Search for brown-dwarf companions of stars*,

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ABSTRACT

Context. The frequency of brown-dwarf companions in close orbit around Sun-like stars is low compared to the frequency of planetary and stellar companions. There is presently no comprehensive explanation of this lack of brown-dwarf companions. Aims. By combining the orbital solutions obtained from stellar radial-velocity curves and Hipparcos astrometric measurements, we attempt to determine the orbit inclinations and therefore the masses of the orbiting companions. By determining the masses of potential brown-dwarf companions, we improve our knowledge of the companion mass-function. Methods. The radial-velocity solutions revealing potential brown-dwarf companions are obtained for stars from the CORALIE and HARPS planet-search surveys or from the literature. The best Keplerian fit to our radial-velocity measurements is found using the Levenberg-Marquardt method. The spectroscopic elements of the radial-velocity solution constrain the fit to the intermediate astrometric data of the new Hipparcos reduction. The astrometric solution and the orbit inclination are found using non-linear 𝜙²-minimisation on a two-parameter search grid. The statistical confidence of the adopted orbital solution is evaluated based on the distribution-free permutation test. Results. The discovery of nine new brown-dwarf candidates orbiting stars in the CORALIE and HARPS radial-velocity surveys is reported. New CORALIE radial velocities yielding accurate orbits of six previously-known hosts of potential brown-dwarf companions are presented. Including the literature targets, 33 hosts of potential brown-dwarf companions are examined. Employing innovative methods, we use the new reduction of the Hipparcos data to fully characterise the astrometric orbits of six objects, revealing M-dwarf companions of masses between 90 ± 55 Jupiters and 0.52 ± 0.05 solar masses. In addition, the masses of two companions can be restricted to the stellar domain. The companion to HD 137510 is found to be a brown dwarf. At 95% confidence, the companion of HD 190228 is also a brown dwarf. Twenty-three companions remain brown-dwarf candidates. On the basis of the CORALIE planet-search sample, we obtain an upper limit of 6.6% for the frequency of brown-dwarf companions around Sun-like stars. We find that the companion-mass distribution function increases toward the lower end of the brown-dwarf mass range, suggesting that we detect the high-mass tail of the planetary distribution. Conclusions. Our findings agree with the results of previous similar studies and confirm the pronounced paucity of brown-dwarf companions around Sun-like stars. They are affected by the Hipparcos astrometric precision and mission duration, which limits the minimum detectable companion mass, and some of the remaining candidates are probably brown-dwarf companions.

Key words. binaries: spectroscopic – brown dwarfs – stars: low-mass – techniques: radial velocities – astrometry

1. Introduction

The mass distribution of close companions to Sun-like stars exhibits a clear deficit of brown dwarfs compared to planets and stellar binaries (Grether & Lineweaver 2006). With a frequency of less than 1% (Marcy & Butler 2000), brown-dwarf companions to solar-type stars at separations below 10 AU are much less common than planetary companions with a frequency of ∼7% (e.g. Udry & Santos 2007) and stellar binaries with a frequency of ∼13% (Duquennoy & Mayor 1991; Halbwachs et al. 2003) at similar separations. Brown dwarfs are substellar objects, massive enough to burn deuterium but too light to permit hydrogen burning, which places them in the mass range of approximately 13–80 Jupiter masses (Burrows et al. 1997; Chabrier & Baraffe 2000; Spiegel et al. 2010). They therefore constitute intermediate objects between planets and stars, though the lower end of the brown-dwarf mass range overlaps with the one of massive planets (e.g., Luhman et al. 2009), and distinguishing planets and brown dwarfs may require us to trace the individual formation processes. A comprehensive description of the formation of low-mass objects, covering the range of planets to very low-mass stars, remains undefined, although the progress in this direction is formidable as reported in the reviews of Whitworth et al. (2007), Luhman et al. (2007), and Burgasser et al. (2007).
A few companions with masses of 18–60 Jupiter masses (\(M_J\)) have been identified by astrometry (e.g., Martíoli et al. 2010; Benedict et al. 2010) or discovered in transiting systems (e.g., Deleuil et al. 2008; Bouchy et al. 2010; Anderson et al., in preparation). Yet, most potential brown dwarf companions are discovered in radial velocity surveys (e.g. Nidever et al. 2002).

In these high-precision programs, designed to achieve planet detection, they are easily spotted and characterised because of their large signatures, by far exceeding the radial velocity measurement precision (e.g., Patel et al. 2007; Bouchy et al. 2009). However, radial-velocity measurements alone do not constrain the orbit inclination. Therefore, they cannot reveal the companion mass, but determine its lower limit. Observations of complementary effects such as the transit light-curve (Deleuil et al. 2008) or the astrometric motion of the host star are required to solve this ambiguity and determine the companion mass.

High-precision astrometry with the Hubble Space Telescope (HST) has been used to find the masses of planetary companions discovered by radial velocity (Benedict et al. 2010, and references therein) and has demonstrated the capability of astrometry to distinguish between planetary, brown-dwarf, or stellar companions. Ground-based optical interferometers achieve the necessary precision to detect orbital motions of binary and multiple systems and to determine the orbital parameters in combination with radial-velocity measurements (Hummel et al. 1993; Muterspaugh et al. 2006; Lane et al. 2007). The astrometric detection of planetary orbits using IR-interferometry is the necessary precision to detect orbital motions of binary and multiple systems and to determine the orbital parameters in combination with radial-velocity measurements (Casertano et al. 2008).

Hipparcos astrometry has been extensively used to constrain the masses of sub-stellar and stellar companions in multiple systems (Perryman et al. 1996; Pourbaix & Jorissen 2000; Zucker & Mazeh 2001; Torres 2007). Han et al. (2001) used Hipparcos astrometry to analyse 30 extrasolar-planet candidates and concluded that half of these companions were rather brown or M dwarfs. With the consequent rectification by Pourbaix & Arenou (2001), showing that the results of Han et al. (2001) had been biased by the fitting procedure, it became clear that the fitting of astrometric data has to be performed very carefully, especially when the size of the orbital signature approaches the measurement precision of the instrument (Pourbaix 2001). All of these past works used the original Hipparcos data (Perryman et al. 1997).

In 2007, the new reduction of the raw Hipparcos data (van Leeuwen 2007) was released, which represents a significant improvement in quality compared to the original data (van Leeuwen 2007). It also includes a revised version of the intermediate astrometric data in a new format that facilitates the search for the signature of orbital motion. We use the new Hipparcos reduction for our study, which aims to determine the masses of companions to stars detected using high-precision radial-velocity measurements and found to have minimum masses in the brown-dwarf mass range of 13–80 \(M_J\). It contains 14 stars from the CORALIE survey (Udry et al. 2000) and one star from the HARPS planet search program (Mayor et al. 2003). To date, the CORALIE planet survey has contributed to the discovery of more than 50 extrasolar planets (e.g. Naef et al. 2001; Mayor et al. 2004; Tamuz et al. 2008). CORALIE is an optical echelle spectrograph mounted on the 1.2 m Swiss Telescope located at the European Southern Observatory in La Silla, Chile. A description of the instrument can be found in Queloz et al. (2000) and the references therein. With its unmatched precision, the HARPS instrument permitted the discovery of more than 70 planetary companions to date, including earth-mass planets (e.g. Pepe et al. 2007; Santos et al. 2007; Mayor et al. 2009; Forveille et al. 2009). The characteristics, radial velocities, and orbital solutions of the 15 CORALIE and HARPS stars are presented and discussed in Sects. 2.1–2.3. In addition, 18 potential brown-dwarf host stars with radial-velocity orbits were selected from the literature and are listed in Sect. 2.4. These include HD 190228, initially a planet-host star (Sivan et al. 2004), whose importance we recognised during the analysis.

2. Brown dwarf candidates

Our target sample consists of 33 stars that exhibit radial velocity variations caused by a companion with minimum mass \(M_J\) \(\sin i\) in the brown-dwarf mass range of 13–80 \(M_J\). It contains 14 stars from the CORALIE survey (Udry et al. 2000) and one star from the HARPS planet search program (Mayor et al. 2003). To date, the CORALIE planet survey has contributed to the discovery of more than 50 extrasolar planets (e.g. Naef et al. 2001; Mayor et al. 2004; Tamuz et al. 2008). CORALIE is an optical echelle spectrograph mounted on the 1.2 m Swiss Telescope located at the European Southern Observatory in La Silla, Chile. A description of the instrument can be found in Queloz et al. (2000) and the references therein. With its unmatched precision, the HARPS instrument permitted the discovery of more than 70 planetary companions to date, including earth-mass planets (e.g. Pepe et al. 2007; Santos et al. 2007; Mayor et al. 2009; Forveille et al. 2009). The characteristics, radial velocities, and orbital solutions of the 15 CORALIE and HARPS stars are presented and discussed in Sects. 2.1–2.3. In addition, 18 potential brown-dwarf host stars with radial-velocity orbits were selected from the literature and are listed in Sect. 2.4. These include HD 190228, initially a planet-host star (Sivan et al. 2004), whose importance we recognised during the analysis.

2.1. Stellar characteristics

The identifiers and basic stellar characteristics of the CORALIE stars and HIP 103019 from the HARPS survey are listed in Table 1. The apparent visual magnitude \(V\), the colour \(B−V\), and the parallax \(\varpi\) are from the new Hipparcos reduction (van Leeuwen 2007), whereas the spectral type is from the original Hipparcos catalogue (Perryman et al. 1997). Table 2 displays the derived stellar parameters. The uncertainty in \(V\) is smaller than 0.002 mag. The absolute magnitude \(MV\) is derived from the Hipparcos magnitude and parallax. The effective temperature \(T_{\text{eff}}\), the surface gravity \(\log g\), and the metallicity [Fe/H] are

<table>
<thead>
<tr>
<th>Nr.</th>
<th>Object HD</th>
<th>HIP</th>
<th>Sp. T.</th>
<th>(B−V)</th>
<th>(V)</th>
<th>(\varpi) (mas)</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>3277</td>
<td>2790</td>
<td>G8</td>
<td>7.59</td>
<td>0.73</td>
<td>34.7 ± 0.7</td>
</tr>
<tr>
<td>2</td>
<td>4747</td>
<td>3850</td>
<td>G8</td>
<td>7.30</td>
<td>0.77</td>
<td>53.5 ± 0.5</td>
</tr>
<tr>
<td>3</td>
<td>17289</td>
<td>12726</td>
<td>G0V</td>
<td>7.56</td>
<td>0.59</td>
<td>20.8 ± 1.2</td>
</tr>
<tr>
<td>4</td>
<td>30501</td>
<td>22122</td>
<td>K0V</td>
<td>7.73</td>
<td>0.88</td>
<td>47.9 ± 0.5</td>
</tr>
<tr>
<td>5</td>
<td>43848</td>
<td>29804</td>
<td>K1IV</td>
<td>8.80</td>
<td>0.93</td>
<td>26.4 ± 0.8</td>
</tr>
<tr>
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<td>52756</td>
<td>34052</td>
<td>K1V</td>
<td>8.60</td>
<td>1.18</td>
<td>57.4 ± 1.2</td>
</tr>
<tr>
<td>7</td>
<td>53680</td>
<td>33736</td>
<td>K5V</td>
<td>8.61</td>
<td>0.90</td>
<td>30.1 ± 0.6</td>
</tr>
<tr>
<td>8</td>
<td>74014</td>
<td>42634</td>
<td>K0</td>
<td>7.73</td>
<td>0.76</td>
<td>29.6 ± 0.7</td>
</tr>
<tr>
<td>9</td>
<td>89707</td>
<td>50671</td>
<td>G1V</td>
<td>7.29</td>
<td>0.55</td>
<td>30.6 ± 0.6</td>
</tr>
<tr>
<td>10</td>
<td>154697</td>
<td>83770</td>
<td>G5</td>
<td>7.97</td>
<td>0.73</td>
<td>30.6 ± 1.0</td>
</tr>
<tr>
<td>11</td>
<td>164427A</td>
<td>88531</td>
<td>G4IV</td>
<td>7.01</td>
<td>0.62</td>
<td>26.3 ± 0.7</td>
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<tr>
<td>12</td>
<td>167665</td>
<td>89620</td>
<td>F8V</td>
<td>6.48</td>
<td>0.54</td>
<td>32.8 ± 0.5</td>
</tr>
<tr>
<td>13</td>
<td>189310</td>
<td>99634</td>
<td>K2V</td>
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<td>0.90</td>
<td>31.0 ± 0.6</td>
</tr>
<tr>
<td>14</td>
<td>211847</td>
<td>110340</td>
<td>G5V</td>
<td>8.78</td>
<td>0.66</td>
<td>19.8 ± 1.3</td>
</tr>
<tr>
<td>15</td>
<td>⋯</td>
<td>103019</td>
<td>K5</td>
<td>10.45</td>
<td>1.33</td>
<td>28.3 ± 2.3</td>
</tr>
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</table>
Table 2. Derived stellar parameters of the surveyed stars.

