Understanding the C$_3$H$_2$ cyclic-to-linear ratio in L1544

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ABSTRACT

Aims. We aim to understand the high cyclic-to-linear C$_3$H$_2$ ratio (32 ± 4) that has been observed toward L1544.

Methods. We combined a gas-grain chemical model with a physical model for L1544 to simulate the column densities of cyclic and linear C$_3$H$_2$ observed toward L1544. The most important reactions for the formation and destruction of both forms of C$_3$H$_2$ were identified, and their relative rate coefficients were varied to find the best match to the observations.

Results. We find that the ratio of the rate coefficients of C$_3$H$_2^+ + e^- \rightarrow$ C$_3$H$_2 + H$ for cyclic and linear C$_3$H$_2$ must be ~20 to reproduce the observations, depending on the branching ratios assumed for the C$_3$H$_2^+ + e^- \rightarrow$ C$_3$H$^+ + H_2$ reaction. In current astrochemical networks it is assumed that cyclic and linear C$_3$H$_2$ are formed in a 1:1 ratio in the aforementioned reactions. Laboratory studies and/or theoretical calculations are needed to confirm the results of our chemical modeling, which is based on observational constraints.

Key words. ISM: abundances – ISM: clouds – ISM: molecules – astrochemistry

1. Introduction

Isomers and isotopologs are excellent tools to probe physical and chemical properties, as well as evolutionary states, in the ISM. Cyclopropenylidene, the cyclic form of C$_3$H$_2$, is one of the most abundant and widespread molecules in our Galaxy (Matthews & Irvine 1985). Propadienyldiene, H$_2$CCC, is a very polar carbene and is the less stable isomer of c-C$_3$H$_2$ by ~10 kcal mol$^{-1}$ (5000 K; Wu et al. 2010). Propadienyldiene was first detected in the laboratory by Vrtilek et al. (1990) by means of millimeter-wave absorption spectroscopy in a DC discharge. It has been observed in prestellar and protostellar cores (Cernicharo et al. 1991; Turner et al. 2000; Kawaguchi et al. 1991; Sakai & Yamamoto 2013), circumstellar envelopes (Cernicharo et al. 1991), translucent clouds (Turner et al. 2000), diffuse clouds (Cernicharo et al. 1999; Kułczak-Iaстрzeńbska et al. 2012; Lisztt et al. 2012), and photodissociation regions (PDRs; Teyssier et al. 2005; Cuadrado et al. 2015).

The cyclic-to-linear (hereafter c/l) ratio of C$_3$H$_2$ tends to increase with increasing $A_V$. Significant variations of this ratio are found in dense cores, ranging from 25 in TMC-1(CP) to 67 in TMC-1C (Spezzano et al. 2016). For completeness, we report in Table A.1 the values of the C$_3$H$_2$ c/l ratio observed in several media, including dense cores. To be able to model the observed c/l ratios use this ratio as an astrochemical tool, the chemistry leading to the formation and destruction of c-C$_3$H$_2$ and H$_2$CCC (hereafter l-C$_3$H$_2$) needs to be better understood, with emphasis on the branching ratios of the reactions involved.

The main formation process of cyclic and linear C$_3$H$_2$ involves the radiative association of C$_3$H$^+$ and H$_2$ to form cyclic and linear C$_3$H$_2^+$, and a subsequent dissociative recombination with electrons that leads to c-C$_3$H$_2$ and l-C$_3$H$_2$, respectively. C$_3$H$^+$ exists in both cyclic and linear forms, but the cyclic isomer lies 17 kcal mol$^{-1}$ (8600 K) above the linear one (Ikuta 1997), hence for our purposes we can assume to deal with l-C$_3$H$_2^+$ only. The radiative association of l-C$_3$H$_2^+$ and H$_2$ will form cyclic and linear C$_3$H$_2^+$ in a 1:1 ratio (Maluendes et al. 1993). Talbi et al. (2009) suggested with their calculations that the dissociative recombination of c-C$_3$H$_2^+$ to form c-C$_3$H$_2$ is more efficient than its linear counterpart. Using an afterglow experiment at 300 K, Adams & Babcock (2005) observed that the cyclic C$_3$H$_2^+$ recombines faster than the linear isomer.

