Dark matter subhaloes as gamma-ray sources and candidates in the first Fermi-LAT catalogue

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Received 7 July 2011 / Accepted 3 November 2011

ABSTRACT

The standard paradigm of hierarchical structure formation in a ΛCDM universe predicts the presence of dark matter subhaloes, hosted by Milky Way-sized galaxies. Anticipated subhalo masses range from 10^{10} down to a cut-off mass between 10^8 and 10^{11} M⊙. If dark matter is composed of heavy self-annihilating or decaying particles, these subhaloes could be visible in the γ-ray band as faint and temporally constant sources without astrophysical counterparts. Based upon realistic subhalo models and current observational constraints on annihilating dark matter scenarios, we predict that one massive Galactic subhalo between 10^8 and 10^9 M⊙ may already be present in the 11-month catalogue of Fermi-LAT. Indeed, at least twelve objects in the first Fermi catalogue qualify as candidates. The most promising object, 1FGL J0030.7+0724, is investigated in detail using a dedicated Swift X-ray follow-up observation and a refined positional analysis of the 24-month Fermi-LAT data. With the new observations, seven point-like X-ray sources have been discovered, of which SWIFT J003119.8+072454, which coincides with a faint radio source (12 mJy at 1.4 GHz), serves as a counterpart candidate of 1FGL J0030.7+0724. The broad-band spectral energy distribution is consistent with a high-energy-peaked blazar. However, flux and extent of 1FGL J0030.7+0724 may also be compatible with a dark matter subhalo. Detection of temporal variability or improved astrometry of 1FGL J0030.7+0724 are necessary to rule out or confirm an astrophysical origin. We discuss strategies to identify γ-ray sources that are associated with self-annihilating dark matter subhaloes.

Key words. dark matter – Galaxy: halo – Galaxy: structure – gamma rays: general

1. Introduction

Several astrophysical observations indicate that in the early as well as in the present Universe a non-baryonic form of dark matter (DM) prevails over the baryonic matter content. Structure formation favours a cold dark matter (CDM) scenario (for recent reviews see, e.g., Bertone et al. 2005; D’Amico et al. 2009; Bertone 2010). However, the nature of DM remains unknown. A class of promising candidates for CDM are stable, weakly interacting, massive particles (WIMPs) with masses between 10 and 10^5 GeV, predicted by theories that extend the standard model of particle physics. The most prominent extensions encompass those based on supersymmetry and universal extra-dimensions, which were invented to solve inconsistencies of the standard model at high energy scales [O(TeV)], and which deliver adequate DM candidates in this way. These particles can self-annihilate or decay, producing detectable signatures in the final states such as energetic photons (γ rays), antimatter, and leptons.

Unravelling the nature of DM remains a challenging problem for astronomy and particle physics, and a variety of attempts to detect signals have been made, using both direct and indirect detection techniques. For instance, multi-wavelength observations of astrophysical targets have constrained the self-annihilation rate of DM, which is related to the thermally averaged annihilation cross section. In particular, regions with high DM densities such as the Galactic Centre (Aharonian et al. 2006a,b; Morselli et al. 2010; Abramowski et al. 2011b), Galactic Ridge (Aharonian et al. 2006d), dwarf spheroidal galaxies (dSphs; Colafrancesco et al. 2007; Albert et al. 2008b; Aharonian et al. 2008, 2009; Aliu et al. 2009; Aharonian et al. 2010; Abdo et al. 2010c; Acciari et al. 2010; Aleksic et al. 2011, H.E.S.S. Collaboration 2011), as well as globular (Colafrancesco et al. 2006; Wagner 2009; Abramowski et al. 2011a) and galaxy clusters (Aleksic et al. 2010; Ackermann et al. 2010) serve as excellent targets. Furthermore, DM annihilation in the entire Galactic halo as well as its subhalo population produces a diffuse γ-ray flux, which contributes to the overall diffuse signal of the Galaxy. Comparatively stringent upper limits on the annihilation cross section have been obtained from the combination of both dedicated observations and bounds obtained from the diffuse γ-ray flux (Abazajian et al. 2010; Cirelli et al. 2010; Meade et al. 2010; Papucci & Strumia 2010; Zaharijas et al. 2010).

Based on the theory of hierarchical structure formation, DM haloes of Milky Way-sized galaxies are anticipated to host numerous DM subhaloes with masses between a cut-off scale 10^{-11}_{-0.3} and 10^{10} M⊙ (e.g., Bringmann 2009), where M⊙ denotes the solar mass. This expectation is a consequence of the early collapse of overdensities in the expanding Universe (Diemand et al. 2005), leading to the formation of initially low-mass haloes, which subsequently serve as building-blocks for larger haloes by merging at later times. Besides analytical calculations (e.g., Berezhnyi et al. 2003, 2006, 2008), recent numerical high-resolution N-body simulations of structure formation in a ΛCDM cosmology (Komatsu et al. 2011), such as the Aquarius Project (Springel et al. 2008a,b) or the Via Lactea II
simulation (Diemand et al. 2008b; Zemp et al. 2009), allow us to study substructures in detail. For a Milky Way-type galaxy, these simulations predict a large number of subhaloes (up to $10^{16}$) with masses distributed following a power law, $dN/dM \propto M^{-\alpha}$, where $\alpha \in [1.9; 2.0]$. The DM density profiles of large-mass subhaloes are found to be similar to the host’s, which results in high central densities (Springel et al. 2008a). The spatial distribution of subhaloes is “anti-biased”, i.e., the dominant fraction is placed far away from the host-halo’s centre.

In self-annihilating DM scenarios, subhaloes are expected to appear as weak point-like or moderately extended γ-ray sources, and a small fraction of them could be detectable with current high- or very high-energy (VHE) γ-ray telescopes (e.g., Pieri et al. 2005, 2008, 2011; Springel et al. 2008b; Kuhlen et al. 2008; Ando 2009; Buckley & Hooper 2010; Brun et al. 2011; Zechlin et al. 2011). Examples for currently operating telescopes are the Fermi-LAT (20 MeV–300 GeV, Atwood et al. 2009) and imaging air Cherenkov telescopes (IACTs; $E \gtrsim 100$ GeV) such as H.E.S.S. (Aharonian et al. 2006c), MAGIC (Albert et al. 2008a; Tridon et al. 2010), and VERITAS (Weeke et al. 2002). In the near future, a significant improvement in the overall sensitivity and lower energy threshold will be achieved by upcoming experiments such as H.E.S.S.–II (Vincent 2005) and CTA (Herrmann 2010; the CTA consortium 2010; Doro 2011). Such instruments are possibly sufficiently sensitive to detect nearby large-mass subhaloes of $O(10^9) M_\odot$ within distances of $O(1)$ kpc.

This paper consists of two separate parts, which can in principle be read independently. The first part (Secs. 2 to 4) investigates the detectability of subhaloes with Fermi-LAT, where the basic framework for predicting the γ-ray properties of subhaloes is laid out in Secs. 2 and 3. In Sec. 4, properties of detectable subhaloes are investigated by means of a fiducial source. In the second part, Sect. 5 discusses the search for DM subhaloes in the first Fermi-LAT point-source catalogue (1FGL) and subsequent multi-wavelength studies of the most promising candidate, 1FGL J0030.7+0724. A discussion of the physical origin of 1FGL J0030.7+0724 and prospects for IACTs are presented in Sect. 6.

Throughout this paper, Hubble’s constant is $H_0 = 73$ km s$^{-1}$ Mpc$^{-1}$, yielding the present value of the Universe’s critical density $\rho_{\text{crit}} = 3H_0^2/(8\pi G_N) \approx 1.48 \times 10^{11} M_\odot\text{Mpc}^{-3}$, where $G_N$ denotes Newton’s gravitational constant (e.g., Spergel et al. 2007).

### 2. Gamma rays from DM subhaloes

In the following, the γ-ray flux from DM subhaloes will be derived, based upon current theoretical models of the corresponding radial density distribution.

With respect to undisturbed, isolated galactic haloes, henceforth field haloes, the general formation history of (embedded) subhaloes differs significantly. Analytical models and numerical N-body simulations of structure formation found their physical properties to depend on particular evolutionary conditions, i.e., formation time, evolution, and orbit (see Diemand et al. 2007, 2008a, and references therein). Tidal interaction with the gravitacional potential of the host halo leads to tidal stripping and heating, and can therefore truncate the outer region of subhaloes. In the following, two different approaches will be discussed. On the one hand, subhaloes are modelled assuming negligible tidal effects and are therefore considered to be in a genuine virialised state. Because this approximation is (at least) valid for field haloes, this model will be tagged as field-halo model (FHM).

On the other hand, a second and more realistic model is considered to account for subhalo evolution, henceforth referred to as subhalo model (SHM).

#### 2.1. Density profile

The subhalo’s DM density profile $\rho(r)$ is assumed to follow

$$
\rho(r) = \frac{\rho_s}{(r/r_s)^\gamma (1 + r/r_s)^{2\gamma}} \begin{cases} 
1 & \text{for } r \leq r_{\text{cut}}, \\
0 & \text{for } r > r_{\text{cut}},
\end{cases}
$$

where $r$ denotes the distance to the subhalo’s centre. In general, the profile cuts at an outer radius $r_{\text{cut}}$, which is the virial or tidal radius ($R_{\text{vir}}$ or $R_{\text{t}}$), respectively. Given $\gamma = 1.0$ for the remainder, the profile follows the universal spherically symmetric Navarro-Frenk-White (NFW) profile, well-fitting haloes in numerical simulations $^1$ (Navarro et al. 1997). The profile is defined by two parameters: a characteristic inner radius $r_s$, where the effective logarithmic slope of the profile is $-2$, and an inner density $\rho_s = 4\rho_c(r_s)$. In case of FHM haloes, which are not subject to tides, both parameters are related to each other by the virial halo mass $M_{\odot}$.

$\rho_c$ is defined as the mass inside the sphere of radius $R_{\text{vir}}$, which encloses a mean density of $\Delta_c$ times the critical density of the Universe at the considered redshift $z$ (Navarro et al. 1997; Bullock et al. 2001), $M_{\odot} := 4\pi/3 \rho_c r_{\text{vir}}^3 R_{\text{vir}}^3$. The virial overdensity at $z = 0$ is $\Delta_c \approx 100$, as suggested by models of the dissipationless spherical top-hat collapse (Eke et al. 1996; Bryan & Norman 1998) and assuming present concordance cosmology. In general, the subhalo mass $M$ is given by a volume integration of Eq. (1), revealing that $M = 4\pi r_s^2 \rho_c f(c)$, where $f(c) \equiv \ln(1 + c) - c/(1 + c)$ and $c$ denotes the concentration parameter of the subhalo. For non-disturbed haloes, the concentration is then given by the virial concentration $c_{\text{vir}} \approx R_{\text{vir}}/r_s$.

