Anatomy of the AGN in NGC 5548

II. The spatial, temporal, and physical nature of the outflow from HST/COS

Observations*


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Received 7 November 2014 / Accepted 2 March 2015

ABSTRACT

Context. AGN outflows are thought to influence the evolution of their host galaxies and of super massive black holes. Our deep multiwavelength campaign on NGC 5548 has revealed a new, unusually strong X-ray obscuration, accompanied by broad UV absorption troughs observed for the first time in this object. The X-ray obscuration caused a dramatic decrease in the incident ionizing flux on the outflow that produces the long-studied narrow UV absorption lines in this AGN. The resulting data allowed us to construct a comprehensive physical, spatial, and temporal picture for this enduring AGN wind.

Aims. We aim to determine the distance of the narrow UV outflow components from the central source, their total column-density, and the mechanism responsible for their observed absorption variability.

Methods. We study the UV spectra acquired during the campaign, as well as from four previous epochs (1998–2011). Our main analysis tools are ionic column-density extraction techniques, photoionization models based on the code CLOUDY, and collisional excitation simulations.

Results. A simple model based on a fixed total column-density absorber, reacting to changes in ionizing illumination, matches the very different ionization states seen in five spectroscopic epochs spanning 16 years. The main component of the enduring outflow is situated at 3.5 ± 1.1 pc from the central source, and its distance and number density are similar to those of the narrow-emitting-line region in this object. Three other components are situated between 5–70 pc and two are farther than 100 pc. The wealth of observational constraints and the anti-correlation between the observed X-ray and UV flux in the 2002 and 2013 epochs make our physical model a leading contender for interpreting trough variability data of quasar outflows.

Conclusions. This campaign, in combination with prior UV and X-ray data, yields the first simple model that can explain the physical characteristics and the substantial variability observed in an AGN outflow.

Key words. galaxies: Seyfert

* Appendix A is available in electronic form at http://www.aanda.org
1. Introduction

AGN outflows are detected as blueshifted absorption troughs with respect to the object systemic redshift. Such outflows in powerful quasars can expel sufficient gas from their host galaxies to halt star formation, limit their growth, and lead to the co-evolution of the size of the host and the mass of its central super massive black holes (e.g., Ostriker et al. 2010; Hopkins & Elvis 2010; Soker & Meiron 2011; Ciotti et al. 2010; Faucher-Giguère et al. 2012; Borguet et al. 2013; Arav et al. 2013). Therefore, deciphering the properties of AGN outflows is crucial for testing their role in galaxy evolution.

Nearby bright AGN are excellent laboratories for studying these outflows because they yield: a) high-resolution UV data, which allow us to study the outflow kinematics and can yield diagnostics for their distance from the central source; and b) high-quality X-ray spectra that give the physical conditions for the bulk of the outflowing material (e.g., Steenbrugge et al. 2005; Gabel et al. 2005; Arav et al. 2007; Costantini et al. 2007; Kastra et al. 2012). Thus, such observations are a vital stepping stone for quantifying outflows from the luminous (but distant) quasars, for which high-quality X-ray data are not available.

For these reasons, we embarked on a deep multiwavelength campaign on the prototypical AGN outflow seen in the intensively studied Seyfert 1 galaxy NGC 5548. For the past 16 years, this outflow has shown six kinematic components in the UV band (labeled in descending order of velocity, following Crenshaw et al. 2003), and their associated X-ray warm absorbers (WA). We note that there is kinematic correspondence between the six UV components and the deduced WA components (see Sect. 5.1). Our 2013 campaign revealed a new X-ray obscurer accompanied by broad UV absorption (analyzed in Kastra et al. 2014). The appearance of the obscurer allows us to derive a comprehensive physical picture of the long-term observed outflow, which we report here.

The plan of the paper is as follows. In Sect. 2 we describe the observations and data reduction; in Sect. 3 we analyze the key component of the outflow; in Sect. 4 we discuss the remaining five components; in Sect. 5 we connect the results of the UV analysis with those of the X-ray warm absorber of the same outflow; and in Sect. 6 we compare our results with previous studies, discuss the implication of our results to the variability of AGN outflow troughs in general, and elaborate on the connection between the X-ray obscurer and the persisting outflow; in Sect. 7 we summarize our results.

2. Observations and data reduction

Our 2013 multiwavelength campaign on NGC 5548 included coordinated observations using XMM-Newton, the Hubble Space Telescope (HST), Swift, INTEGRAL and NuSTAR. The overall structure of the campaign was described in Kastra et al. (2014). A full log of all the observations is given by Mehdipour et al. (2014). Here we present a detailed analysis of the UV observations we obtained using the Cosmic Origins Spectrograph (COS) (Green et al. 2012) onboard HST. We obtained five COS observations simultaneously with five of the XMM-Newton observations between 2013 June 22 and 2013 August 01. Each two-orbit observation used gratings G130M and G160M at multiple central wavelength settings and multiple focal-plane positions (FP-POS) to cover the wavelength range from 1132 Å to 1801 Å at a resolving power of ~15000. Table A.1 lists the observation dates of the individual visits, the exposure times, and the continuum flux measured at 1350 Å in the rest frame, as well as corresponding information for archival HST observations of NGC 5548 that are also used in this analysis. The five observations from June to August 2013 were optimally weighted to produce an average spectrum that we use for our analysis. Kastra et al. (2014) describe the data reduction process from the calibration of the data to the production of this average spectrum.

The 2013 average HST/COS spectrum with all identified absorption features is shown in Fig. A.1. As described in Kastra et al. (2014), we modeled the emission from NGC 5548 using a reddened power law (with extinction fixed at E(B-V) = 0.02, Schlafly & Finkbeiner 2011), weak Fe II emission longward of 1550 Å in the rest frame, broad and narrow emission lines modeled with several Gaussian components, blueshifted broad absorption on all permitted transitions in NGC 5548, and a Galactic damped Ly α absorption line. Using this emission model, we normalized the average spectrum to facilitate our analysis of the narrow intrinsic absorption lines in NGC 5548. Figure 1 shows normalized spectra for absorption lines produced by Si III λ1206, Si IV λλ1394, 1403, C IV λλ1548, 1550, N V λλ1238, 1242, and Ly α as a function of rest-frame velocity relative to a systemic redshift of z = 0.017175 (de Vaucouleurs et al. 1991) via the NASA/IPAC extragalactic database (NED).

As shown in Table A.1, high-resolution UV spectra of NGC 5548 using HST cover an additional four epochs stretching back to 1998. We use the calibrated data sets for each of these observations as obtained from the Mikulski Archive for Space Telescopes (MAST) at the Space Telescope Science Institute (STScI). We compare the strengths of the narrow UV absorption troughs for each of these epochs with our new data set from 2013 in Figs. A.2 and A.3.