The chemistry of c and l C$_3$H$_2$ and that of other associated species has previously been modeled for instance by Turner et al. (2000) and Fossé et al. (2001). Both of these works pointed out that neutral-neutral reactions can also play a part in determining the C$_3$H$_2$ c/l ratio. The rate coefficients of the reactions most relevant in the present context have changed a little since these studies.

A key point of the dissociative recombination of C$_3$H$_2^+$ is its branching ratio, which unfortunately has not yet been investigated either experimentally or theoretically. Chabot et al. (2013) studied the dissociative recombination of cyclic and linear C$_3$H$_2^+$. Their results show that while c-C$_3$H is the main product of the dissociative recombination of c-C$_3$H$_2^+$ with electrons, l-C$_3$H$_2$ is not the main product of the recombination of l-C$_3$H$_2^+$. In this work we study the effect of the branching ratios of the dissociative recombination of both cyclic and linear C$_3$H$_2^+$ on the cyclic-to-linear ratio of C$_3$H$_2$. We use the well-studied prestellar core L1544 to compare the results of our model with the observational results.

2. Formation and destruction of C$_3$H$_2$

In physical conditions with high visual extinction ($A_V >$ a few mag) where photoprocesses do not play a significant role, the formation and destruction of (cyclic and linear) C$_3$H$_2$ follow a relatively simple scheme. The main formation pathway is the dissociative recombination C$_3$H$_2^+$ + e$^-$. In this paper, we use the latest rate coefficient data from the KIDA database (Wakelam et al. 2015; see also below), where three different branches for this reaction are included. These are tabulated in Table 1. The
third tabulated reaction for both forms of C\(_3\)H\(_2\), whose rate coefficient we label here as \(k_{\text{DR,3}}\), is the most important for the C\(_3\)H\(_2\) formation because it produces C\(_3\)H\(_2\) directly (see also Sect. 4).

The main destruction partners of C\(_3\)H\(_2\) change with time, but at the time when C\(_3\)H\(_2\) is the most abundant (see Sects. 3 and 4), the main destruction reactions are ion-molecule reactions that convert C\(_3\)H\(_2\) back to C\(_3\)H\(_2^+\); C\(_3\)H\(_2\) + HCO\(^+\)/H\(_2\)O\(^+\)/H\(_2\) → C\(_3\)H\(_2^+\) + CO/H\(_2\)/H\(_2\) with rate coefficients \(k_1\), \(k_2\), and \(k_3\), respectively. Thus, the abundance of C\(_3\)H\(_2\) in steady state would be given by

\[
\text{C}_3\text{H}_2 = \frac{k_{\text{DR,3}} \left[ \text{C}_3\text{H}_2^+ \right] [\text{e}]}{k_1 [\text{HCO}^+] + k_2 [\text{H}_2\text{O}^+] + k_3 [\text{H}_3^+]},
\]

which leads to an expression of the c/l ratio of C\(_3\)H\(_2\)

\[
\frac{[\text{c-C}_3\text{H}_2]}{[\text{l-C}_3\text{H}_2]} = \frac{k'_{\text{DR,3}}}{k_{\text{DR,3}}} \left[ \frac{\text{c-C}_3\text{H}_2^+}{\text{l-C}_3\text{H}_2^+} \right] \left[ \frac{\text{e}^-}{\text{e}^-} \right] \times \frac{k_1 [\text{HCO}^+] + k_2 [\text{H}_2\text{O}^+] + k_3 [\text{H}_3^+]}{k_1 [\text{HCO}^+] + k_2 [\text{H}_2\text{O}^+] + k_3 [\text{H}_3^+]}.
\]

