

LETTER TO THE EDITOR

Candidate hypervelocity stars of spectral type G and K revisited[★]

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ABSTRACT

Hypervelocity stars (HVS) move so fast that they are not bound to the Galaxy. When they were first discovered in 2005, dynamical ejection from the supermassive black hole (SMBH) in the Galactic centre (GC) was suggested as their origin. The two dozen HVSs known today are young massive B stars, mostly of 3–4 solar masses. Recently, 20 HVS candidates of low mass were discovered in the Segue G and K dwarf sample, but none of them originates from the GC. We embarked on a kinematic analysis of the Segue HVS candidate sample using the full 6D phase space information based on new proper motion measurements. Their orbital properties can then be derived by tracing back their trajectories in different mass models of our Galaxy. We present the results for 14 candidate HVSs for which proper motion measurements were possible. Significantly lower proper motions than found in the previous study were derived. Considering three different Galactic mass models, we found that all stars are bound to the Galaxy. We confirm that the stars do not originate from the GC. The distribution of their proper motions and radial velocities is consistent with predictions for runaway stars ejected from the Galactic disk by the binary supernova mechanism. However, their kinematics are also consistent with old disk membership. Moreover, most stars have rather low metallicities and strong α -element enrichment, as is typical for thick disk and halo stars, whereas the metallicity of the three most metal-rich stars might indicate that they are runaway stars from the thin disk. One star shows halo kinematics.

Key words. stars: kinematics and dynamics – stars: low-mass – stars: late-type – stars: abundances – stars: Population II

1. Introduction

The high space velocities and large distances of hypervelocity stars (HVS) that are not bound to the Galaxy make them important probes for mapping the Galactic dark matter halo. When HVSs were first discovered (Brown et al. 2005; Hirsch et al. 2005; Edelmann et al. 2005), the tidal disruption of a binary by the supermassive black hole (SMBH) in the Galactic centre (GC) was suggested as their origin (Hills 1988). Brown et al. (2014) carried out a systematic survey for B-type stars in the halo and found about two dozen HVSs with intermediate masses in the range of 3 to 4 M_{\odot} . Because such stars are luminous, their survey covered a large volume (out to 100 kpc from the GC). Low-mass stars, on the other hand, can be accelerated more easily and may gain higher ejection velocities (Tauris 2015). Since they are long-lived, they can travel very large distances during their main-sequence lifetime. However, they are less luminous than the B-type HVSs and can only be detected in a smaller volume by flux-limited surveys (<10 kpc), such as the Sloan Extension for Galactic Understanding and Exploration (SEGUE). Moreover, a photometric pre-selection of low-mass main-sequence stars is very difficult because of the overwhelmingly large number of red stars in the halo. Therefore attempts have been made to isolate HVS candidates of low mass from the SEGUE G and K Dwarf Sample (Palladino et al. 2014, hereafter P14), LAMOST (Zhong et al. 2014), and RAVE (Hawkins et al. 2015) surveys using proper motion criteria.

P14 carried out a search for G and K candidate-HVS from the SEGUE sample. High proper motion stars were selected for a detailed analysis of the 6D phase space information, and 20 candidates were found to be probably unbound, four of which at 3σ and six at 2σ significance levels. From calculating the stellar trajectories in the Galactic potential, P14 derived possible places of origin in the Galactic disk. Of the seven stars with the highest probability of being unbound (>98%), none crossed the disk near the GC, but at distances of 5 to 10 kpc away from it. Hence, an origin in the GC was excluded for these stars, which challenges the SMBH slingshot mechanism. Other ejection mechanisms were discussed, including the classical scenario of dynamical interaction in star clusters and the binary supernova scenario (Blaauw 1961). The latter has been revisited by Tauris (2015) to derive the maximum speed of HVS stars ejected from binaries. The simulations indicate that Galactic rest-frame velocities of up to 1280 km s⁻¹ are possible. Such high velocities can explain many, if not all, of the G/K-dwarf HVSs in the SEGUE sample.

