A new SNR with TeV shell-type morphology: HESS J1731-347


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ABSTRACT

Aims. The recent discovery of the radio shell-type supernova remnant (SNR), G353.6-0.7, in spatial coincidence with the unidentified TeV source HESS J1731–347 has motivated further observations of the source with the High Energy Stereoscopic System (HESS) Cherenkov telescope array to test a possible association of the γ-ray emission with the SNR.

Methods. With a total of 59 h of observation, representing about four times the initial exposure available in the discovery paper of HESS J1731–347, the γ-ray morphology is investigated and compared with the radio morphology. An estimate of the distance is derived by comparing the interstellar absorption derived from X-rays and the one obtained from 12CO and HI observations.

Results. The deeper γ-ray observation of the source has revealed a large shell-type structure with similar position and extension (r ∼ 0.25°) as the radio SNR, thus confirming their association. By accounting for the HESS angular resolution and projection effects within a simple shell model, the radial profile is compatible with a thin, spatially unresolved, rim. Together with RX J1713.7–3946, RX J0852.0–4022 and SN 1006, HESS J1731–347 is now the fourth SNR with a significant shell morphology at TeV energies. The derived lower limit on the distance of the SNR of ∼0.06° for the analysis presented in this paper). The largest number of conclusive identifications so far can be attributed to pulsar wind nebulae (PWNe) as presented in e.g. Gallant et al. (2008). Recently, a new radio SNR, catalogued as G353.6-0.7, was discovered by Tian et al. (2008) to be unidentified to date (e.g. Aharonian et al. 2008). Most of the sources are extended beyond the point spread function (PSF) of the HESS experiment (−0.06°) for the analysis presented in this paper).

Key words. astroparticle physics – ISM: supernova remnants – cosmic rays

1. Introduction

In the survey of the Galactic plane carried out by the HESS experiment, many sources emitting at TeV energies remain 

unidentified to date (e.g. Aharonian et al. 2008). Most of the sources are extended beyond the point spread function (PSF) of the HESS experiment (−0.06°) for the analysis presented in this paper). The largest number of conclusive identifications so far can be attributed to pulsar wind nebulae (PWNe) as presented in e.g. Gallant et al. (2008). Recently, a new radio SNR, catalogued as G353.6-0.7, was discovered by Tian et al. (2008) to be
in spatial coincidence with HESS J1731−347, one of the unidentified sources presented in Aharonian et al. (2008). The diameter of the radio shell is nearly 0.5\degree which allows, given the brightness of the source and the HESS angular resolution of ∼0.06\degree, for a morphological comparison of the γ-ray source with the shell observed in radio. Moreover, at least up to the current date, no radio pulsar or X-ray PWN candidate was found that might alternatively explain the TeV emission. This situation should be compared to other γ-ray sources like HESS J1813−178 or HESS J11640−465 (Aharonian et al. 2006), which are also in spatial coincidence with radio SNR shells. However, for these latter sources a morphological identification with the radio shells is not possible, and the emission can plausibly originate from a PWN seen in X-rays as discussed in Funk et al. (2007b) for HESS J11640−465 and in Funk et al. (2007a); Gotthelf & Halpern (2009) for HESS J1813−178.

Observations of the north-eastern part of HESS J1731−347 with the X-ray satellites XMM-Newton, Chandra, and Suzaku have confirmed an X-ray counterpart found in archival ROSAT data (presented in Aharonian et al. 2008; Tian et al. 2008). An X-ray shell partly matching the radio morphology was found and the spectral analysis has revealed that the X-ray emission is of synchrotron origin, indicating that the shock wave of the SNR has accelerated electrons up to TeV energies (Acero et al. 2009b; Tian et al. 2010). A compact (unresolved) X-ray source XMMU J173203.3−344518 (Halpern & Gotthelf 2010) was observed towards the geometrical centre of the remnant and has spectral properties reminiscent of central compact objects (CCOs) found in several other supernova shells (e.g. Pavlov et al. 2004). A search for pulsations using the EPIC PN cameras onboard XMM-Newton shows only marginal evidence of a 1 s period (Halpern & Gotthelf 2010).

Given the recent discovery of G353.6-0.7, little is known about its age and distance. Tian et al. (2008) suggested a distance of 3.2 ± 0.8 kpc assuming that the SNR is at the same distance as the Hr region G353.42-0.37.