Because of gas-grain chemical interaction, a steady-state solution is not attained in our models, but the above expressions turn out to give a rather accurate representation of the C\(_3\)H\(_2\) peak abundances. Inspection of the KIDA data reveals that the rate coefficients of the three C\(_3\)H\(_2\)-destroying ion-molecule reactions discussed above are almost identical (slightly higher for \(k_{\text{DR,2}}\) than for \(k_{\text{DR,3}}\)), so that the latter fraction in Eq. (2) is \(~1\). Furthermore, our chemical model (see below) indicates that the c-C\(_3\)H\(_2^+\)/l-C\(_3\)H\(_2^+\) ratio is usually also very close to unity, so that the c-C\(_3\)H\(_2^+\)/l-C\(_3\)H\(_2^+\) ratio is determined to a good approximation by the rate coefficient ratio \(k'_{\text{DR,3}}/k_{\text{DR,3}}\) at the time when HCO\(^+\), H\(_2\)O\(^+\), and H\(_3^+\) are abundant (i.e., before the onset of freeze-out of neutrals onto grain surfaces). This ratio is unity in the KIDA data, which implies c-C\(_3\)H\(_2^+\)/l-C\(_3\)H\(_2^+\) \sim 1 according to the approximating formula presented above. However, the observed c/l abundance ratio toward the dust peak in L1544 is 32 \pm 4 (Spezzano et al. 2016). We therefore set out to modify the various \(k_{\text{DR}}\) branching ratios to explain the observations. The assumption that the two isomers are formed in the same fashion and in similar conditions is strengthened by our recent maps of c and l C\(_3\)H\(_2\) in the inner 2.5' \times 2.5' of L1544, which demonstrate that c and l C\(_3\)H\(_2\) trace the same region (Spezzano et al., in prep.). A spatial association of c and l C\(_3\)H\(_2\) has also been observed in other sources (e.g., Fossé et al. 2001).

Table 1. Branching ratios for (c/l)-C\(_3\)H\(_2^+\)+e\(^-\) from the KIDA database.

<table>
<thead>
<tr>
<th>Reaction</th>
<th>Branching ratio</th>
</tr>
</thead>
<tbody>
<tr>
<td>c-C(_3)H(_2^+)+e(^-) → CH + C(_2)H(_2)</td>
<td>10% (k_{\text{DR,1}})</td>
</tr>
<tr>
<td>c-C(_3)H(_2^+)+e(^-) → c-C(_3)H + H(_2)</td>
<td>45% (k_{\text{DR,2}})</td>
</tr>
<tr>
<td>c-C(_3)H(_2^+)+e(^-) → c-C(_3)H + H</td>
<td>45% (k_{\text{DR,3}})</td>
</tr>
<tr>
<td>l-C(_3)H(_2^+)+e(^-) → CH + C(_2)H(_2)</td>
<td>10% (k'_{\text{DR,1}})</td>
</tr>
<tr>
<td>l-C(_3)H(_2^+)+e(^-) → c-C(_3)H + H(_2)</td>
<td>45% (k'_{\text{DR,2}})</td>
</tr>
<tr>
<td>l-C(_3)H(_2^+)+e(^-) → l-C(_3)H + H</td>
<td>45% (k'_{\text{DR,3}})</td>
</tr>
</tbody>
</table>

Notes. The third column gives the denominations of the rate coefficients of the various branches discussed in the text. The total rate coefficient is \(k_{\text{tot}} = 7.0 \times 10^{-7} (T/300\text{ K})^{-0.3}\text{ cm}^3\text{ s}^{-1}\) in both cases.

As a guide to the modifications, we analyzed the branching ratios calculated for the C\(_3\)H\(_2^+\)+e\(^-\) reaction by Chabot et al. (2013). These branching ratios, which are included in the KIDA database, are shown in Table 2. Calculations for the C\(_3\)H\(_2^+\)+e\(^-\) system are planned, but it is unclear when the results will be available (M. Chabot, priv. comm.). Therefore, we used the C\(_3\)H\(_2^+\)+e\(^-\) system as a template for the modifications. In the dissociation of C\(_3\)H\(_2^+\), there are pathways that remove either one or both protons from C\(_3\)H\(_2^+\) and pathways that break C-C bonds, that is, there are three groups of pathways. Inspection of the C\(_3\)H\(_2^+\) pathways (Table 1) reveals a similar grouping. As a first approximation we combined the branching ratios belonging to each group in the C\(_3\)H\(_2^+\) data and translated them into the C\(_3\)H\(_2^+\) system. This leads to the branching ratios given in Table 3 (mod_ch). In this modified model, the \(k'_{\text{DR,3}}/k_{\text{DR,3}}\) ratio is \(~5\), which is still not high enough in light of the observations. We then proceeded to vary the branching ratios to attain a fit to the observed C\(_3\)H\(_2\) c/l ratio in L1544. The results of our analysis are presented in Sect. 3.

To simulate the C\(_3\)H\(_2\) abundances in L1544, we used the physical model for L1544 presented by Keto & Caselli (2010) and updated by Keto et al. (2014), which gives us the density, (gas and dust) temperature, and \(A_V\) profiles as functions of distance away from the dust peak. We separated the model core to concentric shells and calculated the chemical evolution separately in each shell. In this way, we can produce simulated chemical abundance gradients and study how the column densities of the various molecules change as functions of time. A similar
The inset in the left panel shows a zoomed-in view of the innermost 10,000 AU in the model core. The black vertical line marks the position of the C\textsubscript{3}H\textsubscript{2} abundance peak (the line is omitted in the left panel for clarity).