Generally, the concentration depends on the subhalo mass and redshift, $c = c(M,z)$, where lighter haloes have higher concentrations (Navarro et al. 1996, 1997; Bullock et al. 2001). Since observational estimates are lacking (see Sect. 2.2.1 in Lallement et al. 2008, and references therein), $c(M)$ is adopted from N-body simulations. For the FHM, the toy model of Bullock et al. (2001)$^2$ is used, where the halo’s (average) virial concentration at redshift $z$ is connected with the density of the Universe at the halo’s (mass-dependent) collapse redshift $z_c(M)$, $c_{\text{vir}} = K(1+z_c)/(1+z)$. The contraction parameter $K$ is constant and independent of cosmology. To determine the mass dependence of $z_c$ and, therefore, of $c_{\text{vir}}$ at $z = 0$, the low-mass extrapolation of the Bullock model by Lallement et al. (2008) is adapted. Because $c_{\text{vir}}$ implicitly depends on $\Delta_c$ (see also Sect. 2.2), a conversion of $c_{\text{vir}}$ to $\Delta_c = 100$ was applied with the relation of Hu & Kravtsov (2003). The concentration-to-mass relation is well-fitted by the polynomial form

$$
\ln(c_{\text{vir}}) = \sum_{i=0}^{4} c_i^{\text{FHM}} \times \left( \ln \left( \frac{M}{M_\odot} \right) \right)^i,
$$

where

$$
c_i^{\text{FHM}} = (4.265, -0.0384, -3.91\times10^{-4}, -2.2\times10^{-6}, -5.5\times10^{-7}).
$$

Note that this model almost equals the relation derived by Pieri et al. (2011) for a cosmology as used in the Aquarius simulation. Regarding the concentration of SHM haloes, the low-mass extrapolation of the Bullock model provides a conservative estimate (cf., Pieri et al. 2008).

$^1$ Note that details on the very inner slope of halo profiles remain to be accounted for.

$^2$ This model extends a proposal by Navarro et al. (1997).

$^3$ Lallement et al. chose $\Delta_c = 81.6$. 
The concentration-to-mass relations are shown in Fig. 1. For the FHM, the scatter is also depicted (68% c.l. of $\log_{10} c$), see Eq. (4). In addition to the concentration of the MW, Fig. 1 contains a selection of eight dSphs that are associated with sufficiently precise stellar data, which allow a conclusive modelling of the DM distribution (see Abdo et al. 2010c). Each dSph is modelled with a NFW profile with parameters chosen to fit measurements of stellar line-of-sight velocities and their distributions (see also Martinez et al. 2009). The dSph’s virial concentration is given by its characteristic density, $\rho_s = \Delta_{\text{half}} c_{\text{vir}}^3 / [3 f(c_{\text{vir}})]$, where tidal effects on the inner system are assumed to be negligible. The SHM is depicted for $D_{\text{gc}} = 74$ kpc, the average galactocentric distance of the dSph subset. Additionally, the models are confronted with direct predictions of the Aquarius simulation, derived from scaling relations fitting subhaloes observed in the simulation. Details are provided in Appendix A.

Within its scatter, the concentration model of FHM haloes consistently describes the DM profile of dSph galaxies and the MW itself. However, the median values $c_{\text{vir}}^{\text{FHM}}$ underpredict dSphs, whereas the subhalo model SHM provides convincing agreement (as expected by $N$-body simulations). The concentration derived directly from the Aquarius simulation confirms the SHM within the validity range, see Fig. 1. Note that the mean distance of subhaloes resolved in Aquarius is 64 kpc.

### 2.2. DM annihilation in subhaloes

For self-annihilating particles, the total rate of photons (or particles) emitted by a DM subhalo with energy $E$ in the interval $[E_1, E_2]$ is

$$L = \frac{\langle \sigma v \rangle_{\text{eff}} N_{\nu}}{2 m_\chi^2} \int dV \rho^2(r) \frac{M^2}{m^2} \gamma(r, \nu, \nu') N_{\nu} = \int \frac{dx}{dx_1} \frac{d N_{\nu}}{dx},$$

(5)

where $\langle \sigma v \rangle_{\text{eff}}$ is the thermally averaged annihilation cross section times the relative velocity, $m_\chi$ the WIMP mass, and $dN_{\nu} / dx$ denotes the differential spectrum of photons per annihilation.

Assuming $r_s \ll D$, the produced photon flux is given by $\phi \propto L/(4 \pi D^2)$. The solution of the integral holds for $y = 1.0$ and $c \gg 1$. In Eq. (5), a small, flat core replacing the unphysical singularity at the halo centre (Berezinsky et al. 1992) is safely neglected (given the NFW profile used here). For $\gamma = 1.0$, Eq. (5) simplifies via $r_c = [3 M / (4 \pi \Delta_{\text{halo}} c_{\text{vir}}^3)]^{1/3}$:

$$L = \frac{\langle \sigma v \rangle_{\text{eff}} N_{\nu} \Delta_{\text{halo}} M c_{\text{vir}}^3}{18 m_\chi^2 f(c_{\text{vir}})^2}.$$

(6)

For a $y = 1.2$ profile (Eq. (1)) the photon rate increases by a factor of $-1.5$ for subhaloes above $10^7 M_\odot$. DM annihilation in subhaloes may be additionally boosted by sub-substructure populations (see Strigari et al. 2007; Kuhlen et al. 2008; Martinez et al. 2009). Conveniently, the value of $\langle \sigma v \rangle_{\text{eff}}$ is normalised to the value $\langle \sigma v \rangle_0 = 3 \times 10^{-26}$ cm$^2$ s$^{-1}$, which leads to the correct relic density. An increase of the annihilation rate, a so-called boost factor $\langle \sigma v \rangle_{\text{eff}} / \langle \sigma v \rangle_0$, could in principle be related to the underlying particle physics framework (e.g., Fornengo et al. 2004) and effects such as Sommerfeld enhancement (e.g., Arkani-Hamed et al. 2009; Kuhlen et al. 2009).

We considered DM to be composed of self-annihilating heavy WIMPs of mass $m_\chi = 500$ GeV and investigated three distinct annihilation models: two of them with total annihilation in heavy quarks or gauge bosons ($b\bar{b}$ and $W^+W^-$) and a...
model with total annihilation in the leptons $\tau^+\tau^-$. Additionally, we considered WIMPs of $m_w = 150$ GeV for annihilation in $\tau^+\tau^-$ final states. Particles of this type are, for instance, provided by supersymmetric theories, e.g., manifested in the neutralino. The WIMP masses chosen are compatible with WIMPs which might explain the recently observed cosmic-ray electron and positron excess, see, e.g., Meade et al. (2010). In general, WIMPs are also supported by collider searches such as the non-detection of supersymmetric particles in the 7 TeV run of the Large Hadron Collider (for an integrated luminosity of 35 pb$^{-1}$) (Aad et al. 2011; CMS Collaboration 2011, and references therein).

Given these final annihilation states, hadronisation and the subsequent decay of $\pi^0$-mesons lead to a continuous $\gamma$-ray spectrum. The resulting photon spectra $dN_\gamma/dE$ were modelled using parametrisations provided by Fornengo et al. (2004), see Fig. 2. That notes photons produced by final state radiation (FSR) and virtual internal bremsstrahlung (VIB) (Bergström et al. 2005a,b; Bringmann et al. 2008) are neglected, because a significant contribution of FSR is only expected for $W^+W^-$ at high energies ($E > 0.6 m_w$) and possible contributions of VIB are highly model-dependent.

### 3. Candidate gamma-ray sources

Given a high WIMP mass, a DM subhalo will show up as steady (very) high-energy $\gamma$-ray source. The differential photon spectrum follows a hard power law (index $\Gamma_\gamma \lesssim 1.5$) that cuts off exponentially at energies beyond 10 GeV, see Fig. 2.

Candidate sources are selected according to their possibility to originate from DM subhaloes based on their observational quantities flux and angular extent. Note that the detailed spectral shape of faint sources is observationally rather unconstrained. Via Eq. (6), the effective self-annihilation cross section $\langle \sigma v \rangle_{\text{eff}}$ required to obtain a given flux $\phi_f$ for the intrinsic source extent $\theta_i$ is determined by $L = 4\pi D^2 \phi_f$, where $\theta_i$ constrains the distance $D$ to the subhalo. For feasible candidate sources the required $\langle \sigma v \rangle_{\text{eff}}$ should comply with current observational constraints.

Conveniently, the characteristic profile radius $r_s$ (see Eq. (1)) traces the intrinsic extent of a DM subhalo, because for an NFW profile 87.5% of the total luminosity is produced within $r_s$ (see Table 1 for values of $r_s$). Therefore, the distance to the subhalo is $D \approx r_s/\theta_i$, where $\theta_i$ denotes the angle corresponding to $r_s$. Owing to the centrally peaked profile, about 68% of the total luminosity is emitted within the angle $\theta_{68} \approx 0.46 \theta_i$. The following relations are given with respect to $\theta_i$ and can easily be adjusted for $\theta_{68}$, which is more convenient for a comparison with observational data. With $r_s = [3M/(4\pi\Delta \rho \sigma v \theta_i^2)]^{1/3}$, the distance to a subhalo with given $\theta_i$ is related to its mass and concentration. In the FHM,

$$D_{\text{FHM}}(M; \theta_i) = 3.8 \left( \frac{M}{10^9 M_\odot} \right)^{1/3} \phi_{\text{FHM}}^{1/3} \left( \frac{\Delta \rho}{37.9} \frac{\theta_i}{\text{deg}} \right)^{-1} \text{kpc}. \quad (7)$$

Note that the concentration depends on the subhalo mass via Eq. (2) as well as, in the SHM, on the object’s position ($l, b$).

For a given WIMP model $(\sigma v)_{\text{eff}}$ is then fully determined by the subhalo mass (Eq. (6)) and the observed quantities flux and extent:

$$\langle \sigma v \rangle_{\text{eff}}(M, \phi, \theta_i) = 96 \pi^2 m_w^2 \left( \frac{3}{4N_p \Delta \rho \sigma v} \right)^{5/3} \frac{\phi}{\theta_i^2} \frac{M_{-1}^{-1/3} f(c_{\text{vir}}^2)}{c_{\text{vir}}^3}. \quad (8)$$

Additional contributions to the DM signal from annihilation in the smooth halo as well as the entire subhalo population were neglected. The required $(\sigma v)_{\text{eff}}$ is highly sensitive to the (observationally unconstrained) concentration, because Eq. (8) roughly depends on $c_{\text{vir}}^{-5}$.

### 4. Interpretation of Fermi sources as DM subhaloes

Based on the study of a fiducial candidate source in Sect. 4.1, the properties of Fermi-LAT detectable subhaloes are investigated in Sect. 4.2.

#### 4.1. A fiducial candidate

##### 4.1.1. Observational properties

In combination with improving (integrated) sensitivity at high energy (Atwood et al. 2009, and cf., Fig. 7), the expected energy spectrum of DM subhaloes (Sects. 2 and 3) favours a detection at the high-energy band of Fermi-LAT. In Appendix B, we investigate the detection sensitivity for faint, moderately extended $(\theta_i \lesssim 1^\circ)$, to corresponding $\theta_{68} \lesssim 0.5^\circ$), and high-latitude ($|b| > 20^\circ$) sources between 10 and 100 GeV in detail. We find that a spectrally hard high-energy source with a flux $\phi(10–100 \text{GeV}) = 1.4 \times 10^{-10} \text{cm}^{-2} \text{s}^{-1}$ and moderate extent $\theta_i = 1^\circ$ can be detected as a point-source with a reconstructed flux $\phi_f(10–100 \text{GeV}) = 0.9 \times 10^{-10} \text{cm}^{-2} \text{s}^{-1}$, with a sky-survey exposure of one year. With respect to the true flux $\phi$ emitted by the entire source, the reconstructed flux $\phi_f$ fitted by the point-source analysis in general decreases with increasing $\theta_i$. To account for this effect, the scaling relation $\phi(\theta_i = h(\theta_i) \phi_f$ is used in Eq. (8), where $h(\theta_i) = 1$ for $\theta_i < 2 \theta_{68}$ and $5$. For the fiducial candidate in Sect. 4 this additional contribution is less than 1%.