As stressed by P14, the HVS nature of the stars strongly depends on the proper motion adopted. Therefore, the data were carefully checked for reliability by simulations. Three stars met all criteria and therefore were characterized as “clean”. The remaining 17 stars were regarded as “reliable”. P14 found that the candidates’ tangential velocities are much higher than their radial velocities, unlike what is expected for an isotropic distribution of stars. The authors therefore cautioned that the high ratio of tangential vs. radial velocity may be characteristic for a sample with large proper motion errors and built a Monte Carlo

[★] Tables 1–3 are available in electronic form at <http://www.aanda.org>

test to estimate the probability of the stars being outliers. All stars show a likelihood of lower than 25%, half of them even lower than 10%. Nevertheless, an independent determination of the proper motions is required.

In this Letter we attempt to determine proper motions for all 20 stars of the sample by combining all astrometric information at hand from digitized plates and modern surveys. We were able to derive proper motions for 14 stars of the sample. Because of the strong implications of the results of P14, we also determined their radial velocities and distances and calculated their trajectories in different Galactic potentials. The confirmation of the candidates as HVSS would demand ejection mechanisms other than the SMBH slingshot to explain their existence.

2. Observations

2.1. Radial velocities and distances

To reanalyse the spectra, we retrieved spectra of all 20 stars from the SDSS data base. In addition, we obtained individual SDSS spectra to search for radial velocity variability so that we could exclude close binaries. No statistically significant radial velocity variations were detected, that is, we found no indication for binarity. The averaged radial velocities were consistent with those given by P14 to within mutual error limits.

P14 made use of the DR9 SEGUE Stellar Parameter Pipeline (SSPP), which provided estimates for effective temperatures T_{eff} , surface gravity $\log g$, iron abundance $[\text{Fe}/\text{H}]$, and α enrichment factor $[\alpha/\text{Fe}]$. They determined the distances of the stars using an isochrone-matching technique. We checked effective temperatures and surface gravities for consistency by comparing *ugriz*-magnitudes from SDSS to synthetic colours from [Castelli \(1999\)](#). Additionally, we compared the synthetic spectra of [Munari et al. \(2005\)](#) to the SDSS spectra. We found no inconsistencies and therefore adopted the atmospheric parameters and distances as given by P14.

2.2. Proper motions

The most problematic information for the kinematic analysis are the proper motions. P14 selected their sample using proper motions provided by the SDSS data base.

To determine the proper motions, we made use of all available images, including SDSS, at different epochs. Early-epoch photographic plates from the Digitised Sky Surveys¹ were combined with those obtained from the data bases of modern digital surveys such as SDSS² and UKIDSS³.

For each HVS candidate, positions were derived from all available images with respect to a set of faint, compact, and well-distributed background galaxies (see [Tillich et al. 2011](#)). This provided a timebase of about 60 years. We were able to derive proper motions for 14 of the 20 P14 candidates. For the low Galactic latitude stars 5, 7, 8, and 10 we failed to identify a sufficiently large set of suitable background galaxies. Stars 1 and 11 suffered from crowding. Therefore we had to dismiss these six stars from further analysis. For the remaining candidates, 16 to 29 background galaxies per field were found to be suitable and distributed favourably around the target. The proper motion components were obtained from the positions of all epochs by

¹ http://archive.stsci.edu/cgi-bin/dss_plate_finder

² <http://skyserver.sdss3.org/public/en/tools/chart/navi.aspx>

³ <http://www.ukidss.org/>

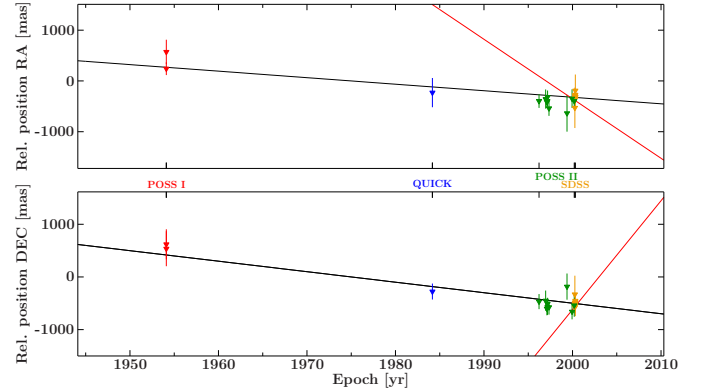


Fig. 1. Proper motion fit through the relative positions of star 4. Red lines indicate the proper motion used by P14.

