Using SDO’s AIA to investigate energy transport from a flare’s energy release site to the chromosphere*

J. W. Brosius1 and G. D. Holman2

1 The Catholic University of America at NASA’s Goddard Space Flight Center, Code 671, Greenbelt, MD 20771, USA
e-mail: Jeffrey.W.Brosius@nasa.gov
2 Code 671, NASA’s Goddard Space Flight Center, Greenbelt, MD 20771, USA
e-mail: Gordon.D.Holman@nasa.gov

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ABSTRACT

Context. Coordinated observations of a GOES B4.8 microflare with SDO’s Atmospheric Imaging Assembly (AIA) and the Ramaty High Energy Solar Spectroscopic Imager (RHESSI) on 2010 July 31 show that emission in all seven of AIA’s EUV channels brightened simultaneously nearly 6 min before RHESSI or GOES detected emission from plasma at temperatures around 10 MK.

Aims. To help interpret these and AIA flare observations in general, we characterized the expected temporal responses of AIA’s 94, 131, 171, 193, 211, and 335 Å channels to solar flare brightenings by combining (1) AIA’s nominal temperature response functions available through SSWIDL with (2) EUV spectral line data observed in a flare loop footprint on 2001 April 24 with the Coronal Diagnostic Spectrometer (CDS) on timescales comparable to AIA’s image cadence.

Methods. The nine emission lines observed by CDS cover a wide range of formation temperature from about 0.05 to 8 MK. Line brightenings observed early during the CDS flare occurred at temperatures less than about 0.7 MK, with the largest values around 0.1 MK. These brightenings were consistent with the flare’s energy transport being dominated by nonthermal particle beams. Because all of AIA’s EUV channels are sensitive to emission from plasma in the 0.1 to 0.7 MK temperature range, we show that all of AIA’s EUV channels will brighten simultaneously during flares like this, in which energy transport is dominated by nonthermal particle beams.

Results. The 2010 July 31 flare observed by AIA and RHESSI displays this behavior, so we conclude that such beams likely dominated the flare’s energy transport early during the event. When thermal conduction from a reconnection-heated, hot (~10 MK) plasma dominates the energy transport, the AIA channels that are sensitive to emission from such temperatures (particularly the 94 and 131 Å channels) will brighten earlier than the channels that are not sensitive to such temperatures (171 and 211 Å).

Conclusions. Thus, based on the differences expected between AIA’s response to flares whose energy transport is dominated by nonthermal particle beams from those whose energy transport is dominated by thermal conduction, AIA can be used to determine the dominant energy transport mechanism for any given event.

Key words. Sun: activity – Sun: corona – Sun: flares – Sun: transition region – Sun: UV radiation – Sun: X-rays, gamma rays

1. Introduction

In the widely accepted “standard” solar flare model, pre-flare energy is stored in non-potential magnetic fields. This energy is released by magnetic reconnection, during which nonthermal energy is stored in non-potential magnetic fields. This energy release site to the chromosphere either by nonthermal beam particles or by thermal conduction from the directly-heated coronal source. When the flux of energy into the chromosphere exceeds what can be shed by radiative losses, the chromosphere heats and expands. Because the density in the overlying corona is less than that in the chromosphere and photosphere, expansion of the heated chromospheric material produces larger upward than downward velocities. This process, now known as “chromospheric evaporation” (see Antonucci et al. 1999 and Bornmann 1999 for reviews) was first described by Neupert (1968). Upward velocities around 100 km s⁻¹ or more in emission lines formed at temperatures T ∼ 10 MK during flares have long been interpreted as evidence for this phenomenon (Antonucci et al. 1982; Antonucci & Dennis 1983; Zarro et al. 1988; Fludra et al. 1989; Canfield et al. 1990; Doschek 1990; Mariska et al. 1993; Mariska 1994; Silva et al. 1997; Brosius 2003, 2009; Brosius & Phillips 2004; Milligan & Dennis 2009). The evaporated material fills magnetic loops with hot, dense, plasma that eventually cools and fades as the flare decays.