In comparison with the point spread function of Fermi-LAT, $\theta_{68} = 1^\circ$ (as implying $\theta_{68} \approx 0.5^\circ$) corresponds to about $3 \theta_{68}$, where $\sigma_{\text{psf}} \approx 0.15^\circ$ for energies beyond 10 GeV (see http://www.glast.slac.stanford.edu/software/IS/glast_lat_performance.htm).
freeze-out (blue-shaded). 

The high-energy flux of the fiducial source above 10 GeV has been chosen to be at the level of the detection sensitivity, \( \phi_F (10–100 \text{ GeV}) = 10^{-10} \text{ cm}^{-2} \text{s}^{-1} \), assuming an extent of \( \theta_{\text{fd}} = 1^\circ \). Given the dependence of the SHM concentration on the galactocentric distance (see Eq. (3)), the fiducial source is placed on a particular line-of-sight chosen to match the location of 1FGL J0030.7–5524 (investigated in Sect. 5). In general, this line-of-sight serves as an appropriate (conservative) benchmark, because it points to (anticentric) positions where the majority of subhaloes is located.

4.1.2. Subhalo interpretation

Adopting the properties of the fiducial source, Fig. 3 depicts the effective enhancement factors \( (\sigma v)_{\text{eff}} / (\sigma v)_0 \) required to generate the emission \( \phi_{\text{fd}} \) by DM annihilation (obtained via Eq. (8)). In the left panel, WIMPs of \( m_\chi = 500 \text{ GeV} \) are considered to totally annihilate in \( b\bar{b} \), while the right panel assumes \( m_\chi = 150 \text{ GeV} \) and annihilation in \( \tau^+ \tau^- \). For a given WIMP model, the resulting enhancement factors of the FHM and SHM are widely different. With respect to the FHM, much less enhancement is required in the SHM, which is manifested in generically higher concentrations of SHM subhaloes. Within the scatter of the concentration intrinsic to the stochastic nature of halo formation (Eq. (4)), which is shown for the SHM, the necessary enhancement spans about one order of magnitude. Only moderate enhancement is required for massive subhaloes between \( 10^6 \) and \( 10^7 M_\odot \), where the lowest \( (\sigma v)_{\text{eff}} \) is needed for \( m_\chi = 150 \text{ GeV} \) and \( \tau^+ \tau^- \) final states (amongst the WIMP models considered here). The lowest possible enhancement factors within the concentration scatter of a \( 10^6 M_\odot \) subhalo are listed in Table 1 for the different subhalo and the WIMP models of Sect. 2.

The distance to the fiducial candidate anticipated in the FHM and the SHM is shown in Fig. 4. The intrinsic concentration scatter implies a corresponding distance scatter for a given halo mass and angular extent. Note that a similar scatter is present for the FHM, but is not shown in the figure. Compared with the FHM, tidal effects lead to higher concentrated subhaloes. This in turn favours a closer distance at the same mass and angular extent for SHM subhaloes than for FHM.

4.1.3. Consistency with observational constraints

The resulting values of \( (\sigma v)_{\text{eff}} \) can now be checked for consistency with the diffuse extragalactic \( \gamma \)-ray background (EGB; Abdo et al. 2010b). Fermi-LAT measurements of the overall diffuse \( \gamma \)-ray flux allow the derivation of the isotropic high-energy EGB, which is shown to be compatible with a featureless power-law spectrum (\( \Gamma = 2.41 \pm 0.05 \)) and integrated diffuse flux \( \phi_{\text{EGB}}(>100 \text{ MeV}) = (1.03 \pm 0.17) \times 10^{-5} \text{ cm}^{-2} \text{s}^{-1} \text{ sr}^{-1} \) (Abdo et al. 2010b). The diffuse \( \gamma \)-ray flux anticipated from DM annihilation in the Galactic halo as well as the entire subhalo population is shown in Fig. 5 (see Appendix C for details), in comparison with the EGB. Both FHM and SHM subhaloes are depicted for \( m_\chi = 150 \text{ GeV} \) and annihilation in \( \tau^+ \tau^- \). In this model, the nearly isotropic diffuse flux from the subhalo population contributes about 1% to the EGB (assuming SHM subhaloes and no sub-substructure) and is fainter than the contribution of the smooth halo (\( \gtrsim 3\% \)). Note that the flux from the extragalactic halo population is lower than the contribution of Galactic subhaloes, see, e.g., Abazajian et al. (2010).
The scatter of the SHM distance is indicated by the red-shaded area. An extent of $\theta_{\text{sh}} \approx 0.5^\circ$ and a source position $(l, b) = (114^\circ, -55^\circ)$ are assumed.

### Table 2. Upper limits on $(\sigma v)_{\text{cut}}/(\sigma v)_0$ from the EGB.

<table>
<thead>
<tr>
<th>WIMP model</th>
<th>Channel</th>
<th>$m_\chi$ [GeV]</th>
<th>Upper limit on $(\sigma v)_{\text{cut}}/(\sigma v)_0$</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>FHM</td>
<td>SHM</td>
<td>FHM</td>
</tr>
<tr>
<td>$b^0 \gamma$</td>
<td>500</td>
<td>1650</td>
<td>530</td>
</tr>
<tr>
<td>$W^+ W^-$</td>
<td>500</td>
<td>2096</td>
<td>673</td>
</tr>
<tr>
<td>$\tau^+ \tau^-$</td>
<td>500</td>
<td>3490</td>
<td>1121</td>
</tr>
<tr>
<td>$\tau^- \tau^-$</td>
<td>150</td>
<td>378</td>
<td>121</td>
</tr>
</tbody>
</table>

$M_{\text{min}} = 10^{-10} M_\odot$, $10^{-6} M_\odot$.

Notes. We assumed a subhalo mass fraction of $f_{\text{sh}} = 15\%$ for a cut-off mass $M_{\text{min}} = 10^{-6} M_\odot$. Upper limits are listed for the cut-off masses bordering a 500 GeV neutralino scenario (see Bringmann 2009) and with respect to $\psi = 180^\circ$ and $E = 40$ GeV. See text and Appendix C for details.

The contribution from the smooth halo peaks at the Galactic Centre, where a high astrophysical foreground is also present, and can therefore not be isotropic. Given that the EGB has been derived assuming isotropy, the most robust upper limits on $(\sigma v)_{\text{cut}}$ are determined by the subhalo contribution and are listed in Table 2, depending on the WIMP model and cut-off mass. The bounds were obtained requiring that the specific intensity of the subhalo population $(I_{\nu}(180^\circ, E))$ does not exceed the EGB, where $(I_{\nu}(\psi, E))$ depends on the angle $\psi$ between the Galactic Centre direction and line-of-sight, and the $\gamma$-ray energy $E$ (see Appendix C).

However, more stringent constraints have been provided by a more detailed modelling of the EGB, including all DM components. To evaluate a possible DM origin of the fiducial source, the results of Abazajian et al. (2010) and Papucci & Strumia (2010)\(^7\) are included in Fig. 3. As stated in the introduction, competitive and similar constraints have been also provided by the non-detection of various objects with high (central) DM densities. Note, for instance, that the constraints used here are consistent with recent bounds from dSph galaxies (Abdo et al. 2010c).

7 In comparison with Papucci & Strumia, the work by Abazajian et al. includes a fore- and background subtraction. Note that the MW halo parameters used by Abazajian et al. are similar to those adopted in this work (see Sect. 2.1).
one year. Given the Poisson distribution of $N$, this means that up to one massive subhalo is expected in the one-year data set (at 95% confidence). Comparable results have been claimed by other authors, e.g., Baltz et al. (2008), Kuhlen et al. (2008), Pieri et al. (2008, 2011), and Anderson et al. (2010).

In general, note that numerical simulations like Aquarius and Via Lactea II neglect the influences of baryonic matter distributed in galactic disks. A recent study by D’Onghia et al. (2010) indicates that a baryonic disk may reduce the number of (massive) subhaloes in the inner galaxy by a factor of 2 to 3.

### 4.2.2. Properties

Given a particular $\gamma$-ray source (such as the fiducial of Sect. 4.1), massive subhaloes between about $10^5$ and $10^8 M_\odot$ require a minimally enhanced annihilation cross section $\langle cr \rangle_{\text{eff}}$, see Fig. 3. Consistently, the probability for these objects to appear in current data sets peaks for high subhalo masses (Fig. 6). Therefore, subhaloes with masses of $10^5$ up to $10^8 M_\odot$ in corresponding distances from 0.5 to 10 kpc (Fig. 4) are favoured for detection with Fermi-LAT as faint and moderately extended sources. Although for the FHM a DM origin of the fiducial source is excluded, within the scatter of the more realistic SHM cross sections required for sources with $\phi(10-100\text{GeV}) \approx 10^{-10} \text{cm}^{-2} \text{s}^{-1}$ and angular extents up to ~1° are well consistent with observational constraints. The presence of sub-structure will even lower $\langle cr \rangle_{\text{eff}}$ by a mass-dependent factor of ~2 to 3 for massive subhaloes (Kuhlen et al. 2008; Martinez et al. 2009). In case of a cuspey profile ($\gamma = 1.2$), the required cross section is lowered by an additional factor of 1.5. At least for WIMPs of $m_\chi = 150$ GeV annihilating to $\tau^+\tau^-$, this leads to a required enhancement of the order of unity within the scatter. Summarising, in optimistic but realistic scenarios a $\gamma$-ray emitter at the (one year) detection level of Fermi-LAT with a measured extent $\theta_{58}$ of about 0.5° might be consistent with a subhalo driven by self-annihilating DM.

In the next years, Fermi-LAT will provide deeper observations with increased observation time $T_{\text{obs}}$. This will improve the sensitivity by roughly $\sqrt{T_{\text{obs}}}$, leading to a factor of about 2 lower values of the minimum detectable flux for the five-year catalogue. This in turn will allow us to detect fainter subhaloes with a correspondingly reduced minimum $\langle cr \rangle_{\text{eff}}$. The average number of detectable subhaloes within five years is about 1.3 (0.8) for the $bb$, $m_\chi = 500$ GeV ($\tau^+\tau^-$, $m_\chi = 150$ GeV) scenario.

Via Eq. (8), a comparison with observational constraints on $\langle cr \rangle_{\text{eff}}$ allows to estimate the maximum subhalo candidates are expected to have. The catalogued flux is $\phi_p \approx \theta_{58}^2 / h(\theta_p)$, which is quadratic for $\theta_p \ll 0.3°$ and linear in the limit of large $\theta_p$. The increase with $\theta_p$ originates from decreasing subhalo distance. Because massive subhaloes require minimum $\langle cr \rangle_{\text{eff}}$, a subhalo of $10^5 M_\odot$ is assumed below. Within the concentration scatter, the maximum flux between 10 and 100 GeV expected for a source with $\theta_p = 1°$ is

$$\phi_p^{\text{max}}(10-100\text{GeV}) \approx 1.2 (2.8) \times 10^{-11} \frac{\langle cr \rangle_{\text{eff}}}{\langle cr \rangle_0} \text{cm}^{-2} \text{s}^{-1}$$

In the $bb$, $m_\chi = 500$ GeV and $\tau^+\tau^-$, $m_\chi = 150$ GeV scenario, respectively. Given the observational constraints of Abazajian et al., the high-energy flux of catalogued candidates should not exceed $\phi_p^{\text{max}} \leq 4.0 (2.8) \times 10^{-10} \text{cm}^{-2} \text{s}^{-1}$.

