, such as hard X-rays, extreme-UV/UV, and microwave emissions, indicating the periodic magnetic reconnection in this flare.

*Conclusions.* To our knowledge, this is the first report about period changes from thermal to nonthermal components in a single flare that occur at almost the same time. This new observational finding could be a challenge to the theory of flare QPPs.

**Key words.** Sun: flares – Sun: oscillations – Sun: radio radiation – Sun: UV radiation – Sun: X-rays, gamma rays

## 1. Introduction

Quasi-periodic pulsations (QPPs) are very common features in flare emissions. They typically exhibit periodic peaks in the light curves. All these peaks have similar lifetimes, which results in regular intervals between the peaks. The observed periods can last from milliseconds through seconds to minutes (e.g., Hoyng et al. 1976; Karlický et al. 2005; Tan et al. 2010; Su et al. 2012; Ning 2014; Li et al. 2015; Zhang et al. 2016). Previous observations indicated that the periods of these flares are strongly dependent on the major radius of flare loops (Aschwanden et al. 1998). On the other hand, QPPs can be detected in intensity (or flux) observations at wide bands, that is, in X-rays (Hoyng et al. 1976; Dolla et al. 2012; Ning 2014; Hayes et al. 2016) from extreme-ultraviolet (EUV; Su et al. 2012) to microwave (Mangeny & Pick 1989; Tan et al. 2010) emissions. QPPs can also be observed in Doppler shifts in spectral observations, such as the hot ( $>6$  MK) lines observed by SOHO/SUMER (Ofman & Wang 2002; Wang et al. 2002), the coronal lines from Hinode/EIS (Tian et al. 2011), and the flare and transition lines from *IRIS* observations (Li et al. 2015; Li & Zhang 2015; Tian et al. 2016). Sometimes, QPPs in the same flare display multiple periods (Inglis & Nakariakov 2009), and these multi-periods are often detected at the same channels, that is, in soft X-ray (SXR; Chowdhury et al. 2015), EUV (Su et al. 2012; Tian et al. 2016), H $\alpha$  (Srivastava et al. 2008; Yang & Xiang 2016), and microwave (Melnikov et al. 2005; Kupriyanova et al. 2014) channels. Multi-periods are also found at various channels

but in different flare loops (Zimovets & Struminsky 2010) or phases (Hayes et al. 2016; Tian et al. 2016). All these period ratios are always deviating from 2.0.

Today, the formation mechanism of QPPs is still debated (e.g., Aschwanden 1987; Nakariakov & Melnikov 2009; Van Doorselaere et al. 2016). Generally, QPPs are thought to be related with waves or energetic particles (electrons). Magnetohydrodynamic (MHD) waves are most frequently used to explain QPPs. These include slow magnetoacoustic, fast kink, and sausage waves (Nakariakov & Melnikov 2009; Tan et al. 2010; Su et al. 2012). QPPs can also be explained by quasi-periodic magnetic reconnection (Nakariakov & Melnikov 2009; Li et al. 2015; Li & Zhang 2015). The periods can be spontaneous (Kliem et al. 2000; Murray et al. 2009) or may be modulated by some MHD waves (Chen & Priest 2006; Nakariakov & Zimovets 2011; Li et al. 2015).

Until now, QPPs in a flare that have periods that change in different channels have rarely been simultaneously observed with different instruments. We here analyze QPPs with changing periods in a solar flare on 2014 October 27, which was detected by The Reuven Ramaty High Energy Solar Spectroscopic Imager (RHESSI; Lin et al. 2002), the Atmospheric Imaging Assembly (AIA; Lemen et al. 2012) onboard Solar Dynamics Observatory (SDO), Nobeyama Radio Polarimeters (NoRP; Nakajima et al. 1985), and the Nobeyama Radioheliograph (NoRH; Hanaoka et al. 1994) in the SXR, HXR, EUV/UV, and microwave channels.