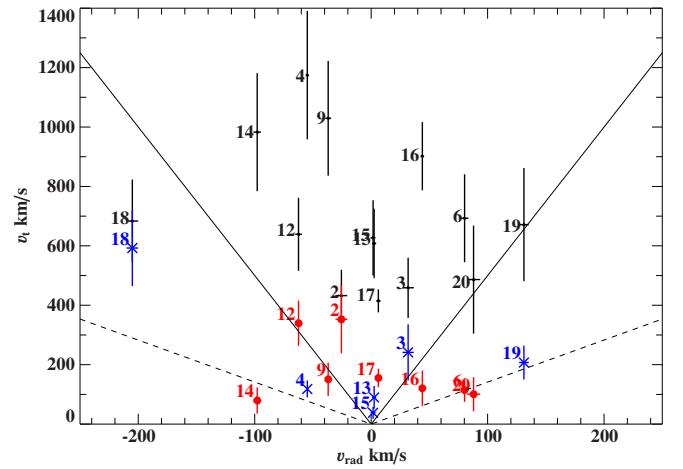


Fig. 2. Comparison of the tangential and radial velocity derived in this work (coloured with error bars, red dots indicate stars with $[\text{Fe}/\text{H}] > -0.7$, blue crosses $[\text{Fe}/\text{H}] < -0.7$ and P14 (black)). Error bars of P14 were derived in the same way as in this work. The dashed lines indicate $v_t = \sqrt{2}v_r$, while the solid lines indicate $v_t = 5v_r$.

linear regression (see Fig. 1). The values only agree with the SDSS measurements for objects 2 and 18 (see Table 1).

The tangential velocities are plotted in Fig. 2 vs. the radial velocities and compared to the results of P14. Except for stars 2 and 18, the values derived here are significantly lower than those given by P14. The ratios of the tangential vs. radial velocity of the stars are much lower than previously found (>5). According to P14, an isotropic distribution of stars would be described by a $\sqrt{2}$ times higher tangential than radial velocity. The new results come close to that expectation.

3. Kinematics

To obtain the dynamical properties of the stars, the orbits were traced back to the Galactic disk. Three different Milky Way mass models, see models I, II, and III in [Irrgang et al. \(2013\)](#), were used for this exercise to estimate possible systematic influences of the applied gravitational potential. Because the escape velocities of model II of [Irrgang et al. \(2013\)](#) are very similar to those calculated by P14, we used this model for the comparison with the results of P14.

By varying the position and the velocity components within their respective errors by applying a Monte Carlo procedure with a depth of 10^6 , we determined the intersection area of

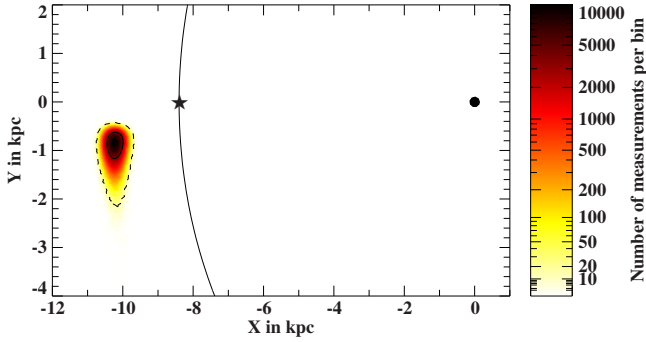


Fig. 3. Disk passages (binned and colour coded) of star 17 obtained by a Monte Carlo simulation with 1 (solid) and 3 (dashed) σ contours. The solid line indicates the solar orbit. The star marks the position of the Sun, the black dot that of the GC.