<table>
<thead>
<tr>
<th>Object</th>
<th>Instr.&quot;</th>
<th>M⊙</th>
<th>Tₐr (K)</th>
<th>log g (cgs)</th>
<th>[Fe/H] (dex)</th>
<th>v sin i (km s⁻¹)</th>
<th>M₁ (M⊙)</th>
<th>Age (Gyr)</th>
</tr>
</thead>
<tbody>
<tr>
<td>HD 3277</td>
<td>C</td>
<td>5.30</td>
<td>5539 ± 50</td>
<td>4.36 ± 0.10</td>
<td>-0.06 ± 0.08</td>
<td>&lt;2.0</td>
<td>0.91 ± 0.03</td>
<td>0.4–8.0</td>
</tr>
<tr>
<td>HD 4747</td>
<td>C</td>
<td>5.94</td>
<td>5316 ± 50</td>
<td>4.48 ± 0.10</td>
<td>-0.21 ± 0.05</td>
<td>0.79 ± 0.06</td>
<td>0.81 ± 0.02</td>
<td>0.1–7.3</td>
</tr>
<tr>
<td>HD 17289</td>
<td>C</td>
<td>4.15</td>
<td>5924 ± 50</td>
<td>4.37 ± 0.10</td>
<td>-0.11 ± 0.06</td>
<td>&lt;2.0</td>
<td>1.01 ± 0.03</td>
<td>6.0–9.3</td>
</tr>
<tr>
<td>HD 30501</td>
<td>C</td>
<td>6.13</td>
<td>5223 ± 50</td>
<td>4.56 ± 0.10</td>
<td>-0.06 ± 0.06</td>
<td>&lt;2.0</td>
<td>0.81 ± 0.02</td>
<td>0.8–7.0</td>
</tr>
<tr>
<td>HD 43848</td>
<td>C</td>
<td>5.91</td>
<td>5534 ± 50</td>
<td>4.56 ± 0.15</td>
<td>0.22 ± 0.06</td>
<td>&lt;2.0</td>
<td>0.89 ± 0.02</td>
<td>0.0–5.1</td>
</tr>
<tr>
<td>HD 52756</td>
<td>C</td>
<td>7.60</td>
<td>5216 ± 65</td>
<td>4.47 ± 0.11</td>
<td>0.13 ± 0.05</td>
<td>&lt;2.0</td>
<td>0.83 ± 0.01</td>
<td>0.2–2.7</td>
</tr>
<tr>
<td>HD 53680</td>
<td>C</td>
<td>6.00</td>
<td>5167 ± 94</td>
<td>5.37 ± 0.29</td>
<td>-0.29 ± 0.08</td>
<td>2.08 ± 0.31</td>
<td>0.79 ± 0.02</td>
<td>0.7–9.4</td>
</tr>
<tr>
<td>HD 74014</td>
<td>C</td>
<td>5.08</td>
<td>5662 ± 55</td>
<td>4.39 ± 0.10</td>
<td>0.26 ± 0.08</td>
<td>&lt;2.0</td>
<td>1.00 ± 0.03</td>
<td>1.3–8.0</td>
</tr>
<tr>
<td>HD 89707</td>
<td>C</td>
<td>4.71</td>
<td>6047 ± 50</td>
<td>4.52 ± 0.10</td>
<td>-0.33 ± 0.06</td>
<td>&lt;2.0</td>
<td>0.96 ± 0.04</td>
<td>0.5–6.0</td>
</tr>
<tr>
<td>HD 154697</td>
<td>C</td>
<td>5.40</td>
<td>5648 ± 50</td>
<td>4.42 ± 0.10</td>
<td>0.13 ± 0.06</td>
<td>&lt;2.0</td>
<td>0.96 ± 0.02</td>
<td>0.3–3.3</td>
</tr>
<tr>
<td>HD 164427A</td>
<td>C</td>
<td>4.11</td>
<td>6003 ± 50</td>
<td>4.35 ± 0.10</td>
<td>0.19 ± 0.06</td>
<td>&lt;2.0</td>
<td>1.15 ± 0.03</td>
<td>1.2–4.1</td>
</tr>
<tr>
<td>HD 167665</td>
<td>C</td>
<td>4.06</td>
<td>6224 ± 50</td>
<td>4.44 ± 0.10</td>
<td>-0.05 ± 0.06</td>
<td>&lt;2.0</td>
<td>1.14 ± 0.03</td>
<td>0.7–3.6</td>
</tr>
<tr>
<td>HD 189310</td>
<td>C</td>
<td>6.07</td>
<td>5188 ± 50</td>
<td>4.49 ± 0.10</td>
<td>-0.01 ± 0.06</td>
<td>&lt;2.0</td>
<td>0.83 ± 0.02</td>
<td>0.0–7.5</td>
</tr>
<tr>
<td>HD 211847</td>
<td>C</td>
<td>5.26</td>
<td>5715 ± 50</td>
<td>4.49 ± 0.10</td>
<td>-0.08 ± 0.06</td>
<td>&lt;2.0</td>
<td>0.94 ± 0.04</td>
<td>0.1–6.0</td>
</tr>
<tr>
<td>HIP 103019</td>
<td>H</td>
<td>7.61</td>
<td>4913 ± 115</td>
<td>4.45 ± 0.28</td>
<td>-0.30 ± 0.06</td>
<td>...</td>
<td>0.70 ± 0.01</td>
<td>0.7–9.9</td>
</tr>
</tbody>
</table>

Notes. (a) Data from Santos et al. (2005). (b) C and H stand for CORALIE and HARPS, respectively.

2.2. Radial-velocity measurements and orbital solutions

Optical high-resolution spectra of the stars in Table 2 were collected with the CORALIE and HARPS spectrographs over time intervals extending to 11 years. Radial velocities were estimated by means of the cross-correlations of the extracted stellar spectra with numerical templates, which depend on the target’s spectral type. During CORALIE observations, a reference thorium-argon spectrum is recorded simultaneously with the stellar spectrum and used to measure and correct for residual zero-point drifts (Baranne et al. 1996). Photon noise limits the obtained precision to typically 2–4 m s⁻¹ per epoch. To account for systematic drifts, an external error of 0.5 m s⁻¹ is quadratically added to the HARPS radial-velocity uncertainties.

The Keplerian-orbit solution is found by imposing the model function Eq. (11) with 6 parameters (K₁, e, ω, P, and γ) to the radial velocities of a given star. The fit determines the values of the period P, the eccentricity e, the time of periastron passage T₀, the longitude of periastron ω, the radial-velocity semi-amplitude K₁, and the systemic velocity γ. In general, measurements with uncertainties larger than 15 m s⁻¹ are not used in the model adjustments.

A single-companion model closely reproduces the data of 11 stars with reduced-chi values of χ² / ν ≤ 1.7. For 4 stars, HD 17289, HD 89707, HD 211847, and HIP 103019, the fit is not as good with χ² / ν = 1.7–2.0 and produces residuals of unusually large amplitude, without contesting the presence of a large orbital motion. Introducing additional model parameters, corresponding to a linear or quadratic drift or a second companion does not significantly improve the fit. Therefore, the moderate fit quality has to be attributed to either underestimated radial velocity errors or to additional noise sources, such as stellar activity. This is discussed for the individual targets in the next section.

The orbital periods of HD 4747 and HD 211847 are very long and not covered by the measurement time-span. Hence, the best-fit period is poorly constrained and realistic confidence intervals have to be derived by a different method. We chose a criterion based on the fit-quality metric χ², and its error σχ² obtained for the best fit. During a second fit, we fix the period during the adjustment, such that the resulting χ²-value increases to χ² + 3σχ². In this way, we obtain upper and lower limits to the orbital parameters, which are listed at the bottom of Tables 3 and 4. The method that searches for the astrometric signature in the Hipparcos data presented in Sect. 3 relies on the precise knowledge of the radial-velocity parameters. Because the periods of HD 4747 and HD 211847 are poorly constrained and the orbital coverage of Hipparcos is very short, we do not perform the combined astrometric analysis for these targets. They are presented as part of the brown-dwarf candidates characterised with CORALIE.

Table 3 displays the orbital elements of the best-fit solutions and Table 4 lists the basic statistics and the derived orbit parameters (see also Fig. 1). The parameter T₀ is given as reduced barycentric Julian date (BJD), that is JD⊙ = BJD – 2 400 000.

1 The web interface is http://stev.oapd.inaf.it/cgi-bin/param
In addition, $N_{RV}$ is the total number of measurements, $\Delta T$ is the observation time-span, $S/N$ is the median $S/N$ per pixel at 550 nm, $\langle \sigma_{\text{RV}} \rangle$ is the mean measurement uncertainty, $\chi^2$ is the reduced-chi value of the fit, G.o.F. is the goodness of fit, $\sigma_{(O-C)}$ gives the root-mean-square dispersion of the fit residuals, $af(m)$ is the mass function, $a_{\text{rel}}$ is the semimajor axis of the relative orbit, and finally $M_2 \sin i$ is the minimum companion mass. The quoted errors correspond to 1-$\sigma$ confidence intervals derived from 5000 Monte Carlo simulations. Figures 4 and 5 show the phase-folded radial velocities and the adopted orbital solution. All measurements are available in electronic form at the CDS as indicated on the title page.

Fig. 1. Orbital characteristics of the 15 brown-dwarf candidates observed with CORALIE and HARPS. Filled symbols indicate the companions with stellar mass identified by this study. The error bars are often smaller than the symbol size, though the orbits of HD 4747 and HD 211847 are not covered and therefore less constrained, see the text for details. Top: minimum companion mass and separation. Bottom: eccentricity and orbital period.

Fig. 2. Radial velocities (top) and residuals (bottom) to the best fit, represented by the solid line, of HD 4747 measured with HIRES (orange diamonds), CORALIE C98 (red circles), and CORALIE C07 (blue squares).

2.3. Details of the surveyed stars

We now present the 15 stars observed with CORALIE and HARPS, and discuss details about their respective radial-velocity solutions:

- HD 3277 (HIP 2790) is found to host a potential brown-dwarf companion. The star was listed as an uncertain spectroscopic binary in Tokovinin et al. (2006), who did not have sufficient radial-velocity measurements to constrain the orbit.

- HD 4747 (HIP 3850) has a potential brown-dwarf companion discovered with radial velocities by Nidever et al. (2002). We combine CORALIE observations with the KECK HIRES radial velocities of Nidever et al. (2002) retrieved from the SB9-catalogue (Pourbaix et al. 2004). However, the orbital period is very long and not covered by the measurements (see Fig. 2). Tables 3 and 4 show the best-fit solution and the solutions with fixed period such that the $\chi^2$-value increases to $\chi^2 + 3\sigma\chi^2 = 1.56$. The resulting period range is $P = 7900–29000$ days with corresponding minimum companion-mass $M_2 \sin i = 39.6–58.1 \text{ M}_J$.

- HD 17289 (HIP 12726) is found to host a potential brown-dwarf companion. Goldin & Makarov (2007) independently found an orbital signature from Hipparcos astrometry alone with a period of $536 \pm 12$ days, which is close to our solution ($562.1 \pm 0.3$ days). However, the orbital parameters derived by Goldin & Makarov (2007) are less precise and less accurate than the radial-velocity solution. The dispersion in the fit residuals is abnormally large for CORALIE. The detailed analysis of the cross-correlation function (CCF) indicates that its bisector velocity span and full-width-at-half-maximum (FWHM) are correlated with the orbital phase (Fig. 6). As we show in Sect. 4, the companion of HD 17289 is a star with mass of $0.52 \text{ M}_\odot$. The two stars are very close (<50 mas, Sect. 4.1) and their estimated intensity ratio in the visible is about 1:180. Therefore, the companion’s light is also injected in the science fibre of
Table 3. Single-companion Keplerian orbital solutions of the surveyed stars.