**Table 1.** Details of the observational instruments and data.

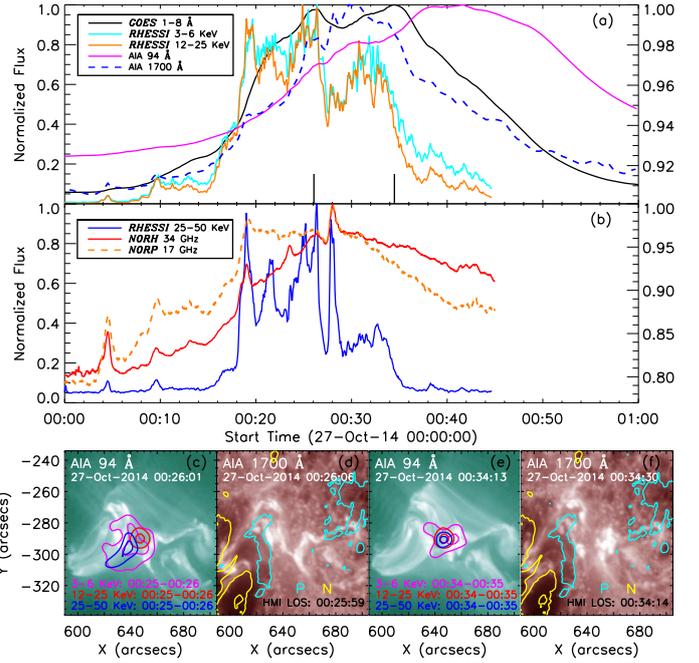
Instruments	Channels	Cadence (s)	Descriptions	<i>c.c.</i>
RHESSI	3–6 keV	4.0	SXR	0.88
	12–25 keV	4.0	SXR/HXR	0.92
	25–50 keV	4.0	HXR	1.0
GOES	1–8 Å	~2.0	SXR	0.82
NoRP	17 GHz	1.0	Microwave	0.81
NoRH	34 GHz	1.0	Microwave	0.82
SDO/AIA	94 Å	12	EUVV	0.57
	1700 Å	24	UV	0.86
SDO/HMI	6173 Å	45	magnetogram	–

## 2. Observations and results

We explore an M7.1 flare that erupted in NOAA AR12192 (S12W42). It started at ~00:06 UT, peaked at ~00:34 UT, and stopped at ~00:44 UT according to the GOES SXR flux. This flare was detected by many instruments over a broad channel with various temporal resolution, as shown in Table 1. Figure 1 displays the normalized light curves and images detected by GOES, RHESSI, SDO, NoRP, and NoRH. The black profile in panel a) shows the GOES SXR flux at 1–8 Å, which clearly shows two peaks at about 00:26 UT and 00:34 UT, indicating two heating processes during the impulsive phase of the flare (see also Chen et al. 2015; Polito et al. 2016). Panel a also displays the light curves from RHESSI 3–6 keV (turquoise) and 12–25 keV (orange), AIA 94 Å (purple), and 1700 Å (blue). Panel b shows the HXR flux at 25–50 keV (blue), microwave emissions from NoRP 17 GHz (orange), and NoRH 34 GHz (red). These three light curves in panel b seem to be well correlated and have similar peaks during the impulsive phase of the flare; see also Table 1. The reference channel is HXR emissions at 25–50 keV. Table 1 gives the correlation coefficients (*c.c.*) between the light curves at HXR 25–50 keV and other channels. Most of the coefficients are higher than 0.8. However, the coefficient between the flux at AIA 94 Å is slightly lower, only about 0.57. The reason might be that the heating process at AIA 94 Å occurs later than that at other channels. Figure 1a also shows that the peak time of AIA 94 Å is later than others, that is, later than GOES 1–8 Å and AIA 1700 Å. In order to clearly show the light curves, the solid profiles in panels a and b work with the left axes, while the dashed profiles correspond to the right axes.