Additional HESS observations, carried out since the discovery paper of HESS J1731−347 (Aharonian et al. 2008), allow to investigate the compatibility of the TeV source with the radio shell SNR G353.6-0.7. The observations and the data analysis are described in Sect. 2 and the morphological and spectral results in Sect. 3. The multi-wavelength counterparts of HESS J1731−347 are described in Sect. 4 and a general discussion is presented in Sect. 5.

2. HESS observations and analysis methods

HESS is an array of four identical imaging atmospheric Cherenkov telescopes (IACTs) located in the Khomas Highland of Namibia 1800 m above sea level (Bernlöff et al. 2003). The survey of the Galactic plane by the HESS collaboration has led to the discovery of the γ-ray source HESS J1731−347, presented as an unidentified extended source in Aharonian et al. (2008). In this first data set, 14 h of observation time were available. Additional dedicated observations were carried out in July 2007 and in July and August 2009 with zenith angles ranging from 9\degree to 42\degree, the mean angle being 16.5\degree. The total HESS observation time for this target is 59 h after data quality cuts.

The data set was analyzed using the Model analysis (de Naurois & Rolland 2009) which exploits the full pixel information by comparing the recorded shower images with a pre-calculated shower model using log-likelihood minimization. In comparison with conventional analysis techniques, no cleaning or parametrization of the image shape is required and the full camera information is used. This method leads to a more accurate reconstruction and better background suppression than more conventional techniques and thus to an improved sensitivity.

Spectral and spatial analyses were carried out using a minimum image intensity of 60 photoelectrons (p.e.) resulting in an energy threshold of 240 GeV and an angular resolution of 0.06\degree (68\% containment radius). All results presented were cross-checked with a multivariate analysis (Ohm et al. 2009) using an independent calibration and gamma/hadron separation, which yielded consistent results. Unless otherwise quoted, the error bars in the following section are given at 1\sigma.

3. TeV γ-rays analysis results

The HESS excess map of the region of HESS J1731−347 is shown in Fig. 1 smoothed with a Gaussian of σ = 0.04\degree. For the background estimation in the image and in the morphology studies, the ring background method presented in Berge et al. (2007) was used. Because of the larger data set and the more sensitive reconstruction technique, the presented image is much more detailed than the one shown in the discovery paper (Aharonian et al. 2008). This reveals a complex region composed of a large and bright structure (HESS J1731−347), detected at 22\sigma, with a suggestive shell-like morphology.

A smaller and fainter structure named HESS J1729−345 (detected at 8\sigma) is also observed, the properties of which are presented separately in Sect. 3.3.

3.1. TeV energy morphology

To further test the hypothesis of a shell morphology for HESS J1731−347 and its association with the radio SNR, radial and azimuthal profiles in radio and γ-rays were extracted centered

![Image](309x561 to 529x767)
The PSF derived for this analysis (sphere of adjustable radius, projected on the sky and then folded sphere and a shell model. The first model is a uniformly emitting γ-ray excess profile restricted to radius $r \leq 0.3\degree$ and using the same center as in Fig. 2. The brightness distribution is compatible with a flat profile.

Fig. 2. The γ-ray excess and radio radial profiles are shown with blue crosses and red squares respectively. The best fits to the γ-ray data of a sphere and a shell model are overlaid. Both radial profiles are centered on the compact central object ($\alpha_{J2000} = 17^h32m03^s, \delta_{J2000} = -34^\circ45'18'\). The flux measured here is lower than what has been derived from a circular region centered on the CCO, illustrated by the large dashed circle ($r_{68\%} = 0.06\degree$). The shell model consists of a uniformly emitting shell of variable outer radius and thickness (defined as $r_{outer} = r_{inner}$) projected on the sky and then folded with the same PSF.

In the morphological test, the best fit statistically favors the shell model and the sphere model is ruled out at 3.9σ (90% confidence level).

To compare the TeV morphology with the shell seen in radio, the radio continuum map from the ATCA southern Galactic plane survey (SGPS) (Haverkorn et al. 2006) was smoothed to match the HESS spatial resolution and a radial profile was extracted (excluding point sources). The radio profile was then scaled by a normalization factor calculated as the ratio of the total number of excess γ-rays over the total radio flux on the whole remnant. The resulting profiles, presented in Fig. 2, show an extended emission in γ-rays similar to that seen in radio.

In contrast with RX J1713.7−3946 which is brighter in the North-West and SN 1006 that exhibits a bipolar morphology, the azimuthal profile of HESS J1731−347 (see Fig. 3) integrated for $r \leq 0.3\degree$ shows no significant deviation from a flat profile ($\chi^2$/d.o.f. = 8.8/9).