Investigations of solar flares based on rapid cadence (7.2 to 9.8 s) stare spectra obtained with the Coronal Diagnostic Spectrometer (CDS; Harrison et al. 1995) aboard the Solar and Heliospheric Observatory (SOHO) spacecraft showed that intensities of emission lines formed in plasma at transition region temperatures (e.g., the O V line at 629.7 Å, formed at T ≈ 0.25 MK) began to increase several minutes earlier than intensities of emission lines formed in plasma at flare temperatures (e.g., the Fe XIX line at 592.2 Å, T = 8 MK; Brosius & Phillips 2004; Brosius 2009; Brosius & Holman 2009, 2010). Statistically significant relative Doppler velocities also typically accompany the intensity increases in the emission lines. This time difference is consistent with the chromosphere being heated predominantly by beams of nonthermal particles; if

* Movie is available in electronic form at http://www.aanda.org
the chromosphere were heated predominantly by thermal conduction from a directly-heated $T \sim 10$ MK source, that source would have appeared first, before the chromosphere had a chance to respond to thermal conduction from it. When possible, CDS stare spectra were combined with X-ray observations from the Ramatyn High Energy Solar Spectroscopic Imager (RHESSI; Lin et al. 2002; Hurford et al. 2002; Smith et al. 2002). RHESSI combines full disk solar imaging, wide energy coverage (3 to 17000 keV), and high sensitivity for investigating both nonthermal and high temperature ($\gtrsim 10$ MK) thermal emission from flares and microflares.

A shortage of the rapid cadence CDS stare studies is their comparatively limited field of view. In previous investigations, this was partly compensated by acquiring full disk 195 Å images at 12 min cadence from the Extreme-ultraviolet Imaging Telescope (EIT; Delaboudinière et al. 1995) aboard SOHO. Nevertheless, there remained some uncertainty as to what was going on outside the CDS field of view. This spatial uncertainty can be eliminated by using the Atmospheric Imaging Assembly (AIA; Lemen et al. 2011) aboard the Solar Dynamics Observatory (SDO) satellite; since AIA obtains snapshots of the entire Sun in each of its EUV channels every $12\, s$, the start of any given flare or microflare is never hidden outside AIA’s field of view, and can always be determined for each channel within the instrumental cadence of $12\, s$.

Launched on 2010 February 11, SDO observes the Sun nearly continuously in an inclined geosynchronous orbit from which its extremely high data rate is enabled through contact with a single dedicated ground station. AIA contains seven EUV channels roughly centered on the following wavelengths, designed to include solar radiation in and around strong emission lines from the following ions: 304 Å (includes emission from He II, formed at $T \approx 0.05$ MK), 171 Å (Fe IX at $T \approx 0.63$ MK), 211 Å (Fe XV at $T \approx 1.9$ MK), 335 Å (Fe XVI at $T \approx 2.7$ MK), 94 Å (Fe XVIII at $T \approx 6.3$ MK), 193 Å (Fe XII at $T \approx 1.4$ MK and Fe XXIV at $T \approx 19$ MK), and 131 Å (Fe VIII at $T \approx 0.37$ MK, Fe XX at $T \approx 9.1$ MK, Fe XXI at $T \approx 10$ MK, and Fe XXIII at $T \approx 14$ MK). Wavelength response functions for each channel were measured in the laboratory prior to SDO’s launch (Boerner et al. 2011), and are available through SSWIDL. Temperature response functions for each channel were derived (Boerner et al. 2011) by combining the measured wavelength response functions with emission line data from the CHIANTI database (Dere et al. 1997, 2009), and are also available through SSWIDL. Based on AIA’s powerful combination of continuous, rapid cadence (12 s), full disk EUV imaging and wide range of temperature sensitivity, its observations are expected to provide new information on solar activity including the temperature dependence of flare start times. The use of timing differences among various EUV channels to assess the roles of thermal conduction and nonthermal particle beams in transporting the energy that drives chromospheric evaporation is one of AIA’s objectives (concept study report; see http://aia.lmsal.com/public/CSR.htm).

Here we present coordinated AIA and RHESSI observations of a GOES B4.8 microflare on 2010 July 31. AIA light curves averaged over selected areas within the active region are expected to reveal whether timing differences exist among the various channels and, hence, the importance of thermal conduction relative to nonthermal particles in transporting the energy that drives chromospheric evaporation. To help with the interpretation of these and additional AIA flare observations, we characterize the expected response of AIA’s EUV channels to the onset of solar flare brightenings by combining (1) AIA’s temperature response functions (Boerner et al. 2011) with (2) existing EUV spectral line flare data obtained with CDS on timescales comparable to AIA’s image cadence (Brosius & Phillips 2004). In Sect. 2 we describe the 2010 July 31 flare observational results, in Sect. 3 we present the expected temporal response of AIA’s EUV channels based on CDS flare spectra obtained on 2001 April 24, in Sect. 4 we discuss implications of this work, and in Sect. 5 we summarize our findings.