the trajectories with the Galactic plane. From these simulations we also derived the median Galactic rest-frame velocities at the present locations and their distributions. These values were compared with the respective local escape velocity. Because we assumed a lower tangential velocity, the stars show a much lower probability of being not bound to the Galaxy in all three applied gravitational potentials. The bound probability is defined as the number of orbits not exceeding the local escape velocity with respect to the number of all calculated orbits. In Table 1 the bound probabilities listed by P14 are compared to the bound probabilities obtained by applying model II of Irrgang et al. (2013) using proper motions of SDSS, as used by P14, and of this work. The bound probabilities calculated using SDSS proper motions are similar to those obtained by P14, as expected. For the Galactic models I (see Table 2) and III (see Table 3), the differences between P14 and our calculations using SDSS proper motions can be explained by the higher Galaxy masses of those models.

Using our revised proper motions, none of the 14 candidates is unbound, irrespective of the Galactic potential model used. In fact, the probability to be unbound is lower than 0.1% for all but three candidates. Star 18 has the highest probability of being unbound, but even in the lightest Galaxy model II, this probability does not exceed 36%. We also calculated the distances to the GC at the time of the disk passage and the shortest distances in all three potentials (Tables 1–3). They do not depend on the choice of the potential except for stars 12, 16, and particularly star 2 (the latter changes by almost a factor of 2), for which the location of disk crossing must be regarded as highly uncertain. The GC is excluded at 3σ level for all stars irrespective of the choice of potential. Among the programme stars the present position of star 17 is closest to the Galactic plane and, hence, its disk-crossing location outside of the solar circle can be well constrained (see Fig. 3).

4. Runaway, disk, or halo stars?

Since all 14 candidates are bound and none comes from anywhere near the GC, these stars have to be dismissed as HVS stars. P14 discussed whether they might be disk runaways or genuine halo stars. In Fig. 4 the Galactic radial velocity component U is plotted against the rotational velocity component V and compared with contour lines denoting the limits for the thin and thick disk by Pauli et al. (2006). However, kinematics cannot easily distinguish between thick disk and runaway origins. Because most runaways are launched at modest velocities from the rotating disk, their kinematic and spatial

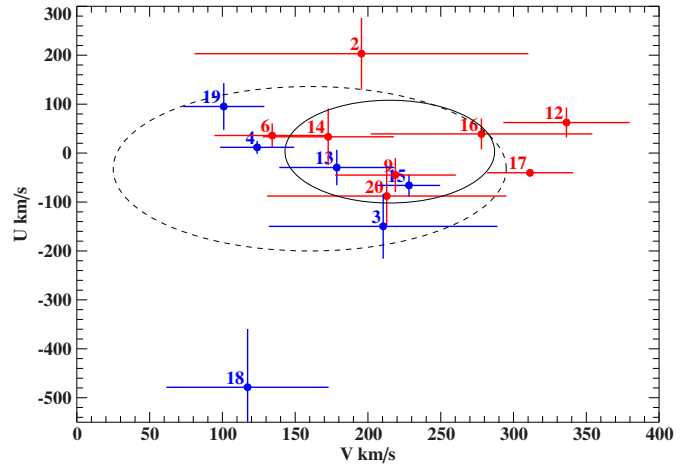


Fig. 4. $U - V$ -diagram, numbered dots with error bars: stars examined in this work (red indicates stars with $[\text{Fe}/\text{H}] > -0.7$, blue $[\text{Fe}/\text{H}] < -0.7$), dashed line: 3σ contour of the thick disk, solid line: 3σ contour of the thin disk.