The energy spectrum of the SNR was obtained by means of a forward-folding maximum likelihood fit (Piron et al. 2001) from a circular region centered on the CCO, illustrated by the large dashed circle ($r = 0.3\degree$) in Fig. 1, chosen to fully enclose the emission of the remnant. The background is estimated using the multiple reflected-regions technique where background events are selected from regions of the same size and shape as the source region and at equal angular distance from the observation position (Berge et al. 2007). The resulting spectrum, shown in Fig. 4, is well described by a power-law model (equivalent $\chi^2$/d.o.f. = 27.7/35) defined as $dN/dE = N_0(E/E_0)^{-\Gamma}$ where $E_0$ is the decorrelation energy (energy at which the correlation between the slope and the normalization vanishes). The best fit parameters, listed in Table 1, result in an integrated 1–10 TeV energy flux of $(6.91 \pm 0.75_{stat} \pm 1.38_{syst}) \times 10^{-12}$ erg cm$^{-2}$ s$^{-1}$.

3.2. Spectral results

### Table 1

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\alpha_{J2000}$</td>
<td>17h32m03s</td>
</tr>
<tr>
<td>$\delta_{J2000}$</td>
<td>-34°45'18&quot;</td>
</tr>
<tr>
<td>$E_0$</td>
<td>$(6.91 \pm 0.75_{stat} \pm 1.38_{syst}) \times 10^{-12}$ erg cm$^{-2}$ s$^{-1}$</td>
</tr>
</tbody>
</table>

1 Position angle 0° corresponds to North and 90° to East.
Table 1. Best fit spectral parameters obtained for different extraction regions in HESS J1731–347.

<table>
<thead>
<tr>
<th>Region</th>
<th>Photon index $\Gamma$</th>
<th>Decorrelation energy $E_0$ (TeV)</th>
<th>Normalization $N_0$ (10$^{-12}$ cm$^{-2}$ s$^{-1}$ TeV$^{-1}$)</th>
<th>1–10 TeV integrated flux (10$^{-12}$ erg cm$^{-2}$ s$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>HESS J1731–347</td>
<td>2.32 ± 0.06$_{\text{stat}}$ 0.783</td>
<td>4.67 ± 0.19$_{\text{stat}}$</td>
<td>0.91 ± 0.73$_{\text{stat}}$</td>
<td></td>
</tr>
<tr>
<td>sub-region of HESS J1731–347$^a$</td>
<td>2.34 ± 0.09$_{\text{stat}}$ 0.780</td>
<td>1.41 ± 0.11$_{\text{stat}}$</td>
<td>2.02 ± 0.36$_{\text{stat}}$</td>
<td></td>
</tr>
<tr>
<td>HESS J1729–345</td>
<td>2.24 ± 0.15$_{\text{stat}}$ 0.861</td>
<td>0.44 ± 0.07$_{\text{stat}}$</td>
<td>0.88 ± 0.29$_{\text{stat}}$</td>
<td></td>
</tr>
</tbody>
</table>

Notes. The model used is a power-law of the form $dN/dE = N_0(E/E_0)^{-\Gamma}$. The systematic errors are conservatively estimated to be ±0.2 on the photon index and 20% on the flux. ($^a$) A spectral analysis corresponding to the FoV of the XMM-Newton data (see Fig. 5, center) has been carried out in order to build a SED.

Initially in Aharonian et al. (2008): (16.2 ± 3.6$_{\text{stat}}$ ± 3.2$_{\text{sys}}$) × 10$^{-12}$ erg cm$^{-2}$ s$^{-1}$ in the same energy band. However, the region of extraction in the discovery paper was much larger ($r = 0.6^\circ$ versus $r = 0.3^\circ$ in this paper), including HESS J1729–345 and possibly some surrounding diffuse emission. A cross-check to derive the flux from the SNR only using the same data set as used in Aharonian et al. (2008) and following the original analysis method gave results consistent with the complete data set presented here thus confirming that the flux difference was mainly due to the choice of the integration region. A power-law model with an exponential cutoff was also tested which did not improve the quality of the fit (equivalent $\chi^2$/d.o.f. = 24.0/34).

3.3. HESS J1729–345

A γ-ray excess of TeV emission was found at the best fit position $\alpha_{2000} = 17^h32^m35.5^s$, $\delta_{2000} = -34^d32^m22^s$ with a statistical error of 0.03$^\circ$ and the source was therefore labeled HESS J1729–345.