2. Observation and analysis of microflare on 2010 July 31

We observed NOAA Active Region 11092 on 2010 July 31 with AIA from 05:00 to 05:45 UT and with RHESSI from 05:00 to 05:40 UT, during which the GOES satellite observed a B4.8 microflare from 05:17 to 05:36 UT. We used AIA’s online “cutout service” (which provides Level 1.5 data; see Lemen et al. 2011) to extract high spatial resolution ($0.6\, \text{pixels}$), rapid cadence (12 s), co-aligned $2.5\times 2.5\, \text{images}$ centered on ($-6463\, \text{,} +113\, \text{s}$) from all seven of AIA’s EUV channels as well as its 1600 Å channel (at $24\, \text{s}$ cadence).

Figure 1 displays a sample of $2.5\times 2.5\, \text{images}$ of AR 11092 in four different channels observed by AIA on 2010 July 31. The source areas A and B in which two successive microflare bursts originate are separated from each other by $30\, \text{r}$ or more but both connected to the sunspot toward the center of the field of view. Movies in all eight of the AIA channels (one movie showing two channels is available online) indicate that the microflare started within the area outlined by the lower dashed rectangle (A), and occurred later within the area outlined by the upper dashed rectangle (B). It is possible that the B4.8 microflare observed by GOES was actually a pair of closely spaced, closely timed sympathetic microflares.

Figure 2 displays a pair of RHESSI $3-25\, \text{keV}$ images that show the two spatially separated sources of brightenings that are seen in the AIA images of the microflare. The time interval for each image is indicated in the figure captions. We integrated over 1-minute intervals because the count rates were low. These images were obtained using the CLEAN technique and RHESSI front detectors 1–9, giving a spatial resolution $\sim 3\, \text{r}$.

Figures 3 and 4 both show $3-12$ and $12-25\, \text{keV}$ RHESSI light curves. Unfortunately, a particle event in Earth’s magnetosphere contaminated these with gradual humps that peak toward the center of RHESSI’s observing window, so while the resulting increased background is not expected to significantly alter the start times or peak times of the flare bursts, it effectively hides any early gradual increase to which RHESSI may have been sensitive. Note the mound in the $12-25\, \text{keV}$ emission atop which bursts starting at 05:17 and 05:27 UT are seen in Figs. 3 and 4. GOES, however, is unaffected by this, and shows a flat background until the start of its first recorded burst at 05:17:32 UT. Figure 3 displays light curves for all seven of AIA’s EUV wavelength channels, plus that of 1600 Â, averaged over area A in Fig. 1 to ensure that the microflare’s starting position is included in the light curves. Solid vertical lines indicate the start of an increase in EUV emission observed by AIA no later than 05:11:34 UT, a peak in the early (precursor?) brightening at 05:13:30 UT, the start of the first impulsive rise of AIA emission at 05:16:44 UT, and the start of the first burst recorded by GOES at 05:17:32 UT. The first impulsive peak recorded by AIA occurred at 05:18:30 UT (simultaneously for all wavebands, within the uncertainty of the cadence), and the rise toward a secondary
peak began at 05:19:34 UT. Figure 4 displays light curves for the same AIA wavelength channels averaged over area B in Fig. 1. The GOES and RHESSI light curves from the start of the microflare component observed in area B are not as pristine as those from the start of the microflare component observed in area A because they are superimposed on declining emission from A, but they show that the rise in the AIA light curves at that location begins before the start of the second burst in the GOES or RHESSI light curves. Solid vertical lines indicate the start of an increase in EUV emission no later than 05:23:45 UT in B, the start of a rapid rise in EUV emission no later than 05:26:15 UT, and the start of the second burst recorded by GOES at 05:27:08 UT.

Eventually we see different temporal behaviors in the different channels. For example, post-impulsive maxima occur at progressively later times in channels that were designed to observe bright lines of successively cooler ions, consistent with expectations of a cooling plasma after the cessation of energy deposition. During this post-impulsive cooling phase, the AIA channel that observes Fe XXI flare emission (131 Å) peaks first (a somewhat broad peak centered around 05:23 UT), followed by 94 Å (05:25 UT), and so on down to 171 Å (05:36 UT). The fact that the 193 Å channel does not show a similar post-impulsive peak prior to that of the 131 Å emission indicates that the maximum flare temperature remained smaller than that required to produce Fe XXIV.