<table>
<thead>
<tr>
<th>Object</th>
<th>$\gamma_{\text{COR}}$ (km s(^{-1}))</th>
<th>$\gamma_{\text{COV}}$ (km s(^{-1}))</th>
<th>$P$ (day)</th>
<th>$e$</th>
<th>$K_1$ (m s(^{-1}))</th>
<th>$T_0$ (JD*)</th>
<th>$\omega$ (deg)</th>
</tr>
</thead>
<tbody>
<tr>
<td>HD 3277</td>
<td>−13.799±0.006</td>
<td>−13.796±0.004</td>
<td>46.151±0.002</td>
<td>0.285±0.001</td>
<td>4076.4±3.5</td>
<td>5432.14±0.03</td>
<td>−39.4±0.3</td>
</tr>
<tr>
<td>HD 172289</td>
<td>35.400±0.009</td>
<td>35.420±0.017</td>
<td>562.1±0.4</td>
<td>0.533±0.004</td>
<td>1414.5±10.0</td>
<td>5376.25±0.9</td>
<td>52.4±0.6</td>
</tr>
<tr>
<td>HD 30501</td>
<td>23.72±0.026</td>
<td>23.710±0.028</td>
<td>2073.6±3.0</td>
<td>0.741±0.004</td>
<td>1703.1±26.0</td>
<td>5385.1±3.0</td>
<td>70.4±0.7</td>
</tr>
<tr>
<td>HD 43848*</td>
<td>68.130±0.045</td>
<td>68.110±0.042</td>
<td>2354.8±6.9</td>
<td>0.709±0.002</td>
<td>517.0±31.7</td>
<td>55858.8±3.3</td>
<td>−131.7±1.3</td>
</tr>
<tr>
<td>HD 43848b</td>
<td>66.845±0.492</td>
<td>66.837±0.496</td>
<td>2349.8±8.8</td>
<td>0.884±0.030</td>
<td>1859.3±570</td>
<td>56564.6±10.8</td>
<td>−151.5±10.0</td>
</tr>
<tr>
<td>HD 52756</td>
<td>56.907±0.003</td>
<td>56.904±0.005</td>
<td>52.8657±0.001</td>
<td>0.6780±0.003</td>
<td>4948.8±3.0</td>
<td>52882.617±0.003</td>
<td>139.9±0.1</td>
</tr>
<tr>
<td>HD 53680</td>
<td>89.065±0.005</td>
<td>89.048±0.004</td>
<td>1688.6±1.1</td>
<td>0.475±0.002</td>
<td>1239.8±4.0</td>
<td>54182.9±1.4</td>
<td>−133.2±0.3</td>
</tr>
<tr>
<td>HD 74014</td>
<td>−16.371±0.002</td>
<td>−16.387±0.008</td>
<td>6936.3±134.9</td>
<td>0.532±0.006</td>
<td>617.5±23.3</td>
<td>53431.9±5.2</td>
<td>−34.7±0.6</td>
</tr>
<tr>
<td>HD 89070</td>
<td>84.360±1.886</td>
<td>84.390±1.699</td>
<td>298.5±0.1</td>
<td>0.900±0.030</td>
<td>4189.5±476.7</td>
<td>53296.4±5.5</td>
<td>74.8±18.7</td>
</tr>
<tr>
<td>HD 154697</td>
<td>−34.870±0.007</td>
<td>−34.858±0.011</td>
<td>1835.9±2.0</td>
<td>0.164±0.002</td>
<td>1230.1±8.1</td>
<td>52579.7±9.5</td>
<td>179.1±1.6</td>
</tr>
<tr>
<td>HD 164427A</td>
<td>5.506±0.015</td>
<td>5.512±0.015</td>
<td>108.537±0.001</td>
<td>0.551±0.002</td>
<td>2231.1±126.0</td>
<td>53026.8±0.1</td>
<td>−3.3±0.4</td>
</tr>
<tr>
<td>HD 167665</td>
<td>7.984±0.004</td>
<td>8.003±0.008</td>
<td>4451.8±27.6</td>
<td>0.340±0.005</td>
<td>6095.3±33.3</td>
<td>56987.6±29.7</td>
<td>−134.3±0.9</td>
</tr>
<tr>
<td>HD 189310</td>
<td>−10.250±0.003</td>
<td>−10.250±0.003</td>
<td>14.186±3.0</td>
<td>0.359±0.001</td>
<td>2606.0±2.8</td>
<td>53542.399±0.004</td>
<td>50.1±0.1</td>
</tr>
<tr>
<td>HIP 103019</td>
<td>−4.296±0.002</td>
<td>−4.292±0.002</td>
<td>917.3±11.1</td>
<td>0.507±0.003</td>
<td>1614.2±18.0</td>
<td>54681.5±0.2</td>
<td>−74.2±8.1</td>
</tr>
</tbody>
</table>

Notes. (\(a\)) Fixed. (\(b\)) HARPS systemic velocity. (\(c\)) KECK HIRES systemic velocity is \(9.823±0.016\) km s\(^{-1}\). (\(d\)) Lower period limit. (\(e\)) Best-fit solution. (\(f\)) Upper period limit. (\(g\)) Final solution at lower eccentricity. The MIKE systemic velocity \(−0.046±0.046\) km s\(^{-1}\). (\(h\)) Formal solution at higher eccentricity. The MIKE systemic velocity \(−1.331±0.493\) km s\(^{-1}\).

CORALIE and its signature appears in the collected spectra. The spectral lines of the companion distort the CCF depending on the orbital phase. As expected, the CCF width is minimal when the radial-velocity curves of both components cross and equal the systemic velocity. The detailed modelling of this double-lined spectroscopic binary is beyond the scope of this paper, but we have identified the binary as the cause of a phase-dependent radial-velocity bias, which explains the excess noise in our measurements.

– HD 30501 (HIP 22122) is found to host a potential brown-dwarf companion.

– HD 43848 (HIP 29804) has a potential brown-dwarf companion discovered with radial velocity by Minniti et al. (2009). An astrometric analysis using Hipparcos and the solution of Minniti et al. (2009) was performed by Sozzetti & Desidera (2010). We combined the published measurements\(^2\) with the CORALIE velocities to adjust the Keplerian model. While the orbital period is well defined, the distributions of $e$, $K_1$, $T_0$, and $\omega$ obtained from the Monte Carlo simulations are bimodal.

\(^2\) By comparing the published radial velocities to our CORALIE measurements, we discovered that the published Julian dates corresponding to the radial velocities of Minniti et al. (2009) are misprinted and offset by $-2000$ days.

This is because the lower velocity turnover is not well sampled by the measurements. The fit-quality in terms of $\chi^2$ is comparable for all Monte Carlo simulations and cannot be used to identify the correct solution. In Fig. 7, we show the eccentricity and semi-amplitude distributions. Because the realisation of the lower-eccentricity orbit is more probable (77%) than the high-eccentricity orbit (23%), we consider only the orbits with $e \leq 0.78$ to find the final solution and perform the astrometric analysis. For completeness, we also list the less probable and higher-eccentricity solution in Tables 3 and 4.

– HD 52756 (HIP 33736) is found to host a potential brown-dwarf companion.

– HD 53680 (HIP 34052) is found to host a potential brown-dwarf companion.

– HD 74014 (HIP 42634) has a potential brown-dwarf companion discovered with radial velocities by Patel et al. (2007). The original orbit is compatible with our solution obtained from 118 CORALIE velocities covering the complete orbit.

– HD 89070 (HIP 50671) has a potential brown-dwarf companion announced by Duquennoy & Mayor (1991) and analysed by Halbwachs et al. (2000) based on radial velocities collected with the CORAVEL and ELODIE instruments. The original orbit is compatible with our solution obtained from 34.
Table 4. Companion parameters for the surveyed stars.

<table>
<thead>
<tr>
<th>Object</th>
<th>$N_{RV}$</th>
<th>$\Delta T$ (yr)</th>
<th>$S/N$</th>
<th>$\langle\sigma_{RV}\rangle$ (m s$^{-1}$)</th>
<th>$\chi_t$</th>
<th>G.o.F.</th>
<th>$\sigma_{(o-C)}$ (m s$^{-1}$)</th>
<th>$a_1 \sin i$ (10$^{-3}$ AU)</th>
<th>$f(m)$ (10$^{-9}$ $M_J$)</th>
<th>$a_{rel}$ (AU)</th>
<th>$M_J \sin i$ (AU)</th>
</tr>
</thead>
<tbody>
<tr>
<td>HD 3277</td>
<td>14</td>
<td>10.4</td>
<td>35.4</td>
<td>8.4</td>
<td>1.10 ± 0.27</td>
<td>0.55</td>
<td>5.38</td>
<td>16.58 ± 0.02</td>
<td>285270 ± 700</td>
<td>0.249 ± 0.004</td>
<td>64.7 ± 2.2</td>
</tr>
<tr>
<td>HD 17289</td>
<td>29</td>
<td>11.3</td>
<td>46.9</td>
<td>6.6</td>
<td>2.02 ± 0.15</td>
<td>6.07</td>
<td>12.46</td>
<td>61.9 ± 0.3</td>
<td>100178 ± 1619</td>
<td>1.357 ± 0.023</td>
<td>48.9 ± 1.7</td>
</tr>
<tr>
<td>HD 30501</td>
<td>46</td>
<td>11.3</td>
<td>32.5</td>
<td>6.6</td>
<td>0.91 ± 0.11</td>
<td>-0.73</td>
<td>6.33</td>
<td>217.9 ± 2.0</td>
<td>320923 ± 4957</td>
<td>0.306 ± 0.049</td>
<td>62.3 ± 2.1</td>
</tr>
<tr>
<td>HD 43848</td>
<td>38°</td>
<td>9.2</td>
<td>30.1</td>
<td>5.6</td>
<td>1.31 ± 0.13</td>
<td>2.40</td>
<td>7.29</td>
<td>87.5 ± 4.8</td>
<td>16285 ± 2711</td>
<td>3.360 ± 0.057</td>
<td>24.5 ± 1.6</td>
</tr>
<tr>
<td>HD 43848e</td>
<td>38°</td>
<td>9.2</td>
<td>30.1</td>
<td>5.6</td>
<td>1.31 ± 0.13</td>
<td>2.40</td>
<td>7.29</td>
<td>179.1 ± 32.4</td>
<td>150526 ± 2638</td>
<td>3.385 ± 0.057</td>
<td>50.2 ± 2.8</td>
</tr>
<tr>
<td>HD 52756</td>
<td>33</td>
<td>10.1</td>
<td>28.8</td>
<td>6.1</td>
<td>1.15 ± 0.14</td>
<td>1.17</td>
<td>7.72</td>
<td>17.68 ± 0.01</td>
<td>263682 ± 483</td>
<td>0.265 ± 0.004</td>
<td>59.3 ± 2.0</td>
</tr>
<tr>
<td>HD 53680</td>
<td>48</td>
<td>9.1</td>
<td>30.5</td>
<td>7.1</td>
<td>1.03 ± 0.11</td>
<td>0.30</td>
<td>7.77</td>
<td>169.4 ± 0.5</td>
<td>227414 ± 2125</td>
<td>2.620 ± 0.043</td>
<td>54.7 ± 1.8</td>
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<tr>
<td>HD 74014</td>
<td>118</td>
<td>11.3</td>
<td>41.8</td>
<td>5.3</td>
<td>0.96 ± 0.07</td>
<td>-0.61</td>
<td>6.29</td>
<td>333.2 ± 4.8</td>
<td>102583 ± 947</td>
<td>7.722 ± 0.153</td>
<td>49.0 ± 1.7</td>
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<tr>
<td>HD 89707</td>
<td>37</td>
<td>9.0</td>
<td>57.5</td>
<td>6.1</td>
<td>1.70 ± 0.13</td>
<td>5.03</td>
<td>11.63</td>
<td>46.0 ± 6.8</td>
<td>158142 ± 6203</td>
<td>0.877 ± 0.015</td>
<td>53.6 ± 2.9</td>
</tr>
<tr>
<td>HD 154697</td>
<td>49</td>
<td>10.8</td>
<td>31.6</td>
<td>7.2</td>
<td>1.09 ± 0.11</td>
<td>0.85</td>
<td>7.82</td>
<td>204.7 ± 1.3</td>
<td>339722 ± 6646</td>
<td>2.960 ± 0.048</td>
<td>71.1 ± 2.4</td>
</tr>
<tr>
<td>HD 164427A</td>
<td>21</td>
<td>10.7</td>
<td>43.1</td>
<td>6.7</td>
<td>1.25 ± 0.19</td>
<td>1.39</td>
<td>7.12</td>
<td>18.6 ± 0.1</td>
<td>72658 ± 991</td>
<td>0.472 ± 0.007</td>
<td>48.0 ± 1.6</td>
</tr>
<tr>
<td>HD 167665</td>
<td>38</td>
<td>10.9</td>
<td>76.3</td>
<td>6.0</td>
<td>1.30 ± 0.13</td>
<td>2.36</td>
<td>8.28</td>
<td>234.6 ± 2.3</td>
<td>86931 ± 1802</td>
<td>5.608 ± 0.097</td>
<td>50.6 ± 1.7</td>
</tr>
<tr>
<td>HD 189310</td>
<td>19</td>
<td>8.3</td>
<td>22.2</td>
<td>7.6</td>
<td>0.70 ± 0.20</td>
<td>-1.43</td>
<td>4.39</td>
<td>3.172 ± 0.003</td>
<td>21156 ± 56</td>
<td>0.109 ± 0.002</td>
<td>25.6 ± 0.9</td>
</tr>
<tr>
<td>HIP 103019</td>
<td>26</td>
<td>2.3</td>
<td>44.0</td>
<td>2.3</td>
<td>1.92 ± 0.16</td>
<td>5.28</td>
<td>3.57</td>
<td>117.7 ± 0.1</td>
<td>258370 ± 9999</td>
<td>1.679 ± 0.027</td>
<td>52.5 ± 1.7</td>
</tr>
</tbody>
</table>

Notes. (a) 21 velocity measurements are from Nidever et al. (2002). (b) Lower period limit. (c) Best-fit solution. (d) Upper period limit. (e) 10 velocity measurements are from Minniti et al. (2009). (f) Final solution at lower eccentricity. (g) Formal solution at higher eccentricity.

37 CORALIE velocities. A combined astrometric analysis using Hipparcos was performed by Zucker & Mazeh (2001). We find slight evidence of re-emission at the bottom of the Ca II H absorption line (Fig. 8) and derive an activity index of log $R'_{HK} = -4.66 ± 0.03$ from a high $S/N$ CORALIE spectrum. Using the relation of Santos et al. (2000), we expect an additional activity-induced radial-velocity noise of 12 m s$^{-1}$, which explains the unusually-large fit residuals encountered for HD 89707. Accounting for the additional noise during the adjustment yields an acceptable fit quality with $\chi_t = 0.91 ± 0.13$.