Fig. 1. Left panel: abundances of c-C\textsubscript{3}H\textsubscript{2} (solid lines) and l-C\textsubscript{3}H\textsubscript{2} (dotted lines) as functions of distance from the core center in the L1544 model at \( t = 10^7 \) yr. Middle panel: C\textsubscript{3}H\textsubscript{2} c/l ratio as a function of distance from the core center in the L1544 model at \( t = 10^7 \) yr. Right panel: density (green solid line), dust temperature (red solid line), and gas temperature (red dashed line) as functions of distance from the core center in the L1544 model of Keto et al. (2014). The different colors in the left and middle panels correspond to \( k\textsubscript{DR,1}/k\textsubscript{DR,2} = 1 \) (red), \( \sim 7 \) (blue), and 25 (green).

The procedure has been used to successfully model the emission and absorption of the ortho and para H\textsubscript{2}D\textsuperscript{+} ground-state rotational lines toward IRAS 16293A (Br"{u}nken et al. 2014). The (gas-grain) chemical code used here is described in detail in Sipil"{a} et al. (2015a). The gas-phase chemical reaction set used in the present work is based on the latest KIDA reaction file (Wakelam et al. 2015), which was deuterated and spin-state separated according to the prescriptions of Sipil"{a} et al. (2015a) and Sipil"{a} et al. (2015b). However, in this paper we do not explicitly consider spin states or deuteration. The physical parameters of the model and the initial chemical abundances are adopted from Tables 1 and 3 in Sipil"{a} et al. (2015a) except for the ratio of diffusion to binding energy (\( E\textsubscript{D}/E\textsubscript{B} \)), for which we adopt here a value of 0.60, in line with the recent results of Minissale et al. (2016).

3. Results

We calculated the abundances and column densities of c-C\textsubscript{3}H\textsubscript{2} and l-C\textsubscript{3}H\textsubscript{2} in the L1544 model with a grid of values for the \( k\textsubscript{DR,1}/k\textsubscript{DR,2} \) ratio, ranging from 1/1 (the KIDA value) up to 50/2 (mod hr, Table 3). We plot in Fig. 1 the abundances of c-C\textsubscript{3}H\textsubscript{2} and l-C\textsubscript{3}H\textsubscript{2} and the C\textsubscript{3}H\textsubscript{2} c/l ratio at \( t = 10^7 \) yr assuming three different values for \( k\textsubscript{DR,1}/k\textsubscript{DR,2} \). We also show the density and temperature structures of the L1544 core model (Keto et al. 2014) for completeness. Additional calculations (not shown) indicate a simple increasing trend of the C\textsubscript{3}H\textsubscript{2} c/l ratio with \( k\textsubscript{DR,1}/k\textsubscript{DR,2} \). Evidently, neither the KIDA model nor our initial modified model (mod_ch) produce a high enough c/l ratio, and a high \( k\textsubscript{DR,1}/k\textsubscript{DR,2} \) is needed to produce a c/l ratio comparable to the observed value.

To supplement the results shown in Fig. 1, we show in Fig. 2 the column densities of cyclic and linear C\textsubscript{3}H\textsubscript{2} and their ratio as functions of time toward the center of the L1544 model. We also show the observed values (Spezzano et al. 2016). The simulated column densities were convolved to a 30" beam so that they correspond to the observations. The peak value of the C\textsubscript{3}H\textsubscript{2} c/l ratio is \( \sim 32 \), which corresponds to \( 8 \times 10^4 \) yr of chemical evolution. The best fit to the observations is attained at \( t = 10^5 \) yr, when the column densities of both species and their ratio are all within the observed limits. The timescale required to attain the best fit depends on the initial chemical abundances, and so our results do not necessarily reflect the “absolute” age of the core. However, it is clear that we can only attain a high enough C\textsubscript{3}H\textsubscript{2} c/l ratio with the model if the production of cyclic C\textsubscript{3}H\textsubscript{2} is strongly boosted over that of linear C\textsubscript{3}H\textsubscript{2} in the dissociation of C\textsubscript{3}H\textsubscript{2}\textsuperscript{+}.