### 5. Searches for DM subhaloes in the 1FGL

In the previous section we demonstrated that DM subhaloes could appear in $\gamma$-ray catalogues of sufficient sensitivity as faint, non-variable, and moderately extended objects without astrophysical counterparts. The 11-month$^8$ point-source catalogue of Fermi-LAT (1FGL, Abdo et al. 2010a) lists 1451 sources significantly detected above 100 MeV (test statistic TS ≥ 25, corresponding to a significance $S \geq 4.1\sigma$), together with the flux in five discrete energy bins (up to 100 GeV), position, significance of variability, and spectral curvature. Source spectra have been fitted with power laws. Among the sources, 630 objects$^9$ are not confidently associated with known sources at other wavelengths.

Although sophisticated methods have been applied to find multi-wavelength associations for unidentified sources, all algorithms suffer from lacking sensitivity or incomplete sky-coverage of current surveys. Therefore, the sample of unassociated high-latitude Fermi-LAT sources is expected to be composed of several source classes, among them faint AGN (Active Galactic Nuclei), galaxy clusters, and new exotic sources like DM subhaloes (Stephen et al. 2010; Mirabal et al. 2010, and references therein). Concerning the 1FGL catalogue, improved association methods recently presented by Mirabal et al. (2010) revealed that ≤20% of all unassociated 1FGL sources with $|b| \geq 15°$ may contain new $\gamma$-ray emitters.

To single out possible subhalo candidates within the sample of unassociated sources (cf., Buckley & Hooper 2010), we searched the sample for non-variable$^{10}$ sources detected between 10 and 100 GeV. Requiring a detection at high energy provides subhalo candidates driven by heavy WIMPs and avoids confusion with high-energy pulsars$^{11}$. Furthermore, the candidate’s
Table 3. Cut efficiencies on the sample of unassociated sources and AGN.

<table>
<thead>
<tr>
<th>Cut</th>
<th>Unassociated</th>
<th>AGN</th>
</tr>
</thead>
<tbody>
<tr>
<td>$</td>
<td>\theta</td>
<td>\geq 20^\circ$</td>
</tr>
<tr>
<td>non-variable</td>
<td>241 (97%)</td>
<td>372 (69%)</td>
</tr>
<tr>
<td>detected between 10–100GeV</td>
<td>12 (5%)</td>
<td>58 (11%)</td>
</tr>
</tbody>
</table>

Notes. The cuts are cumulative, i.e., each number includes all cuts listed by previous rows. See text for details.

location was constrained to galactic latitudes $|\theta| \geq 20^\circ$ to avoid a general confusion with Galactic sources. Additionally, the lower Galactic foreground improves the detection sensitivity of Fermi-LAT at high latitudes (Atwood et al. 2009).

Applying all cuts, twelve unassociated sources remain. The twelve sources are listed in Table 4 together with additional information from the catalogue. With the exception of 1FGL J0614.1-3328, the sample consists of sources at the faint end of the entire 1FGL sample. Given the result of Mirabal et al. (2010), the sample should statistically contain two to three subhaloes at most, consistent with the estimate discussed previously (Fig. 6). The expectation of the sample consisting mostly of AGN is met by applying the same cuts to all AGN detected by Fermi-LAT. A comparison with the sample of unassociated sources reveals similar cut efficiencies (5% vs. 11%, see Table 3), indicating that the two populations share common properties. Note that for the AGN the variability cut has subdominant influence as well, see Table 3. Except for three, all AGN that passed the cuts have been classified as BL Lac.

Even though the twelve candidate objects are listed in the 1FGL catalogue as unassociated, we extended the counterpart search to a wider choice of astronomical catalogues. Table 4 lists the classifications of counterpart candidates in the 68% confidence regions around the most likely 1FGL positions, retrieved from the NASA/IPAC Extragalactic Database (NED). In particular, possible associations are provided by radio and X-ray sources, since most of the selected γ-ray sources are expected to be AGN. Given that no detailed association study was conducted, some of the tabulated sources might be by-chance associations.

Governed by lacking association, faintness, and spectral shape, this study focusses on the most promising candidate, 1FGL J0030.7+0724. Within the errors, its high-energy flux and spectral index are well-compatible with a self-annihilating DM scenario. The source has only been detected between 10 and 100 GeV, see Fig. 7.

5.1. Multi-wavelength properties of 1FGL J0030.7+0724

5.1.1. Catalogued data

No counterpart candidate was found within the positional uncertainty of the γ-ray source at a 68% confidence level (Table 4). In the 95% confidence region, the faint radio object NVSS J003030+072132 is located (f1400mJy = (3.5 ± 0.4)mJy; Condon et al. 1998). However, no conclusive infrared (2MASS, Skrutskie et al. 2006) or optical (USNO B1.0, Monet et al. 2003) association of the NVSS source is known so far12. Note that

12 Within the $2\sigma$ positional uncertainty of NVSS J003030+072132, a very faint optical SDSS (Sloan Digital Sky Survey) source is located no dSph galaxy is located in the source region (NED). ROSAT (0.1–2.4 keV; Voges et al. 1999) observations of the region with an exposure of about 107 revealed no X-ray source down to an energy-flux level of $\sim 10^{-12}$ erg cm$^{-2}$ s$^{-1}$ (Borm 2010).

We emphasise that the orphaned faint radio source is likely located in the uncertainty region of 1FGL J0030.7+0724 by chance, because about 0.7 NVSS sources are expected by statistics.

5.1.2. Fermi-LAT data

By analysing the 24-month public archival data between 10 and 100 GeV, updated results on 1FGL J0030.7+0724 will be provided. For the same energy range, a reanalysis of the 11-month data is presented for comparison. Particular focus will be drawn on positional properties, the high-energy flux, and the photon distribution, which allows us to investigate possible counterparts, temporal variability, and the angular extent.

The data analysis was performed with the latest public version of the Fermi science tools (v9r18p6)13 along with recommended options and the set of instrument-response functions $P6\_V3\_DIFFUSE$ (Rando 2009). Throughout the analysis, the optimiser MINUIT was used. For reliable results, photons of event class 3 (Diffuse) and 4 (DataClean) within a radius of 10°, centred on the nominal position of 1FGL J0030.7+0724, were selected. Given that most of the photons are at highest energies (Fig. 7), only events between 10 and 100 GeV were selected to minimise the background and ensure a narrow PSF. The data were processed using gtselect, gmktime, gtcube, gtexpmap, and unbinned gtlike. To compute the most likely position and its corresponding uncertainty on basis of the 10–100 GeV photon sample, we used gfindsrc. For the purpose of detailed counterpart searches, the two-dimensional likelihood function $L(RA, Dec)$ was computed, which provides the 95% uncertainty contour by $2 \Delta(log L) = 6.18$ (2 degrees of freedom).

(26.0°) – SDSS J003031.22+072132.2 (SDSS DR7, Abazajian et al. 2009). However, this object was observed with the edge of the plated SDSS camera. Therefore, this detection is probably spurious.

Table 4. Unassociated, non-variable 1FGL sources at high galactic latitudes.

<table>
<thead>
<tr>
<th>Name</th>
<th>$\sigma_{\alpha/\delta}$ [arcmin]</th>
<th>$S$ [$]</th>
<th>$f_\nu(0.1–100 \text{ GeV})$ [10$^{-11}$ erg cm$^{-2}$ s$^{-1}$]</th>
<th>$\Gamma$</th>
<th>$\phi_\nu(10–100 \text{ GeV})$ [10$^{-10}$ cm$^{-2}$ s$^{-1}$]</th>
<th>Possible associations\textsuperscript{a}</th>
<th>Remarks</th>
</tr>
</thead>
<tbody>
<tr>
<td>0022.2-1850</td>
<td>6.0/9.6</td>
<td>9.4</td>
<td>1.3(4)</td>
<td>1.6(1)</td>
<td>1.6(7)</td>
<td>RadioSs (4, 21, 22), Gs (20–18)</td>
<td>68% c.l.</td>
</tr>
<tr>
<td>0030.7+0724</td>
<td>3.0/5.1</td>
<td>5.8</td>
<td>1.0(4)</td>
<td>1.7(4)</td>
<td>1.5(7)</td>
<td>–</td>
<td></td>
</tr>
<tr>
<td>0051.4-6242</td>
<td>2.4/4.2</td>
<td>12.0</td>
<td>1.8(5)</td>
<td>1.7(1)</td>
<td>1.7(8)</td>
<td>Gs (20), XrayS (3.8)</td>
<td>c</td>
</tr>
<tr>
<td>0143.9-5845</td>
<td>3.0/4.7</td>
<td>9.0</td>
<td>1.4(4)</td>
<td>2.0(2)</td>
<td>2.0(9)</td>
<td>RadioS (28'), Gs (20–13)</td>
<td>c</td>
</tr>
<tr>
<td>0335.5-4501</td>
<td>2.4/4.0</td>
<td>8.6</td>
<td>1.5(4)</td>
<td>2.1(2)</td>
<td>1.6(8)</td>
<td>Gs (19, 18)</td>
<td>c</td>
</tr>
<tr>
<td>0614.1-3328</td>
<td>1.2/1.7</td>
<td>54.4</td>
<td>11.2(6)</td>
<td>1.93(3)</td>
<td>3(1)</td>
<td>GrayS</td>
<td>b</td>
</tr>
<tr>
<td>0848.6+0504</td>
<td>5.4/8.6</td>
<td>5.4</td>
<td>1.0(5)</td>
<td>1.2(3)</td>
<td>1.6(8)</td>
<td>RadioSs (2, 3, 5), Gs &amp; *s, XrayS (4.4)</td>
<td>c</td>
</tr>
<tr>
<td>1331.3+2942</td>
<td>1.8/2.7</td>
<td>11.9</td>
<td>2.0(1)</td>
<td>2.0(1)</td>
<td>2.1(8)</td>
<td>RadioSs (2.8, 263, 724), Gs &amp; *s</td>
<td>c</td>
</tr>
<tr>
<td>1754.3+3212</td>
<td>2.4/4.1</td>
<td>15.6</td>
<td>2.6(4)</td>
<td>2.0(9)</td>
<td>1.4(7)</td>
<td>RadioS (38')</td>
<td></td>
</tr>
<tr>
<td>2134.5-2130</td>
<td>3.0/5.1</td>
<td>6.7</td>
<td>1.1(3)</td>
<td>1.9(2)</td>
<td>1.4(7)</td>
<td>RadioS (22), Gs (20)</td>
<td></td>
</tr>
<tr>
<td>2146.6-1345</td>
<td>3.0/4.4</td>
<td>9.8</td>
<td>1.5(5)</td>
<td>1.8(2)</td>
<td>1.8(8)</td>
<td>RadioS (23), Gs (20), XrayS (1.9)</td>
<td>c</td>
</tr>
<tr>
<td>2329.2+3755</td>
<td>1.2/1.9</td>
<td>10.4</td>
<td>1.7(5)</td>
<td>1.6(2)</td>
<td>2.4(9)</td>
<td>G (14)</td>
<td>c</td>
</tr>
</tbody>
</table>