3. Physical and temporal characteristics of component 1

The key for building a coherent picture of the long-seen outflow is component 1: the strongest and highest velocity outflow component (centered at ~1160 km s⁻¹). Due to the strong suppression of incident ionizing flux by the obscurer, the 2013 HST/COS data of component 1 show a wealth of absorption troughs from ions never before observed in the NGC 5548 outflow. These data allow us to decipher the physical characteristics of this component. In Sect. 3.1 we use the column-density measurements of P V, P III, Fe III, and Si II as input in photoionization models, and derive the total hydrogen column-density (N H) for component 1, and its ionization parameter (U H). In Sect. 3.2 we use the column-density measurements of C III* and Si III* to infer the electron number density (n e), which combined with the value of the incident U H yields a distance (R) between component 1 and the central source. In Sect. 3.3 we construct a simple model based on a fixed total column-density absorber, reacting to changes in ionizing illumination, that matches the very different ionization states seen in five HST high-resolution spectroscopic epochs spanning 16 years.

3.1. Total column-density (N H) and ionization parameter (U H)

In the Appendix we describe the methods we use to derive the ion column-densities (N ion) from the outflow absorption troughs. Table A.2 gives the N ion measurements for all observed troughs of all six outflow components in the five HST epochs (spanning 16 years) of high-resolution UV spectroscopy. The upper and lower limits were derived using the apparent optical depth (AOD) method. All the reported measurements were done...
using the partial covering (PC) method. We used the power-law
(PL) method only on the C\textsc{iii} using the partial overing (PC) method. We used the power-law

To find the pair of \((U_{\text{H}}, N_{\text{H}})\) that best predicts the set of observed column-densities, we vary \(U_{\text{H}}\) and \(N_{\text{H}}\) in 0.1 dex steps to generate a grid of models (following the same approach described in Borguet et al. 2012b) and perform a minimization of the function

\[
\chi^2 = \sum_i \left( \frac{\log(N_{\text{obs}}) - \log(N_{\text{mod}})}{\sigma_{\text{obs}}} \right)^2
\]

where, for ion \(i\), \(N_{\text{obs}}\) and \(N_{\text{mod}}\) are the observed and modeled column-densities, respectively, and \(\sigma_{\text{obs}}\) is the error in the observed column-density. The measurement errors are not symmetric. We use the positive error \((+\sigma_{\text{obs}})\) when \(\log(N_{\text{mod}}) > \log(N_{\text{obs}})\) and the negative error \((-\sigma_{\text{obs}})\) when \(\log(N_{\text{mod}}) < \log(N_{\text{obs}})\).

The ionization solution for component 1 at the 2013 epoch is shown in Fig. 3. We only show constraints from \(N_{\text{H}}\) measurements, and note that all the lower limits reported in Table 2 are satisfied by this solution. We find \(log(N_{\text{H}}) = 21.5_{-0.2}^{+0.4} \) cm\(^{-2}\), and an ionization parameter of \(log(U_{\text{H}}) = -1.5_{-0.2}^{+0.4}\) where the errors are strongly correlated as illustrated by the \(1\sigma\) \(\chi^2\) contour.
3.2. Number density and distance

As shown in Fig. 4, we detect absorption troughs from the C**III** 1175 Å multiplet, arising from the metastable 3P levels of the 2s2p term. As detailed in Gabel et al. (2005), the excited C**III** 1175 Å multiplet comprises six lines arising from three J levels. The J = 0 and J = 2 levels have significantly lower radiative transition probabilities to the ground state than the J = 1 level and are thus populated at much lower densities than the latter. In particular, Fig. 5 in Borguet et al. (2012a) shows that the relative populations of the three levels are a sensitive probe to the age of our CLOUDY model for component 1 during the 2013 epoch.

In addition to the C**III** and Si**III** troughs, the 2013 COS spectra of component 1 show several additional troughs from excited states: C**II**, Si**II**, P**III**, S**III** and Fe**III** (as well as their associated resonance transitions). Careful measurement of those troughs show that in these cases, the deduced n_e is a lower limit that is larger than the critical density of the involved excited states. From these measurements we infer log(n_e) = 4.8 ± 0.1 cm⁻³ (see Fig. 4). We also detect two shallow troughs from the same metastable level of Si**III** (see panel b of Fig. 4), which we use to measure an independent and consistent value for n_e (see panel c of Fig. 4). The collisional excitation simulations shown in Figure 4c were performed using version 7.1.3 of CHIANTI (Dere et al. 1997; Landi et al. 2013), with a temperature of 10 000K (similar to the predicted temperature of our CLOUDY model for component 1 during the 2013 epoch).

As can be seen from the definition of the ionization parameter U_H (Eq. (1)), knowledge of the hydrogen number density n_H for a given U_H and N_H allows us to derive the distance R. Our photoionization models show that for component 1, log(U_H) = −1.5 and since n_e ≈ 1.2 n_H (as is the case for highly ionized plasma), n_e = 5.3 × 10⁹ cm⁻³. To determine the Q_H that affects component 1, we first calculate the bolometric luminosity using...
the average flux at 1350 Å for visits 1–5 in 2013, the redshift of the object and the obscured SED (see Fig. 2). We find \( L_{\text{bol}} = 2.6 \times 10^{44} \text{ erg s}^{-1} \) and from it, \( Q_{\text{II}} = 6.9 \times 10^{43} \text{ s}^{-1} \). Therefore, Eq. (1) yields \( R = 3.5 \pm 1.2 \text{ pc} \), where the error is determined from propagating the errors of the contributing quantities. We note that the distance and number density of component 1 are similar to those of the narrow-emitting-line-region (NELR) in this object. Using the variability of the narrow [O\( \text{III} \)]\( \lambda4959,5007 \) emission-line fluxes, **Peterson et al. (2013)** derive a radius of 1–3 pc and \( n_4 \sim 10^6 \text{ cm}^{-3} \) for the NELR.

Roughly 50% of all Seyfert 1 galaxies show evidence for absorption outflows (Crenshaw et al. 1999). From this observational evidence we assume that all Seyfert 1 galaxies have outflows that cover 50% of the solid angle around the central source. Using Eq. (1) in Arav et al. (2013) we find that the mass flux associated with the UV manifestation of component 1 is \( 1.0^{+0.6}_{-0.5} \) solar masses per year, and that the kinetic luminosity is \( 4^{+8}_{-3} \times 10^{41} \text{ erg s}^{-1} \). We note that most of the \( N_{\text{II}} \) in the various outflow components is associated with the higher ionization X-ray phase of the outflow. Therefore, we defer a full discussion of the total mass flux and kinetic luminosity of the outflow to a future paper that will present a combined analysis of the UV and X-ray data sets.