4. Discussion and conclusions

The C\textsubscript{3}H\textsubscript{2}\textsuperscript{+} + e\textsuperscript{−} branches (Tables 1 and 3) show that the branch listed second (C\textsubscript{3}H\textsubscript{2}\textsuperscript{+} + e\textsuperscript{−} \( \rightarrow \) C\textsubscript{3}H + H\textsubscript{2}) forms C\textsubscript{3}H, which can be converted back into C\textsubscript{3}H\textsubscript{2} through the sequence C\textsubscript{3}H + H\textsubscript{2} \( \rightarrow \) C\textsubscript{3}H\textsuperscript{+} + H\textsubscript{2} \( \rightarrow \) C\textsubscript{3}H\textsuperscript{+} + H \( \rightarrow \) C\textsubscript{3}H\textsubscript{2}. Therefore it seems reasonable to assume that increasing \( k\textsubscript{DR,2}/k\textsubscript{DR,1} \) would further increase the C\textsubscript{3}H\textsubscript{2} c/l ratio. We tested this by modifying the \( k\textsubscript{DR,2}/k\textsubscript{DR,2} \) and \( k\textsubscript{DR,1}/k\textsubscript{DR,1} \) ratios, keeping \( k\textsubscript{DR,1}/k\textsubscript{DR,1} \) fixed at 50/2. We found that increasing \( k\textsubscript{DR,2} \) and decreasing \( k\textsubscript{DR,1} \) by 10%\textsuperscript{1} further increases the peak C\textsubscript{3}H\textsubscript{2} (solid line), dust temperature (red solid line), and gas temperature (red dashed line) as functions of distance from the core center in the L1544 model of Keto et al. (2014). The different colors in the left and middle panels correspond to \( k\textsubscript{DR,1}/k\textsubscript{DR,2} = 1 \) (red), \( \sim 7 \) (blue), and 25 (green).

\[ \text{L1, page 3 of 5} \]
c/l ratio to ~40, while decreasing $k_{DR,3}$ and increasing $k_{DR,3}$ by 10\% leads to a peak $C_{3}H_{2}$ c/l ratio of ~26. Therefore we can attain a good fit to the observations with several combinations of the $k_{DR,2}$ and $k_{DR,3}$ values, and it is difficult to derive strict upper limits to acceptable values of the branching ratios based on the present analysis. However, our calculations indicate that the effect of $CH_{3}^{+} + e^{-} \rightarrow C_{3}H_{2} + H$ is greater than that of $CH_{3}^{+} + e^{-} \rightarrow CH_{2}H + H$ in determining the $C_{3}H_{2}$ c/l ratio. This is because the former pathway is a direct source of $C_{3}H_{2}$ while the latter pathway depends on multiple reactions (see above) that can also form species other than $C_{3}H_{2}$. We require $k_{DR,3}/k_{DR,3}$ clearly higher than unity. We note that the $C_{3}H/c/l$ ratio is in our models nearly independent of the various $k_{DR,3}$ values, and is ~3 at the time of the $C_{3}H_{2}$ peak, increasing to ~10 at later times (as the $C_{3}H_{2}$ c/l ratio decreases). These values agree with observations (Turner et al. 2000; Fossé et al. 2001; Cuadrado et al. 2015). However, because of its weak dependence on $k_{DR,3}$, the $C_{3}H/c/l$ ratio cannot be used to constrain the branching ratios of the $CH_{3}^{+} + e^{-}$ process.

Figure 2 demonstrates that the $C_{3}H_{2}$ c/l ratio is not constant in time, meaning that the approximation of Eq. (2) does not hold universally. For the highest $k_{DR,3}/k_{DR,3}$ ratio tested here (50/2, Table 3), the $C_{3}H_{2}$ c/l ratio presents a maximum for a relatively brief period of time and finally settles to ~15. Our chemical model shows that at late times the slow neutral-neutral reaction $H + CH_{2}CH_{2} \rightarrow H_{2} + C_{3}H_{2}$ is a significant source of $C_{3}H_{2}$. Because the $k_{DR,3}/k_{DR,3}$ ratio is so high, this neutral-neutral reaction is the most important formation pathway for $l$-$C_{3}H_{2}$ (and the second most important pathway for c-$C_{3}H_{2}$), so that the $C_{3}H_{2}$ c/l ratio is no longer controlled solely by $C_{3}H_{2}^{+} + e^{-}$. This finding is in line with previous discussions by Turner et al. (2000) and Fossé et al. (2001).