- HD 154697 (HIP 83770) is found to host a potential brown-dwarf companion.
- HD 164427A (HIP 88531) has a potential brown-dwarf companion announced by Tinney et al. (2001), whose original orbit is compatible with our solution. A combined astrometric analysis using Hipparcos was performed by Zucker & Mazeh (2001).
- HD 189310 (HIP 99634) is found to host a potential brown-dwarf companion in a short-period orbit of 14.2 days.
- HD 167665 (HIP 89620) has a potential brown-dwarf companion announced by Patel et al. (2007), whose original orbit is compatible with our solution.
- HD 211847 (HIP 110340) is found to host a potential brown-dwarf companion. The long-period orbit is not completely covered by our measurements (Fig. 3) and the CORALIE upgrade occurred at minimum radial velocity. To obtain a reasonable solution, we therefore fix the offset between the

![Fig. 3. Radial velocities (top) and residuals (bottom) to the best fit, represented by the solid line, of HD 211847 measured with CORALIE C98 (red circles) and CORALIE C07 (blue squares).](image)
Fig. 4. Phase-folded radial velocities of 8 stars with potential brown-dwarf companions. Red circles and blue squares indicate measurements with CORALIE C98 and C07, respectively. For HD 43848, the orange diamonds show the MIKE measurements. The solid lines correspond to the best-fit solutions. The error bars are smaller than the symbol size.
parameters $\gamma_{\text{CORALIE}}$ and $\gamma_{\text{HARPS}}$ to zero, which is the mean offset value found for three other stars of the same spectral type. As for HD 4747, we derive realistic confidence intervals by fixing the period during the adjustment, such that the resulting $\chi$-value increases to $\chi_r + 3\sigma_{\chi} = 2.32$. The period lower limit is $P = 3750$ days and we set the upper limit to $P = 100000$ days, because we cannot derive an upper limit using the defined criterion. The corresponding minimum companion mass is $M_2 = 17.0 - 23.3 M_J$. The dispersion in the best-fit residuals is abnormally large for CORALIE. We find evidence of re-emission at the bottom of CaII H and CaII K (see Fig. 9) and derive an activity index of $\log R'_{\text{HK}} = -4.74 \pm 0.04$ from a high $S/N$ CORALIE spectrum. Using the relation of Santos et al. (2000), we expect an additional activity-induced radial-velocity noise of $11 \text{ m s}^{-1}$, which explains the fit residuals encountered for HD 211847. Accounting for the additional noise during the adjustment yields an acceptable fit quality with $\chi_r = 1.15 \pm 0.15$.

- HIP 103019 is discovered to host a potential brown-dwarf companion using HARPS. The star is included in the search for planets in a volume-limited sample (e.g. Naef et al. 2010). Twenty-six radial velocities were collected with HARPS with a mean uncertainty of $2.3 \text{ m s}^{-1}$. The dispersion in the fit residuals is abnormally large for HARPS ($\sigma_{\text{HARPS}} = 3.57 \text{ m s}^{-1}$). Because of its faintness, we were not able to measure the activity level of HIP 103019. We therefore cannot exclude an additional radial-velocity jitter induced by the presence of spots on the surface of the star.

Fig. 5. Phase-folded radial velocities of 5 stars with potential brown-dwarf companions. Error bars are smaller than the symbol size.
In addition to the stars from the CORALIE and HARPS surveys, we selected 18 stars from the literature using the following criteria. They have announced radial-velocity companions with minimum masses of 13–80 $M_\odot$, and are not flagged as binary stars in the new Hipparcos reduction. The selected stars are listed in Table 5 and details are given below:

- HD 13189 is a K2 II star with a very uncertain mass of 2–7 $M_\odot$ and a companion announced by Hatzes et al. (2005) with $M_2 \sin i = 8–20 M_\oplus$.
- HD 30339 (F8, 1.1 $M_\odot$), HD 65430 (K0V, 0.78 $M_\odot$), HD 140913 (G0V, 0.98 $M_\odot$), and HD 174457 (F8, 1.19 $M_\odot$) have companions announced by Nidever et al. (2002) with $M_2 \sin i$ ranging between 42–78 $M_\oplus$.
- HD 38529 is a G4 IV star of mass 1.48 $M_\odot$. Companion announced by Fischer et al. (2003) with $M_2 \sin i = 12.7 M_\oplus$. Refined orbits are given by Wright et al. (2009) and Benedict et al. (2010). We use the latest orbital elements for analysis.
- HD 91669 is a KO dwarf of mass 0.88 ± 0.04 $M_\odot$. Companion announced by Wittnemeyer et al. (2009) with $M_2 \sin i = 30.6 M_\oplus$.
- HD 107583 is a G8 III giant star of mass 2.7 ± 0.3 $M_\odot$ and a companion announced by Liu et al. (2008) with $M_2 \sin i = 19.4 M_\oplus$.
- HD 119445 is a G6 III giant star of mass 3.9 ± 0.4 $M_\odot$ and a companion announced by Omiya et al. (2009) with $M_2 \sin i = 37.6 M_\oplus$.
- HD 131664 is a G3V star of mass 1.10 ± 0.03 $M_\odot$. Companion announced by Moutou et al. (2009) with $M_2 \sin i = 18.15 M_\oplus$.
- HD 137510 is a G0 IV star of mass 1.42 $M_\odot$ and a companion announced by Endl et al. (2004). The updated data by Butler et al. (2006) lists $M_2 \sin i = 22.7 M_\oplus$.
- HD 180777 is a A9V star of mass 1.7 ± 0.1 $M_\odot$ and a companion announced by Galland et al. (2006) with $M_2 \sin i = 25 M_\oplus$.
- HD 190228 is a G5 IV star of mass 0.83 $M_\odot$ and a planetary-mass companion with $M_2 \sin i = 3.58 M_\oplus$ (Perrier et al. 2003). Zucker & Mazeh (2001) performed an astrometric analysis using Hipparcos and derived $M_2 = 67 ± 29 M_\oplus$ at 95% confidence based on the spectroscopic elements of Sivan et al. (2004); see also Sect. 3.8.2).
- HD 191760 is a G3 IV star of mass 1.26 ± 0.06 $M_\odot$. Companion announced by Jenkins et al. (2009) with $M_2 \sin i = 38.17 M_\oplus$.
- HD 202206 is a G6 V star of mass 1.15 $M_\odot$ and has two companions with minimum masses in the planetary (HD 202206c) and brown dwarf (HD 202206b, $M_2 \sin i = 17.4 M_\oplus$) discovered by Udry et al. (2002) and Correia et al. (2005).

We identified our target sample of stars with potential brown-dwarf companions and obtained the orbital parameters accessible from radial-velocity measurements. As the next step, we search the intermediate astrometric data of the new Hipparcos reduction for signatures of the orbital motion corresponding to the radial-velocity orbit.

3. Combined astrometric analysis

We describe a method to search the intermediate astrometric data (IAD) of the new Hipparcos reduction for the orbital signatures of stars, whose spectroscopic elements are known from a reliable radial-velocity solution. The significance of the derived orbit is inferred from the distribution-free permutation test (Good 1994; Zucker & Mazeh 2001).
Fig. 7. Eccentricities (left) and velocity semi-amplitudes (right) of 10000 Monte Carlo simulations performed to find the radial-velocity orbit of HD 43848. The bi-modal structure is clearly seen in both histograms. The final solution is found by considering only the lower eccentricity realisations, which is done by separating the distributions at $\hat{K}_1 = 910 \text{ m s}^{-1}$ and correspondingly at $\hat{e} \sim 0.78$. The probability of obtaining a high- or low-eccentricity solution is calculated by counting the number of realisations above and below this threshold value and equals 77% and 23%, respectively.

Fig. 8. CaII H emission region for HD 89707 showing slight re-emission at the bottom of this absorption line.

Fig. 9. CaII H emission region for HD 211847 showing re-emission at the bottom of the absorption line.

Techniques to perform the simultaneous fitting to radial-velocity and Hipparcos astrometry data for substellar companions and verify the statistical confidence of the solution have been developed by Zucker & Mazeh (2000), Pourbaix (2001), Halbwachs et al. (2000), and Zucker & Mazeh (2001). Reffert & Quirrenbach (2006) and Sozzetti & Desidera (2010) used similar approaches to claim the detection of three brown-dwarf companions, although with a significance not exceeding 2-$\sigma$. Independently, successful detections of the astrometric signature of exoplanet candidates using the HST fine guidance sensor (HST FGS) optical interferometer are reported for a few stars with spectroscopic elements known from radial velocity (e.g. Bean et al. 2007; Martioli et al. 2010). Even though the orbital phase is not always fully covered by HST observations, for instance the orbital coverage for HD 33636 is 20% (Bean et al. 2007), their superior precision makes detection possible.

Our analysis is applied to stars with an orbital solution obtained from radial velocities that provides the spectroscopic elements $P, e, T_0, \omega, a$, and $K_1$. The system is modelled as a primary mass $M_1$ orbited by one invisible companion of mass $M_2$. Using Hipparcos astrometric measurements, we seek to determine the two remaining unknown parameters that characterise the orbit, which are the inclination $i$ and the longitude of the ascending node $\Omega$. The analysis is performed in three steps. First, the Hipparcos abscissa, possibly containing the companion signature, is constructed from the catalogue data. This is important, because the IAD contains only abscissa residuals, which are the astrometric residuals after subtraction of the model adopted by the new Hipparcos reduction. Second, the best-fit values for $i$ and $\Omega$ are found by determining the global $\chi^2$ minimum on a two-dimensional search grid. Finally, the significance of the obtained solution is evaluated.

3.1. Constructing the Hipparcos abscissa

On the basis of the Hipparcos identifier of a given star, the astrometric data is retrieved from the new reduction catalogue. It includes the astrometric parameters of the solution, the solution type, and the goodness of fit. The IAD is read from the resrec folder in the catalogue DVD of van Leeuwen (2007) and contains the satellite orbit number, the epoch $t$, the parallax factor $\Pi$, the scan angle orientation $\psi$, the abscissa residual $\delta \Lambda$, and the abscissa error $\sigma_\Lambda$ for every satellite scan.

For the stars in our sample, we encounter the solution types “1”, “5”, “7”, and “9”. The solution type “5” indicates that the five standard astrometric parameters fit the data reasonably well, whereas a solution type “1” is termed a stochastic solution and is adopted when the five-parameter fit is not satisfactory and neither orbital nor acceleration models improve the solution in terms of $\chi^2$. The solution types “7” and “9” are given, when the solution has to include acceleration parameters (proper-motion derivatives of first and second order) to obtain a reasonable fit. The solution types “1”, “7”, and “9” may be the first indications of a possible astrometric perturbation by an unseen companion.
The Hipparcos abscissa $\Delta_{\text{HIP}}$ is then given by:

$$
\Delta_{\text{HIP}} = (a^* + \mu_\alpha \cdot t + \Delta \gamma a^* + \Delta \delta a^*) \cos \psi \\
+ (\delta + \mu_\delta \cdot t + \Delta \gamma \delta + \Delta \delta \delta) \sin \psi + \epsilon \Pi + \delta \Delta, 
$$

(1)

where the modified right ascension is given by $a^* = \alpha \cos \delta$, the declination $\delta$, the associated proper motions $\mu_\alpha$ and $\mu_\delta$, and the parallax $\epsilon$. The epoch $t$ is given as time in years since $t_{\text{HIP}} = 1991.25$ A.D. = 2 448 348.8125 U and the correction factors for solution types “7” and “9” are given by

$$
\Delta_7 p = \mu_\alpha \left( t^2 - 0.81 \right) / 2
$$

(2)

and

$$
\Delta_9 p = \mu_\delta \left( t^2 - 1.69 \right) / 6,
$$

(3)

where $p$ is either $a^*$ or $\delta$ (van Leeuwen 2007), and $\mu_\alpha$ and $\mu_\delta$ are the acceleration and the change in acceleration in the coordinate direction $p$, respectively. For solution types “1” and “5”, $\Delta_7 p$ and $\Delta_9 p$ vanish and the five parameter model holds.

Because of the linear nature of Eq. (1), it is usually convenient to decompose each of the five standard parameters into a constant term and an offset, e.g. $a^* = a^*_0 + \Delta a^*$, and to set the constant terms to zero to compute the abscissa. As explained in Sect. 3.4, however, our formalism requires the value of the parallax $\epsilon$ to be present in Eq. (1), whereas we can safely set $a^*_0 = \delta_0 = \mu_\alpha r_0 = \mu_\delta \delta_0 = 0$ for the calculation of the abscissa.