The source is extended beyond the size of the PSF (Gaussian width $\sigma = 0.12^\circ$ ± 0.03$^\circ$) and the region used to derive the spectral parameters is shown by the small dashed circle ($r = 0.14^\circ$) in Fig. 1. The spectrum obtained is well modeled by a power-law model (see Fig. 4) and the best fit parameters are listed in Table 1. The integrated flux in the 1–10 TeV energy band is (0.88 ± 0.29$_{\text{stat}}$ ± 0.18$_{\text{sys}}$) × 10$^{-12}$ erg cm$^{-2}$ s$^{-1}$.

4. Multi-wavelength counterparts

One of the interesting characteristics of HESS J1731–347 is that non-thermal emission is clearly identified in radio, X-rays and at TeV energies. In X-rays however, the access to the spectral properties is limited to a subregion of the SNR as the coverage with the XMM-Newton, Chandra and Suzaku satellites is only partial, and the statistics in the ROSAT All Sky Survey data are too low. In order to study the spectral energy distribution (SED) of the source, the radio flux and the TeV spectral properties were extracted only from the region observed in X-rays (see region definition in Fig. 5, right). The multi-wavelength counterparts of HESS J1729–345 are discussed later in Sect. 4.5.

4.1. Radio continuum

The shell observed in radio is spatially coincident with the γ-ray shell and has a similar extent (see radial profile in Fig. 2). The flux obtained (excluding point sources) from the SGPS ATCA data in the region observed by the XMM-Newton pointing is 0.8 ± 0.3 Jy at 1420 MHz. The total radio flux for the SNR measured by Tian et al. (2008) is of 2.2 ± 0.9 Jy. The compact HII region (G353.42-0.37) located to the West of the remnant at a distance of 3.2 ± 0.8 kpc (Tian et al. 2008) is indicated in Fig. 5 (left).

4.2. X-rays

In order to derive spectral information from the X-ray emission from the remnant, the XMM-Newton pointing obtained as a follow up of the HESS source (ObsId: 0405680201; PI: G. Pühlhofer) was analyzed. To clean the proton flare...
contamination during the observation, a histogram of the 10–12 keV count rates of each camera was built. A Gaussian fit was then performed in order to remove time intervals where the count rates were beyond 3σ from the mean value (Pratt & Arnaud 2002). The remaining exposure time after flare screening is 22 ks out of the 25 ks of observation for MOS and 15 ks for PN. For the image generation, the instrumental background was derived from the compilation of blank sky observations by Carter & Read (2007) and renormalized in the 10–12 keV band over the whole FoV. The image resulting from the combination of the two MOS instruments is presented in Fig. 5 (middle). For this mosaic, the data from the PN instrument were not used because of straylight contamination to the North-East (photons singly reflected by the mirrors) from a bright X-ray source located outside the FoV. This results in some spurious arc features near the border of the FoV in the North-East.

The X-ray emission is characterized by extended emission which is concentrated in arc-like features, similar to broken shell seen from many shell-type SNRs. Some of the arcs partly coincide with the radio and γ-ray shell (see Fig. 5). Some of the structures could hint at an additional, smaller shell, but might also come from irregular SNR expansion in an inhomogeneous and/or dense medium (Blondin et al. 2001). A double-shell structure is also observed in RX J1713.7–3946 in X-rays (Lazendic et al. 2004; Cassam-Chenaï et al. 2004; Acero et al. 2009a).

The spectral analysis of the diffuse X-ray emission was carried out using the Extended Source Analysis Software (ESAS) provided in the XMM-Newton Science Analysis System (SAS v9.0) to model the particle and instrumental backgrounds. The error bars in this section are quoted at 90% level confidence. For this analysis, the three instruments PN+MOS1+MOS2 were used and the regions were selected to avoid the straylight features.

The spectrum derived from the region covered by the FoV of XMM-Newton that is used for the SED is shown in Fig. 6. The emission is well represented by an absorbed power-law model and no emission lines were found (see also Tian et al. 2010). The best fit parameters obtained from a joint fit of MOS1, MOS2 and PN spectra are that is used for the SED is shown in Fig. 6. The spectrum derived from the region covered by the FoV of XMM-Newton that is used for the SED is shown in Fig. 6. The comparison of the absorption along the line of sight derived from X-ray data, 12CO and HI observations can be used to constrain the distance to the SNR. The velocity spectra of the 12CO emission (using data from the CfA survey, Dame et al. 2001) and the HI emission (using data from the SGPS survey, Haverkorn et al. 2006) derived from the region of highest X-ray absorption (1J2000 = 17h31m43.0s, δJ2000 = −34◦34′58″) are shown in Fig. 8 (bottom).