Figure 1 shows that the $C_{3}H_{2}$ c/l abundance ratio presents a minimum about 10 000 AU away from the core center and increases again at larger radii. This behavior is a result of changing physical conditions, particularly the visual extinction. A detailed analysis of our results reveals that around 10 000 AU at $t = 10^{6}$ yr, the formation of linear $C_{3}H_{2}$ is governed by reactions other than $C_{3}H_{2}^{+} + e^{-}$ because of its low rate coefficient, while at still higher radii photochemistry comes into play and the $C_{3}H_{2}^{+} + e^{-}$ reaction is again the dominant formation pathway for both cyclic and linear $C_{3}H_{2}$. However, the abundances are very low in the outer core; by far the largest part of the cyclic and linear $C_{3}H_{2}$ column densities comes from the interval ~3000–7000 AU.

The total rate coefficient for $C_{3}H_{2}^{+} + e^{-}$ assumed in the earlier modeling works of Turner et al. (2000) and Fossé et al. (2001) is equal to that adopted here (Table 1) within a factor of two. Although only a limited amount of rate coefficient data is quoted in Turner et al. (2000) and Fossé et al. (2001), it seems that the rate coefficients for the main reactions in the present context have not significantly changed in the past two decades. Still, it is clear from all the models that the observed c/l ratios can only be reproduced if the relative rate coefficients for the formation and destruction of the various $c$ and $l$ species have different values.

As stated in the Introduction and expanded upon in Appendix A, the $C_{3}H_{2}$ c/l ratio has been observed toward various environments (e.g., protostellar cores and diffuse clouds) with values ranging from 3 to ~70. The observations show a trend of increasing $C_{3}H_{2}$ c/l ratio at increasing $A_V$. The analysis presented above demonstrates that various reactions can be responsible for the formation of $C_{3}H_{2}$ (and hydrocarbons in general; see Alata et al. 2014, 2015; Duley et al. 2015). We cannot, based on the present models, draw quantitative conclusions on the expected $C_{3}H_{2}$ c/l ratio in objects other than L1544 without detailed (physical) modeling, particularly toward environments with low gas density and low $A_V$. Such an undertaking is beyond the scope of this Letter. As an example of the interpretational difficulty, we note that the $C_{3}H_{2}$ c/l ratio in the outer parts of the L1544 model, which is roughly consistent with a PDR in terms of physical conditions, follows the $k_{DR,2}/k_{DR,3}$ ratio (Fig. 1), and we checked that the $C_{3}H_{2}$ c/l column density ratio is indeed high (~17–35 depending on the time) even 100″ away from the core center. This agrees with the observations (Spezzano et al., in prep.).

We conclude that to explain the $C_{3}H_{2}$ c/l ratio observed toward L1544 (32 ± 4; Spezzano et al. 2016), the branching ratios of the various outcomes of the $C_{3}H_{2}^{+} + e^{-}$ reaction should be adjusted to favor cyclic $C_{3}H_{2}$ over linear $C_{3}H_{2}$. We will discuss the relative importance and isotopic features of various reactions and the observed $C_{3}H_{2}$ c/l ratios in a future publication.

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References

Kulczak-Jastrzębska, M., Lis, D., & Gerin, M. 2012, Acta Astron., 60, 32
Appendix A: Additional data

As noted in the introduction, the C$_3$H$_2$ c/l ratio tends to increase with increasing $A_V$. Table A.1 reports the values of the c/l ratio observed in several media. The values range from 3 in diffuse clouds to over 50 in dense cores. However, the density of the medium is not the only quantity that has an effect on this ratio. Teyssier et al. (2005) showed that inside the Horsehead Nebula PDR the C$_3$H$_2$ c/l ratio does not exceed 15 even though the extinction is on the order of 10–20, which is comparable with what is found in TMC-1(CP), where the c/l ratio has been observed to be 25 (Turner et al. 2000). In the Orion Bar a c/l ratio of 34 has been recently observed (Cuadrado et al. 2015), which is higher than the value of ~4 observed toward the Horsehead PDR (Teyssier et al. 2005). Furthermore, the c/l ratio is also different among dense cores, ranging from 25 in TMC-1(CP) to 67 in TMC-1C (Spezzano et al. 2016).