### 3.2. Model function

The model function describes the astrometric signature along the scan axis $\phi$ of an isolated star, orbited by one invisible companion. To develop our approach, we list the standard formulae (e.g. Hilditch 2001) applicable to this system. The astrometric motion of the star is characterised by $P$, $e$, $T_0$, $i$, $\omega$, $\Omega$, and the astrometric semimajor axis $a$, expressed in angular units. We write the Thiele-Innes constants as

$$
A = a c_A, \quad \alpha = \cos \Omega \cos \omega - \sin \Omega \sin \omega \cos i, \\
B = a c_B, \quad \beta = \sin \Omega \cos \omega + \cos \Omega \sin \omega \cos i, \\
F = a c_F, \quad \gamma = - \cos \Omega \sin \omega - \sin \Omega \cos \omega \cos i, \\
G = a c_G, \quad \delta = - \sin \Omega \sin \omega + \cos \Omega \cos \omega \cos i,
$$

(4)

where $c_{A,B,F,G}$ are functions of $i$, $\Omega$, and $\omega$, and $a$ is expressed in milli-arcseconds (mas). The elliptical rectangular coordinates $X$ and $Y$ are functions of eccentric anomaly $E$ and eccentricity

$$
E - e \sin E = \frac{2\pi}{P} (t - T_0),
$$

(5)

$$
X = \cos E - e, \\
Y = \sqrt{1 - e^2} \sin E.
$$

(7)

If the single-companion assumption is true, the Hipparcos abscissa is described by a linear equation with 9 parameters, which are the 5 standard astrometric parameters plus the 4 Thiele-Innes constants

$$
\Delta_{\text{HIP}} = (a^* + \mu_\alpha \cdot t) \cos \psi + (\delta + \mu_\delta \cdot t) \sin \psi + \epsilon \Pi + \delta \Delta,
$$

(8)

The relationship between astrometrically and spectroscopically derived parameters is (Pourbaix 2001)

$$
a \sin i = c_p K_1 P \sqrt{1 - e^2} \sigma, \quad c_p = 3.35729138 \times 10^{-5},
$$

(9)

where $K_1$, $P$, and $\sigma$ are expressed in m s$^{-1}$, years, and mas, respectively. Hence, the astrometric model function can also be written as a non-linear function of 12 independent parameters

$$
\Delta_{\text{HIP}} = \Delta_{\text{HIP}} \left( a^*, \delta, \omega, P, T_0, i, \epsilon, \Pi, \Omega, K_1 \right).
$$

(10)

The observed radial velocity $v_{\text{rad}}$ of this orbit is given by

$$
v_{\text{rad}} = K_1 \cos \theta \cos \epsilon + e \cos \omega \sin \gamma,
$$

(11)

where $\gamma$ is the systemic velocity. The true anomaly $\theta$ and the semi-amplitude $K_1$ are constrained by

$$
\tan \left( \frac{\theta}{2} \right) = \frac{1 + e}{1 - e} \tan \left( \frac{E}{2} \right),
$$

(12)

$$
K_1 = \sqrt{2} \left( \frac{P}{1 - e^2} \right)^{1/2}.
$$

(13)

### 3.3. Definition of the search grid

One technique to find a possible orbital signature consists of defining a two-dimensional search grid in inclination $i$ and ascending node $\Omega$ to solve for the remaining parameters of the model function (Zucker & Mazeh 2000; Reffert & Quirrenbach 2006). In this way, the goodness-of-fit is obtained for each point on the grid and the optimal solution can be identified as the global $\chi^2$ minimum. The search grid is defined by $i = 0 – 180^\circ$ and $\Omega = 0 – 360^\circ$, to avoid negative semimajor axes and sample the complete range of values of the cosine function (cf. Eqs. (4) and (9)). We use a grid size of $30 \times 30$, which provides a reasonable compromise between computation time and grid resolution.

### 3.4. Joint confidence intervals

For each point on the $i$-$\Omega$-grid, we find the least squares solution by solving a linear equation using matrix inversion, where $P$, $e$, $T_0$, $i$, $\Omega$, and $K_1$ are fixed. The $\chi^2$ value hence obtained for each grid-point allows us to derive joint confidence intervals on the $i$-$\Omega$-grid. The parallax dependence of the Thiele-Innes constants is removed by writing

$$
a = a_i \sigma, \quad a_i = c_p K_1 P \sqrt{1 - e^2} \frac{1}{\sin i},
$$

(14)

and defining the new constant

$$
\Upsilon = a_1 \left[ (c_B X + c_G Y) \cos \psi + (c_A X + c_F Y) \sin \psi \right].
$$

(15)

We note that $a_1$ is expressed in AU and corresponds to the angular semimajor axis $a$ of the astrometric orbit. We can now rewrite Eq. (8) as

$$
\Delta_{\text{HIP}} = (a^* + \mu_\alpha \cdot t) \cos \psi + (\delta + \mu_\delta \cdot t) \sin \psi + \epsilon \Pi + \delta \Delta
$$

(16)

This relation is linear in the five remaining free parameters ($a^*$, $\delta$, $\epsilon$, $\Pi$, $\Omega$) and can easily be solved analytically. The companion signature is solely contained in an additive modulation $\Upsilon(i, \Omega, e, P, T_0, \omega, K_1)$ of the parallax factor. It is important to realise that the presence of $\Upsilon$ forces the parallax (instead of a parallax offset) to be present in the model function, which originates in the abscissa reconstruction from Eq. (1). For the remaining parameters, it is sufficient to consider the offsets $\Delta a^*$, $\Delta \delta$, $\Delta \epsilon$, $\Delta \Pi$, and $\Delta \Omega$ to the catalogue values because the abscissa was constructed with $a^*_0 = \delta_0 = \mu_\alpha r_0 = \mu_\delta \delta_0 = 0$.

For every combination of $i$ and $\Omega$, the linear equation Eq. (16) is solved, yielding the five astrometric parameters and the corresponding $\chi^2$. The best-fit parameters are identified by

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the minimum $\chi^2$ value on the $i$-$\Omega$-grid. The results of the linear adjustment are not directly used in the quoted final solution, but they serve as a consistency check of the results from the non-linear fitting and are used for graphical illustration of the joint confidence intervals shown in Fig. 12. Low-significance orbits can show several local $\chi^2$-minima corresponding to approximately opposite orbit orientations (see also e.g. Zucker & Mazeh 2001; Reffert & Quirrenbach 2006). We find that moderate- and high-significance orbits have one global $\chi^2$-minimum and the confidence contours cover a small area of the $i$-$\Omega$-space (similarly to HD 53680 in Fig. 12). The solution parameters corresponding to an opposite-orientation orbit are always beyond the 4-$\sigma$ contour, i.e. all orbital parameters of significant solutions are unambiguously determined.

3.5. Determining the final solution

While the solution method discussed in the previous section is favourable because the model function is linear, its accuracy is limited by the resolution of the $i$-$\Omega$-grid. For instance, the resolution of a square grid with 900 points is $6^\circ$ and $12^\circ$ in $i$ and $\Omega$, respectively. To avoid this limitation, we determine the optimal solution by means of $\chi^2$-minimisation of a model function, where $i$ and $\Omega$ are free parameters and thus can adopt continuous values. For this purpose, the $i$-$\Omega$-grid is used to define the starting values of a non-linear least squares fit by the Levenberg-Marquardt method. The model function is given by Eq. (10), but is limited to seven free parameters: the five astrometric parameters $\alpha^e$, $\delta$, $\epsilon$, $\mu_x$, $\mu_y$, plus $i$ and $\Omega$. We select the non-linear solution yielding the smallest $\chi^2$ and perform 1000 Monte Carlo simulations. Every Monte Carlo realisation consists of generating a set of Hipparcos abscissa measurements and consequent $\chi^2$-minimisation of the non-linear model and therefore provides 1000 sets of solution elements.

To incorporate the uncertainties in the spectroscopic parameters, we perform the complete analysis of 100 sets of spectroscopic parameters, where each set is randomly drawn from a Gaussian distributions with mean and standard deviation given by the radial-velocity solution and its error, respectively (as in Sozzetti & Desidera 2010). The final solution accounts for all 1000 Monte Carlo solutions, which are obtained for each of the 100 sets of spectroscopic parameters. To achieve this, we combine all solutions to yield distributions of 100,000 values for each parameter in Eq. (10). The final parameter value and its error is derived from the mean and the 1-$\sigma$ confidence interval of the associated distribution, respectively. The values and errors of derived quantities, such as the companion mass, are obtained in the same way, thus take into account possible correlations between individual parameters.

3.6. Outliers in the astrometric data

Outlying astrometric measurements are examined if they deviate by more than 4-$\sigma$ from the orbital solution. Here, $\sigma$ is calculated as the root-mean-square of the fit residuals. The datapoint is discarded if there are at least three measurements of the same satellite orbit number and if it is more than 3-$\sigma$ away from these datapoints. If fewer than three datapoints of the same orbit number are present, they are all removed. The analysis is then iterated. This procedure is similar to the one applied by Pourbaix & Jorissen (2000) and in some cases also agrees with the outlier rejection, which is applied by the new Hipparcos reduction itself. Outliers are removed for the objects GJ 595, HD 3277, HD 43848, HD 74842, HD 154697, HD 164427A, HD 167665, HD 191760, and HIP 103019. This never requires more than one iteration.

3.7. Statistical significance of the orbit

Orbital solutions can be found for any set of astrometric data and the crucial step in the analysis is to determine the credibility of the derived orbits. We accomplish this by applying two statistical tests: the F-test for its simplicity and the permutation test because of its lack of underlying assumptions.

The F-test is extensively used for the statistical analysis of astrometric orbit signatures in Hipparcos data (Pourbaix & Jorissen 2000; Pourbaix & Arenou 2001; Pourbaix 2001; Reffert & Quirrenbach 2006), although it relies on the assumption that the measurement errors are Gaussian. That this assumption is not necessarily fulfilled for Hipparcos data is mentioned by Zucker & Mazeh (2001), who do not apply the F-test to their analysis. The F-test yields the probability that the null assumption that no orbital motion present is true, by comparing the $\chi^2$-values of two models with differing number of parameters. In our case, we consider the $\chi^2$-value of the 7-parameter solution by non-linear minimisation for each of the 100 draws from the spectroscopic elements. The F-test is applied to the median $\chi^2$-value with respect to the $\chi^2$ value, obtained from the standard 5-parameter solution, and results in the null probability, which is tabulated in Table 8 and gives the probability that the simpler model is valid. We use the F-test as an additional indicator of the statistical confidence of our solution, but the conclusive argument is derived from the permutation test.

The permutation test has the important advantage of being part of the distribution-free tests (Zucker & Mazeh 2001) and therefore does not rely on the assumption of Gaussian error distribution. For our purpose, it can be summarised by the idea that a periodic signal present in a given data set will be destroyed by a random permutation of the individual datapoints in all but a few cases. In the present analysis, the periodicity is defined by the radial velocity solution. To derive a significance, the semimajor axis of the solution orbit, derived with the original astrometric data, is compared to the semimajor axes of pseudo-orbits, which are derived from randomly permuted astrometric data. If a signal is present in the astrometric data, it will be destroyed by the random permutations and the solution orbit will be easily distinguishable in the pseudo-orbits because of its large amplitude.

Because of the pre-reduction present in the original Hipparcos IAD (ESA 1997), its permutation had to be performed with great care (Zucker & Mazeh 2001). Fortunately, this became considerably easier with the availability of the new Hipparcos reduction IAD. The permutation is applied only to the list of abscissa residuals $\Delta\alpha$ while taking the possible contribution of the 7- and 9-parameter solution into account. The remaining data, i.e. epochs, parallax factors, and scan angle orientations are left unchanged. To obtain the pseudo-orbits, 1000 permutations are performed and the analysis is run on the synthesised data, using the nominal spectroscopic elements. The statistical significance of the solution orbit is equal to the percentage of pseudo-orbits, which have a smaller semimajor axis than the solution orbit (Zucker & Mazeh 2001). Examples of high- and low-significance orbits are shown in Fig. 13. In contrast to Halbwachs et al. (2000), the analysis of the pseudo-orbits is performed with spectroscopic parameters identical to the nominal ones, hence the method of comparing the distribution of $a/\sigma_a$ to the expected Rayleigh-Rice law is not applicable.
3.8. Validation of the method

The combined analysis of radial-velocity and astrometric data can lead to erroneous conclusions, if the results are not carefully checked (Han et al. 2001; Pourbaix & Arenou 2001). In particular, the relationship between the inclination and the semimajor axis has to be investigated during the interpretation, since over- and underestimation of these parameters during the fitting process affects the derived companion masses. To permit us to have high confidence in our results, we validate our method by both simulation and comparison with published work.

3.8.1. Simulation

Simulations are employed to investigate the effect of Hipparcos measurement precision on the derived orbital parameters and the orbit significance. Of particular interest is the behaviour of the measurement precision on the derived orbital parameters and the orbit significance. Of particular interest is the behaviour of the measurement precision on the derived orbital parameters and the orbit significance.

Instead of modelling a complete dataset, we chose the case of HD 17289, whose orbit is discussed in Sect. 4 and for which the Hipparcos observations cover twice the orbital period. To inject the astrometric signature of a companion, we use the satellite configuration for this star, i.e. the scan angles, the parallax factors, and the measurement precision. Noise is generated by adding a random bias term to each measurement, which is drawn from a Gaussian distribution with a standard deviation given by the satellite measurement precision. Configurations with different $S/N$ are generated by altering the radial-velocity semi-amplitude $K_1$ of the simulated orbit, while keeping all other parameters fixed. The emulated Hipparcos errors in this simulation are consequently Gaussian, which is a simplifying assumption and may not be fulfilled in reality. However, the simulation is valid for investigating the biases of the method.