3. Expected temporal response of AIA based on CDS flare spectra from 2001 April 24

To quantify the rapid cadence (9.8 s) temporal behavior of the solar flare atmosphere over a wide range of temperature
Fig. 3. Light curves (log of the count rate for AIA, arbitrary units for GOES, log of the count rate for RHESSI, but shifted downward by 2.2 for the 3–12 keV light curve, and 1.65 for the 12–25 keV light curve) from eight wavelength channels on AIA averaged over area A (see Fig. 1) in which the brightening first occurred during the B4.8 microflare on July 31. Solid vertical lines indicate the start of an increase in EUV emission observed by AIA no later than 05:11:34 UT, a peak in the early (precursor?) brightening at 05:13:30 UT, the start of an impulsive rise in AIA emission at 05:16:44 UT, the start of the first burst recorded by GOES at 05:17:32 UT, the first impulsive peak detected by AIA at 05:18:30 UT (simultaneously for all wavebands, within the uncertainty of the cadence), the rise toward a secondary peak at 05:19:34 UT, and the start of the second burst recorded by GOES at 05:27:08 UT.

Fig. 4. Light curves (log of the count rate for AIA, arbitrary units for GOES, log of the count rate for RHESSI, but shifted downward by 2.2 for the 3–12 keV light curve, and 1.65 for the 12–25 keV light curve) from eight wavelength channels on AIA averaged over area B (see Fig. 1) in which the brightening started later during the B4.8 microflare on July 31. Solid vertical lines indicate the start of an increase in EUV emission observed by AIA no later than 05:11:34 UT, a peak in the early (precursor?) brightening at 05:13:30 UT, the start of an impulsive rise in AIA emission at 05:16:44 UT, the start of the first burst recorded by GOES at 05:17:32 UT, the first impulsive peak detected by AIA at 05:18:30 UT (simultaneously for all wavebands, within the uncertainty of the cadence), the rise toward a secondary peak at 05:19:34 UT, and the start of the second burst recorded by GOES at 05:27:08 UT.

Fig. 5. Light curves of O V 629.7 Å, Al XI 550.0 Å, and Fe XIX 592.2 Å obtained with SOHO’s CDS in rapid cadence (9.8 s) stare mode during the 2001 April 24 GOES M2.3 flare reported by Brosius & Phillips (2004). We use these and light curves from six additional spectral lines to measure average preflare intensities and flare intensity enhancement factors as described in the text. Solid horizontal lines indicate ±1σ from the average preflare Fe XIX noise. The long vertical tick mark at 18:02:00 UT indicates the time at which the Fe XIX intensity exceeds and remains continuously greater than its pre-flare noise level by more than 1σ, and the long vertical tick mark at 18:07:40 UT indicates the time at which Yohkoh hard X-ray count rates became statistically significant (with a peak at 18:10:40 UT).

(4.67 \leq \log T \leq 6.89), we use the integrated intensities (ergs cm\(^{-2}\) s\(^{-1}\) sr\(^{-1}\)) derived by Brosius & Phillips (2004) for the following nine emission lines measured at the commencement site (an area in and around one of the flare loop’s footpoints) of a GOES M2.3 flare in Active Region 9433 on 2001 April 24: He II 303.78 Å (log T = 4.67), O III 599.60 Å (log T = 4.96), O IV 554 Å (a self-blend of four lines at 553.33, 554.08, 554.51, and 555.26 Å, formed at log T = 5.24), O V 629.76 Å (log T = 5.37), Ne VI 558.59 Å (log T = 5.63), Ca X 557.76 Å (log T = 5.87), Mg X 609.79 Å (log T = 6.05, from which its O IV blend was removed based on the density-insensitive intensity ratio O IV 609.83(553.33+554.08+554.51+555.26) = 0.1645 ± 0.0022 derived with CHIANTI v. 6.01; Dere et al. 1997, 2009), Al XI 550.04 Å (log T = 6.17), and Fe XIX 592.23 Å (log T = 6.89). See Fig. 5. Using the 51 exposures whose start times occurred between 17:36:05 and 17:44:18 UT, we derived pre-flare average intensities and their corresponding 1-sigma scatters. All of the lines except that of Fe XIX at 592.2 Å were observed during this pre-flare interval; for Fe XIX the pre-flare average intensity is defined by the noise. We consider the detection of Fe XIX emission to be reliable only when the intensity derived from fits to the line profile exceeds and continuously remains greater than the pre-flare average noise level by more than 1σ (which occurs at 18:02:00 UT; see Fig. 5). Upward velocities approaching –40 km s\(^{-1}\) in those same lines during the impulsive phase, when Fe XIX showed maximum upward velocities of ~64 km s\(^{-1}\), indicated explosive evaporation.