Table A.1. Observed C$_3$H$_2$ c/l ratios toward various objects, and the associated physical conditions (visual extinction $A_V$, medium density $n$, (rotational) temperature $T_{rot}$) when available in the literature.

<table>
<thead>
<tr>
<th>Source</th>
<th>$A_V$ [mag]</th>
<th>$n$ [cm$^{-3}$]</th>
<th>$T_{rot}$ [K]</th>
<th>c/l ratio</th>
<th>Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>Prestellar cores</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>TMC-1(CP$^b$)</td>
<td>$3 \times 10^4$</td>
<td>7</td>
<td>70</td>
<td>25</td>
<td>Cernicharo et al. (1991)</td>
</tr>
<tr>
<td>TMC-1(CP)</td>
<td>$3 \times 10^4$</td>
<td>7</td>
<td>20–40</td>
<td>25</td>
<td>Turner et al. (2000)</td>
</tr>
<tr>
<td>TMC-1(CP)</td>
<td>$3 \times 10^4$</td>
<td>7</td>
<td>28</td>
<td>25</td>
<td>Fossé et al. (2001)</td>
</tr>
<tr>
<td>TMC-1(edge$^c$)</td>
<td>$3 \times 10^4$</td>
<td>7</td>
<td>10</td>
<td>10</td>
<td>Fossé et al. (2001)</td>
</tr>
<tr>
<td>L183</td>
<td></td>
<td></td>
<td></td>
<td>50</td>
<td>Turner et al. (2000)</td>
</tr>
<tr>
<td>TMC-1C</td>
<td>7 (c), 4 (l)</td>
<td></td>
<td>67</td>
<td>67</td>
<td>Spezzano et al. (2016)</td>
</tr>
<tr>
<td>L1544</td>
<td>6 (c), 4 (l)</td>
<td></td>
<td>32</td>
<td>32</td>
<td>Spezzano et al. (2016)</td>
</tr>
<tr>
<td>Protostellar cores</td>
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<td>L1527</td>
<td>$7 \times 10^5$</td>
<td>12</td>
<td>12</td>
<td>12</td>
<td>Sakai &amp; Yamamoto (2013)</td>
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<tr>
<td>Translucent clouds</td>
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<tr>
<td>CB17, CB24, CB228</td>
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<td>Circumstellar envelopes</td>
<td></td>
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<td>IRC+10216</td>
<td>25</td>
<td>30</td>
<td></td>
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<td>Cernicharo et al. (1991)</td>
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<td>Diffuse clouds</td>
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<td>W49N</td>
<td>$T_{CMB}^d$</td>
<td>4</td>
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<td>Kulczak-Jastrzębska et al. (2012)</td>
</tr>
<tr>
<td>W51</td>
<td>$T_{CMB}$</td>
<td>13</td>
<td></td>
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<td>Kulczak-Jastrzębska et al. (2012)</td>
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<tr>
<td>B0355+508, B0415+379,</td>
<td>$T_{CMB}$</td>
<td>15–40</td>
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<td>Liszt et al. (2012)</td>
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<td>B2200+420, B2251+158</td>
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<td>PDRs</td>
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<td>Horsehead Nebula (edge$^e$)</td>
<td>$&lt;10^4$</td>
<td>3–5</td>
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<td>Teyssier et al. (2005)</td>
</tr>
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<td>Horsehead Nebula (dense cloud$^e$)</td>
<td>$2 \times 10^5$</td>
<td>15</td>
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<td>Teyssier et al. (2005)</td>
</tr>
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<td>Orion Bar</td>
<td>17 (c), 26 (l)</td>
<td>34</td>
<td></td>
<td></td>
<td>Cuadrado et al. (2015)</td>
</tr>
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</table>

Notes. (a) The rotational temperature of c and/or l C$_3$H$_2$, which may not be equal to the kinetic temperature. (b) CP stands for cyanopolyyne peak. (c) The edge position is (−40, 0) offset from the CP position $\alpha_{2000} = 4^h 41^m 42.5^s$, $\delta_{2000} = 25^\circ 41' 27''$ (Kawaguchi et al. 1991). (d) $T_{CMB}$ stands for the temperature of the cosmic microwave background (2.7 K). (e) The edge and cloud positions are (by estimation) (−20, −15) and (−72, −62) offset from the (0, 0) position (Teyssier et al. 2005). (f) Density corresponds to the model presented by Pety et al. (2012).