Figures 1c–f show the AIA images with a size of about 115"×105" at 94 Å and 1700 Å from two peak times marked by the vertical ticks in panel a. The light curves at 94 and 1700 Å in panel a are integrated from the all flare regions. The color contours in panels c and e represent the RHESSI SXR/HXR images using the CLEAN algorithm from detectors 3, 4, 5, 6, and 8, while the color contours in panels d and f are from the SDO/HMI line-of-sight (LOS) magnetograms at a scale of ±800 G for positive (turquoise) and negative (yellow) fields, respectively. The imaging observations show that the flare has complex ribbons and the HXR sources are located between two of them. This M7.1 flare is also associated with strong magnetic fields that were connected by the positive and negative poles.

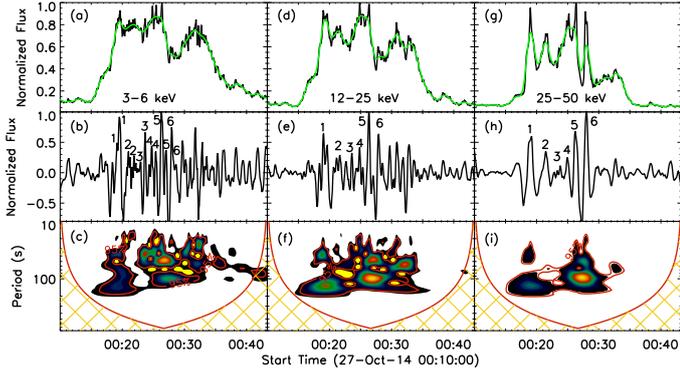
Figure 1b shows several peaks during the impulsive phase of the flare in HXR and microwave light curves between 00:18 UT and 00:28 UT. These peaks appear to be regular and periodic, but they are superposed by the gradual background emissions, which



**Fig. 1.** Upper panel: normalized light curves of the flare on 2014 October 27 from GOES 1–8 Å, RHESSI 3–6 keV and 12–25 keV, SDO/AIA 94 Å and 1700 Å. The vertical ticks indicate the time of two GOES SXR peaks. Middle panel: normalized light curves of the same flare from RHESSI 25–50 keV, NoRP 17 GHz, and NoRH 34 GHz. Lower panels: SDO/AIA images at 94 Å c, e) and 1700 Å d, f) at two SXR peak times. The X-ray contour levels from the RHESSI images are 95%, 80%, while the HMI contour levels are at ±800 G.

makes them difficult to distinguished. Therefore, these light curves were decomposed into the slowly varying and rapidly varying signals with the Fourier transform. First, the power spectrum of the light curve was calculated using the fast Fourier transform (FFT). Second, the power spectrum was split into domains with lower and higher frequency. Third, the inverse FFT was applied to these two domains to obtain the slowly varying (background) and rapidly varying (quasi-periodic peaks) signals (see also Ning 2014; Ning & Guo 2014). Here, a uniform timescale of 120 s was used to distinguish these two signals for all the observation data.

Figures 2a–d, g show the slowly varying signals (green profiles) that overlap the X-ray light curves at three energy channels from RHESSI. The middle panels display their rapidly varying signals. To better show the rapidly varying signal, it was normalized by its slowly varying signal (see also Kupriyanova et al. 2010, 2013). The period is clearly identified in the higher energy channel, that is, 25–50 keV. Panel h clearly shows 6 peaks (as marked by the numbers) from around 00:18 UT to 00:28 UT (~600 s), which means that the period is about 100 s at this channel. Then we applied a wavelet analysis to the rapidly varying signals, and we show the results in the bottom panels. The wavelet results further confirm the period of about 100 s at the higher energy, see panel i. However, at the lower energy (3–6 keV), there are more peaks and their intervals are shorter, that is, 12 peaks from about 00:18 UT to 00:28 UT, and some peaks are weaker than others. The wavelet results show a period of ~50 s, as shown in panel c, which is twice lower than the period at the higher energy. It is interesting that both the periods of ~50 s and ~100 s are detected in the medium-energy channel, that is, 12–25 keV, as shown in panel f. Panel e also shows that there are 6 main peaks in the rapidly varying signals