The results are illustrated in Figs. 10 and 11, where the accuracy of the determined values for inclination and semimajor axis are displayed as a function of statistical significance, derived from the permutation test. The accuracy is expressed in terms of the difference between determined and true value, i.e. for the inclination it is computed as $i_{\text{derived}} - i_{\text{true}}$. In general, the accuracy of the derived $a$ and $i$ increases with $S/N$ and significance. For significances above 2-$\sigma$, the accuracy is comparable to the precision. However, the semimajor axis retains a small bias even at high $S/N$, whose magnitude and sign is influenced by the satellite configuration concerning orbit orientation and period. The determination of the orbit inclination greatly improves with significance. To investigate whether our analysis is biased even when exceeding the 3-$\sigma$ significance, we repeated the simulation for 3 different satellite configuration, i.e. three different stars. We find no sign of a systematic bias, favouring large or small inclinations and semimajor axes. Above a 3-$\sigma$ significance, derived and true values always agree within the error bars. This simulation shows that a large bias can be introduced in the solution, if an orbit significance of 1-$\sigma$ is used to qualify a detection. In contrast, the biases are small if higher significance is requested.

We find that an orbit is detected at greater than 3-$\sigma$ significance, when its semimajor axis amounts to $\sim$70% of the measurement uncertainty. This result may depend on the assumption of Gaussian errors and the specific satellite configuration, but it confirms the warning of Pourbaix (2001) that derived orbits with sizes comparable to the instrument precision have to be evaluated very carefully.

![Fig. 10. Simulated difference between the derived semimajor axis $a_{\text{derived}}$ and its true value $a_{\text{true}}$. Top: $a_{\text{derived}}$ as function of $a_{\text{true}}$. Filled circles, open circles, and diamonds designate orbits of 1–2$\sigma$, 2–3$\sigma$, and >3$\sigma$ significance, respectively. The solid line has unity slope and the dashed line shows the single-measurement precision of $\sigma_a = 7.2$ mas. Bottom: the relative deviation calculated as $(a_{\text{derived}} - a_{\text{true}})/a_{\text{true}}$ as a function of significance. The accuracy and relative precision of the solution improve with significance. At the detection limit of 3-$\sigma$, the deviation from the true value is smaller than 0.5 mas.](image1)

![Fig. 11. Simulated difference between the derived inclination $i$ and its true value. The accuracy and the precision of the solution improve with significance. At the detection limit of 3-$\sigma$, the deviation from the true value is smaller than 0.02°.](image2)

3.8.2. Comparison sample

For independent validation of our method and to investigate the impact of the new reduction, we screened the literature for objects for which a similar analysis was performed. We chose the three stars HR 6046, GJ 1069, and HD 190228 with companions found by Torres (2007), Halbwachs et al. (2000),...
and Zucker & Mazeh (2001)\(^3\), respectively, that exhibit diverse orbit inclinations and companion masses. All these works use the original IAD, whereas we use the new reduction IAD.

For all three stars, good agreement is found between our results and the published works, both in terms of value and 1-σ results and the published works, both in terms of value and 1-σ uncertainty in semimajor axis and companion mass, as shown in Table 6. The orbit significances derived from pseudo-orbits given by Zucker & Mazeh (2001) for two objects match our results. For HD 190228, we find the same solution as Zucker & Mazeh (2001) and a 2-σ significance (see Sect. 4.2 for further discussion of this object). We confirm the detection of the GJ 19832 orbit, caused by a low mass star (Zucker & Mazeh 2001; Halbwachs et al. 2000), and the binary orbit of HR 190228 has also been analysed by Pourbaix & Arenou (2001). Halbwachs et al. 2000), and the binary orbit of HR 6046 (Torres 2007) at greater than 3-σ significance. In particular, the case of GJ 1069, which is now analysed by three independent teams finding essentially identical results, attests that our analysis is robust. Both the simulation results and the analysis of a sample of comparison stars confirm that our method for the combined analysis of radial-velocity orbits and Hipparcos astrometry is correct. The consideration of the orbit significance ensures that the derived result is bias-free to a reasonable extent.

### 3.9. Companion mass limits

The combination of astrometric and spectroscopic data makes it possible to find a complete orbital solution. However, the result

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**Table 6.** Results for the comparison sample.

<table>
<thead>
<tr>
<th>Object</th>
<th>HIP</th>
<th>(a) (mas)</th>
<th>(M_2) ((M_\odot))</th>
<th>Significance</th>
<th>Ref.</th>
</tr>
</thead>
<tbody>
<tr>
<td>HD 190228</td>
<td>98714</td>
<td>(1.8^{+0.5}_{-0.5}) [1.8 ± 0.8]</td>
<td>(65.7^{+20.4}<em>{-20.4}) [67.1^{+30.4}</em>{-28.3}]</td>
<td>96.3</td>
<td>[95]</td>
</tr>
<tr>
<td>GJ 1069</td>
<td>19832</td>
<td>(17.7^{+0.8}_{-0.8}) [17.1 ± 0.8]</td>
<td>(263.2^{+18.7}<em>{-18.7}) [253.0^{+14.7}</em>{-15.3}]</td>
<td>&gt;99.9</td>
<td>[99.9]</td>
</tr>
<tr>
<td>GJ 1069</td>
<td>19832</td>
<td>(17.7^{+0.8}_{-0.8}) [17.2 ± 0.8]</td>
<td>(263.2^{+18.7}_{-18.7}) [256.7 ± 14.7]</td>
<td>&gt;99.9</td>
<td>[...]</td>
</tr>
<tr>
<td>HD 145849</td>
<td>79358</td>
<td>(11.1^{+0.7}_{-0.7}) [9.69 ± 0.85]</td>
<td>(1443.5^{+18.5}<em>{-20.6}) [1424.8^{+73.3}</em>{-31.4}]</td>
<td>&gt;99.9</td>
<td>[...]</td>
</tr>
</tbody>
</table>

**Notes.** The values in square brackets are the corresponding values from the reference paper. Our error in the companion mass does not include the uncertainty in the primary mass. \(^{\text{a}}\) Primary mass from Sivan et al. (2004).

**References.** (1) Zucker & Mazeh (2001); (2) Halbwachs et al. (2000); (3) Torres (2007).
can be misleading or wrong, if the significance of the solution is not verified. For significant orbits, the derived parameters, e.g., the companion mass, and their errors are valid. For low-significance orbits, the derived parameters and their errors are very probably false. Therefore, the distribution function for the companion mass of a low-significance orbit cannot be used to derive upper and lower mass limits.

The question is then turned to the definition of a significant orbit. In the literature, various significance estimators (e.g. F-test or permutation test) and detection limits are employed. We use a 3-$\sigma$ detection limit based on the permutation test and deem orbits at less than 2-$\sigma$ unreliable. For a 2-$\sigma$ orbit, we adopt a reduced validity and expect that the mass limits do hold.

In principle, it is possible to derive the upper mass limit for any companion by employing simulations similar to the one described in Sect. 3.8.1. This simulation is extremely elaborate and its application to the list of host stars with undetected astrometric motion is beyond the scope of this paper. However, we have demonstrated that an orbit is detected at greater than 3-$\sigma$ significance when its semimajor axis amounts to $\sim$70% of the measurement precision, if the Hipparcos observations cover one complete orbit. Including an additional confidence margin, we thus claim that the signature of a fully covered orbit would have been detected, if its semimajor axis equaled the single-measurement precision. Hence, an upper limit to the companion mass can be set by enforcing the equation $a = \sigma_A$ for orbits with significances smaller than 2-$\sigma$ and $N_{\text{rob}} > 1$.

4. Results

For all objects in our sample, we are able to obtain an orbital solution from the combination of the radial-velocity orbit and the Hipparcos astrometric data, but not all of these objects are credible. To distinguish a real orbit from meaningless output of the fitting procedure, we use the permutation test to evaluate the orbit significance. This results in the detection of 6 orbits having high significance above 3-$\sigma$ (Tables 7 and 8). The orbits of 5 objects have moderate significance between 2-$\sigma$ and 3-$\sigma$. Three orbits have low significance between 1-$\sigma$ and 2-$\sigma$ and the orbits of 17 objects have very-low significance (<1-$\sigma$, Table 10)².

The Tables 3, 9, and 10 list the spectroscopic and astrometric elements of the adopted orbital solution together with the orbit significance and additional information: $N_{\text{rob}}$ is the number of orbits covered by the Hipparcos observations, $\sigma_A$ is the median Hipparcos single-measurement precision, $a \sin i$ is the minimum astrometric semimajor axis, $M_2$ (3-$\sigma$) are upper and lower companion-mass limits at 3-$\sigma$, $\Delta q_{\text{red}}$ is the semimajor axis of the relative orbit, $x_{\text{red}}$ is the reduced chi-square value of the adopted 7-parameter solution, and the null probability gives the result of the F-test. The quoted errors correspond to Monte Carlo based 1-$\sigma$ confidence intervals.

For orbits with significances below 2-$\sigma$ and where the orbital period is fully covered by Hipparcos observations, we derived the upper limit to the companion mass $M_{\text{2,up-lim}}$ by enforcing the equation $a = \sigma_A$ (cf. Sect. 3.9).

4.1. Orbits with high significance exceeding 3-$\sigma$

Clear orbital signatures are detected for six target stars. The orbit significances exceed 99.7%, which we adopt as a criterion for doubtful detection. The orbits show a large variety in size ($a = 2.3–35$ mas) and inclination ($i = 12–173^\circ$) and reveal companion masses ranging from 90 $M_J$ to 0.52 $M_\odot$. The obtained relative precision of the companion mass is typically of the order of 10% and does not include any contribution by the primary-mass uncertainty. The visualisation of Hipparcos astrometric data is not very intuitive, because positions are measured in one dimension ² As motivated in Sect. 2.3, we do not present an astrometric analysis for HD 4747 and HD 211847 and therefore the number of objects listed here adds up to 31 instead of 33.

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References. (1) Minniti et al. (2009); (2) Nidever et al. (2002); (3) Perrier et al. (2003); (4) Wright et al. (2009); (5) Jenkins et al. (2009).

Table 8. Solution parameters and companion masses for stars with orbits that have significances above 1-σ.

<table>
<thead>
<tr>
<th>Object</th>
<th>Sol. type</th>
<th>$N_{orb}$</th>
<th>$\sigma_{\alpha}$ (mas)</th>
<th>$a_{sin i}$ (mas)</th>
<th>$a$ (mas)</th>
<th>$M_2$ ($M_\odot$)</th>
<th>$M_2$ (3-σ) ($M_\odot$)</th>
<th>$a_{rel}$ (mas)</th>
<th>$\chi^2_{red}$</th>
<th>Null prob. (%)</th>
<th>Significance</th>
</tr>
</thead>
<tbody>
<tr>
<td>HD 17289</td>
<td></td>
<td>1</td>
<td>2.1</td>
<td>7.2</td>
<td>1.29</td>
<td>10.8 $^{+7.0}_{-4.7}$</td>
<td>547.4 $^{+277.4}_{-173.5}$</td>
<td>(386, 700)</td>
<td>32.0</td>
<td>0.32</td>
<td>7e-29</td>
</tr>
<tr>
<td>HD 30501</td>
<td></td>
<td>5</td>
<td>0.6</td>
<td>3.3</td>
<td>10.44</td>
<td>14.0 $^{+1.8}_{-1.9}$</td>
<td>89.6 $^{+12.3}_{-12.5}$</td>
<td>(65, 132)</td>
<td>147.7</td>
<td>0.79</td>
<td>0.005</td>
</tr>
<tr>
<td>HD 43848*</td>
<td></td>
<td>9</td>
<td>0.5</td>
<td>4.2</td>
<td>9.9 $^{+1.2}_{-1.3}$</td>
<td>101.8 $^{+10.4}_{-15.0}$</td>
<td>(59, 150)</td>
<td>90.8</td>
<td>1.09</td>
<td>4e-8</td>
<td>&gt;999</td>
</tr>
<tr>
<td>HD 43848b</td>
<td></td>
<td>9</td>
<td>0.5</td>
<td>4.2</td>
<td>9.3</td>
<td>10.5 $^{+13.1}_{-13.3}$</td>
<td>120.2 $^{+17.1}_{-17.2}$</td>
<td>(73, 179)</td>
<td>91.7</td>
<td>1.07</td>
<td>1e-8</td>
</tr>
<tr>
<td>HD 53680</td>
<td></td>
<td>1</td>
<td>0.7</td>
<td>5.7</td>
<td>9.73</td>
<td>34.9 $^{+3.2}_{-3.2}$</td>
<td>226.7 $^{+20.4}_{-20.9}$</td>
<td>(159, 301)</td>
<td>160.7</td>
<td>0.82</td>
<td>3e-16</td>
</tr>
<tr>
<td>HD 164427A</td>
<td></td>
<td>5</td>
<td>1.1</td>
<td>2.4</td>
<td>0.94</td>
<td>23 $^{+0.5}_{-0.5}$</td>
<td>269.9 $^{+63.2}_{-65.2}$</td>
<td>(101, 482)</td>
<td>12.8</td>
<td>1.15</td>
<td>0.003</td>
</tr>
<tr>
<td>HIP 103019</td>
<td></td>
<td>5</td>
<td>1.2</td>
<td>7.1</td>
<td>3.33</td>
<td>10.8 $^{+1.3}_{-1.3}$</td>
<td>188.1 $^{+26.5}_{-26.4}$</td>
<td>(118, 280)</td>
<td>53.4</td>
<td>1.13</td>
<td>4e-10</td>
</tr>
<tr>
<td>HD 3277</td>
<td></td>
<td>5</td>
<td>24.8</td>
<td>3.6</td>
<td>0.58</td>
<td>2.5 $^{+0.5}_{-0.5}$</td>
<td>344.2 $^{+76.1}_{-76.1}$</td>
<td>(143, 600)</td>
<td>9.4</td>
<td>1.0</td>
<td>5e-4</td>
</tr>
<tr>
<td>HD 131664</td>
<td></td>
<td>5</td>
<td>0.6</td>
<td>4.7</td>
<td>0.88</td>
<td>4.3 $^{+0.6}_{-0.7}$</td>
<td>89.6 $^{+16.4}_{-16.6}$</td>
<td>(44, 153)</td>
<td>60.3</td>
<td>0.77</td>
<td>2e-5</td>
</tr>
<tr>
<td>HD 154697</td>
<td></td>
<td>7</td>
<td>0.6</td>
<td>3.1</td>
<td>6.26</td>
<td>11.7 $^{+17}_{-17}$</td>
<td>151.9 $^{+24.9}_{-25.0}$</td>
<td>(81, 235)</td>
<td>89.2</td>
<td>1.60</td>
<td>1e-6</td>
</tr>
<tr>
<td>HD 174457</td>
<td></td>
<td>7</td>
<td>1.3</td>
<td>3.2</td>
<td>1.77</td>
<td>2.8 $^{+0.6}_{-0.6}$</td>
<td>107.8 $^{+23.8}_{-24.1}$</td>
<td>(65, 190)</td>
<td>35.2</td>
<td>1.22</td>
<td>12.3</td>
</tr>
<tr>
<td>HD 190228</td>
<td></td>
<td>5</td>
<td>1.0</td>
<td>3.5</td>
<td>0.13</td>
<td>1.8 $^{+0.5}_{-0.5}$</td>
<td>49.4 $^{+14.3}_{-14.3}$</td>
<td>(10, 99)</td>
<td>32.9</td>
<td>1.08</td>
<td>0.3</td>
</tr>
<tr>
<td>HD 74014</td>
<td></td>
<td>5</td>
<td>0.1</td>
<td>2.5</td>
<td>9.84</td>
<td>64.0 $^{+32.3}_{-33.4}$</td>
<td>421.3 $^{+25.3}_{-25.2}$</td>
<td>(51, 1568)</td>
<td>235.4</td>
<td>1.34</td>
<td>19.9</td>
</tr>
<tr>
<td>HD 168443</td>
<td></td>
<td>5</td>
<td>0.6</td>
<td>1.9</td>
<td>1.25</td>
<td>2.4 $^{+0.6}_{-0.6}$</td>
<td>34.3 $^{+31}_{-31}$</td>
<td>(18.7, 67)</td>
<td>75.9</td>
<td>1.53</td>
<td>3.6</td>
</tr>
<tr>
<td>HD 191760</td>
<td></td>
<td>5</td>
<td>1.6</td>
<td>3.3</td>
<td>0.46</td>
<td>1.9 $^{+0.8}_{-0.8}$</td>
<td>185.3 $^{+94.1}_{-94.1}$</td>
<td>(39, 659)</td>
<td>16.0</td>
<td>1.39</td>
<td>27.6</td>
</tr>
</tbody>
</table>