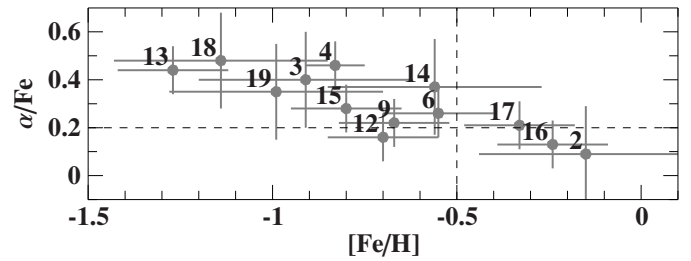


Fig. 5. Position of the programme stars in the $([\text{Fe}/\text{H}], [\alpha/\text{Fe}])$ diagram. The dashed lines separate thick-disk (top left) from thin-disk stars (bottom right). Error bars were estimated from the signal-to-noise ratio of the spectra using the prescriptions of Allende Prieto et al. (2008) and Lee et al. (2011).

distributions will naturally look quite similar to those of the thick disk (Bromley et al. 2009). Population membership can also be assigned by chemical tagging, using $[\text{Fe}/\text{H}]$ and $[\alpha/\text{Fe}]$ as indicators. According to Fuhrmann (2011), thin-disk stars have $[\text{Fe}/\text{H}] > -0.5$ and $[\alpha/\text{Fe}] < 0.2$, whereas thick-disk stars have $[\text{Fe}/\text{H}] < -0.3$ and $[\alpha/\text{Fe}] > 0.2$ (see Fig. 15 of Fuhrmann 2011). However, some halo stars may have similar chemical abundances as thick-disk stars. In Fig. 5 we plot $[\alpha/\text{Fe}]$ vs. $[\text{Fe}/\text{H}]$ for the 14 stars and conclude that most stars have chemical characteristics of the thick disk, whereas stars 2, 16, and 17 are more likely to be thin-disk stars. The chemical composition of star 18 strengthens the conclusion that it is a halo star.

Dynamical ejection in dense star clusters or binary supernova ejection may result in disk runaway stars. Kenyon et al. (2014) simulated the Galactic populations of (i) HVS stars ejected from the GC via the SMBH slingshot mechanism; (ii) runaway stars from the disk ejected at distances of 3 to 30 kpc from the GC by either the binary supernova mechanism or dynamical ejection from clusters. They compared their predictions to the observed properties of B-type HVS (Brown et al. 2014) and B-type runaway stars (Silva & Napiwotzki 2011) as well as to the P14 sample of G- and K-type HVS candidates. Kenyon et al. (2014) concluded that the runaway sample of B-type stars can be matched by their simulation for binary supernova ejection or dynamical ejection, but not by the SMBH slingshot. The HVS sample of B-type stars (Brown et al. 2014), on the other hand,

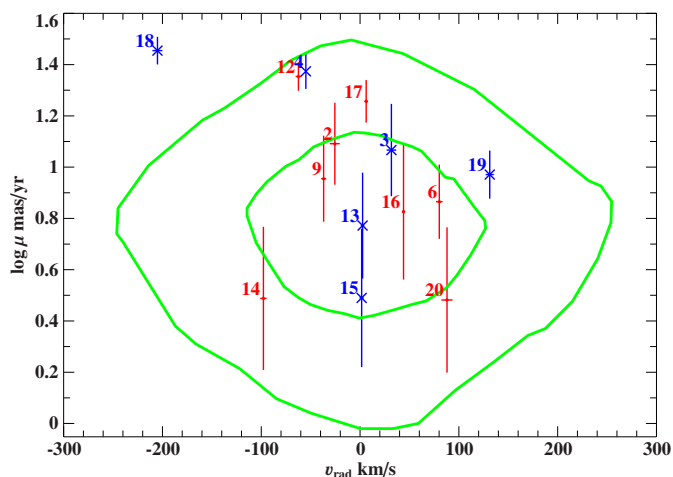


Fig. 6. Comparison of the HVS candidates from P14 with proper motions derived by us (red indicates stars with $[\text{Fe}/\text{H}] > -0.7$, blue crosses $[\text{Fe}/\text{H}] < -0.7$) in the proper motion vs. radial velocity plane with predictions of simulations for a sample of $1 M_{\odot}$ runaways (green curves) generated by the binary supernova mechanism (Kenyon et al. 2014). The inner contour includes 50% of the stars in the distance-limited sample, while the outer one includes 90% (adapted from Fig. 25 of Kenyon et al. 2014).

is well reproduced by the prediction from the SMBH slingshot, but the data cannot be explained by runaway models.