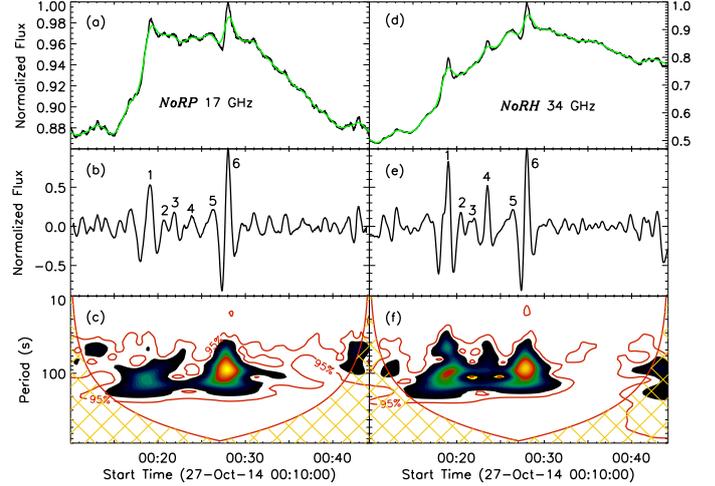
**Fig. 2.** Upper panels: RHESSI light curves (black) and their slowly varying signals (green). Middle panels: rapidly varying signals. Lower panels: wavelet power spectra.

at 12–25 keV, as indicated by the numbers. Meanwhile, there are double subpeaks at almost every main peak. This is possible because the emissions from this channel include both thermal and nonthermal components. Therefore, periods of both  $\sim 50$  s and  $\sim 100$  s are displayed in this channel. To our knowledge, this is the first time that changing periods in a flare are reported that change from the thermal components (SXR emissions) to the nonthermal components (HXR emissions) at almost the same time.

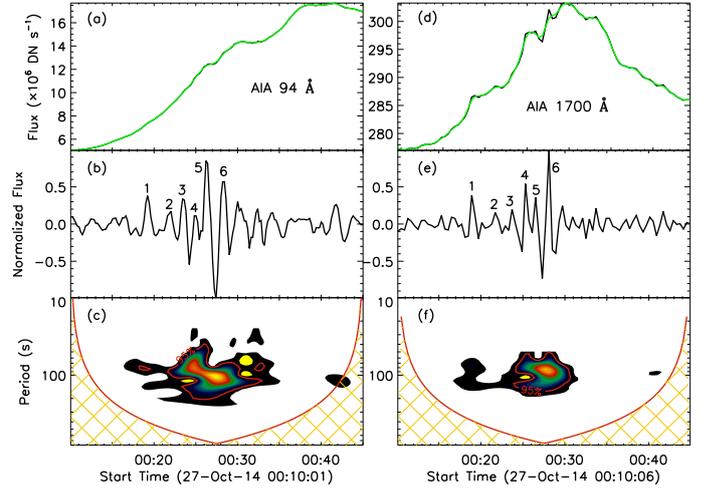
Except for the RHESSI HXR channels, the  $\sim 100$  s QPPs are also detected in microwave and EUV/UV emissions using the same method. Figure 3 shows the wavelet results from the microwave emissions at frequencies of 17 GHz and 34 GHz. The green profiles that overlap on the black profiles are the slowly varying signals using the Fourier transform (Ning 2014; Ning & Guo 2014). The rapidly varying signals and their wavelet power spectra are shown in the middle and bottom panels. Similar to the RHESSI HXR emissions, the microwave emissions at 17 GHz and 34 GHz exhibit the  $\sim 100$  s QPPs. All these emissions at HXR and microwave bands could be considered as the nonthermal components. Figure 4 shows the wavelet analysis results from the EUV/UV wavelengths at AIA 94 Å and 1700 Å. In the same way as in Fig. 3, the upper panels show the integrated intensity flux (black profiles) and their slowly varying signals (green profiles), and the integral region is shown in Fig. 1c, while the rapidly varying signals and their wavelet power spectra are shown in the middle and bottom panels. The same  $\sim 100$  s QPPs are found using the wavelet method. Figures 3 and 4 show that there are more than 6 peaks in some channels, and some peaks are out of the time between 00:18 UT and 00:28 UT; these are AIA 94 Å and RHESSI 12–25 keV. However, only 6 main peaks are marked by the numbers from about 00:18 UT to 00:28 UT, which is exactly the HXR peak time during the impulsive phase of the flare. More important, these 6 peaks appear to match well in the HXR, microwave, and EUV/UV channels, as indicated by the number ticks.