For each exposure beginning with the start of the first precursor, we calculate an intensity enhancement factor for each emission line as the ratio of the line’s intensity in the given exposure to its pre-flare average value. Since the Fe XIX line emission is considered to be noise until its intensity continuously exceeds the pre-flare average noise by more than 1σ, we define its enhancement factor as 1 until that time. Thus for each 9.8 s interval we calculate 9 enhancement factors, and we
assume that each corresponds to its respective emission line’s formation temperature (listed above). We use a spline fit to place these onto the finer temperature grid (0.05 dex) of the AIA temperature response functions (Boerner et al. 2011). In this way we obtain enhancement factors as functions of temperature for 4.65 ≤ log T ≤ 6.90. See Fig. 6, which shows normalized AIA temperature response functions for all of the EUV channels except 304 Å, along with normalized enhancement factor curves for six CDS exposures up to and including the first peak during the flare’s first precursor. In Fig. 6 and subsequent figures, the AIA temperature response functions are normalized to the maximum value within their respective channels. The EUV spectral line enhancement factors are normalized to the largest overall value during the flare (325, which occurs for Fe XIX during its peak at 18:15:18 UT), multiplied by 10 in Fig. 6 to improve visibility. The point of Fig. 6 is that the majority of the intensity enhancements observed early during the first precursor occur at temperatures less than about log T = 5.8 (0.6 MK), with the largest values occurring around log T = 5.0 (0.1 MK).

Figure 7 is similar to Fig. 6, and shows normalized enhancement factors (in this case not multiplied by a factor of 10) for four CDS exposures during the flare. These include the first peak during the first precursor (17:45:56 UT, which is also shown in Fig. 6), the start of the impulsive increase of the Fe XIX emission at 18:08:24 UT, the peak intensity of the cool, transition region emission lines at 18:10:32 UT, and the peak intensity of the Fe XIX line observed by CDS at 18:15:18 UT (the maximum to which all of the CDS enhancements are normalized). It is evident that although intensity enhancements due to low temperature emission dominate early during the flare, by the time of the maximum O V intensity the contributions from hot flare emission (approaching 10 MK) are significant while at the time of the maximum Fe XIX intensity they dominate the low temperature contributions.

Figure 8 shows fractional contributions to the total brightness increase in AIA’s 94, 131, 171, 193, 211, and 335 Å channels due to plasma with log T ≤ 5.85. Because the temperature range for this calculation is limited to that available to CDS (4.67 ≤ log T ≤ 6.89), we define the fractional contribution for each frame in this figure as the ratio of the enhancement summed over the interval from log T = 4.65 to 5.85 to that summed over the interval from log T = 4.65 to 6.90. The curves show that contributions to the total brightness increase in each channel are greatest for low-temperature emission until about 18:10 UT, when contributions from emission formed at larger temperatures begin to dominate. The lower, darker curve in Figure 8a was calculated with the 94 Å channel’s nominal temperature response function (Boerner et al. 2011) obtained through SSWIDL, while the upper, lighter curve was calculated with a “corrected” version of this function in which the temperature response was multiplied by 5 for log T ≤ 5.85 (see discussion).

Based on our combination of the nominal AIA temperature response curves available through SSWIDL (Boerner et al. 2011) with the (temperature-dependent) intensity enhancement factors derived from the CDS spectral line data of the M2.3 flare on 2001 April 24 (Brosius & Phillips 2004), we estimate that AIA would have observed count rate increases in its 94 Å channel by factors of 1.87 at the first peak of the flare’s first precursor, 6.94 by the start of the impulsive rise of the Fe XIX emission, 40.6 at the overall maximum in the “cool” (e.g., O V) line emission, and 186.5 at the time of the overall maximum in the Fe XIX line emission. These and corresponding values at the same times in the 131, 171, 193, 211, and 335 Å channels are listed in Table 1. By way of comparison, AIA channel enhancement factors observed on 2010 July 31 at the peak (05:13:30 UT) of the earliest flare (precursor?) brightening and the first impulsive peak (05:18:30 UT) of the light curves from area A (see Fig. 1) are given in Table 1.