### 3.3. Modeling the temporal behavior of the outflow

The absorption troughs of component 1 change drastically between the five HST high-resolution spectroscopic epochs spanning 16 years (see Figs. A.2 and A.3). After finding the location and physical characteristics of component 1 using the 2013 data, the next step is to derive a self-consistent temporal picture for this component. There are two general models that explain trough variability in AGN outflows (e.g., Barlow et al. 1992; Gabel et al. 2003; Capellupo et al. 2012; Arav et al. 2012; Filiz Ak et al. 2013, and references therein). One model attributes the trough variability to changes of the ionizing flux experienced by the outflowing gas. In its simplest form, this model assumes that the total \( N_{\text{II}} \) along the line of sight does not change as a function of time. A second model invokes material moving across the line of sight, which in general causes changes of \( N_{\text{II}} \) along the line of sight as a function of time to explain the observed trough changes.

In the case of component 1, we have enough constraints to exclude the model of material moving across our line of sight. The outflow is situated at 3.5 pc from the central source, which combined with the estimated mass of the black hole in NGC 5548 (\( 4 \times 10^7 \) solar masses; Pancoast et al. 2014), yields a Keplerian speed of \( 1.9 \times 10^3 \text{ cm s}^{-1} \) at that distance. As can be seen from Fig. A.1, 2/3 of the emission at the wavelength of component 1 arises from the C IV broad emission line (BEL). Therefore, the transverse motion model crucially depends on the ability of gas clouds to cross most of the projected size of the broad line region (BLR) in the time spanning the observations epochs. Reverberation studies (Korista et al. 1995) give the diameter of the C IV BLR as 15 light days or 3.9 \( \times 10^{16} \text{ cm} \), which for \( v_\perp = 1.9 \times 10^4 \text{ cm s}^{-1} \), yields a crossing time of \( 2.0 \times 10^9 \text{ s} \), i.e., 65 years.

Thus, in the 16 years between our epochs, material that moves at the Keplerian velocity, 3.5 pc away from the NGC 5548 black hole, will cross only about 25% of the projected size of the C IV BLR. Therefore, the much larger change in the residual intensity of the component 1 C IV trough cannot be attributed to new material appearing due to transverse motion at this distance. We note that 25% motion across the projected size of the C IV BLR is a highly conservative limit for two reasons: 1) at certain velocities there are changes of 50% in the residual intensity in the component 1 C IV trough between the 2011 and 2013 epochs; and in the elapsing 2 years, transverse motion will only cover 3% of the projected size of the C IV BLR; 2) material that moves away from the central source under the influence of radial forces should conserve its angular momentum. Therefore, if it moved to distances that are much larger compared with its initial distance, its \( v_\perp \ll v_{\text{kep}} \) at its current distance. We conclude that even under favorable conditions, the transverse motion model of gas into or out of the line of sight cannot explain the observed behavior of component 1 over the five observed epochs.

Can changes of the ionizing flux experienced by the outflowing gas explain the observed trough changes? We construct such a model under the simplest and restrictive assumption that the \( N_{\text{II}} \) of component 1 did not change over the 16 years spanning the five high resolution UV spectral epochs. Furthermore, for \( \log (n_4) = 4.8 \text{ cm}^{-3} \), the absorber should react to changes in incident ionizing flux on time-scales of 5 days (see Eq. (3) here, and discussion in Arav et al. 2012). Therefore, for component 1 we use the restrictive assumption of a simple photoionization equilibrium, determined by the flux level of the specific observation.

In 1998 the AGN was in a high flux level of \( F = 6 \) (measured at 1350 Å rest-frame and given in units of \( 10^{14} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ Å}^{-1} \)). At that epoch the absorber only showed a \( \text{Ly}\alpha \) trough necessitating \( \log (U_{\text{II}}) > 0.1 \) (otherwise a N V trough would be detected, see Fig. 5). In 2002 the AGN was in a medium flux level with \( F = 2 \), at which time the absorber showed C IV and N V troughs in addition to Ly\( \alpha \). In 2004 the AGN was at a historically low flux of \( F = 0.25 \).
In that epoch a Si\textsc{iii} trough appeared in addition to the C\textsc{iv}, N\textsc{v} and Ly\textalpha{} troughs; however, a C\textsc{ii} trough did not appear. The combination of the Si\textsc{iii} and C\textsc{ii} constraints necessitates \(-1.3 < \log(U_{H}) < -1.15\) (see Fig. 5). The change in \(\log(U_{H})\) required by the photoionization models agrees remarkably well with the change in flux between the 1998 and 2004 epochs as \(\log(F_{\text{2004}}/F_{\text{1998}}) = -1.4\). Thus, a constant \(N_{H}\) absorber yields an excellent fit for the absorption features from two epochs with the same SED (high SED in Fig. 5), but with very different \(U_{H}\) values. Comparison of several key troughs between the five epochs is shown in Figs. A.2 and A.3. The 1530 Å flux measurements, plus observation details are given in Table A.1 and the derived column-densities for all the outflow features are given in Table A.2.

In 2013 the AGN flux was \(F = 3\). With this flux level and assuming the same SED, the \(U_{H}\) value should have been 50% higher than in 2002, and we would expect to see only C\textsc{iv}, N\textsc{v} and Ly\textalpha{} troughs. Instead we also detect Si\textsc{iii}, C\textsc{ii}, Si\textsc{ii} and Al\textsc{ii}. Therefore, the incident SED for component 1 must have changed, and indeed the 2013 soft X-ray flux is 25 times lower compared to that of the 2002 epoch (see Fig. 1 in Kaastra et al. 2014). This drop is caused by the newly observed absorber close to the AGN, which does not fully cover the source (Kaastra et al. 2014). We found a good match to the UV absorption and soft X-ray flux with an SED that is similar to that of the high flux one. A 1 Ry, but abruptly drops to 6% of that flux between X-ray flux with an SED that is similar to that of the high flux one (see discussion below), this \(n_{e}\) range yields distances of \(5 < R < 15\) pc. Similarly, component 6 shows C\textsc{iv} and N\textsc{v} troughs in 2011 but not in 2002. This nine years timescale yields \(R < 100\) pc. Component 2 and component 4 do not show changes in the UV absorption between any of the epochs. Therefore, we can derive a lower limit for their distance of \(R > 130\) pc.

4.2. Constraining \(N_{H}\) and \(U_{H}\)

It is not feasible to put physically interesting constraints on components 2, 4 and 6. First, they only show troughs from C\textsc{iv}, N\textsc{v} and Ly\textalpha{}, which (based on our analysis of component 1) are probably highly saturated. Second, the ionization time-scales of components 2 and 4 are larger than 16 years. Therefore, even if a \(U_{H}\) can be deduced from the measurements, it will only be a representative average value for a period of time larger than 16 years.