### 3. Conclusions and discussions

We reported  $\sim 100$  s QPPs in the HXR, microwave, and EUV/UV channels and the  $\sim 50$  s QPPs in the SXR bands in the M7.1 flare on 2014 October 27. The result was obtained using the multi-instrument observations from RHESSI, NoRP, NoRH, and SDO/AIA. To our knowledge, this is the first report about periods in the same flare that change at almost the same time from nonthermal components to thermal components. We identified



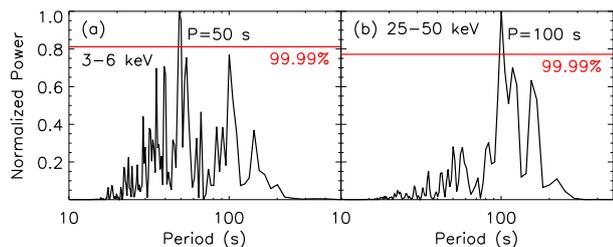
**Fig. 3.** Upper panels: microwave emissions (black) and their slowly varying signals (green). Middle panels: rapidly varying signals. Lower panels: wavelet power spectra.



**Fig. 4.** Upper panels: EUV and UV flux (black) and their slowly varying signals (green). Middle panels: rapidly varying signals. Lower panels: wavelet power spectra.

QPPs based on the rapidly varying signals that were obtained from the original light curves with the FFT. Then the  $\sim 100$  s and  $\sim 50$  s QPPs were derived with the wavelet method. Based on this, the ratio between these two periods is 2.0. This is well consistent with the period ratio (2.0) between the fundamental and harmonic MHD waves.

It is very interesting that the changing periods are detected in the same flare at almost the same time and that the periods change from  $\sim 50$  s at the thermal components to  $\sim 100$  s at the nonthermal components. There have been some reports about multi-period QPPs in the same flare. However, these periods have always been detected in the same channels, that is, in SXR light curves (Chowdhury et al. 2015), EUV flux/spectra (Su et al. 2012; Tian et al. 2016),  $H\alpha$  images (Srivastava et al. 2008; Yang & Xiang 2016), or microwave emissions (Melnikov et al. 2005; Kupriyanova et al. 2014). On the other hand, multiple periods have been detected in various channels, but these periods typically appeared at different times or loops, such as the impulsive and decay phases of the flare (Zimovets & Struminsky 2010; Huang et al. 2014; Hayes et al. 2016; Tian et al. 2016). It is interesting that the period ratio



**Fig. 5.** FFT spectra of the rapidly varying signals at 3–6 keV, and 25–50 keV. The red line indicates the 99.99% confidence level.

is exactly 2.0. This could be interpreted as modulations in the magnetic reconnection rate that are caused by the fundamental and harmonic modes of the MHD waves. However, previous multi-period observations usually showed that the period ratio deviated from 2.0 (e.g., [Srivastava et al. 2008](#); [Chowdhury et al. 2015](#); [Yang & Xiang 2016](#)), which may be due to the longitudinal density stratification ([Andries et al. 2005](#)), the loop expansion ([Verth & Erdélyi 2008](#)), the siphon flow ([Li et al. 2013](#)), or the flare phase/loop differences ([Zimovets & Struminsky 2010](#); [Tian et al. 2016](#); [Hayes et al. 2016](#)).