Notes. The top 6 orbits have high significance >99.7% and are therefore fully characterised. The bottom 3 orbits are of low significance (1-σ–2-σ) and represent the formal solution. The remaining 5 orbits have moderate significance (2-σ–3-σ) and provide lower and upper mass-limits for the companions. (a) Final solution at lower eccentricity. (b) Formal solution at higher eccentricity.

Table 9. Spectroscopic elements of literature targets with orbits of greater than 1-σ significance.

<table>
<thead>
<tr>
<th>Object</th>
<th>HIP</th>
<th>$M_1$ ($M_\odot$)</th>
<th>$M_2$ sin $i$ ($M_\odot$)</th>
<th>$P$ (day)</th>
<th>$e$</th>
<th>$K_1$ (m s$^{-1}$)</th>
<th>$T_0$ (JD$^+$)</th>
<th>$\omega$ (deg)</th>
<th>Ref.</th>
</tr>
</thead>
<tbody>
<tr>
<td>HD 131664</td>
<td>73408</td>
<td>1.1</td>
<td>18.2</td>
<td>19510.0 ± 41.0</td>
<td>0.64 ± 0.02</td>
<td>359.5 ± 22.3</td>
<td>52 060.0 ± 41.0</td>
<td>1497.0 ± 1.0</td>
<td>1</td>
</tr>
<tr>
<td>HD 174457</td>
<td>92418</td>
<td>1.2</td>
<td>65.8</td>
<td>840.8 ± 0.1</td>
<td>0.23 ± 0.01</td>
<td>12500.0 ± 10.0</td>
<td>52 020.0 ± 41.0</td>
<td>1390.0 ± 1.0</td>
<td>2</td>
</tr>
<tr>
<td>HD 190228</td>
<td>98714</td>
<td>0.8</td>
<td>11.6</td>
<td>9160.0 ± 16.0</td>
<td>0.50 ± 0.04</td>
<td>910.0 ± 5.0</td>
<td>51 236.0 ± 25.0</td>
<td>100.7 ± 3.2</td>
<td>3</td>
</tr>
<tr>
<td>HD 168443</td>
<td>89844</td>
<td>1.0</td>
<td>18.1</td>
<td>1748.2 ± 1.0</td>
<td>0.212 ± 0.001</td>
<td>298.1 ± 0.6</td>
<td>53 769.8 ± 0.0</td>
<td>64.7 ± 0.5</td>
<td>4</td>
</tr>
<tr>
<td>HD 191760</td>
<td>99661</td>
<td>1.2</td>
<td>38.0</td>
<td>505.6 ± 0.4</td>
<td>0.63 ± 0.01</td>
<td>10478.0 ± 38.7</td>
<td>54 835.7 ± 2.1</td>
<td>200.4 ± 0.3</td>
<td>5</td>
</tr>
</tbody>
</table>

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References. (1) This work; (2) Nidever et al. (2002); (3) Hatzes et al. (2005); (4) Benedict et al. (2010); (5) Wittenmyer et al. (2009); (6) Liu et al. (2008); (7) Omiya et al. (2009); (8) Butler et al. (2006); (9) Udry et al. (2002); (10) Galland et al. (2006); (11) Correia et al. (2005).
only, i.e. along the scan-angle orientation $\psi$. We choose a representation similar to Torres (2007) and the resulting stellar astrometric orbits are shown in Fig. 20. For a clearer representation, we compute normal points for each satellite orbit number. Their values and errors are given by the mean and standard deviation of the Hipparcos abscissae obtained during one given satellite orbit, respectively. Normal points are used only for the visualization and not during the orbit adjustment. We prove that orbit detections are also possible in apparently difficult conditions, such as incomplete orbit coverage by the satellite ($N_{\text{tot}} = 0.5$ for HD 43848) and an instrument precision comparable to the the orbit size (HD 164427A). The final 7-parameter fit for the detected orbits has an acceptable reduced chi-square value ($\chi^2_{\text{red}} = 0.8–1.2$) in all but one case. For HD 17289, this value is very small at 0.3 and the Hipparcos precision is unusually large at 7.2 mas, which may indicate that the astrometric errors for this star are overestimated. As for the radial-velocity measurements, the light emitted by the stellar companion of this star may have disturbed the Hipparcos observations.

Three out of these six stars have non-standard Hipparcos solutions. Thus, first indications for the astrometric disturbance induced by the companion were already detected during the standard Hipparcos data analysis. An additional outcome of the companion solution are refined positions, distances, and proper motions of the objects (Table 7). For five stars, the parallax estimate becomes more precise, whereas its value is compatible within the error bars. Only the parallax of HIP 103019 is 1.9 mas larger than the value given by the new reduction. Changes in position and proper motion are of the order of 1 mas and 1 mas yr$^{-1}$, respectively, with extrema of 41 mas and 22 mas yr$^{-1}$, respectively. Our results for the individual objects are described in the following:

- HD 17289: The orbit inclination is $173.2 \pm 0.5^\circ$ and the companion mass is $0.52 \pm 0.05 M_{\odot}$, which makes it the most massive companion detected in this study. HD17289 is thus identified as a binary star with mass ratio $q = 1.97\pm 0.06$. Because the orbit is completely characterised, we can calculate the projected angular separation between the two components. Along the orbit, their separation on the sky ranges between 15–49 mas. The mass-luminosity relations of Delfosse et al. (2000) estimate an apparent $V$-magnitude of the secondary of $V_2 = 13.18$. The primary magnitude obtained from Hipparcos is $V = 7.56$ (Table 1) and the expected intensity contrast between the secondary and primary is $\sim 1:180$ in the visible. The combination of close separation and moderate intensity contrast causes a small and periodic bias in the radial-velocity measurement (see Sect. 2.3).

- HD 30501: The orbit inclination is $49^\circ \pm 10^\circ$ deg and the companion mass is $90 \pm 12 M_J$, which is the lightest companion detected in this study. This companion is just at the threshold between a very low-mass star and a brown dwarf.

- HD 43848: The orbit inclination is $165^{+3}_{-2}^\circ$ deg and the companion is a very low mass star with mass $0.10 \pm 0.01 M_{\odot}$. This value is slightly lower and more precise than the one derived by Sozzetti & Desidera (2010), who obtain $0.11 \pm 0.04 M_{\odot}$ based on the radial velocities of Minniti et al. (2009). This solution is found with the lower-eccentricity orbit (cf. Sect. 2.2). For completeness, we include the formal astrometric solution for the less probable and higher-eccentricity radial-velocity orbit in Tables 7 and 8. The resulting companion mass obtained in this case would be higher at $M_2 = 0.11 \pm 0.02 M_{\odot}$. Hipparcos astrometry cannot be used to determine the true eccentricity, because the fit quality in terms of $\chi^2$ is comparable in both cases.

- HD 53680: The orbit inclination is $163.6^{+1.2}_{-1.3}$ deg and the companion is a low mass star with mass $0.22 \pm 0.02 M_{\odot}$.

- HD 164427A: The orbit inclination is $11.8^{+3.1}_{-1.4}$ deg and the companion is a low mass star with mass $0.26 \pm 0.06 M_{\odot}$. The stellar companion has also been detected by Zucker & Mazeh (2001), who derived a slightly less inclined orbit and therefore obtained the higher mass of $0.35 \pm 0.09 M_{\odot}$.

- HD 103019: The orbit inclination is $160.9^{+2.7}_{-2.6}$ deg and the companion is a very low mass star with mass $0.18 \pm 0.03 M_{\odot}$.

### 4.2. Orbits with moderate significances between 2-$\sigma$ and 3-$\sigma$

The orbits of 5 stars are found to have significances of 95.4–99.7%. Although we do not consider that these solutions accurately characterise the stellar orbits, the moderate level of significance indicates that the orbital motion is present in the astrometric data. The formal orbits are shown in Fig. 21 and allow us to obtain a visual impression of the fit quality. Four out of five orbits have sizes comparable to or smaller than the instrument precision and the astrometric signal is therefore hidden. The permutation test allows us to assess the credibility of the derived orbit especially in this difficult regime. Since there is signal in the astrometric data, the orbit inclination is constrained and we expect the companion mass to be confined within the 3-$\sigma$ limits given in Table 8. The individual objects are discussed below:

- HD 3277: The companion is of stellar nature and its mass is confined to within $0.14–0.57 M_{\odot}$. The formal companion mass at moderate significance is $0.33 \pm 0.07 M_{\odot}$. The upper mass-limit that we independently derived from $a = \sigma_n = M_{\text{up-lim}} = 0.51 M_{\odot}$, which is compatible with our solution.

- HD 131664: The companion is a very low-mass star or a brown dwarf with a mass of $44–153 M_I$. The available data is not sufficient to distinguish between these two cases. The formal companion mass at moderate significance is $90 \pm 16 M_I$. Although the orbit significance agrees with the result of Sozzetti & Desidera (2010), who obtain a companion mass of $23^{+26}_{-20} M_I$ with 95% confidence, the mass range derived from our analysis is significantly higher.

- HD 154697: The companion is a low-mass star with a mass confined within $0.08–0.22 M_{\odot}$. The formal companion mass at moderate significance is $0.15 \pm 0.02 M_{\odot}$. The longitude of the ascending node $\Omega$ of this orbit is hardly constrained and is correlated with the position offsets $\Delta \Omega^*$ and $\Delta \theta$. The $\Omega–\Delta \Omega^*$-correlation is shown in Fig. 14. Because the adopted solution is derived from the mean of the complete distribution, it does not represent a valid solution with a valid orbit, as is also visible in Fig. 21, where the fit residuals are large. However, this does not affect the validity of the derived companion mass, because that is independent of $\Omega$.

- HD 174457: The companion is a very low-mass star or a brown dwarf. Its mass is confined to within $0.06–0.18 M_{\odot}$. The formal companion mass at moderate significance is $0.10 \pm 0.02 M_{\odot}$. The upper mass-limit that we independently derived from $a = \sigma_n = M_{\text{up-lim}} = 0.12 M_{\odot}$ is compatible with our solution.