In order to derive a lower limit on the integration distance required to match the Nh derived from X-rays, all the material is assumed to be at the near distance allowed by the Galactic rotation curve. Under this hypothesis, the cumulative absorption column derived from the atomic and molecular hydrogen shown in Fig. 8 (top) is similar to the one observed in X-rays, Nh = (2.23 ± 0.21) × 1022 cm−2 when integrating up to a radial velocity relative to the local standard of rest (LSR) of ~25 km s−1. The CO-to-H2 mass conversion factor and the HI brightness temperature to column density used are respectively of 1.8 × 1020 cm−2 K−1 km−1 s−1 (Dame et al. 2001) and 1.82 × 1018 cm−2 K−1 km−1 s (Dickey & Lockman 1990).

When integrating up to the same velocity, the map of Nh derived from the atomic and molecular hydrogen shown in Fig. 7 (right) exhibits an increase of absorption towards the Galactic plane similar to that in the X-ray absorption map in Fig. 7 (left). The peak of 12CO emission at a LSR velocity of ~18 km s−1 is thus in the foreground of the SNR and is likely to be the cause of the gradient of absorption seen in X-rays. Using the circular Galactic rotation model of Hou et al. (2009) with a distance to the Galactic center of 8.0 kpc, the nearest distance corresponding to the LSR velocity of ~18 km s−1 is 3.2 kpc thus setting a lower limit for the distance of the remnant.

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Fig. 6. X-ray spectrum, using the PN camera, extracted from the sub-region of HESS J1731–347 shown in Fig. 5 (right). The non-thermal emission from the SNR is described by an absorbed power-law model (dashed line). The local astrophysical background, fitted to an off region outside the SNR, is modeled by two components (dotted lines). The low energy component is an APEC model (astrophysical plasma emission code, see http://hea-www.harvard.edu/APEC) representing the background from the Local Bubble and the high-energy component is an absorbed power-law representing the hard X-ray background (unresolved AGNs, cataclysmic variables, etc.). The residuals of the total model (SNR+local astrophysical background) are shown in the lower panel and the χ/s.d.o.f. is 1921/1569.

4.3. 12CO (J = 1–0) and HI

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2 http://xmm2.esac.esa.int/external/xmm_sw_cal/background/epic_esas.shtml
Fig. 7. Left: X-ray absorption map derived from a spectral fit to XMM-Newton data assuming an absorbed power-law model. A significant increase of $N_H$ towards the Galactic plane is observed. Right: absorption column map derived from atomic and molecular hydrogen when integrating over radial velocities from 0 km s$^{-1}$ to $-25$ km s$^{-1}$ (see Sect. 4.3 for more details). The Galactic plane is represented by the white dashed line. In both panels, the XMM-Newton field of view is represented by a dashed circle and the X-ray contours obtained from Fig. 5 (center) are overlaid.

Fig. 8. Top: cumulative absorbing column density (solid line) as a function of radial velocity at the position of highest X-ray absorption (see Sect. 4.3). The relative contributions from the atomic and molecular hydrogen are represented by the dashed and dash-dotted lines respectively. Middle: rotation curve towards the same direction as derived from the model of Galactic rotation of Hou et al. (2009). Bottom: $^{12}$CO (dashed line) and HI (dash-dotted line) spectra obtained the region highest X-ray absorption.

4.4. GeV $\gamma$-rays

In the Fermi-LAT first year catalog (Abdo et al. 2010) the source 1FGL J1729.1–3452c is found in the neighborhood of HESS J1731–347 as shown in Fig. 9. The Fermi source has an analysis flag that indicates that the source position moved beyond its 95% error ellipse when changing the model of diffuse emission. The $^{12}$CO map in Fig. 9 shows that the Fermi source is located near a small scale gas clump that could be not well represented in the diffuse emission model. The position of the source presented in the catalog is to be used with caution and is therefore possibly not incompatible with HESS J1729–345.

The Fermi source has a photon spectral slope of 2.26±0.08 and shows neither indication for spectral curvature nor time variability on a time scale of months (the catalog does not address shorter or longer time variations). This source is the closest Fermi detection near the newly discovered SNR and the flux derived in the Fermi catalog is used as an upper limit in the SED of the SNR in Fig. 10.