Notes. The columns list the positional uncertainty $\sigma_{\alpha/\delta}$ [68% (95%) c.l., semimajor axis], detection significance $S$ in Gaussian sigma, integrated energy flux $f_\nu(0.1–100 \text{ GeV})$, spectral index $\Gamma$, and the photon flux $\phi_\nu(10–100 \text{ GeV})$. Here, parentheses indicate the corresponding error on the last decimal(s). Furthermore, the type of classifications sources found in astronomical catalogues within the 68% uncertainty region of the Fermi-LAT position are listed. \textsuperscript{a} Classification refers to RadioS (radio source), G (galaxy), \textsuperscript{b} S (star), XrayS (X-ray source), and GrayS (gamma-ray source). For radio, optical, and X-ray sources corresponding fluxes are given in mJy (at 1.4 GHz [($E_{\gamma}$)], position are listed. \textsuperscript{c} Hartman et al. (1999) catalogue, respectively. \textsuperscript{d} The spectrum is probably curved. \textsuperscript{e} The second column lists the analysed energy range. Their parameters were taken as catalogued set \textsuperscript{f} of 0.5\degree \textsuperscript{g} of 0.5\degree. The unabsorbed X-ray flux was derived from the catalogued count-rate, assuming a power law with index 2.0 (with WebPIMMS, http://heasarc.gsfc.nasa.gov/Tools/w3pimms.html). The hydrogen column density was obtained from the LAB survey, see Sect. 5.1. Sources refer to are listed in the FIRST (Becker et al. 1995), JVASS/CLASS (Jackson et al. 2007), NVSS, SUMSS (Mauch et al. 2003), 2MASS, APMUKS (Maddox et al. 1990), SDSS (Abazajian et al. 2009), ROSAT, or EGRET (Hartman et al. 1999) catalogue, respectively. \textsuperscript{f} The spectrum is probably curved. \textsuperscript{g} The $\gamma$-ray source has been associated by a cross-correlation of unidentified Fermi-LAT sources with the ROSAT All Sky Survey Bright Source Catalogue (see Stephen et al. 2010 for details).

Table 5. Positional and spectral properties of 1FGL J0030.7+0724 as given in the catalogue (11 months) and by the (re-)analysis of the first 11 and 24-month data sets.

<table>
<thead>
<tr>
<th>Data set</th>
<th>$E$ [GeV]</th>
<th>RA (J2000)</th>
<th>Dec (J2000)</th>
<th>$f_\nu(10–100 \text{ GeV})$ [10$^{-10}$ cm$^{-2}$ s$^{-1}$]</th>
<th>$\phi_\nu(10–100 \text{ GeV})$ [10$^{-10}$ cm$^{-2}$ s$^{-1}$]</th>
<th>$\Delta t$ [s]</th>
</tr>
</thead>
<tbody>
<tr>
<td>1FGL</td>
<td>0.1–100</td>
<td>00 30 42.6</td>
<td>+07 24 09</td>
<td>1.5 ± 0.7</td>
<td>6.6</td>
<td>3</td>
</tr>
<tr>
<td>11</td>
<td>10–100</td>
<td>00 30 37.6</td>
<td>+07 24 15</td>
<td>1.4 ± 0.7</td>
<td>6.5</td>
<td>3</td>
</tr>
<tr>
<td>24</td>
<td>10–100</td>
<td>00 30 47.6</td>
<td>+07 24 20</td>
<td>0.9 ± 0.4</td>
<td>6.6</td>
<td>3</td>
</tr>
</tbody>
</table>

Notes. The second column lists the analysed energy range. The 11 (24)-month analysis focuses on the high-energy flux $\phi_\nu(10–100 \text{ GeV})$ only. In all cases, the significance $S$ of the high-energy bin is well above 6$\sigma$.

The source model for the data analysis contains all 1FGL sources within the region of interest (ROI, radius 10\degree). Their parameters were taken as catalogued and we used the latest Galactic (\texttt{gll_iem_v02.fit}) and extragalactic (\texttt{isotropic_iem_v02.txt}) diffuse background models. All parameters but those of 1FGL J0030.7+0724 were kept fixed. Furthermore, the catalogued power-law index of 1FGL J0030.7+0724 was used while fitting the flux between 10 and 100 GeV. Although the exposure of the 24-month data has almost doubled with respect to the catalogue, the use of the catalogued properties for sources within the ROI will not affect the analysis between 10 and 100 GeV. The three nearby sources, i.e., 1FGL J0022.5+0607, 1FGL J0030.4+0451, and 1FGL J0023.5+0930, are not only more than 2\degree away from 1FGL J0030.7+0724, but they are also not significantly detected between 10 and 100 GeV. Furthermore, visual inspection does not reveal any other relevant source within this nearby region.

The analysis of the 11-month data reproduces the catalogued values well (Table 5). After 11 (24) months, five (six) photons between 10 and 100 GeV have been detected within a radius of 0.5\degree around the nominal position, listed in Table 6. Except one, all photons are classified as class 4 events and are therefore very likely signal events. The Galactic foreground and the extragalactic background at the source position are negligible with respect to the signal, with an expected total number of background photons $N_{bg} = 0.6 (1.2)$ within the considered region of 0.5\degree. For comparison, the predicted number of signal events is $N_{sig} = 4.9 (5.8)$ after 11 (24) months. According to the 11-month data set, the (10–100 GeV) best-fit position shifts by about 2.5\degree. The small positional error of the sixth photon also accounts for the increase of the source’s positional uncertainty, see Fig. 8.

The average flux over the entire data set has decreased by a factor of roughly 1.5 with respect to the first 11 months (Table 5). To judge on the variability of 1FGL J0030.7+0724, its temporal photon distribution (Table 6) was tested for compatibility with a constant flux, using an unbinned Kolmogorov-Smirnov (KS) test (Press et al. 2007). The KS test is already valid for low photon counts, unlike the binned chi-square method used by the catalogue. The KS test confirms the null-hypothesis.
of non-variability with a probability of about 0.7 (0.5) for the 11 (24) months data set. The varying exposure on the region was taken into account by examining the photon distribution of the bright pulsar nearby (1FGL J0030.4+0451).

The analysis of the (intrinsic) spatial extent of the source is based upon a likelihood-ratio test, using all photons listed in Table 6. The (corresponding) statistical measure is given by $L(\theta_0) = -2 \sum_{i=1}^{N} \ln[p_{\text{det}}(x_i; \theta_0) + b]$, where $p_{\text{det}}(x; \theta_0)$ is the probability distribution function for a photon detected at $x$, $\theta_0$ denotes the best-fit position (Table 5), and $b$ incorporates the flat background. For a spatially extended 1FGL J0030.7+0724, the number of photons detected $p_{\text{det}} = \text{psf} \times \text{int}$, the (two dimensional) convolution of the Fermi-LAT PSF (P6_v3, diffuse class) with the intensity profile of the emitter. In the subhalo case, the intensity profile follows the line-of-sight integral of the squared NFW profile (Eq. (1)). The quantity $\Delta L = L - L_{\text{min}}$ follows a chi-square distribution with one degree of freedom, with additional terms of the order of $1/N^{1/2}$, which are important for a small number of counts (Wilks 1938; Cash 1979). The likelihood is minimised ($L_{\text{min}}$) for the intrinsic extension parameter fitting the photon distribution best. Examining the 11-month data set, the test shows the source to be consistent with a point source, implying that the intrinsic extent is smaller than the (average) PSF (about 0.15°). The 24-month data favour a moderate extent $\theta_0 \approx 0.14^{+0.20}_{-0.12}$ deg, which is, however, not significant. Upper limits on the intrinsic parameter are $\theta_0 < 0.54(0.72)$ deg at 95% confidence level, derived from the 11 (24) months data set. Since the low statistics affect the chi-square distribution, note that the confidence level is not precisely defined (Cash 1979; Mattax et al. 1996). Furthermore, we point out that the PSF of Fermi-LAT (P6_v3) may be underestimated and changes will have an impact on the fitted extent.

5.1.3. Swift-XRT data

The field was successfully proposed for observation with the X-ray telescope (XRT, 0.2–10 keV) onboard the Swift satellite (Gehrels et al. 2004, 2005). The observations

\[ \text{http://fermi.gsfc.nasa.gov/ssc/data/analysis/LAT} \]

were carried out on 10 November, 2010, between 00:23:46 and 19:52:56 UT with a total effective exposure of 10.1 ks. Observations with the XRT were performed in photon-counting (PC) mode. The XRT data were calibrated and selected with standard screening criteria (xrtpipeline), using the HEAsoft 6.10 package for data reduction with the current version of calibration files available (release 2010-09-30). For the analysis, events with grades 0–12 (Burrows et al. 2005) were used. The spectral analysis was carried out with xspec (12.6.0, Arnaud 1996), using the PC grade 0–12 response matrix swxpc0to12_126_200709011011.rmf with the ancillary response function generated by xrtmkarf for PSF correction and the position of the source considered. The on-source region was selected to contain about 90% of the PSF ($\pm 47''$). For background subtraction, an off-source region with radius of about 4'' was used. To ensure a spectral fit of sufficient quality, the spectra were binned to a minimum of 10 events per bin (with grapha). Owing to the low statistics accumulated, the C-statistic was used for spectral fitting.

In the field-of-view (FoV) of XRT, seven new X-ray sources were discovered with a probability of being background fluctuations smaller than $10^{-6}$. We show them in Fig. 8. Their positional properties, measured flux, and the flux corrected for photoelectric absorption between 0.2 and 2 keV are listed in Table 7. The spectra of the two brightest sources are well-fitted by an absorption corrected power-law model, fixing the hydrogen column density $N_H$ to the nominal Galactic value. The power-law index for the faint sources was fixed to 2.0. The Galactic hydrogen column density was obtained from the LAB HI survey (Kalberla et al. 2005) for the corresponding celestial positions.

Associations. The newly discovered X-ray sources were studied for possible associations in other accessible wavelengths. Multi-wavelength surveys covering the region are the NVSS in the radio, the 2MASS in the infrared, and the USNO B1.0 and SDSS DR7 catalogues for the optical band. For every Swift source we found at least one SDSS source be positionally coincident (Table 8), with apparent magnitudes between 21$m$ and 17$m$. Owing to insufficient sensitivity, the very faint SDSS sources have not been detected by USNO.

6. Discussion

6.1. An AGN origin of 1FGL J0030.7+0724

The γ-ray signal of 1FGL J0030.7+0724 can be explained by a conventional AGN. With respect to the unified scheme for the spectral energy distribution (SED) of AGN (namely Flat Spectrum Radio Quasars (FSRQs) and blazars), see, e.g., Donato et al. (2001), the hard spectral index of 1FGL J0030.7+0724 ($\Gamma \approx 1.7$) is compatible with a high-energy-peaked blazar (HBL). Within the updated positional uncertainty of 1FGL J0030.7+0724 derived from the 24-month data set, the most likely radio counterpart is NVSS J003119+072456 ($f_{1.4\text{GHz}} = (11.6 \pm 0.6)\text{mJy}$), which positionally coincides with the newly discovered hard X-ray source SWIFT J003119.8+072454 ($\Gamma \approx 1.6$). Note that corresponding to the notation of Table 7, the Swift source is flagged with a $G$ in Fig. 8. The energy flux observed between 0.2 and 2 keV is $\sim 2 \times 10^{-13}$ erg cm$^{-2}$ s$^{-1}$ (Table 7). Additionally, an optical counterpart of the radio and X-ray source is listed in the SDSS catalogue ($\sim 17.4$), see Table 8. In Fig. 9 we show an empirical model for the average SED of HBLs, which is based on the bolometric luminosity distribution of FSRQs and

Fig. 8. Celestial region of 1FGL J0030.7+0724, 25' × 13'. The catalogued position is indicated by the red +, while the dashed red ellipse borders its uncertainty (95% c.l.). The black × marks the 24-month position, the solid black contour its uncertainty at 95% confidence. Photons detected by Swift-XRT (10.1 ks) are mapped by the back-image, which is smoothed with a Gaussian (7′). The positions of NVSS radio sources are given by the dark green boxes, the two NVSS sources discussed in the text are named. In this region, seven X-ray sources have been discovered, indicated by the blue arrows (see Table 7). Note that the boxes’ size is chosen arbitrarily.