Figure 7 shows the \(N_{H}/U_{H}\) phase plot for component 3 based on the \(N_{H}\) reported in Table A.2 (the Ly\textalpha{} and C\textsc{iv} \(N_{H}\) lower limits are trivially satisfied by the lower limit shown for the N\textsc{v} \(N_{H}\)). The phase plot constraints given by the \(N_{H}\) measurements are mostly parallel to each other. Therefore, the \(N_{H}/U_{H}\) constraints are rather loose, allowing a narrow strip from \(\log(N_{H}) = 19.6\) cm\(^{-2}\) and \(\log(U_{H}) = -2\), to \(\log(N_{H}) = 21.5\) cm\(^{-2}\) and \(\log(U_{H}) = -1.1\). If we take the most probable value of \(\log(U_{H}) = -1.3\), the distance estimate for component 3 will drop by 30% compared with the estimate of \(5 < R < 15\) pc, which used the \(\log(U_{H}) = -1.5\) of component 1. For component 5, the situation is rather similar, as the detected Si\textsc{iii} and Si\textsc{iv} allow a narrow strip from \(\log(N_{H}) = 19.2\) cm\(^{-2}\) and \(\log(U_{H}) = -1.8\), to \(\log(N_{H}) = 20.7\) cm\(^{-2}\) and \(\log(U_{H}) = -1.2\), while the lowest \(\chi^{2}\) is achieved at \(\log(N_{H}) = 20.7\) cm\(^{-2}\) and \(\log(U_{H}) = -1.3\).

5. Comparison with the warm absorber analysis

Here we compare the physical characteristics inferred from the outflows’ UV diagnostics to the properties of the X-ray manifestation of the outflow known as the warm absorber (WA). Since the soft X-ray flux in our 2013 data is very low due to the appearance of the obscurer, we cannot characterize the WA that is connected with the six UV outflow components at that epoch. Our main inferences about the WA are from the 2002 epoch when we obtained simultaneous X-ray and UV spectra of the outflow (when no obscurer was present) that gave a much higher soft X-ray flux (compared with the 2013 observations) and allowed a detail modeling of the WA in that epoch (Steenbrugge et al. 2005; Kaastra et al. 2014). Due to the inherent complications of comparing analyses on different spectral regions (X-ray
and UV) separated by 11 years (2002 and 2013), of a clearly time-dependent phenomenon, we defer a full comparison to a later paper (Ebrero et al., in prep.). Here we outline some of the main points in such a comparison, based on the analysis presented here and the published analysis of the WA (Steenbrugge et al. 2005; Kaastra et al. 2014) and discuss some of the similarities and challenges of such a combined analysis.

5.1. Kinematic similarity

There is kinematic correspondence between the UV absorption troughs in components 1–5 and the six ionization components (A–F) of the X-ray WA (Kaastra et al. 2014, see Table 1 here). X-ray components F and C span the width of UV component 1. X-ray component E matches UV component 2. The lowest ionization X-ray components, A and B comprise the full width of the blended UV troughs of components 3 and 4. Finally, X-ray component D kinematically matches UV component 5. However, as we show in point 2 below, this kinematic matching is physically problematic as the ionization parameters of WA components C, D, E and F are too high to produce observed troughs from C IV and N V that are observed in all the UV components.

5.2. Comparing similar ionization phases

We note that the X-ray analysis of the WA in the Chandra 2002 observations uses a different ionization parameter ($\xi$) than the $U_{H}$ we use here; where $\xi \equiv L/(n_H r^2)$ (erg cm s$^{-1}$) with $n_H$ being the hydrogen number density, $L$ the ionizing luminosity between 13.6 eV and 13.6 keV and $r$ the distance from the central source. For the high SED, $\log(U_{H}) = \log(\xi) + 1.6$. In Table 1, we give the $\log(U_{H})$ for the WA components. From the WA analysis and Fig. 5 here, we deduce that 90% of the WA material (components C, D, E and F in Table S2 of Kaastra et al. 2014) is in a too high ionization stage to produce measurable lines from the UV observed ions (e.g., C IV, N V). Only component A and B of the WA are at low enough ionization states to give rise to...
the UV observed material. The issue of pressure equilibrium between the different ionization phases is beyond the scope of this paper, and will be discussed in a forthcoming paper.

5.3. Assuming constant \( N_{\text{H}} \) for the UV components and components A and B of the WA

Our temporal model for component 1 has a constant \( N_{\text{H}} \) value in all the observed epochs, including 2002. The model also predicts the \( U_{\text{H}} \) of the 2002 epoch (see Fig. 5). We can therefore compare the predictions of this model to the results of the re-analysis of the 2002 WA (Kaastra et al. 2014), provided that the \( N_{\text{H}} \) for components A and B of the WA also did not change over the 11 years between the epochs. We note that since UV component 1 is the closest to the central source, the assumption of constant \( N_{\text{H}} \) for the other UV components, over this 11 years time period, is reasonable (see discussion in Sect. 3.3). Therefore, we use the same ionization assumptions for UV components 3 and 5 as for component 1. That is, their \( N_{\text{H}} \) is fixed to the 2013 value and their \( \log(U_{\text{H}})_{2002} = \log(U_{\text{H}})_{2013} + 1.1 \), which are the values we list in Table 1. We do not have empirical constraints on the distances of WA components A-F from the central source.

5.4. Comparing UV components 1 and 3 to components A and B of the WA

Using proto-Solar abundances (Lodders et al. 2009), our 2002 model prediction for UV component 1 has \( \log(N_{\text{H}}) = 21.5_{-0.2}^{+0.4} \) cm\(^{-2} \), and an ionization parameter of \( \log(U_{\text{H}}) = -0.4_{-0.5}^{+0.2} \) (see Sect. 3.1 and Fig. 5). This model gives a good match for the UV data of that epoch (2002) and its \( \log(U_{\text{H}}) \) is in between those of WA components A and B (\( \log(U_{\text{H}}) = -0.8 \) and \( \log(U_{\text{H}}) = -0.1 \)). However, there are two inconsistencies between the models. First, components A and B have a total \( \log(N_{\text{H}}) = 20.95 \pm 0.1 \) cm\(^{-2} \) or about 2\( \sigma \) disagreement with that of UV component 1. This discrepancy is mainly due to the limit on the O\( \text{vii} \) \( N_{\text{ion}} \), which arises from the bound-free edge of this ion in the X-ray data. In the WA model, about 95% of the O\( \text{vii} \) \( N_{\text{ion}} \) arises from components A and B. Second, the reported velocity centroids for WA components A and B (\( \sigma_{\text{H}} \pm 34 \) km s\(^{-1} \) and \( \sigma_{\text{H}} \pm 31 \) km s\(^{-1} \), respectively) are in disagreement with the velocity centroids of UV component 1 (\( \sigma_{\text{H}} = 1160 \) km s\(^{-1} \)) and its 300 km s\(^{-1} \) width.