4.5. Multi-wavelength counterparts for HESS J1729–345

At radio wavelengths, the $\gamma$-ray contours of HESS J1729–345 lie near the HII region G353.381-0.114. Using HI radio recombination line data, the LSR velocity corresponding to this source is either $-54$ km s$^{-1}$ or $-82$ km s$^{-1}$ (Caswell & Haynes 1987). In the latter case this HII region could be associated with the molecular cloud observed around velocities of $\sim-80$ km s$^{-1}$ (see Fig. 9). At X-ray energies, no archival dedicated observations were found, and no emission is detected in the ROSAT all sky survey, probably due to the high absorption in the line of sight. As discussed in the previous section, a Fermi source is found to lie close to HESS J1729–345.

5. Discussion

The newly discovered SNR HESS J1731–347 is in several ways comparable to RX J1713.7–3946 and RX J0852.0–4622. Those objects are X-ray synchrotron emitters and exhibit no thermal emission lines. A CCO is also found within those three SNRs indicating a core collapse SN. Moreover at a distance of 3.2 kpc (see Sect. 4.3), the TeV luminosity of HESS J1731–347 in the 1–30 TeV energy band is $1.07 \times (d/3.2\text{ kpc})^2 \times 10^{34}$ erg s$^{-1}$ which is similar to the luminosity of RX J1713.7–3946 (the brightest TeV shell SNR detected until now), of $0.81 \times 10^{34}$ erg s$^{-1}$.
second and third $^{12}$CO peak shown in Fig. 8. The HESS significance
LSR velocity
are overlaid. The linear scale is in units of K km s$^{-1}$.

95% position confidence level contours presented in the 1 year catalog
contours from Fig. 5 together with the Galactic plane and the Fermi
using a distance of 1 kpc (Fukui et al. 2003; Cassam-Chenaï
Fig. 9. $^{12}$CO map of the vicinity of HESS J1731–347 integrated from
LSR velocity −13 km s$^{-1}$ to −25 km s$^{-1}$ (top) and from −75 km s$^{-1}$ to
−87 km s$^{-1}$ (bottom) respectively corresponding to the intervals of the
second and third $^{12}$CO peak shown in Fig. 8. The HESS significance
contours from Fig. 5 together with the Galactic plane and the Fermi
95% position confidence level contours presented in the 1 year catalog
are overlaid. The linear scale is in units of K km s$^{-1}$.

using a distance of 1 kpc (Fukui et al. 2003; Cassam-Chenaï
et al. 2004; Moriguchi et al. 2005) and slightly higher than
that of RX J0852.0−4622 with 0.65 × 10$^{52}$ erg s$^{-1}$ at a
distance of 0.75 kpc (Katsuda et al. 2008). A difference with RX J1713.7−3946 is that the flat γ-ray azimuthal profile of
HESS J1731–347 (see Fig. 3) suggests that the remnant is evolving
in a relatively uniform ambient medium and that it is not in
interaction with the cloud (shown in Fig. 9, top) used to derive a
lower limit to the distance of the SNR. This significantly differs
from the case of RX J1713.7−3946 which exhibits much brighter γ-ray emission in the North-West where the shock is thought to interact with denser material.

The distance used for the luminosity is derived from the
absorption in the foreground and provides only a lower limit of 3.2 kpc. However, as it is believed that supernova explosions are more likely to occur in the spiral arms of the Galaxy where the density of massive stars (i.e. SNR progenitors) is higher (Russell 2003; Hou et al. 2009), it is likely that HESS J1731–347 could be located within the Scutum-Crux or Norma
arms, which cross the line of sight at $l = 353.5^\circ$ at ≈3.0 and

~4.5 kpc respectively (Hou et al. 2009). The next arm in the
same line of sight is the Sagittarius arm lying at a distance of
12 kpc. This latter possibility for the location of the SNR would
lead to a much higher γ-ray luminosity, an order of magnitude
higher than RX J1713.7−3946. Also at such a distance, the physical
size of the remnant would exceed 50 pc, substantially larger
than other TeV shell SNRs whose physical size is ≤15 pc. As
a result it is reasonable to believe that the real distance to the
SNR should not be much larger than the derived lower limit of
3.2 kpc.