Table 7. X-ray sources detected with the Swift-XRT.

<table>
<thead>
<tr>
<th>ID</th>
<th>Name</th>
<th>SWIFT J</th>
<th>$\sigma_0$</th>
<th>$S/N$</th>
<th>$N_H$</th>
<th>$\phi_0$</th>
<th>$\Gamma$</th>
<th>$f_{\text{obs}}$(0.2–2 keV)</th>
</tr>
</thead>
<tbody>
<tr>
<td>A</td>
<td>003000.3+072301</td>
<td>6</td>
<td>3.5</td>
<td>3.5$^{+1.1}_{-1.0}$</td>
<td>3.98</td>
<td>1.4 ± 0.4</td>
<td>2.0</td>
<td>5.2 ± 1.5</td>
</tr>
<tr>
<td>B</td>
<td>003017.8+072142</td>
<td>5</td>
<td>5.4</td>
<td>5.0$^{+3.0}_{-2.1}$</td>
<td>3.71</td>
<td>2.2$^{+0.6}_{-0.4}$</td>
<td>1.4 ± 0.3</td>
<td>6.7$^{+2.3}_{-1.8}$</td>
</tr>
<tr>
<td>C</td>
<td>003022.1+072623</td>
<td>6</td>
<td>3.1</td>
<td>1.7$^{+0.5}_{-0.4}$</td>
<td>3.10</td>
<td>0.6$^{+0.3}_{-0.2}$</td>
<td>2.0</td>
<td>2.7$^{+1.1}_{-0.9}$</td>
</tr>
<tr>
<td>D</td>
<td>003030.0+072013</td>
<td>5</td>
<td>5.1</td>
<td>5.2$^{+2.4}_{-2.0}$</td>
<td>3.71</td>
<td>2.0$^{+0.6}_{-0.4}$</td>
<td>2.0</td>
<td>7.4$^{+1.2}_{-0.5}$</td>
</tr>
<tr>
<td>E</td>
<td>003049.8+072316</td>
<td>5</td>
<td>3.0</td>
<td>3.1$^{+1.0}_{-0.9}$</td>
<td>3.10</td>
<td>1.2$^{+0.4}_{-0.3}$</td>
<td>2.0</td>
<td>4.4$^{+1.5}_{-1.1}$</td>
</tr>
<tr>
<td>F</td>
<td>003054.9+072328</td>
<td>6</td>
<td>2.8</td>
<td>2.0$^{+0.8}_{-0.6}$</td>
<td>3.10</td>
<td>0.8$^{+0.4}_{-0.3}$</td>
<td>2.0</td>
<td>3.0$^{+1.5}_{-1.1}$</td>
</tr>
<tr>
<td>G</td>
<td>003119.8+072454</td>
<td>5</td>
<td>6.5</td>
<td>15.9$^{+4.5}_{-3.0}$</td>
<td>3.10</td>
<td>6.5$^{+1.1}_{-0.9}$</td>
<td>1.6 ± 0.3</td>
<td>20.7$^{+8.8}_{-4.7}$</td>
</tr>
</tbody>
</table>

Notes. The FoV is centred on (RA, Dec) = (7.6315, 7.4211) deg with a radius of 13’. We give an internal ID, the position (SWIFT JHHMMSS.x±DDMMSS) and its corresponding error at 90% confidence level $\sigma_0$ (determined with xrtcentroid), and the signal-to-noise ratio $S/N$ (Ximage) of the observed flux $f_{\text{obs}}$. If constraining, a power-law model corrected for photoelectric absorption was fitted to the spectrum, $dy/dE = \phi_0 (E/\text{keV})^{-\Gamma}$. The hydrogen column density $N_H$ was fixed during the fit. The unabsorbed flux $f_{\text{obs}}$ was derived from the power-law fit. Due to a low $S/N$ a two-parameter power-law fit is not constraining. The fluxes were derived assuming the index $\Gamma = 2.0$.

Table 8. Likely counterparts of the X-ray sources listed in Table 7.

<table>
<thead>
<tr>
<th>ID</th>
<th>Name</th>
<th>2MASS Name</th>
<th>NVSS DR7 Name</th>
<th>R</th>
<th>$\gamma$</th>
<th>Type</th>
</tr>
</thead>
<tbody>
<tr>
<td>A</td>
<td>003000.3+072255</td>
<td>47(2)</td>
<td>–</td>
<td>–</td>
<td>0.0</td>
<td>6</td>
</tr>
<tr>
<td>B</td>
<td>003000.3+072255</td>
<td>–</td>
<td>003000.5315</td>
<td>20.0</td>
<td>0.0</td>
<td>6</td>
</tr>
<tr>
<td>C</td>
<td>003000.3+072255</td>
<td>–</td>
<td>003000.5428</td>
<td>20.4</td>
<td>0.0</td>
<td>6</td>
</tr>
<tr>
<td>D</td>
<td>003000.3+072255</td>
<td>–</td>
<td>003000.5481</td>
<td>18.9</td>
<td>0.0</td>
<td>6</td>
</tr>
<tr>
<td>E</td>
<td>003000.3+072255</td>
<td>–</td>
<td>003000.5484</td>
<td>19.3</td>
<td>0.0</td>
<td>6</td>
</tr>
<tr>
<td>F</td>
<td>003000.3+072255</td>
<td>–</td>
<td>003000.5560</td>
<td>18.2</td>
<td>0.0</td>
<td>6</td>
</tr>
<tr>
<td>G</td>
<td>003000.3+072255</td>
<td>–</td>
<td>003000.5567</td>
<td>18.6</td>
<td>0.0</td>
<td>6</td>
</tr>
</tbody>
</table>

Notes. Scans ranging from radio (NVSS) to infrared (2MASS) and optical (USNO and SDSS) wavelength bands are given. The table lists the object’s name and the catalogued flux or apparent magnitude. Here, parentheses indicate the corresponding error on the last decimal. Based on photometric morphology, SDSS provides a separation between galaxy-like (3) and star-like objects (6), see Lupton et al. (2001). Frequency $\nu = 1.4 \text{ GHz}$. The column lists R2. If not available, R1 or B1 is given instead (see Monet et al. 2003, and references therein).

blazars (Fosati et al. 1997, 1998; Donato et al. 2001). The SED is normalised to the radio flux of NVSS J003119+072456 (at 5 GHz). For comparison, the spectral measurements of the optical and X-ray counterparts are presented as well. Within the observational errors and assuming temporal variability, the $\gamma$-ray spectrum of 1FGL J0030.7+0724 is consistent with the model prediction. Furthermore, the spectral index of the X-ray source agrees with an HBL, while its flux is fainter than predicted for an average-concentrated SHM subhalo with a corresponding distance of 4 kpc, while the upper limit on the extent is $\theta_{68} \leq 0.7''$, corresponding to $\theta_{68} \leq 0.3''$. As shown in Sect. 4, in realistic WIMP scenarios the high effective self-annihilation cross section required to explain the source with DM annihilation in a FHM subhalo is hardly compatible with current observational constraints (see Fig. 3 and Table 1). However, given the more realistic SHM, flux and extent of 1FGL J0030.7+0724 are consistent with a subhalo of mass between $10^8$ and $10^9 M_{\odot}$. Assuming a DM subhalo of $10^8 M_{\odot}$, the resulting distance would be $2.4^{+1.0}_{-0.7}$ kpc, given the concentration scatter of the SHM model. For a WIMP of 500 GeV annihilating to $b\bar{b}$, the required minimum effective enhancement is 7 for a high-concentrated SHM subhalo with a corresponding distance of 1.7 kpc, while it increases to 31 for an average-concentrated subhalo with a corresponding distance of 2.4 kpc. Note that $h(0.7'') \approx 1.4$. An even lower boost factor is required for a lighter WIMP of 150 GeV which predominantly annihilates to $\tau^+\tau^-$: 3 (13) for a high-concentrated (average-concentrated) SHM subhalo.

6.2. A DM subhalo origin of 1FGL J0030.7+0724

Without a clear indication for variability, it remains plausible that the $\gamma$-ray emission of 1FGL J0030.7+0724 originates from a DM subhalo. The analysis of the arrival times of the source photons (Sect. 5.1.2) is consistent with a temporally constant source of moderate spatial extent. The reconstructed high-energy flux within the statistical errors is $\phi_0(10–100 \text{ GeV}) \geq 5 \times 10^{-11} \text{ cm}^{-2} \text{s}^{-1}$, while the upper limit on the extent is $\theta_{68} \leq 0.7''$, corresponding to $\theta_{68} \leq 0.3''$. As shown in Sect. 4, in realistic WIMP scenarios the high effective self-annihilation cross section required to explain the source with DM annihilation in a FHM subhalo is hardly compatible with current observational constraints (see Fig. 3 and Table 1). However, given the more realistic SHM, flux and extent of 1FGL J0030.7+0724 are consistent with a subhalo of mass between $10^8$ and $10^9 M_{\odot}$. Assuming a DM subhalo of $10^8 M_{\odot}$, the resulting distance would be $2.4^{+1.0}_{-0.7}$ kpc, given the concentration scatter of the SHM model. For a WIMP of 500 GeV annihilating to $b\bar{b}$, the required minimum effective enhancement is 7 for a high-concentrated SHM subhalo with a corresponding distance of 1.7 kpc, while it increases to 31 for an average-concentrated subhalo with a corresponding distance of 2.4 kpc. Note that $h(0.7'') \approx 1.4$. An even lower boost factor is required for a lighter WIMP of 150 GeV which predominantly annihilates to $\tau^+\tau^-$: 3 (13) for a high-concentrated (average-concentrated) SHM subhalo.

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Further decrease of the necessary boost may be provided by sub-substructure and cuspy profiles (Sect. 4).

In addition to theoretical uncertainties on halo properties and their expected scatter (Sect. 2.1), observational uncertainties affect the distance and boost factor estimates. The uncertainties on the flux directly change the boost, while uncertainties on the most crucial measurement, the angular extent on the flux directly change the boost, while uncertainties on the effect the distance and boost factor estimates. The uncertainties on the X-ray flux at the nominal position of 1FGL J0030.7+0724 be-

Further decrease of the necessary boost may be provided by sub-substructure and cuspy profiles (Sect. 4).

In addition to theoretical uncertainties on halo properties and their expected scatter (Sect. 2.1), observational uncertainties affect the distance and boost factor estimates. The uncertainties on the flux directly change the boost, while uncertainties on the most crucial measurement, the angular extent, affect both the required boost and the distance estimate (Sect. 3). The discussed object 1FGL J0030.7+0724 serves as an appropriate benchmark, because the corresponding uncertainties are representative for a typical DM subhalo source. The observational uncertainties are of similar magnitude as the theoretical ones.