UV component 3 has a velocity centroid at \( \sigma_{\text{H}} = 640 \) km s\(^{-1} \) and a width of \( \sigma_{\text{H}} = 200 \) km s\(^{-1} \), and therefore is a better velocity match with WA components A and B. The large uncertainties in the inferred \( N_{\text{H}} \) and \( U_{\text{H}} \) for UV component 3 (see Fig. 7 and Table 1), make these values consistent with the \( N_{\text{H}} \) and \( U_{\text{H}} \) deduced for WA components A and B. However, the uncertainties also allow UV component 3 to have a negligible \( N_{\text{H}} \) compared to WA components A and B.

We note that with the current analyses, the better the agreement between UV component 3 and WA components A and B, the worse is the disagreement between UV component 1 and WA components A and B. This is because UV component 1 already predicts higher values of \( N_{\text{H}} \) and O\( \text{vii} \) \( N_{\text{ion}} \) than are measured in WA components A and B, and the kinematics of the deduced O\( \text{vii} \) \( N_{\text{ion}} \) disagree considerably. In points 5 and 6 below we identify two possible ways to alleviate and even eliminate these apparent discrepancies.

5.5. Existence of considerable O\( \text{vii} \) \( N_{\text{ion}} \) at the velocity of UV component 1

The 2002 X-ray spectra presented by Steenbrugge et al. (2005) consist of two different data sets that were acquired in the same week: 170 ks observations taken with the High Energy Transmission Grating Spectrometer (HETGS) and 340 ks observations with the Low Energy Transmission Grating Spectrometer (LETGS). Figure 2 in Steenbrugge et al. (2005) shows some of the low ionization WA troughs in velocity presentation, where the dotted lines show the position of the UV components (with somewhat different velocity values than we use here due to the use of a slightly different systemic redshift for the object). The LETGS data of the O\( \text{vii} \) and O\( \text{v} \) troughs are consistent with one main kinematic component matching the velocity of UV component 3. However, the more noisy but higher resolution HETGS data for the same transitions, show two subtroughs: one corresponding to UV component 1 and one
to UV component 3. Therefore, it is possible that much of the O\textsc{vii} \textit{N}_{\text{ion}} is associated with UV component 1.

5.6. Abundances considerations

As mentioned in Sect. 3, for the UV analysis, we use pure proto-Solar abundances (Lodders et al. 2009), which well-match the measured \textit{N}_{\text{ion}} from the UV data. But these models produce considerably more O\textsc{vii} \textit{N}_{\text{ion}} in the 2002 epoch, than the measured O\textsc{vii} \textit{N}_{\text{ion}} in the warm absorber data. However, the \textit{N}_\text{H} (and therefore also the O\textsc{vii} \textit{N}_{\text{ion}}) of UV component 1 is critically dependent on the assumed phosphorus abundance. Ionization models with all elements having proto-Solar abundances except phosphorus, for which we assume twice proto-Solar abundance, preserve the fit to the UV data (at 1/3 the \textit{N}_\text{H} value) and at the same time match the O\textsc{vii} \textit{N}_{\text{ion}} measured in the X-ray warm absorber at the 2002 epoch. Larger over-abundances of phosphorus further reduce the \textit{N}_\text{H} value and therefore the predicted O\textsc{vii} \textit{N}_{\text{ion}} for the 2002 epoch.

But are such assumed abundances physically reasonable? AGN outflows are known to have abundances higher by factors of two (Arav et al. 2007) and even ten (Gabel et al. 2006) compared with the proto-Solar values. In particular, phosphorus abundance in AGN outflows, relative to other metals, can be a factor of several higher than in proto-Solar abundances (see Sect. 4.1 in Arav et al. 2001). Furthermore, the theoretical expectations for the value of chemical abundances in an AGN environment as a function of metallicity are highly varied. The leading models can differ about relative abundances values by factors of three or more (e.g., comparing the values of Hamann & Ferland 1993; Ballero et al. 2008).

We conclude that if roughly half of the observed O\textsc{vii} \textit{N}_{\text{ion}} is associated with UV component 1 (as discussed in point 5), and if the relative abundance of phosphorus to oxygen is twice solar or larger, than the \textit{N}_\text{H}, \textit{U}_\text{H} and velocity distribution of the WA and UV outflowing gas can be consistent.

6. Discussion

6.1. Comparison with Crenshaw et al. (2009)

Here we compare our results with previous work by Crenshaw et al. (2009, hereafter C09) on the enduring outflow in NGC 5548. For the first time a simple model of a constant \textit{N}_\text{H} absorber yields a physical picture that is consistent with the substantial trough variations seen in all epochs of high-resolution UV spectral observations. The trough changes are explained solely by the observed differences in the incident ionizing flux. In addition, we determine robust distances for (or limits on) all six kinematic components.

Our results differ considerably from those previously found for this outflow in C09. The HST epochs of 1998, 2002 and 2004 were analyzed by C09. In Table 2 we show for component 1 the deduced log(\textit{N}_\text{H}) and log(\textit{U}_\text{H}) values for these 3 epochs, for both C09 (see their Table 6) and this work (see Fig. 5 here). It is clear that we derive very different parameters than those of C09.

The main reason for these differences is the reliability of the measured ionic column densities (\textit{N}_{\text{ion}}), which are used as input in photoionization models. Both works use CLOUDY, assume the same abundances and use a similar SED. C09 base their photoionization solution on the \textit{N}_{\text{ion}} of C\textsc{iv} and N\textsc{v} (see C09 Table 5). We obtain similar C\textsc{iv} and N\textsc{v} \textit{N}_{\text{ion}} values in these epochs (see Table A.2 here), but we only use them as lower limits. The new diagnostic troughs seen during the 2013 campaign (especially P\textsc{v}) demonstrate this crucial point. The existence of similar depth C\textsc{iv} and P\textsc{v} troughs in component 1 is evidence that the actual C\textsc{iv} \textit{N}_{\text{ion}} is at least a factor of 100 larger than the lower limit (see discussion in Borguet et al. 2012a). A similar situation occurs for N\textsc{v}. This assertion is confirmed quantitatively by our full photoionization models. With such large saturation factors, photoionization solutions for this outflow based on the trough-derived C\textsc{iv} and N\textsc{v} \textit{N}_{\text{ion}} are unreliable.

Examining the results for the 2004 epoch demonstrates these issues. Using the C\textsc{iv} and N\textsc{v} \textit{N}_{\text{ion}} values of C09 for component 1, we can reproduce their log(\textit{U}_\text{H}) and log(\textit{N}_\text{H}) parameters. However, the 2004 HST spectrum shows clear Si\textsc{iii} troughs in components 1 and 3 (see Fig. A.3f here) that are unaccounted for in C09. The C09 solution shown in Table 2 is incompitable with the Si\textsc{iii} \textit{N}_{\text{ion}} measurement, even if that trough is unsaturated. To reproduce the Si\textsc{iii} measurement using the C09 log(\textit{N}_\text{H}) = 20.35 cm$^{-2}$, log(\textit{U}_\text{H}) need to be $-1.7 \pm 0.1$ instead of the C09 value of log(\textit{U}_\text{H}) = $-0.55$. However, at log(\textit{N}_\text{H}) = 20.35 cm$^{-2}$ and log(\textit{U}_\text{H}) = $-1.7$, CLOUDY yields a C\textsc{iv} \textit{N}_{\text{ion}} that is 100 times larger than the value C09 deduce from the observations (see their Table 5).