The radio flux and the X- and γ-ray spectra derived in Sect. 4
from the sub-region of HESS J1731–347 that is covered by the
FoV of XMM-Newton were combined in the SED presented in
Fig. 10. The X-ray data were corrected for the interstellar ab-
sorption with $N_H = 1.08 \times 10^{22}$ cm$^{-2}$. To model the SED, a
simple one-zone stationary model (presented in Acero et al. 2010)
was used. In this model, the spectrum of electrons and protons is
represented by a power-law of slope $c$ with exponential cutoffs
at energies $E_{c,e}$ and $E_{c,p}$ for the electrons and protons respec-
tively. For the modeling of the object, it is assumed that the mea-
sured multi-wavelength emission from the sub-region of HESS
J1731–347 is entirely coming from the SNR located at a distance
of 3.2 kpc. As this distance is only a lower limit, the total
energy of accelerated particles ($W_e$ and $W_p$) in the SNR should
also be viewed as lower limits.
In the pure leptonic scenario, the slope of the electrons is constrained by the radio and the X-ray synchrotron emission between 1.9 and 2.1 and the strength of the magnetic field required to reproduce the ratio of observed synchrotron and IC emission lies between 20 and 30 μG for 15 \( \leq E_{\text{el}} \leq 25 \) TeV. Although the relative ratio of radio, X- and TeV fluxes can be fairly well reproduced by this leptonic scenario, the model is inadequate to account for the X-ray and the γ-ray spectral slope as illustrated in Fig. 10 (top). The corresponding parameters for the latter model are summarized in Table 2.

This limitation no longer occurs in a scenario where the TeV emission is dominated by hadronic processes as the X- and γ-ray emission are now independent and both spectral slopes can be reproduced as shown in Fig. 10 (bottom). Moreover, the strength of the magnetic field can be increased as it is no longer fixed by the X/γ ratio. In order to reproduce the observed TeV flux, the total energy in high-energy protons (\( E \geq 1 \) GeV) assuming a spectral slope of 2.0 is \( W_p = 2 \times 10^{50} \text{ (n/1 cm}^{-3})^{-1} \text{ (d/3.2 kpc})^{-3} \text{ erg} \). It should be noted that this energy content only represents a sub-region of the SNR accounting for \( \sim 1/3 \) of the total TeV flux (see Table 1) implying that the total energy transferred to accelerated protons in the whole SNR is a substantial fraction of the energy available in the remnant for \( n \sim 1 \) cm\(^{-3} \). For this energetic reason, gas densities much below this value appear incompatible with the hadronic emission scenario.

Although it is not possible to measure the density of the ambient medium surrounding the SNR as no X-ray thermal emission is detected, an upper limit on the density can be derived. In order to do so, a thermal component, whose normalization is fixed for a given density using the method presented in Acero et al. (2007, Sect. 3.1), is added to the X-ray spectrum. The shocked ambient medium is assumed to be in a non-equilibrium ionization state with an ionization timescale parameter \( \tau = 10^8 \text{ cm}^{-3} \text{ s} \) and an electron plasma temperature \( kT_e = 1 \) keV. Such values are commonly observed in other young SNRs for which the X-ray emission of the shocked ambient medium has been studied as in e.g. RCW86 (see Vink et al. 2006). For the given parameters, the derived upper limit (90% confidence level) on the ambient medium density is \( 10^{-2} \) cm\(^{-3} \). In the case of a lower temperature (\( kT_e = 0.15 \) keV), an upper limit of 1 cm\(^{-3}\) is reached.

For a density of 1 cm\(^{-3}\) the corresponding shock speed and age of the SNR would be \( \sim 410 \text{ km s}^{-1} \) and 14,000 yrs in order to match a physical radius of \( R_{\text{shock}} = 15 \text{ pc} \) (0.27\(^{+} \) at 3.2 kpc), for a SN explosion of \( E_{\text{SN}} = 1 \times 10^{51} \text{ erg} \) with a mass of ejecta of 5 M\(_{\odot}\) using equations from Truelove & McKee (1999). However, this shock speed is an order of magnitude lower than what has been measured in other bright synchrotron emitting SNRs like SN 1006 (\( V_{\text{sh}} = 5000 \pm 400 \text{ km s}^{-1} \) at a distance of 2.2 kpc; Katsuda et al. 2009), RCW 86 (\( V_{\text{sh}} = 6000 \pm 3000 \text{ km s}^{-1} \); Helder et al. 2009), CasA (\( V_{\text{sh}} = 4900 \text{ km s}^{-1} \); Patnaude & Fesen 2009) or Tycho (\( V_{\text{sh}} = 3000 \pm 1000 \text{ km s}^{-1} \) at a distance of 2.3 kpc; Katsuda et al. 2010). As a rough estimate, the required density to reproduce a canonical shock speed of 3000 km s\(^{-1}\) using the aforementioned SN parameters is of the order of 0.01 cm\(^{-3}\) (compatible with the upper limit derived from the lack of thermal X-ray emission in the previous paragraph) for a corresponding age of \( \sim 2500 \) yrs.