It is also interesting that the QPPs exhibit the same period of  $\sim 100$  s at a broad wavelength from HXR through EUV/UV to microwave emissions. Their radiation sources spread from the chromosphere through the transition region to the corona. It is still an open question what modulates the 100 s QPPs at such broad wavelengths. Three models have been presented in previous documents (e.g., [Aschwanden 1987](#)): (i) MHD flux tube oscillations modulated by waves; (ii) periodic self-organizing systems of plasma instabilities of wave-particle or wave-wave interactions; and (iii) modulation of periodic particle acceleration. We here observed QPPs in the impulsive phase, which gives us a clue about how to explain the QPPs. It is generally thought that the magnetic reconnection process occurs in the impulsive phase. Thus, it is possible that the QPPs are produced by periodic magnetic reconnection in this flare. Based on the standard flare model, magnetic reconnection could simultaneously accelerate the bidirectional electron beams ([Aschwanden et al. 1995](#); [Su et al. 2013](#); [Li et al. 2016](#)). The downward beam produces the microwave emission peak in its trajectory propagation, it also produces one peak of the HXR emission when it is injected into the chromosphere, and this process will heat the local plasmas to produce one peak in the EUV/UV emission. In this case, the periodic magnetic reconnection can produce the QPPs peaks at HXR, EUV/UV, and microwave emissions. The period of  $\sim 100$  s could be modulated by certain MHD waves in the solar corona, such as the fast kink mode ([Nakariakov et al. 2006](#); [Nakariakov & Melnikov 2009](#); [Huang et al. 2014](#)), or the slow sausage mode ([Van Doorselaere et al. 2011](#)), which could modulate the period of around 100 s. However, it is difficult to determine the MHD waves in this flare because of the observation limitation, such as the absence of a three-dimensional magnetic configuration of the active regions.

The wavelet spectra in Figs. 2–4 exhibit broad amplitudes with brightest cores and weaker borders, which may result in a fuzzy result for the QPP periods. To unambiguously conclude on the periods, we performed a periodogram analysis of the rapidly varying signals using the Lomb-Scargle periodogram method ([Scargle 1982](#)). Figure 5 shows the FFT spectra at RHESSI 3–6 keV and 25–50 keV, and the red line indicates the 99.99% confidence level defined by [Horne & Baliunas \(1986\)](#). We were easily able to detect the periods of  $\sim 50$  s and  $\sim 100$  s in the FFT spectra in the SXR and HXR channels, respectively.

This additionally confirms that the periods change depending on the thermal or nonthermal emission.