However, SMBH slingshot models fail to account for the observations of the P14 SEGUE HVS candidates of spectral type G and K. In Fig. 6 we compare the radial velocities and our new proper motions of the 14 P14 stars analysed here to the predictions from simulations for $1 M_{\odot}$ runaway stars generated by the binary supernova mechanism. The distribution of stars match the prediction of Kenyon et al. (2014) very well. However, the stellar chemical composition is inconsistent with the disk runaway scenario, except possibly for the three most metal-rich ones. Such stars could be ejected as the surviving donor of an exploding white dwarf in a SN Ia event, as is discussed for Tycho's supernova (Ruiz-Lapuente et al. 2004).

5. Summary and conclusions

We presented a revised analysis of the 20 candidate HVSS introduced by P14. Our work was motivated by the warning of P14 that the proper motions they extracted from the SDSS data base were abnormally high and that therefore the stars have to be considered HVS candidates only until their proper motions are confirmed. Therefore, we independently determined proper

motions from all astrometric images using grids of distant galaxies as frames of reference for 14 of these candidates. We repeated the analysis of P14 and confirmed the radial velocities, atmospheric parameters, and distances. However, except for two stars, all measured proper motions are significantly lower than those used by P14. We carried out a kinematic analysis of the stars in three different Milky Way mass models and confirmed that no star originated from the GC. All stars are bound to the Galaxy and therefore have to be dismissed as HVS. We considered the possibilities that the stars are either old stars of the halo or thick-disk stars, as suggested by their metallicities and α element enhancements from the (U, V) diagram or disk-runaway stars. The stellar kinematics as well as their chemical composition indicate that they mostly belong to the thick-disk population; accordingly, star 18 might belong to the halo. However, the disk-runaway option would also be consistent with model predictions by Kenyon et al. (2014) for all stars (except for star 18). However, this is at variance with the stellar chemical composition except for the three most metal-rich stars, which might be runaway stars from the thin-disk population.

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Table 1. Kinematic properties of the programme stars.