All of AIA’s EUV channels are sensitive to emission around log T ≈ 5.5 (see Figs. 6 and 7). This means that any flare (like that observed by Brosius & Phillips 2004) in which emission from plasma at log T ≈ 5.5 brightens significantly earlier (by more than AIA’s 12 s cadence) than emission from plasma at “flare” temperatures (log T ≈ 7) will brighten simultaneously in AIA’s 94, 131, 171, 193, 211, and 335 (and 304) Å channels. As can be seen from AIA’s temperature response functions shown in Figs. 6 and 7, all of those channels are also sensitive to emission from plasma at log T ≈ 6. Thus, while a simultaneous brightening in all of AIA’s EUV channels would be consistent with previous spectroscopic observations in which emission from plasma at log T ≈ 5.5 is the first to brighten during a flare (e.g., Brosius & Phillips 2004; Brosius 2009; Brosius & Holman 2009, 2010), the possibility would remain that the actual temperature of the emitting plasma is close to log T = 6 or perhaps spread over a range from log T = 5 to log T ≈ 6. Thus for AIA flare observations in which all of AIA’s EUV channels brighten simultaneously we encounter some ambiguity regarding the temperature distribution of the emitting plasma. Nevertheless, a simultaneous brightening in all of AIA’s EUV channels is consistent with chromospheric heating and evaporation being driven by beams of nonthermal particles.

If, on the other hand, emission from plasma at log T ≈ 7 (the result of direct coronal heating at the reconnection site) brightened significantly before emission from plasma at log T ≈ 5.5, then only the AIA channels that are sensitive to emission from such hot plasma (94 and 131 Å) will brighten significantly while others (193 and 335 Å) may show a smaller brightening and the remainder (171 and 211 Å) will show no effect. This enables AIA to distinguish which mechanism (thermal conduction or nonthermal particle beam) dominates the transport of flare energy from its release site (typically thought to be in the corona) to the chromosphere. By way of quantifying AIA’s flare response for a case in which reconnection-heated, high-temperature (10 MK) plasma appears before the brightening of plasma at lower temperatures, we consider a synthetic case in which only the intensity of the hot flare plasma (6.8 ≤ log T ≤ 6.9) is enhanced early during the flare. Based on calculations similar to those described above for the combination of AIA temperature response curves and CDS flare spectra, we find that the 94 Å channel would brighten by a factor of 4.51 when the intensity enhancement factor is 10, and by a factor of 12.3 when it is 30. Similarly, the 131 Å channel would brighten by factors of 1.18 and 1.57 for the same intensity enhancement factors. For the 171 Å channel we get 1.00 and 1.00, for the 193 Å channel we get 1.03 and 1.09, for the 211 Å channel we get 1.02 and 1.08, and for the 335 Å channel we get 1.18 and 1.57.

4. Discussion

Figure 3 shows that all of the AIA light curves from area A in Fig. 1 started to rise at the same time, nearly 6 min before RHESSI or GOES detected any increase due to the microflare, while Fig. 4 shows that all of the AIA light curves from area B started their rise more than 3 min before RHESSI or GOES detected an increase. This is consistent with expectations presented in Sect. 3, in which the simultaneous increase observed in AIA’s
Fig. 6. Nominal, normalized AIA temperature response functions (Boerner et al. 2011), with CDS enhancement factors (normalized to the largest value which occurs for Fe XIX, then multiplied by 10 to improve visibility in these frames) for six of the first ten CDS exposures, beginning with the first exposure during the first precursor and ending with the first peak during the first precursor of the M2.3 flare. The interval between exposures is 9.8 s. The majority of the intensity enhancements occur at temperatures less than log \( T = 5.8 \), with the largest values at log \( T = 5.0 \) (for the last four CDS exposures in this series). The greatest contributions to AIA’s brightenings occur where the product of the intensity enhancements and the response functions are largest.

Fig. 7. Nominal, normalized AIA temperature response functions, with CDS enhancement factors (normalized to the largest value which occurs for Fe XIX) for four times during the M2.3 flare, including (1) the first peak of the first precursor at 17:45:56 UT (the top-most CDS curve in Fig. 6, but here not multiplied by 10), (2) the start of the impulsive increase of the Fe XIX emission at 18:08:24 UT, (3) the peak intensity of the cool, transition region emission lines at 18:10:32 UT, and (4) the peak intensity of the Fe XIX line observed by CDS at 18:15:18 UT (the maximum to which all of the CDS enhancements are normalized). The interval between exposures is 9.8 s.