To summarize, the presented static one-zone model suffers from limitations in both scenarios. In the leptonic case, the model allows to estimate the average B-field \(( \sim 25 \mu\text{G}) \) and the total energy in accelerated electrons present in the shell of the SNR but fails to reproduce the observed X-ray and γ-ray spectral slope. In the hadronic model, the high medium density required to reproduce the observed TeV flux is hardly compatible with the hydrodynamics of the SNR. More detailed models using non-linear diffuse shock acceleration theory have been developed (e.g. Zirakashvili & Aharonian 2010; Ellison et al. 2010) and would provide more accurate predictions than the simple model presented here. It should be noted that the considered model does not take into account evolution related to radiative cooling which could yield a steeper gamma-ray spectrum, in better agreement with the data. Also, the presented scenarios do not cover possible non-homogeneous surroundings such as wind bubble blown by the progenitor. Such detailed spectral and evolutionary modeling depends on many poorly known parameters and is therefore beyond the scope of the present discussion.

Concerning the source HESS J1729−345, detected in the vicinity of the SNR, the presented multi-wavelength data do not provide a clear understanding of the nature of the object. The closest structures located near the γ-ray emission are the HII region G353.381-0.114 (seen in the radio in Fig. 5, left) and a molecular gas clump observed in 12CO (see Fig. 9, bottom) when integrating around a LSR velocity of \( \sim 80 \text{ km s}^{-1} \) (corresponding to near and far kinematic distances of \( \sim 6 \) and \( \sim 10 \text{ kpc} \) respectively). If the γ-ray source HESS J1729−345 is associated with those gas structures, it would therefore not be associated with the SNR HESS J1731-347 thought to lie at a closer distance.

### 6. Conclusion

The newly discovered SNR HESS J1731–347 exhibits a significant shell morphology spatially resolved by HESS, similar to the one observed in radio. Together with RX J1713.7−3946, RX J0852.0−4622 and SN 1006, HESS J1731–347 is now the fourth TeV γ-ray source to join this small but growing class. A lower limit to the distance of the SNR of 3.2 kpc was obtained by comparing the absorption derived from the X-rays and from HI and 12CO observations.

The multi-wavelength emission from the SNR, detected in radio, X-rays and γ-rays, was combined in anSED to investigate the origin of the γ-ray emission assuming that the broadband emission stems from the same region (one-zone model). While the measured fluxes can be accounted for in a purely leptonic model with a magnetic field of the order of 25 μG, this simple model fails to reproduce the spectral shape of the X- and γ-ray emission. A second model that assumes that the TeV emission is produced by hadronic processes is able to reproduce the spectral slopes in X- and γ-rays at the cost of requiring that a large fraction of the kinetic energy of the explosion must be transferred to the accelerated protons and a high ambient medium density of \( n \sim 1 \) cm\(^{-3}\) for \( d \geq 3.2 \) kpc. Moreover for such a density, the corresponding shock speed of the SNR would be an order of magnitude lower than in other SNRs exhibiting bright synchrotron emission.

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**Table 2.** List of the parameters used for the spectral energy distribution modeling presented in Fig. 10.

<table>
<thead>
<tr>
<th>Model</th>
<th>( E_{\text{TeV}} )</th>
<th>( E_{\text{p}} )</th>
<th>( W_p )</th>
<th>( W_{\gamma} )</th>
<th>B</th>
</tr>
</thead>
<tbody>
<tr>
<td>leptonic</td>
<td>18</td>
<td>1.1</td>
<td>10(^{52})</td>
<td>10(^{50})</td>
<td>25</td>
</tr>
<tr>
<td>hadronic</td>
<td>16</td>
<td>100</td>
<td>0.25</td>
<td>2.0</td>
<td>50</td>
</tr>
</tbody>
</table>

**Notes.** The spectral slope are fixed at 2.0 for the electron and the proton distribution. The density of the ambient medium was set to 1 cm\(^{-3}\) in the case of the hadronic model.