P14 No.	This work		SDSS		$v_{\text{GRF},1}$ (km s ⁻¹)	$v_{\text{GRF},2}$ (km s ⁻¹)	z (kpc)	r (kpc)	r_{min} (kpc)	Bound probabilities (%)		
	$\mu_{\alpha} \cos \delta$ (mas/yr)	μ_{δ} (mas/yr)	$\mu_{\alpha} \cos \delta$ (mas/yr)	μ_{δ} (mas/yr)						P14 PM: SDSS	model II	this work
2	8.9 ± 4.7	8.6 ± 3.7	-2.6 ± 3.0	15.5 ± 3.0	423.9 ± 100.7	635.2 ± 86.7	-2.9 ± 0.2	60.1	3.4	7.43	10.7	84.7
3	5.5 ± 4.0	-10.4 ± 4.7	23.6 ± 3.0	-1.6 ± 3.0	303.7 ± 75.9	645.9 ± 96.7	2.5 ± 0.5	8.7	1.5	34.88	33.3	99.9
4	-12.9 ± 2.4	-20.0 ± 4.1	-117.2 ± 5.9	206.8 ± 5.9	126.5 ± 23.2	1304.9 ± 213.5	0.8 ± 0.1	6.7	3.9	0.00	0.0	100
6	5.8 ± 2.2	-4.5 ± 2.5	-41.6 ± 5.5	20.1 ± 5.5	145.4 ± 38.3	920.9 ± 148.2	-2.1 ± 0.4	6.8	0.8	0.07	0.3	100
9	-3.1 ± 3.1	-8.5 ± 3.2	4.8 ± 2.9	-65.3 ± 2.9	254.1 ± 42.7	937.9 ± 187.3	-1.3 ± 0.2	11.1	7.5	1.20	1.9	100
12	7.3 ± 2.6	21.4 ± 2.9	19.4 ± 2.8	38.1 ± 2.8	482.4 ± 69.9	735.7 ± 117.4	-1.2 ± 0.2	27.7	16.2	3.77	5.5	81.1
13	-2.0 ± 2.5	-5.6 ± 2.7	39.6 ± 3.0	-19.1 ± 3.0	185.7 ± 35.8	716.7 ± 110.0	2.8 ± 0.5	7.6	3.6	4.42	8.9	100
14	-0.9 ± 2.9	-1.3 ± 2.4	-5.7 ± 3.9	-44.5 ± 3.9	200.2 ± 50.9	840.6 ± 194.6	-3.3 ± 0.6	7.4	1.0	5.86	7.4	100
15	1.1 ± 2.2	-2.2 ± 2.3	-58.6 ± 5.4	8.1 ± 5.4	239.5 ± 18.1	656.7 ± 117.5	1.5 ± 0.3	11.2	7.1	15.98	21.7	100
16	-5.2 ± 4.6	1.1 ± 5.1	0.8 ± 5.7	-58.1 ± 5.7	293.3 ± 74.6	676.8 ± 113.2	1.1 ± 0.1	19.3	2.4	19.70	12.0	99.9
17	12.5 ± 3.2	13.0 ± 3.5	25.2 ± 2.5	42.1 ± 2.5	346.8 ± 30.9	603.0 ± 37.3	0.4 ± 0.1	10.3	9.8	20.01	10.9	100
18	1.8 ± 2.3	-28.6 ± 3.5	7.2 ± 3.1	-32.2 ± 3.1	537.4 ± 123.3	651.1 ± 134.8	2.7 ± 0.5	3.6	0.5	21.30	29.8	63.6
19	-9.0 ± 2.0	-2.8 ± 1.7	-28.8 ± 6.9	8.7 ± 6.9	150.2 ± 41.6	644.5 ± 188.0	2.9 ± 0.5	3.6	0.9	23.69	34.6	100
20	-0.4 ± 2.4	-1.0 ± 3.1	15.4 ± 5.8	-5.8 ± 5.8	255.7 ± 69.0	630.2 ± 178.2	4.6 ± 0.8	7.8	1.9	43.24	43.2	100

Notes. Comparison of our measured proper motions and those from the SDSS database as used by P14 (Cols. 2–5). Galactic rest-frame velocities in Cols. 6 and 7 are based on our proper motions $v_{\text{GRF},1}$ or SDSS proper motions $v_{\text{GRF},2}$. The present height above the Galactic plane is denoted with z . The quantities r and r_{min} give the average and shortest distance to the GC at the 3σ level of the disk passage using new proper motions and model II of Irrgang et al. (2010). The last three columns give the bound probabilities as listed by P14 (Col. 11), our results based on model II, and the proper motions from SDSS as given in Cols. 4 and 5 (Col. 12) and based on our revised proper motions from Cols. 2 and 3 in the last column.

Table 2. Kinematic properties of the programme stars using model I.