In conclusion, our improved physical modeling is due to powerful diagnostics that were revealed during the 2013 campaign, establishing that many of the observed troughs are highly saturated; and also because C09 did not account for the Si\textsc{iii} troughs associated with components 1 and 3 in the 2004 data.

6.2. Implications for BALQSO variability studies

Our multiwavelength campaign has significant consequences for studies of absorption trough variability in quasar outflows, and in particular for the intensive studies of trough variability in broad absorption line (BAL) quasars (e.g., Barlow et al. 1992; Capellupo et al. 2012; Filiz Ak et al. 2013, and references therein). AGN outflow absorption systems with different velocity width, and different AGN luminosities, share similar physical characteristics (e.g., ionization states, trough variability time scales, trough variability pattern, solar or larger metallicity, non-black saturation in the troughs). Therefore, we can make a comparison between the mechanisms in the narrow outflow of NGC 5548 to that of BALs in QSOs.

As discussed in Sect. 3.3, the two main proposed mechanisms for trough variability are (1) reaction to changes in ionizing flux of a constant absorber (which is the model we successfully use to explain the NGC 5548 outflow trough changes); and (2) absorber motion across the line of sight (e.g., Gabel et al. 2003), which as we demonstrated, cannot explain the variability of the enduring, narrow-trough outflow in NGC 5548 (see Sect. 3.3).

In some BALQSO cases the rest-frame UV flux around 1350 Å does not change appreciably between the studied epochs while significant trough variability is observed. This behavior is taken as an argument against mechanism (1) (e.g., Barlow et al. 1992; Filiz Ak et al. 2013) as the ionizing flux (below 912 Å)

### Table 2. Comparing the photoionization solutions for component 1.

<table>
<thead>
<tr>
<th>Epoch</th>
<th>\textit{This work}</th>
<th>\textit{C09}</th>
<th>\textit{This work}</th>
<th>\textit{C09}</th>
</tr>
</thead>
<tbody>
<tr>
<td>1998</td>
<td>21.50</td>
<td>19.23</td>
<td>0.14</td>
<td>$-0.50$</td>
</tr>
<tr>
<td>2002</td>
<td>21.50</td>
<td>19.55</td>
<td>$-0.41$</td>
<td>$-0.74$</td>
</tr>
<tr>
<td>2004</td>
<td>21.50</td>
<td>20.35</td>
<td>$-1.27$</td>
<td>$-0.55$</td>
</tr>
</tbody>
</table>
is assumed to correlate with the longer wavelength UV flux. However, our simultaneous UV/X-ray observations show a clear case where the ionizing photon flux drops by a factor of 25 between the 2002 and 2013 epochs, while the 1350 Å UV flux actually increases by 50%.

Mechanism (1) is a simpler explanation for cases where velocity-separated outflow troughs change in the same way between different epochs (i.e., when all troughs become either shallower or deeper: Capellupo et al. 2012; Filiz Ak et al. 2013). As we demonstrated, the objection to this model based on the assumed correlation between the rest-frame UV flux around 1350 Å and the ionizing flux might not be valid. Therefore, due to its simplicity, mechanism (1) is a strong contender for interpreting trough variability data of quasar outflows.

We note that the above discussion does not exclude the possibility of trough variability due to absorber motion across the line of sight, in some cases. Indeed it is possible that this mechanism explains the variability in the newly discovered broad UV absorption in NGC 5548, as argued in Kaastra et al. (2014).

The data set of NGC 5548 (spanning five high-resolution spectroscopic epochs between 1998 and 2013) gives high S/N spectra of UV absorption troughs that arise from many different ions, and simultaneously yields the crucial soft X-ray flux spectra of UV absorption troughs that arise from many different ions (Peterson et al. 2013). See Sect. 3.2.

The distance and number density of component 1 are similar to those of the narrow-emitting-line-region in this object (Peterson et al. 2013). See Sect. 3.2.

The wealth of observational constraints makes our changes-of-ionization model a leading contender for interpreting trough variability data of quasar outflows, in particular BAL variability. See Sect. 6.2.

Components 3 and 5 are situated between 5–15 pc from the central source, component 6 is closer than 100 pc and components 2 and 4 are further out than 130 pc. See Sect. 4.

A detailed comparison of the physical characteristics inferred from the outflows’ UV diagnostics with those of the X-ray Warm Absorber is deferred to a future paper. Here we outline some of the main points in such a comparison, and discuss some of the similarities and challenges of such a combined analysis. See Sect. 5.

### 6.3. Implications for the X-ray obscurer

Our results about the enduring outflow have implications for the X-ray obscurer and the broad UV absorption discovered by Kaastra et al. (2014). The derived transmission for the X-ray obscurer is consistent with the SED required for the 2013 spectrum, thus showing that the obscurer is closer to the super massive black holes than 3.5 pc, and that its shadow influences the conditions in the more distant narrow UV absorbers.

## 7. Summary

In 2013 we executed the most comprehensive multiwavelength spectroscopic campaign on any AGN to date, directed at NGC 5548. This paper presents the analysis’ results from our HST/COS data of the enduring UV outflow, which is detected in six distinct kinematic components. Our campaign revealed an unusually strong X-ray obscuration. The resulting dramatic decrease in incident ionizing flux on the outflow allowed us to construct a comprehensive physical, spatial and temporal picture for the well-studied AGN wind in this object. Our main findings are listed below (see Fig. 6 for a graphic illustration of our results):

1. Our best constraints are obtained for component 1 (the highest velocity component). It is situated at $R = 3.5 \pm 1.2$ pc from the central source, has a total hydrogen column-density of $\log(N_H) = 21.5^{+0.4}_{-0.2}$ cm$^{-2}$, an ionization parameter of $\log(U_0) = -1.5^{+0.2}_{-0.2}$, and an electron number density of $\log(n_e) = 4.8 \pm 0.1$ cm$^{-3}$. This component probably carries the largest $N_H$ associated with the UV outflow. See Sects. 3.1 and 3.2.

2. For component 1 a simple model based on a fixed total column-density absorber, reacting to changes in ionizing illumination, matches the very different ionization states seen at five spectroscopic epochs spanning 16 years. See Sect. 3.3.