94, 131, 171, 193, 211, and 335 Å channels is due to emission from plasma at temperatures between 0.1 and 0.7 MK. This is consistent with nonthermal electrons starting to heat the chromosphere nearly 6 min before RHESSI or GOES detected microflare emission. If the chromosphere were heated by thermal conduction, an earlier increase in high-temperature (≥10 MK) emission is expected to have been seen by RHESSI, GOES, or the AIA channels that were designed with the greatest sensitivity to emission from high-temperature flare plasma (Fe XVIII at 6.3 MK in the 94 Å channel, Fe XXI at 10 MK in the 131 Å
Fig. 8. Fractional contribution to total brightness increase in six AIA EUV channels attributed to plasma with log $T \leq 5.85$. For this investigation our temperature range is constrained to that available to CDS, namely, $4.67 \leq \log T \leq 6.89$. We increased AIA’s 94 Å channel’s low-temperature ($4.65 \leq \log T \leq 5.85$) sensitivity by a factor of 5 to compensate for its known underestimate of its low-temperature sensitivity; this is plotted as the upper (lighter color) curve in frame a), as described further in the text.

Table 1. AIA intensity enhancement factors at selected times during flares.

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<th>Channel</th>
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Notes. (a) Expected AIA intensity enhancement factor based on nominal temperature response functions (Boerner et al. 2011) combined with intensity enhancement factors measured with CDS spectra during the M2.3 flare on 2001 April 24 (Brosius & Phillips 2004). (b) Values observed by AIA in area A of Fig. 1 during the microflare on 2010 July 31. (c) At the first peak of the flare’s first precursor (17:45:56 UT). (d) At the impulsive rise of the Fe XIX emission (18:08:24 UT). (e) At the maximum intensity observed in the cool (e.g., O V) emission lines (18:10:32 UT). (f) At the maximum of the Fe XIX emission (18:15:18 UT). (g) At the first peak after the microflare’s onset (05:13:30 UT). (h) At the microflare’s first impulsive peak (05:18:30 UT).

channel, Fe XXIV at 19 MK in the 193 Å channel, and Ni XXII at 10 MK and Fe XXI at 10 MK in the 335 Å channel). Since this was not the case, we find no observational evidence to suggest that thermal conduction heated the chromosphere or drove evaporation at the start of the microflare (in either area A or area B).

The RHESSI data show no X-ray emission before about 15:16 UT. After this time, the 12–25 keV light curve (see Fig. 3) displays the impulsive character of bremsstrahlung from nonthermal electrons. The first X-ray peak, beginning between 5:17 UT and 5:18 UT, may in fact contain significant emission from nonthermal electrons down to the lowest observed photon energy of 3 keV. Spectral analysis to investigate the likelihood that the higher energy photons (and a significant fraction of the lower energy X-ray photons in the first peak) were from nonthermal electrons was inconclusive because of the low level of emission above 12 keV and the uncertainty in the background subtraction. Despite RHESSI’s observation of no hard X-ray emission before 5:16 UT, nonthermal electrons could still have provided the power input required to heat the chromospheric plasma at these earlier times (e.g., Siarkowski et al. 2009; Falewicz et al. 2011).

In addition to the light curves from areas A and B, we also obtained light curves averaged over the entire $2.5' \times 2.5'$ field of view that we had selected from the AIA cutout service. Because these light curves are averaged over large areas that do not participate in the flare brightenings, features in them appear relatively weakened or suppressed when compared with their counterparts from areas A and B. Although this makes the starting times of some of the features indicated in Figs. 3 and 4 more difficult to determine, the relative timings remain unchanged. Further, and
perhaps most important, is the fact that the light curves averaged over this large area yield no evidence of a “hot” \((T \approx 10 \text{ MK})\) start that may have occurred outside areas A and B above. This provides further support for our conclusion that flare energy was transported by nonthermal particle beams rather than by thermal conduction during the 2010 July 31 B4.8 microflare.