P14 No.	r (kpc)	r_{min} (kpc)	$v_{\text{GRF},1}$ (km s ⁻¹)	$v_{\text{GRF},2}$ (km s ⁻¹)	Bound probabilities (%)		
					P14 PM: SDSS	model I	this work
2	50.3	3.8	424.5 ± 100.5	636.4 ± 86.7	7.43	23.4	92.3
3	8.5	1.3	304.6 ± 75.8	647.2 ± 96.7	34.88	49.3	100
4	6.7	3.9	128.1 ± 23.2	1305.8 ± 213.4	0.00	0.0	100
6	6.8	0.9	146.7 ± 38.4	922.6 ± 147.9	0.07	0.8	100
9	11.1	7.5	255.4 ± 42.6	937.5 ± 187.1	1.20	3.3	100
12	24.2	16.0	483.5 ± 69.9	736.4 ± 117.4	3.77	11.0	91.7
13	7.6	3.6	187.2 ± 35.8	717.7 ± 109.8	4.42	16.9	100
14	7.4	1.0	201.5 ± 50.9	839.7 ± 194.4	5.86	11.3	100
15	11.1	7.1	241.1 ± 18.1	657.0 ± 117.3	15.98	33.8	100
16	18.6	2.5	294.5 ± 74.4	675.9 ± 113.3	19.70	22.2	100
17	10.3	9.8	348.2 ± 30.9	604.4 ± 37.3	20.01	44.4	100
18	3.5	0.5	537.1 ± 123.5	651.1 ± 134.9	21.30	41.4	74.6
19	3.7	1.0	150.7 ± 41.3	644.6 ± 187.9	23.69	43.4	100
20	7.9	2.0	257.0 ± 69.0	631.7 ± 178.0	43.24	52.0	100

Notes. The quantities r and r_{min} give the average and shortest distance to the GC at the 3σ level of the disk passage using new proper motions and model I of Irrgang et al. (2010). Galactic rest-frame velocities v_{GRF} based on our proper motions (1) or SDSS proper motions (2). The last three columns give the bound probabilities as listed by P14 (Col. 6), our results based on model I, and the proper motions from SDSS as given in Cols. 4 and 5 (Table 1) (Col. 7) and based on our revised proper motions from Cols. 2 and 3 (Table 1) in the last column.

Table 3. Kinematic properties of the programme stars using model III.

P14 No.	r (kpc)	r_{min} (kpc)	$v_{\text{GRF},1}$ (km s ⁻¹)	$v_{\text{GRF},2}$ (km s ⁻¹)	Bound probabilities (%)		
					P14 PM: SDSS	model III	this work
2	34.1	4.0	423.3 ± 100.5	634.4 ± 86.7	7.43	94.4	99.9
3	7.9	1.1	303.2 ± 75.8	645.3 ± 96.9	34.88	96.3	100
4	6.7	3.9	126.5 ± 23.1	1304.1 ± 213.8	0.00	1.0	100
6	6.8	0.9	144.7 ± 38.4	920.5 ± 147.8	0.07	19.8	100
9	11.1	7.5	253.5 ± 42.7	937.8 ± 187.4	1.20	22.6	100
12	18.9	14.2	481.8 ± 69.9	735.1 ± 117.6	3.77	68.4	100
13	7.6	3.7	185.1 ± 35.8	716.6 ± 110.0	4.42	79.4	100
14	7.3	1.1	199.6 ± 51.0	841.0 ± 194.5	5.86	43.3	100
15	11.2	7.2	239.5 ± 18.1	656.5 ± 117.5	15.98	88.3	100
16	16.3	2.5	292.5 ± 74.5	677.7 ± 113.3	19.70	83.5	100
17	10.3	9.8	346.2 ± 30.9	602.5 ± 37.3	20.01	100.0	100
18	3.6	0.5	537.4 ± 123.3	651.1 ± 134.8	21.30	87.7	97.9
19	3.7	1.0	149.8 ± 41.7	644.1 ± 187.9	23.69	80.0	100
20	8.0	2.1	255.1 ± 68.8	630.2 ± 178.4	43.24	85.4	100

Notes. The quantities r and r_{min} give the average and shortest distance to the GC at the 3σ level of the disk passage using new proper motions and model III of Irrgang et al. (2010). Galactic rest-frame velocities v_{GRF} are based on our proper motions (1) or SDSS proper motions (2). The last three columns give the bound probabilities as listed by P14 (Col. 6), our results based on model III, and the proper motions from SDSS as given in Cols. 4 and 5 (Table 1) (Col. 7) and based on our revised proper motions from Cols. 2 and 3 (Table 1) in the last column.