3. The distance and number density of component 1 are similar to those of the narrow-emitting-line-region in this object (Peterson et al. 2013). See Sect. 3.2.

### Acknowledgements

This work was supported by NASA through grants for HST program number 13184 from the Space Telescope Science Institute, which is operated by the Association of Universities for Research in Astronomy, Incorporated, under NASA contract NAS5-26555. SRON is supported financially by NWO, the Netherlands Organization for Scientific Research. M.M. acknowledges the support of a Studentship Enhancement Programme awarded by the UK Science & Technology Facilities Council (STFC). P.-O.P. and F.U. thank financial support from the CNES and from the CNRS/PICS. F.U. acknowledges Ph.D. funding from the VINCI program of the French-Italian University. K.C.S. acknowledges financial support from the Fondo Fortalecimiento de la Productividad Científica VRIDT 2013. E.B. is supported by grants from Israel’s MoST, ISF (1163/10), and I-CORE program (1937/12). J.M. acknowledges funding from CNRS/PNHE and CNRS/PICS in France. G.M. and F.U. acknowledges financial support from the Italian Space Agency under grant ASI/INAF I/037/12/0-011/13. B.M.P. acknowledges support from the US NSF through grant AST-1008882. M.C., S.B., G.M. and A.D.R. acknowledge INAF/PICS support. G.P. acknowledges support via an EU Marie Curie Intra-European fellowship under contract no. FP-PEOPLE-2012-IEF-331095. M.W. acknowledges the support of a Ph.D. studentship awarded by the UK Science & Technology Facilities Council (STFC). The data used in this research are stored in the public archives of the satellites that are involved. We thank the International Space Science Institute (ISSI) in Bern for support. This work is based on observations obtained with XMM-Newton, an ESA science mission with instruments and contributions directly funded by ESA Member States and the USA (NASA). It is also based on observations with INTEGRAL, an ESA project with instrument and science data center funded by ESA member states (especially the PI countries: Denmark, France, Germany, Italy, Switzerland, Spain), Czech Republic, and Poland and with the participation of Russia and the USA. This work made use of data supplied by the UK Swift Science Data Centre at the University of Leicester. This research made use of the Chandra Transmission Grating Catalog and archive (http://tgcat.mit.edu). This research has made use of data obtained with the NuSTAR mission, a project led by the California Institute of Technology (Caltech), managed by the Jet Propulsion Laboratory (JPL) and funded by NASA, and has utilized the NuSTAR Data Analysis Software (NUSTARDAS) jointly developed by the ASI Science Data Center (ASDC, Italy) and Caltech (USA). Figure 6 was created by Ann Feild from STScI.

### References


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Appendix A: Ionic column-density measurements

The column-density associated with a given ion as a function of the radial velocity \( v \) is defined as:

\[
N_{\text{ion}}(v) = \frac{3.8 \times 10^{14}}{f_j \lambda_j} \langle \tau_j(v) \rangle \quad \text{(cm}^{-2} \text{km}^{-1} \text{s}) \tag{A.1}
\]

where \( f_j \), \( \lambda_j \) and \( \langle \tau_j(v) \rangle \) are respectively the oscillator strength, the rest wavelength and the average optical depth across the emission source of the line \( j \) for which the optical depth solution is derived (see Edmonds et al. 2011). The optical depth solution across a trough is found for a given ion by assuming an absorber model. As shown in Edmonds et al. (2011), the major uncertainty on the derived column-densities comes from the choice of absorption model. In this study we investigate the outflow properties using column-densities derived from three common absorber models.

Assuming a single, homogeneous emission source of intensity \( F_0 \), the simplest absorber model is the one where a homogeneous absorber parameterized by a single optical depth fully covers the photon source. In that case, known as the apparent optical depth scenario (AOD), the optical depth of a line \( j \) as a function of the radial velocity \( v \) in the trough is simply derived by the inversion of the Beer-Lambert law: \( \tau_j(v) = -\ln(F_j(v)/F_0(v)) \), where \( F_j(v) \) is the observed intensity of the line.

Early studies of AGN outflows pointed out the inadequacy of such an absorber model, specifically its inability to account for the observed departure of measured optical depth ratio between the components of typical doublet lines from the expected laboratory line strength ratio \( R = \lambda_j f_j / \lambda_j f_j \). Two parameter absorber models have been developed to explain such discrepancies.

The partial covering model (e.g. Hamann et al. 1997; Arav et al. 1999, 2002, 2005) assumes that only a fraction \( C \) of the emission source is covered by absorbing material with constant optical depth \( \tau \). In that case, the intensity observed for a line \( j \) of a chosen ion can be expressed as

\[
F_j(v) = F_0(v)(1 + C(\tau)(e^{-\tau j(v)} - 1)). \tag{A.2}
\]

Our third choice are inhomogeneous absorber models. In that scenario, the emission source is totally covered by a smooth distribution of absorbing material across its spatial dimension \( x \). Assuming the typical power law distribution of the optical depth \( \tau(x) = \tau_{\text{max}} x^q \) (de Kool et al. 2002; Arav et al. 2005, 2008), the observed intensity observed for a line \( j \) of a chosen ion is given by

\[
F_j(v) = F_0(v) \int_0^1 e^{-\tau_{\text{max}} x^q} dx. \tag{A.3}
\]

Once the line profiles have been binned on a common velocity scale (we choose a resolution \( \Delta v = 20 \text{ km} \text{s}^{-1} \), slightly lower than the resolution of COS), a velocity dependent solution can be obtained for the couple of parameters \( (C, \tau) \) or \( (\alpha, \tau_{\text{max}}) \) of both absorber models as long as one observes at least two lines from a given ion, sharing the same lower energy level. Once the velocity dependent solution is computed, the corresponding column density is derived using Eq. (A.1) where \( \tau_j(v) = C_{\text{ion}}(v) \tau_j(v) \) for the partial covering model and \( \tau_j(v) = \tau_{\text{max}}(v)/(\tau_{\text{max}}(v) + 1) \) for the power law distribution. Note that the AOD solution can be computed for any line (singlet or multiplet), without further assumption on the model, but will essentially give a lower limit on the column-density when the expected line strength ratio observed is different from the laboratory value.