6.3. Remarks and prospects for IACTs

The 28-month data of Fermi-LAT contains no additional photon detected around the nominal position. This lowers the probability of steadiness to $\sim 25\%$ and may indicate variability, which supports a BL Lac scenario. Vice versa, such a behaviour would also be anticipated by a selection bias: If the true flux is lower than the value found in the discovery data set, the discovery condition would only allow for the detection of sources where the photon number has been fluctuating upwards. Poisson fluctuations of this faint signal could have accounted for a detection with the LAT even if the true flux had remained below the detection sensitivity.

It is instructive to note that with regard to a definite identification of a counterpart (or ruling out a candidate) from observations in other wavelength regimes the limiting factor is the accuracy of the Fermi-LAT source position ($\mathcal{O}(5')$, cf., Table 4) and PSF. With just six detected photons, probably including one background photon, the source is close to the confusion limit. This situation can only be resolved by future instruments with much larger effective areas, such as the proposed CTA, which will probe deep into the expected population of subhaloes. The much larger number of photons would help to infer significantly improved source positions. Furthermore, for detecting a spectral cut-off and in case of heavy WIMPs ($m_\chi > 1\, \text{TeV}$), observations in the VHE range with IACTs are favoured. For the particular DM scenarios proposed for 1FGL J0030.7+0724, fluxes anticipated in the energy ranges accessible for MAGIC and H.E.S.S., $\phi(>50\, \text{GeV})$ and $\phi(>300\, \text{GeV})$, respectively, are listed in Table 9 (given by $\phi(>E) \propto N_c(>E)/N_c(100\, \text{GeV})$, see Eq. (8)). Additionally, flux estimates for WIMPs with $m_\chi = 1\, \text{TeV}$ were derived. Note that the required effective cross sections (see Sect. 6.2) increase by a factor of 2.3, because $\langle \sigma v \rangle_{\text{eff}} \propto m_\chi^2 N_c(10\, \text{GeV}\Rightarrow 10\, \text{GeV})^{-1}$. Also note that fluxes expected for VERITAS are comparable to those for H.E.S.S.

The low energy threshold of MAGIC leads to comparatively high integrated VHE fluxes for $m_\chi < 1\, \text{TeV}$. The flux prediction for MAGIC is of $\mathcal{O}(1\%)$ of the Crab Nebula’s for the favoured $\gamma^+\gamma^-$, $m_\chi = 150\, \text{GeV}$ and $\gamma^+\gamma^-$, $m_\chi = 1\, \text{TeV}$ model. With MAGIC, 50 h of observation are necessary to detect this source with more than 5$\sigma$. For comparison, predicted fluxes for H.E.S.S. are not higher than 0.2$\%$ Crab for these models, which requires a few hundred hours of observation (Aharonian et al. 2006c). We remark that advanced analysis methods improve the sensitivity of H.E.S.S. by a factor of 2 (de Naurois & Rolland 2009). In the near future, an additional telescope (H.E.S.S.-II) will lower the energy threshold of H.E.S.S. to about 25–50 GeV. For the corresponding flux level of 1 Crab, the required observation time for H.E.S.S.-II and MAGIC will be similar. Furthermore, the planned CTA observatory will be able to detect such a source in about 50 h (the CTA consortium 2010).

7. Summary and conclusions

Hierarchical structure formation predicts Milky Way-sized galaxies to host numerous DM subhaloes with masses ranging from $10^{10}$ down to a cut-off scale of $10^{-3}$–$10^{-11} M_\odot$. Given standard WIMP scenarios, e.g., motivated by supersymmetry, we have demonstrated that DM subhaloes are detectable with the currently operating $\gamma$-ray telescope Fermi-LAT. Based upon state-of-the-art models, detectable subhaloes would observationally appear as faint high-energy $\gamma$-ray sources between 10 and 100 GeV with a flux at the sensitivity level of Fermi-LAT ($\sim 10^{-10}\, \text{cm}^{-2}\, \text{s}^{-1}$) between 10 and 100 GeV for one
year). The observable γ-ray emission exhibits a moderate spatial extent below ~0.5◦. Subhaloes favoured for detection are massive (10^6–10^8 M⊙) at distances of O(kpc), while low-mass subhaloes are not detectable. Within the intrinsic halo-to-halo scatter, only a moderate enhancement of the self-annihilation subhaloes are not detectable. Within the intrinsic halo-to-halo mass ratio below ∼10^5, 3 between 3 and 12 is necessary (dependent on the WIMP model), which is consistent with current observational constraints. Increasing sensitivity for a data-taking period of five years will allow us to resolve subhaloes requiring a cross section enhanced by a factor between 1.3 and 5. Additional substructure within a subhalo may lower the required enhancement. Within statistics, one massive subhalo could be detectable with Fermi-LAT in one year and might appear in the first-year catalogue (1FGL), assuming a subhalo population predicted by numerical N-body simulations. Regarding the 1FGL, the high-energy flux (10–100 GeV) of subhalo candidates should be fainter than ~4 × 10^{-10} cm^{-2} s^{-1} (for the WIMP scenarios considered here).

Intensive searches for subhaloes in the 1FGL reveal twelve candidates, which are unassociated, non-variable, high-latitude sources detected above 10 GeV. The physical origin of the most promising object selected by lacking association, faintness, and spectral index, 1FGL J0030.7+0724, was investigated by analysing the 24-month data set of Fermi-LAT. With dedicated Swift-XRT observations (10.1 ks), seven X-ray sources were discovered around 1FGL J0030.7+0724. Located within the positional uncertainty of the γ-ray source, a radio source positionally coincides with a newly discovered X-ray source hints at a conventional HBL origin of 1FGL J0030.7+0724. However, owing to a large positional uncertainty and the lacking detection of variability, the possibility of a dark nature remains. The measured high-energy flux and spatial extent of the source is compatible with a DM subhalo between 10^7 and 10^8 M⊙ in a distance of about 2 kpc, driven by a 500 (150) GeV WIMP self-annihilating to b̅b (τ+τ−). In this case, the required enhancement of (σv)0 is 7 (3) within the intrinsic scatter of the subhalo model, given a subhalo of 10^6 M⊙.

Establishing the – probably more likely – HBL scenario of 1FGL J0030.7+0724 requires a significant detection of γ-ray variability and a confirmation of the radio as well as X-ray counterparts. Vice versa, a steady γ-ray flux with a spectral shape predicted by self-annihilating WIMPs would hint at a DM nature of the object. This validates the necessity of additional intense and long multi-wavelength observations. In particular, IACTs offer a unique capability to reduce the positional uncertainty of faint LAT sources and to detect a spectral cut-off in the VHE range. A detection of the subhalo candidate 1FGL J0030.7+0724 may be possible with telescope systems like H.E.S.S.-II, MAGIC, and CTA within about 50 h of observation.

Our results encourage the search for more subhalo candidates in current and upcoming (very-)high-energy data releases. However, even in optimistic scenarios the expected number of LAT-detectable subhaloes is small. Furthermore, a longer exposure time – while certainly helpful with regard to the single candidate discussed in this work – will not necessarily remedy the general problem of the γ-ray photon count that limits the positional accuracy and therefore the chance of identifying counterparts. Given m_γ < 1 TeV, acquiring a sufficiently large number of detections which may solve the subhalo problem requires higher sensitivity in the high-energy range.

Acknowledgements. We kindly acknowledge helpful discussions with our colleagues Katharina Born, Torsten Bringmann, Wilfried Buchmüller, Frederike Jäger, Andrei Lobanov, and Martin Rau. We kindly acknowledge the Swift PI Neil Gehrels and his team for the prompt response to our ToO request and the corresponding observations. The help of the Fermi HelpDesk is kindly acknowledged. We kindly thank the anonymous referee for useful comments.

This work was supported through the collaborative research center (SFB) 676 “Particles, Strings, and the Early Universe” at the University of Hamburg. This publication makes use of data products from the Two Micron All Sky Survey, which is a joint project of the University of Massachusetts and the Infrared Processing and Analysis Center/California Institute of Technology, funded by the National Aeronautics and Space Administration and the National Science Foundation. This publication makes use of data from the Sloan Digital Sky Survey (SDSS). Funding for the SDSS and SDSS-II has been provided by the Alfred P. Sloan Foundation, the Participating Institutions, the National Science Foundation, the US Department of Energy, the National Aeronautics and Space Administration, the Japanese Monbukagakusho, the Max Planck Society, and the Higher Education Funding Council for England. The SDSS Web Site is http://www.sdss.org/. The SDSS is managed by the Astrophysical Research Consortium for the Participating Institutions. The Participating Institutions are the American Museum of Natural History, Astrophysical Institute Potsdam, University of Basel, University of Cambridge, Case Western Reserve University, University of Chicago, Drexel University, Fermilab, the Institute for Advanced Study, the Japan Participation Group, Johns Hopkins University, the Joint Institute for Nuclear Astrophysics, the Kavli Institute for Particle Astrophysics and Cosmology, the Korean Scientist Group, the Chinese Academy of Sciences (LAMOST), Los Alamos National Laboratory, the Max-Planck-Institute for Astronomy (MPIA), the Max-Planck-Institute for Astrophysics (MPA), New Mexico State University, Ohio State University, University of Pittsburgh, University of Portsmouth, Princeton University, the United States Naval Observatory, and the University of Washington. This research has made use of the NASA/IPAC Extragalactic Database (NED) which is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration.

Appendix A: Concentration of Aquarius subhaloes

The Aquarius simulation provides results on the profile parameters of resolved subhaloes, taking tidal interaction into account (Springel et al. 2008a). These results are used to derive the distance-averaged virial concentration of subhaloes to confront it with the toy-model predictions used here.

Following up on Eq. (1), the tidal concentration c(t) ≡ R_t/r_s is introduced (cf., Ando 2009), where R_t denotes the tidal and therefore physical radius of a subhalo. For an NFW-type mass density profile, c(t) = exp[(W(e−r) + a) − 1], a ≡ 1 + M/(4πρ_s r_s^2), where W(x) denotes Lambert’s W-function and M_t the tidal subhalo mass. In numerical simulations, the directly “observable” quantities of (sub)haloes are related to the dynamics of the halo system, including the maximum velocity V_max and the distance r_max where V_max is reached. To recover the canonical parameters r_s and ρ_s, related to the density profile, we use approximate relations 2[V_max/(H_0 max)]^2 = 5.80 × 10^4 [M/(10^8 M⊙)]^{0.18} and M_t ≈ 3.37 × 10^4 [V_max/(10 km s^{-1})]^2 M⊙, fitting the results of the simulation (Springel et al. 2008a; Ando 2009). Given analytical equations between (r_max, V_max) and (r_s, ρ_s) for the NFW profile (e.g., Eq. (8) in Kuhlen et al. 2008), this yields r_s(Aq) = 0.094 [M/(10^8 M⊙)]^{0.38} kpc and ρ_s(Aq) = 9.6 × 10^4 [M/(10^8 M⊙)]^{0.18}. Therefore, the tidal concentration c(Aq) is determined via a ≡ c ≈ 1 + 0.66 [M/(10^8 M⊙)]^{0.04}, which is valid for masses above the resolution limit of the simulation, M_t ≥ 3.2 × 10^4 M⊙.