In a recent theoretical investigation O’Dwyer et al. (2010) computed the contributions of spectral lines and continuum emission to AIA’s EUV channels for different regions of the solar atmosphere (coronal hole, quiet Sun, active region, and flare plasma), and concluded that no waveband is strictly isothermal: emission from plasma at a variety of temperatures contributes significantly to each. Although Table 1 of O’Dwyer et al. (2010) indicates that AIA’s 94, 131, and 193 Å channels are dominated by Fe XVIII, Fe XXI, and Fe XXIV, respectively, during a flare, those calculations were performed with a differential emission measure obtained during the decay of an M2 flare (Dere & Cook 1979). This DEM did not include the flare impulsive phase, when bursts of emission from plasma at transition region temperatures \((0.1–0.5 \text{ MK})\) are particularly bright (see, e.g., Fisher et al. 1985a,b,c; and CDS light curves presented by Brosius 2003; Brosius & Phillips 2004 and Brosius & Holman 2007). Thus, as we have demonstrated in Sect. 3 above, emission from transition region plasma will dominate in all of AIA’s EUV channels during a flare’s precursors and early impulsive rise, while emission from plasma at temperatures around 10 MK will dominate later during the event. In an even more recent comparison of AIA images of a non-flaring loop (Boerner et al. 2011) with (2) EUV spectral line flare data obtained with CDS on 2001 April 24 (Brosius & Phillips 2004) on timescales comparable to AIA’s image cadence. The line brightenings that were observed early during the CDS flare occurred at temperatures less than about 0.7 MK, with the largest values around 0.1 MK; this is consistent with the flare’s energy transport being dominated by nonthermal particle beams. Because all of AIA’s EUV channels are sensitive to emission from plasma in the 0.1 to 0.7 MK temperature range, we show that all of AIA’s EUV channels will brighten simultaneously during flares like this, in which energy transport is dominated by nonthermal particle beams. The 2010 July 31 flare observed by AIA and RHESSI displays this behavior, and is consistent with the idea that particle beams transported the energy from the energy release site to the chromosphere, where it began to drive chromospheric evaporation nearly 6 min before flare temperatures around 10 MK were reached. When thermal conduction from a directly-heated, hot (\sim 10 \text{ MK}) plasma is the dominant energy transport mechanism, the AIA channels that are sensitive to emission from such temperatures (particularly the 94 and 131 Å channels) will brighten earlier than the channels that are not sensitive to such temperatures \((171 \text{ and } 211 \text{ Å})\). Thus, based on the differences between AIA’s response to flares whose energy transport is dominated by nonthermal particle beams from those whose energy transport is dominated by thermal conduction, AIA can be used to determine the dominant energy transport mechanism for any given event.

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5. Summary

Coordinated AIA and RHESSI observations of a GOES B4.8 microflare on 2010 July 31 show that emission in all seven of AIA’s EUV channels brightened simultaneously nearly 6 minutes before RHESSI or GOES detected emission from plasma at temperatures around 10 MK. To help interpret these and AIA flare observations in general, we characterized the expected temporal responses of AIA’s 94, 131, 171, 193, 211, and 335 Å channels to solar flare brightenings by combining (1) AIA’s nominal temperature response functions available through SSWIDL (Boerner et al. 2011) with (2) EUV spectral line flare data obtained with CDS on 2001 April 24 (Brosius & Phillips 2004) on timescales comparable to AIA’s image cadence. The line brightenings that we observed early during the CDS flare occurred at temperatures less than about 0.7 MK, with the largest values around 0.1 MK; this is consistent with the flare’s energy transport being dominated by nonthermal particle beams. Because all of AIA’s EUV channels are sensitive to emission from plasma in the 0.1 to 0.7 MK temperature range, we show that all of AIA’s EUV channels will brighten simultaneously during flares like this, in which energy transport is dominated by nonthermal particle beams. The 2010 July 31 flare observed by AIA and RHESSI displays this behavior, and is consistent with the idea that particle beams transported the energy from the energy release site to the chromosphere, where it began to drive chromospheric evaporation nearly 6 min before flare temperatures around 10 MK were reached. When thermal conduction from a directly-heated, hot (\sim 10 \text{ MK}) plasma is the dominant energy transport mechanism, the AIA channels that are sensitive to emission from such temperatures (particularly the 94 and 131 Å channels) will brighten earlier than the channels that are not sensitive to such temperatures \((171 \text{ and } 211 \text{ Å})\). Thus, based on the differences between AIA’s response to flares whose energy transport is dominated by nonthermal particle beams from those whose energy transport is dominated by thermal conduction, AIA can be used to determine the dominant energy transport mechanism for any given event.

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