### Table A.1. Observations and flux values for all epochs.

<table>
<thead>
<tr>
<th>Epoch</th>
<th>Observing dates</th>
<th>Instrument</th>
<th>Grating</th>
<th>Exposure</th>
<th>avg ( F_{1350} )</th>
<th>Fit ( F_{1350} )</th>
<th>( \alpha' )</th>
</tr>
</thead>
<tbody>
<tr>
<td>2013_v1</td>
<td>2013 June 22</td>
<td>HST:COS</td>
<td>G130M</td>
<td>1.8 ks</td>
<td>1.95 ± 0.36</td>
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<td></td>
</tr>
<tr>
<td>2013_v2</td>
<td>2013 July 12</td>
<td>HST:COS</td>
<td>G160M</td>
<td>2.1 ks</td>
<td>2.22 ± 0.34</td>
<td></td>
<td></td>
</tr>
<tr>
<td>2013_v3</td>
<td>2013 July 24</td>
<td>HST:COS</td>
<td>G160M</td>
<td>2.2 ks</td>
<td>3.70 ± 0.44</td>
<td></td>
<td></td>
</tr>
<tr>
<td>2013_v4</td>
<td>2013 July 30</td>
<td>HST:COS</td>
<td>G160M</td>
<td>2.0 ks</td>
<td>3.51 ± 0.44</td>
<td></td>
<td></td>
</tr>
<tr>
<td>2013_v5</td>
<td>2013 August 01</td>
<td>HST:COS</td>
<td>G130M</td>
<td>2.0 ks</td>
<td>3.28 ± 0.40</td>
<td></td>
<td></td>
</tr>
<tr>
<td>2013_v6</td>
<td>2013 December 20</td>
<td>HST:COS</td>
<td>G160M</td>
<td>2.2 ks</td>
<td>3.26 ± 0.40</td>
<td>3.14 ± 0.02</td>
<td>-0.79 ± 0.03</td>
</tr>
<tr>
<td>2013_v345</td>
<td>2013 July 24 to 2013 August 01</td>
<td>HST:COS</td>
<td>G130M</td>
<td>6.0 ks</td>
<td>3.51 ± 0.24</td>
<td>3.44 ± 0.03</td>
<td>-0.776 ± 0.014</td>
</tr>
<tr>
<td>2013_v12345</td>
<td>2013 June 22 to 2013 August 01</td>
<td>HST:COS</td>
<td>G130M</td>
<td>12.8 ks</td>
<td>3.11 ± 0.19</td>
<td>3.051 ± 0.015</td>
<td>-0.736 ± 0.026</td>
</tr>
<tr>
<td>2011</td>
<td>2011 June 16</td>
<td>HST:COS</td>
<td>G130M</td>
<td>1.9 ks</td>
<td>6.17 ± 0.52</td>
<td>6.13 ± 0.05</td>
<td>-0.86 ± 0.07</td>
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<tr>
<td>2004</td>
<td>2004 February 10</td>
<td>HST:STIS</td>
<td>E140M</td>
<td>52.2 ks</td>
<td>0.25 ± 0.07</td>
<td>0.203 ± 0.017</td>
<td>-1.86 ± 0.29</td>
</tr>
<tr>
<td>2002</td>
<td>2002 January 22</td>
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<td>E140M</td>
<td>15.3 ks</td>
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<td>-1.46 ± 0.06</td>
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<td>1998 March 11</td>
<td>HST:STIS</td>
<td>E140M</td>
<td>4.75 ks</td>
<td>6.41 ± 0.75</td>
<td>6.43 ± 0.12</td>
<td>-1.51 ± 0.13</td>
</tr>
</tbody>
</table>

**Notes.** (a) Flux at rest-frame 1350 Å in units of \( 10^{-14} \text{ erg} \text{ s}^{-1} \text{ cm}^{-2} \text{ Å}^{-1} \). (b) Flux from a power-law fit of the form \( F(\lambda) = F_{1350}(\lambda/1350)^{\alpha} \). (c) Spectral index from the above power-law fit.
Table A.2. UV column-densities for the outflow components in NGC 5548.

<table>
<thead>
<tr>
<th>Ion</th>
<th>$v_{t1}$</th>
<th>$v_{t2}$</th>
<th>$v_{t3}$</th>
<th>$v_{t4}$</th>
<th>$v_{t5}$</th>
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<td>Epoch 2013</td>
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<tr>
<td>H I</td>
<td>$&gt;14.39^{b}$</td>
<td>$&gt;13.86$</td>
<td>$&gt;14.28$</td>
<td>$&gt;14.40$</td>
<td>$&gt;14.00$</td>
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</tr>
<tr>
<td></td>
<td>$&gt;14.52$</td>
<td>$&lt;12.81^{c}$</td>
<td>$13.67^{+0.1}_{-0.1}$</td>
<td>$&lt;13.24$</td>
<td>$&lt;13.40$</td>
<td>$&lt;13.00$</td>
</tr>
<tr>
<td>C II</td>
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<td></td>
<td></td>
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<tr>
<td>C III</td>
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<td></td>
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<td></td>
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</tr>
<tr>
<td>C IV</td>
<td>$&gt;14.60$</td>
<td>$14.00^{+0.3}_{-0.3}$</td>
<td>$&gt;14.30$</td>
<td>$&gt;14.30$</td>
<td>$&gt;14.10$</td>
<td>$&gt;13.70$</td>
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<tr>
<td>N V</td>
<td>$&gt;14.90$</td>
<td>$13.90^{+0.2}_{-0.2}$</td>
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<td>$&gt;14.80$</td>
<td>$&gt;14.60$</td>
<td>$&gt;14.10$</td>
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<tr>
<td>Al II</td>
<td></td>
<td></td>
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Notes. (a) Integration limits in km s$^{-1}$. (b) Lower limit log column-densities given in units of cm$^{-2}$ are shown in blue. (c) Upper limits are likewise shown in red.
Fig. A.1. 2013 spectrum of NGC 5548. The vertical axis is the flux in units of $10^{-14}$ erg s$^{-1}$ cm$^{-2}$ Å$^{-1}$, and the quasar-rest-frame and observer-frame wavelengths are given in Angstroms on the top and bottom of each sub-plot, respectively. Each of the six kinematic components of the outflow shows absorption troughs from several ions. We place a vertical mark at the expected center of each absorption trough (following the velocity template of Si IV and N V) and state the ion, rest-wavelength and component number (C1–C6). We also assign a color to each component number that ranges from blue (C1) to red (C6). Absorption lines from the ISM are likewise marked in black with dashed lines.

Fig. A.1. continued.
Fig. A.1. continued.
Fig. A.1. continued.
Fig. A.1. continued.
Fig. A.1. continued.
Fig. A.1. continued.

Fig. A.1. continued.
Fig. A.1. continued.
Fig. A.2. Spectrum of NGC 5548 during the five epochs of observation. The 2013 spectrum is obtained by co-adding visits 1 through 5. Spectral regions where absorption troughs from five ions are shown in sub-plots a) through e) and the six kinematic components associated with such absorption are labelled $C_1$ through $C_6$.

Fig. A.2. continued.
Fig. A.2. continued.
Fig. A.3. Normalized spectrum of NGC 5548 during the five epochs of observation, plotted in the velocity rest-frame of the quasar (same annotation as Fig. A.2).
Fig. A.3. continued.
Fig. A.3. continued.

Fig. A.3. continued.
Fig. A.3. continued.

Fig. A.3. continued.
Fig. A.3. continued.