The virial concentration of Aquarius subhaloes is given by c(Aq)(M_{vir}) = [3 M_{vir}/(4πΔ_c ρ_{vir})]^{1/3} r_s(Aq)(M_{vir}), where the characteristic radius as function of the virial subhalo mass is obtained from an empirical relation mapping M_{vir} to M_{t}. Based on c(Aq) and assuming the FHM virial concentration-to-mass relation (Eq. (2)), the relative tidal mass is M_t/M_{vir} ≈ f(c(Aq)(M_{vir}))/f(c). Hereby, we
have assumed that tidal effects on inner subhalo parts are negligible: \( \rho_s(M_{\text{vir}}) r_s(M_{\text{vir}})^3 \approx \rho_s(M_{\text{vir}}) r_s(M_{\text{vir}})^3 \). With \( f(c_{\text{vir}}^2) = c_{\text{vir}}^2 - 1 \), the distance-averaged \( M_s-M_{\text{vir}} \) relation is

\[
M_s(M_{\text{vir}}) \approx \left( \frac{712.6 \text{ kpc}^{-3}}{4 \pi \rho_{\text{vir}}^\text{PSF}} \right)^{1.04} M_{\text{vir}}^{1.04}.
\]

For massive subhaloes (\( \geq 10^4 M_{\odot} \)), \( f(c_{\text{vir}}^2) \) is well fit by a power law, \( f(c_{\text{vir}}^2) \) as function of the intrinsic angular extent \( \theta_s \). The intensity profile was assumed to follow the line-of-sight integral over the (squared) NFW profile as function of the intrinsic angular extent \( \theta_s \).

\[
\sigma_{\text{eff}}(M_{\text{vir}}) = 46.8 \left( \frac{M_{\text{vir}}}{10^9 M_{\odot}} \right)^{-0.07}. \tag{A.2}
\]

for \( M_{\text{vir}} \in [1.5 \times 10^5; \sim 10^{10}] M_{\odot} \).

**Appendix B: Moderately extended Fermi sources**

For \( \gamma \)-raycatalogues as for 1FGL, instrument data have been analysed assuming sources to be point-like. Given that detectable subhaloes would appear as moderately extended according to the PSF of Fermi-LAT (see Sect. 4.1.1, \( \sigma_{\text{PSF}} \approx 0.15^\circ \) for \( E = 10 \text{ GeV} \)), we investigated the effect of the 1FGL point-source analysis framework on extended sources.

To study the high-energy flux \( \phi_p(10-100 \text{ GeV}) \) reconstructed from the point-source analysis for a given intrinsic (subhalo) extent \( \theta_s \), a Monte-Carlo (MC) simulation study is restricted to the signal-dominated regime chosen to be \( \theta_s \approx 200 \text{ MW} \), where \( \theta_s \approx 1^\circ \) given the low background \( N_{\text{bg}} \). Since \( \theta_{\text{bg}} \approx 0.46^\circ \), this corresponds to \( \sim 3 \sigma_{\text{PSF}} \). Justified by the low background, all photons were treated as signal events.

The \( \theta_s \) dependence of the sample-averaged reconstructed flux \( \phi_p(10-100 \text{ GeV}) \) and the corresponding test statistic TS is shown in the two upper panels of Fig. B.1. For large \( \theta_s \), the probability of photons to be located far away from their central position increases. Therefore, both \( \phi_p \) and TS decrease because of a minor contribution of outer photons to the point-source region (defined by the PSF). For \( \theta_s \approx 1^\circ \), the average significance drops below the detection criterion (TS \( \geq 25 \)). Note that TS \( \geq 25 \) still holds for about 35% of the simulated samples.

In terms of Eq. (8), appropriate investigations of candidates provided by point-source catalogues is therefore admitted by a scaling \( h(\theta_s) \), which allows us to map the catalogued flux \( \phi_p \) to the true flux \( \phi \) of the entire source. The angular dependence of \( h \) is shown in the lower panel of Fig. B.1. Given \( \phi = h(\theta_s) \phi_p \), the factor was derived by defining \( h(0^\circ) = 1 \). Conservatively, the complete MC sample was used to derive \( h(\theta_s) \), including realisations with TS \( < 25 \). As expected in the signal-dominated regime, the increase of \( h \) with increasing \( \theta_s \) is comparatively slight, while it is fairly linear in the background-dominated regime. Note again that this result holds for sources similar to 1FGL J0030.7+0724 at high galactic latitudes only, while in general \( h(\theta_s) \) is well fit by

\[
\phi_p(10-100 \text{ GeV}) \approx 10^{-10} \text{ cm}^{-2} \text{ s}^{-1} \text{ TeV}^{-1}\text{ sr}^{-1}, \tag{B.1}
\]

and \( \sim 3 \sigma_{\text{PSF}} \) for the signal-dominated regime.

**Appendix C: Subhalo-induced diffuse flux**

In the following, the diffuse flux of the subhalo population is derived using a prescription by \( \text{Ando} \) (2009), which is extended to include the probability distribution of the concentration parameter \( c \) (see Eq. (4)). Numerical \( N \)-body simulations have demonstrated that the differential subhalo number density \( dN_{\text{sh}} = N(r,M) dM \) follows a power-law in subhalo mass \( M \). Following standard assumptions, the number density \( N(r,M) \) factorises such that \( N(r,M) \approx n_{\text{sh}}(r) \cdot M^{-\alpha} \), where \( \alpha = 1.9 \) and \( r \) is the distance to the host’s centre. In simulations, the spatial density distribution \( n_{\text{sh}}(r) \) is consistently found to be "anti-biased" and, e.g., \( n_{\text{sh}}(r) \approx r_{\text{Ein}}(r) \) (Springel et al. 2008a), where \( r_{\text{Ein}}(r) \) denotes the Einasto profile (Einasto 1965)

\[
r_{\text{Ein}}(r) \propto \exp \left\{ -\frac{2}{\alpha} \left[ \left( \frac{r}{r_{\text{Ein}}} \right)^{\alpha - 1} \right] \right\}. \tag{C.1}
\]

For a Milky Way-sized halo, the best-fit parameters for the subhaloes’ spatial distribution \( r_{\text{Ein}}(r) \) have been found to be \( \alpha = 0.68 \) and \( r_{\text{Ein}} \approx 0.81 c_{\text{MW}}^{\text{MW}} < 200 \text{ kpc} \) (Springel et al. 2008a), where \( c_{\text{MW}} \approx 15 \) (Catena & Ullio 2010). Using \( N(r,M) \) normalised

\[\text{Given the selection bias of the 1FGL catalogue, TS \( \geq 25 \), a more stringent deduction of } h(\theta_s) \text{ should include realisations with TS \( \geq 25 \) only. This lowers the effective scaling factor } h(\theta_s).\]
to represent a probability density function in $M$, the differential density is
\[
\frac{dn_M(r,M)}{dM} = \frac{n_M}{M_{\text{min}}^{\alpha-1}} \left( \frac{M}{M_{\text{min}}} \right)^{-\alpha}, \tag{C.2}
\]
where $M_{\text{min}} \ll M_{\text{max}}$ are the minimum and maximum mass of Galactic subhaloes, respectively. The normalisation of the subhalo number density $n_M(r)$ is chosen such that the fraction of the host’s mass distributed in subhaloes $f_{\text{sh}} \equiv M_{\text{sh}}/M_{\text{vir}} \approx 15\%$ for the cut-off scale $M_{\text{vir}} = 10^{15} M_{\odot}$, where $M_{\text{vir}} = (1.49 \pm 0.17) \times 10^{11} M_{\odot}$ (Catena & Ullio 2010). The chosen value of $f_{\text{sh}}$ is consistent with recent estimates $f_{\text{sh}} \approx 10-50\%$ (Diemand et al. 2005; Diemand & Moore 2011; Diemand et al. 2008b; Springel et al. 2008b).

The total mass contained in subhaloes is given by
\[
f_{\text{sh}} M_{\text{vir}} \approx 4\pi \int_0^{M_{\text{vir}}} dr r^2 \int_{M_{\text{max}}}^{M_{\text{vir}}} dM \frac{dn_M(r,M)}{dM}. \tag{C.3}
\]
Solving for an Einstain-type profile and $\alpha \neq 2$ yields
\[
n_M(r) = \frac{f_{\text{sh}} M_{\text{vir}}}{2\pi r^3_{\text{min}}} \left( \frac{2}{\alpha \Gamma} \right)^{\alpha-1} \Gamma \left[ \frac{3}{\alpha \Gamma} \left( \frac{r_{\text{vir}}}{r} \right)^{\alpha} \right]^{-1} \times \exp \left[ - \frac{2}{\alpha \Gamma} \left( \frac{r}{r_{\text{min}}} \right)^{\alpha} \right]. \tag{C.4}
\]
where $\Gamma(a,x)$ is the lower incomplete gamma function and $\Lambda = M_{\text{min}}/M_{\text{vir}}$.

The minimum mass $M_{\text{min}}$ of subhaloes is governed by the details of kinetic decoupling of WIMPs in the early Universe (Berezhiani et al. 2003, 2006; Green et al. 2005; Bringmann 2009). Depending upon the mass and composition of, e.g., the neutralino, a wide range of minimal subhalo masses has been considered in the literature, namely $M_{\text{min}} \in [10^{11}, 10^{13}] M_{\odot}$. Here, two benchmark cases are considered for $M_{\text{min}}$, i.e., $10^{10}$ and $10^6 M_{\odot}$ bracketing the 500 GeV neutralino scenario discussed by Bringmann (2009). The upper mass limit was fixed to $M_{\text{max}} = 10^{7} M_{\text{vir}} \approx 10^{10} M_{\odot}$. Results do not depend on the exact value chosen for $M_{\text{max}}$.

Using $\mathcal{L}(M,D)$ (Sects. 2.1 and 2.2), the average specific intensity from a subhalo population with extended, isotropic emissivity profiles is given towards a galactic radius $\tilde{h}$ by
\[
\frac{d\mathcal{L}}{d\tilde{h}} = \int_{M_{\text{min}}}^{M_{\text{max}}} \int_{L(M)} dM \int_{s(M)}^{s_{\text{max}}(M)} ds \frac{dn_M}{dM} (r, s, M) \int_{-\infty}^{\infty} dc P(c, \tilde{h}) \mathcal{L}(M, c, \tilde{h}) \frac{4\pi}{4\pi}. \tag{C.5}
\]
assumed that the spatial extent of each subhalo is much smaller than the scale on which the subhalo distribution changes significantly. The total photon rate $\mathcal{L}(\tilde{h})$ is given by Eq. (6) with the substitution $N_e \rightarrow E dN_e/DE$. Furthermore, $L(M)$ is required to be one-to-one. The galactocentric radius corresponding to the position $\tilde{h}$ is $r(\tilde{h}, \psi) = (R_0^2 + s^2 - 2R_0 s \cos \psi)^{1/2}$, where $\psi$ denotes the angle between $\tilde{n}$ and $R_0 (\cos \psi = \langle \tilde{n}, R_0 \rangle)$. Subhaloes bright enough to be detected as individual sources are not considered to contribute to the diffuse emission. Therefore, the lower limit of the line-of-sight integral is set by the detection criterion $\mathcal{L} \gtrsim 4\pi x_2 \phi_{\text{sh}},$ where $\phi_{\text{sh}}$ denotes the flux sensitivity for a detection in one year with Fermi-LAT, see Sect. 4. Since $R_0 \ll R_{\text{vir}}$ the upper bound of the $s$-integral $s_{\text{max}}(\tilde{h}) \approx R_{\text{vir}}^\phi$ is $R_{\text{vir}}^\phi$. The SHM photon rate is a function of both $M$ and $s$, and therefore $s$ also depends slightly on $s$. Conservatively, $s_{\text{max}}(M, s) = s_{\text{max}}(M, 3), s = R_{\text{vir}}^\phi$, revealing a lower bound on $s$.