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## References

- Andries, J., Arregui, I., & Goossens, M. 2005, *ApJ*, 624, L57  
 Aschwanden, M. J. 1987, *Sol. Phys.*, 111, 113  
 Aschwanden, M. J., Benz, A. O., Dennis, B. R., & Schwartz, R. A. 1995, *ApJ*, 455, 347  
 Aschwanden, M. J., Kliem, B., Schwarz, U., et al. 1998, *ApJ*, 505, 941  
 Chen, P. F., & Priest, E. R. 2006, *Sol. Phys.*, 238, 313  
 Chen, H., Zhang, J., Ma, S., et al. 2015, *ApJ*, 808, L24  
 Chowdhury, P., Srivastava, A. K., Dwivedi, B. N., Sych, R., & Moon, Y.-J. 2015, *Adv. Space Res.*, 56, 2769  
 Dolla, L., Marqué, C., Seaton, D. B., et al. 2012, *ApJ*, 749, L16  
 Hanaoka, Y., Shibasaki, K., Nishio, M., et al. 1994, Proceedings of Kofu Symposium, 35  
 Hayes, L. A., Gallagher, P. T., Dennis, B. R., et al. 2016, *ApJ*, 827, L30  
 Horne, J. H., & Baliunas, S. L. 1986, *ApJ*, 302, 757  
 Hoynig, P., Brown, J. C., & van Beek, H. F. 1976, *Sol. Phys.*, 48, 197  
 Huang, J., Tan, B., Zhang, Y., Karlický, M., & Mészárosová, H. 2014, *ApJ*, 791, 44  
 Inglis, A. R., & Nakariakov, V. M. 2009, *A&A*, 493, 259  
 Karlický, M., Bárta, M., Mészárosová, H., & Zlobec, P. 2005, *A&A*, 432, 705  
 Kliem, B., Karlický, M., & Benz, A. O. 2000, *A&A*, 360, 715  
 Kupriyanova, E. G., Melnikov, V. F., Nakariakov, V. M., & Shibasaki, K. 2010, *Sol. Phys.*, 267, 329  
 Kupriyanova, E. G., Melnikov, V. F., & Shibasaki, K. 2013, *Sol. Phys.*, 284, 559  
 Kupriyanova, E. G., Melnikov, V. F., Puzynya, V. M., Shibasaki, K., & Ji, H. S. 2014, *Astron. Rep.*, 58, 573  
 Lemen, J. R., Title, A. M., & Akin, D. J., et al. 2012, *Sol. Phys.*, 275, 17  
 Li, T., & Zhang, J. 2015, *ApJ*, 804, L8  
 Li, B., Habbal, S. R., & Chen, Y. 2013, *ApJ*, 767, 169  
 Li, D., Ning, Z. J., & Zhang, Q. M. 2015, *ApJ*, 807, 72  
 Li, L., Zhang, J., Peter, H., et al. 2016, *Nature Physics*, 12, 847  
 Lin, R. P., Dennis, B. R., Hurford, G. J., et al. 2002, *Sol. Phys.*, 210, 3  
 Mangeney, A., & Pick, M. 1989, *A&A*, 224, 242  
 Melnikov, V. F., Reznikova, V. E., Shibasaki, K., & Nakariakov, V. M. 2005, *A&A*, 439, 727  
 Murray, M. J., van Driel-Gesztelyi, L., & Baker, D. 2009, *A&A*, 494, 329  
 Nakajima, H., Sekiguchi, H., Sawa, M., Kai, K., & Kawashima, S. 1985, *PASJ*, 37, 163  
 Nakariakov, V. M., & Melnikov, V. F. 2009, *Space Sci. Rev.*, 149, 119  
 Nakariakov, V. M., & Zimovets, I. V. 2011, *ApJ*, 730, L27  
 Nakariakov, V. M., Foullon, C., Verwichte, E., & Young, N. P. 2006, *A&A*, 452, 343  
 Ning, Z. J. 2014, *Sol. Phys.*, 289, 1239  
 Ning, Z., & Guo, Y. 2014, *ApJ*, 794, 79  
 Ofman, L., & Wang, T. J. 2002, *ApJ*, 580, L85  
 Polito, V., Reep, J. W., Reeves, K. K., et al. 2016, *ApJ*, 816, 89  
 Scargle, J. D. 1982, *ApJ*, 263, 835  
 Srivastava, A. K., Zaqarashvili, T. V., Uddin, W., Dwivedi, B. N., & Kumar, P. 2008, *MNRAS*, 388, 1899  
 Su, J. T., Shen, Y. D., & Liu, Y., et al. 2012, *ApJ*, 755, 113  
 Su, Y., Veronig, A. M., Holman, G. D., et al. 2013, *Nature Physics*, 9, 489  
 Tan, B. L., Zhang, Y., Tan, C. M., & Liu, Y. Y. 2010, *ApJ*, 723, 25  
 Tian, H., McIntosh, S. W., & De Pontieu, B. 2011, *ApJ*, 727, L37  
 Tian, H., Young, P. R., Reeves, K. K., et al. 2016, *ApJ*, 823, L16  
 Van Doorselaere, T., De Groof, A., Zender, J., Berghmans, D., & Goossens, M. 2011, *ApJ*, 740, 90  
 Van Doorselaere, T., Kupriyanova, E. G., & Yuan, D. 2016, *Sol. Phys.*, 291, 3143  
 Verth, G., & Erdélyi, R. 2008, *A&A*, 486, 1015  
 Wang, T., Solanki, S. K., Curdt, W., Innes, D. E., & Dammasch, I. E. 2002, *ApJ*, 574, L101  
 Yang, S., & Xiang, Y. 2016, *ApJ*, 819, L24  
 Zhang, Q. M., Li, D., & Ning, Z. J. 2016, *ApJ*, 832, 65  
 Zimovets, I. V., & Struminsky, A. B. 2010, *Sol. Phys.*, 263, 163