- HD 190228: The companion is probably a brown dwarf, although we cannot exclude a very low-mass star. Assuming a primary mass of $0.83 M_{\odot}$ (Perrier et al. 2003), we find the
The analysis of 17 objects yields orbits with significances below 1−σ. Because the astrometric data does not contain orbital information, the derived result is physically meaningless and only a mathematical solution. Therefore we do not list the solution parameters. Table 10 shows the targets and their basic parameters. Two characteristics can be observed that impede the astrometric orbit detection. In the one case, applicable to both HD 65430 and HD 167665, the orbital coverage is poor and the orbit is not detected, although the expected signal is large compared to the instrument precision. In the other case, for instance for GJ 595, HD 13189, and HD 140913, the minimum orbital signature given by \( a \sin i \) is very small compared to the instrument precision.

The companions of six stars have maximum masses of \( M_{\text{2,up-lim}} \) above 0.63 \( M_\odot \), which does not represent a considerable constraint of the system. Nine companions have upper mass-limits of \( M_{\text{2,up-lim}} = 0.11-0.63 M_\odot \), thus are in the M-dwarf mass range. Finally, the mass of the companion of HD 137510 is confined to be between 22.7 \( M_\odot \) and \( M_{\text{2,up-lim}} = 64.4 M_\odot \). Therefore, we find that HD 137510b has to be a brown dwarf. This detection was possible because of the confidence we gained in the capability of Hipparcos to detect astrometric orbital motion. Using a simpler argument, Endl et al. (2004) derived an upper mass-limit of 94 \( M_\odot \), which was not stringent enough to prove the substellar nature of the companion.

This sample also contains HD 38529, for which Benedict et al. (2010) found the orbit inclination of the outer companion HD 38529c and derived its mass of \( M_2 = 17.6 M_\odot \) using HST FGS astrometry. Our analysis made use of the radial velocity orbit given by Benedict et al. (2010), but failed to detect an orbit reliably, although the Hipparcos precision of 1.6 mas is close to the HST orbit size (\( a = 1.05 \pm 0.06 \) mas). However, the Hipparcos orbit coverage is poor at 0.4, which may explain the non-detection.

4.3. Orbits with low significances between 1−σ and 2−σ

The orbits of three stars are found at low significance of 68.3–95.4\%, i.e. the astrometric data contains very little or no orbital signal. As we have shown by simulation (cf. Sect. 3.8.1), orbital solutions at this significance level are prone to large biases. Therefore, the solution parameters shall not be considered valid in these cases. The formal solution of the fitting procedure is given for completeness in Tables 7 and 8. We discuss the individual objects below:

- HD 74014: The poor Hipparcos coverage of 10\% of this long-period orbit impedes the detection of the astrometric orbit. With a minimum semimajor axis of \( a \sin i = 9.8 \) mas, the orbit would have been detected on a longer timespan, given the instrument precision of 2.5 mas. Because the orbit is not fully covered, we cannot set an upper companion-mass limit \( M_{\text{2,up-lim}} \).

- HD 168443: The formal solution for the companion mass of 34 ± 9 \( M_\odot \) is in agreement with the result of Reffert & Quirrenbach (2006), who derive 34 ± 12 \( M_\odot \) using the radial velocities of Marcy et al. (2001) and the original Hipparcos data. From our analysis, using the CORALIE radial velocities and the new Hipparcos reduction IAD, we conclude that the mass obtained for the outer companion HD 168443c is of low confidence. Because the orbit is not fully covered, we cannot set an upper companion-mass limit \( M_{\text{2,up-lim}} \).

- HD 191760: The minimum orbit size of this star is very small (\( a \sin i = 0.46 \) mas) and Hipparcos astrometry does not reveal the signature of the companion discovered by Jenkins et al. (2009). However, we can set an upper mass-limit for the companion of \( M_{\text{2,up-lim}} = 0.28 M_\odot \).

4.4. Orbits with very-low significances and the detection of one brown-dwarf companion

The analysis of 17 objects yields orbits with significances below 1−σ. Because the astrometric data does not contain orbital information, the derived result is physically meaningless and only a mathematical solution. Therefore we do not list the solution parameters. Table 10 shows the targets and their basic parameters. Two characteristics can be observed that impede the astrometric orbit detection. In the one case, applicable to both HD 65430 and HD 167665, the orbital coverage is poor and the orbit is not detected, although the expected signal is large compared to the instrument precision. In the other case, for instance for GJ 595, HD 13189, and HD 140913, the minimum orbital signature given by \( a \sin i \) is very small compared to the instrument precision.

The companions of six stars have maximum masses of \( M_{\text{2,up-lim}} \) above 0.63 \( M_\odot \), which does not represent a considerable constraint of the system. Nine companions have upper mass-limits of \( M_{\text{2,up-lim}} = 0.11-0.63 M_\odot \), thus are in the M-dwarf mass range. Finally, the mass of the companion of HD 137510 is confined to be between 22.7 \( M_\odot \) and \( M_{\text{2,up-lim}} = 64.4 M_\odot \). Therefore, we find that HD 137510b has to be a brown dwarf. This detection was possible because of the confidence we gained in the capability of Hipparcos to detect astrometric orbital motion. Using a simpler argument, Endl et al. (2004) derived an upper mass-limit of 94 \( M_\odot \), which was not stringent enough to prove the substellar nature of the companion.

This sample also contains HD 38529, for which Benedict et al. (2010) found the orbit inclination of the outer companion HD 38529c and derived its mass of \( M_2 = 17.6 M_\odot \) using HST FGS astrometry. Our analysis made use of the radial velocity orbit given by Benedict et al. (2010), but failed to detect an orbit reliably, although the Hipparcos precision of 1.6 mas is close to the HST orbit size (\( a = 1.05 \pm 0.06 \) mas). However, the Hipparcos orbit coverage is poor at 0.4, which may explain the non-detection.

5. Discussion

The compilation of close companions to Sun-like stars with minimum masses of 10–80 \( M_\odot \) of Sozzetti & Desidera (2010) lists 39 objects. With the discovery of 9 new companions, we have increased the known number of these objects by about 20\%. Ten companions in the Sozzetti & Desidera (2010) list are probably stars. We confirm the stellar nature in two cases (HD 43848B and HD 164427B) and detect 6 new stellar companions (HD 3277B, HD 17289B, HD 30501B, HD 53680B, HD 154697B, and HIP 103019B) at the 3−σ significance level. All of these companions are M dwarfs with masses between 90 \( M_\odot \) and 0.52 \( M_\odot \). For 3 companions, we derive mass limits that enclose the boundary between low-mass stars and brown dwarfs: HD 131664 (\( M_2 = 44-153 M_\odot \)), HD 174457 (\( M_2 = 65-190 M_\odot \)), and HD 190228 (\( M_2 = 10-99 M_\odot \)). HD 190228 was initially a planet-host star discovered by Sivan et al. (2004). We analysed it to validate our method and consequently recognised its importance as a promising brown dwarf candidate.

Furthermore, we set upper mass limits below 0.63 \( M_\odot \) on the companions of nine stars. These companions are therefore M-dwarfs or brown dwarfs. The upper mass-limits of 12 companions within our sample remain unconstrained because either the derived limit is very high (>0.63 \( M_\odot \)) or the orbital period is not covered by Hipparcos measurements. Finally, we found that the companion of HD 137510 is a brown dwarf with a mass range of 22.7–64.4 \( M_\odot \).

Our results are summarised in Fig. 15, where we show the companion-mass constraints obtained from Hipparcos astrometry for all targets in our analysis sample.
5.1. Combining radial velocities and Hipparcos astrometry

We have used the astrometric data of the new Hipparcos reduction to place constraints on the masses of potential brown-dwarf companions detected in radial-velocity surveys. Radial-velocity measurements of a stellar orbit yield five of the seven Keplerian elements but do not constrain the orbit’s inclination and longitude of the ascending node. We have demonstrated how these two remaining elements can be determined, if the orbital signature is present in the Hipparcos astrometry, and how to distinguish a significant orbit.

We used the permutation test to define the criterion for orbit detection. The implementation of this powerful test is simplified considerably by the use of the new Hipparcos reduction, because the individual measurements of the IAD are independent. A source of possible errors and complications in the original Hipparcos data is therefore avoided. We emphasize the importance of eliminating outliers of the IAD, because otherwise the final result can be falsified. On a couple of comparison targets drawn from the literature, we have verified that our analysis method is robust. For these targets, the performance of the new and the original Hipparcos IAD is comparable and no major difference in the results was observed.

In addition to the permutation test, we simultaneously performed the F-test and can thus compare the performance of both significance indicators, keeping in mind that the F-test is only valid if Gaussian errors are assumed. Figure 16 displays the respective values listed in Tables 8 and 10. We find a general agreement between the two indicators, i.e. the null probability decreases with increasing permutation significance. However, the permutation test is more stringent. If we had chosen to rely on
the F-test and imposed a null probability of 0.3% as detection limit, we would have detected 10 orbits. These orbits overlap with the 6 orbits that we did detect with the criterion based on the permutation test (cf. Table 8).

A problem of similar studies in the past was the appearance of biases towards small orbit inclinations, thus towards high companion masses (Han et al. 2001; Pourbaix 2001). Using simulations, we showed that our method is not affected by this bias when using an appropriate detection limit in terms of orbit significance. Figure 17 shows the distributions of orbital inclination and semimajor axis for all orbits with greater than 2-$\sigma$ significance. No preference for very small inclinations can be observed, as it would be expected from biased solutions (Pourbaix 2001). However, we note that no orbit is detected at inclinations above 50°. This is explainable because in this study, we have selected stars with minimum companion masses of $M_2 \sin i = 13–80 M_J$. As shown by Grether & Lineweaver (2006), the frequency of brown-dwarf companions is very low compared to the stellar companions and therefore small inclinations are more likely to be detected.

We showed that our analysis is ultimately limited by the astrometric precision of the satellite and therefore cannot constrain orbits of size typically smaller than $a = 2$ mas, which in most cases will impede the detection of a true brown dwarf companion. In addition, the Hipparcos mission duration of about 1000 days limits the periods of detectable orbits to $\sim 2000$ days.

In summary, we have developed a new method to combine a radial velocity solution with Hipparcos astrometric data that has several advantages compared to previous similar approaches: it uses the new Hipparcos reduction, which simplifies both the mathematical formulation and the performance of the permutation test in contrast to the implementation of e.g. Zucker & Mazeh (2001) for the original Hipparcos data. The new method defines rigorous detection limits in terms of orbit significance and yields companion mass limits for low-significance orbits, provided they are covered by astrometric measurements. It is stringent enough to avoid the output of biased results (an example is the case of HD 38529, see Reffert & Quirrenbach 2006; and Benedict et al. 2010). Because of the general format of the new reduction IAD, our method is universal and applicable to any astrometric data.

5.2. Brown-dwarf companions in the CORALIE survey

The CORALIE planet-search sample contains about 1600 Sun-like stars within 50 pc. After a duration of 12 years, 21 potential brown companions have been observed in this survey. Eighteen have been discussed in this work and three additional hosts of potential brown-dwarf companions are listed by Sozzetti & Desidera (2010) and contained in the CORALIE sample. Assuming that most candidates contained in this survey have already been found, we can derive a typical frequency of 1.3% for potential brown dwarf companions in close orbits ($<10$ AU) around Sun-like stars. Seven of these candidates are identified as stellar companions by this study. In addition, Zucker & Mazeh (2001) showed that the companions of HD 18445, HD 112758, and HD 217580 are of stellar nature. Thus, at least 10 out of 21 brown-dwarf candidates are low-mass stars, which corresponds to a rate of 48%. Based on the uniform stellar sample surveyed by the CORALIE planet search, we thus obtain an upper limit of 0.6% for the frequency of close-in brown-dwarf companions around Sun-like stars within 50 pc, which is close to the frequency of $<0.5%$ initially quoted by Marcy & Butler (2000).

Figure 18 shows the distribution of $M_2 \sin i$ in the brown-dwarf mass range for the CORALIE sample. While the cumulative distribution (blue dashed line) of all 21 candidates is reasonably compatible with a linear increase, i.e. a flat

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5 This is true provided that the implicit approximations are valid, i.e. the astrometric precision is larger than about 1 mas, see e.g. Konacki et al. (2002).

6 Three more stars from this list were not included because their minimum mass is $M_2 \sin i < 13 M_J$. The frequency of bias is reasonably compatible with a linear increase, i.e. a flat
distribution function, the situation changes significantly after removal of the 10 stellar companions. As illustrated in Fig. 19, most stellar companions have $M_2 \sin{i} > 45 M_J$, and after their removal, a cumulative distribution with particular shape emerges (black solid line in Fig. 18). It exhibits a steep ascent in the range of 13–25 $M_J$, which contains one half of the objects, followed by

\begin{figure}
\centering
\includegraphics[width=\textwidth]{Fig_20.png}
\caption{Visualisation of the high-significance orbits. \textit{Top panels:} astrometric stellar orbits projected on the sky. North is up and east is left. The solid red line shows the orbital solution and open circles mark the individual Hipparcos measurements. Dashed lines with orientation along the scan angle $\psi$ and length given by the O–C residual of the orbital solution connect the measurements with the predicted location from our model. The blue solid circles show the normal points for each satellite orbit number. The curl at the lower left corner indicates the orientation of orbital motion. \textit{Bottom panels:} O–C residuals for the normal points of the orbital solution (filled blue circles) and of the standard 5-parameter model without companion (black crosses). The error bars of the normal points correspond to the dispersion of Hipparcos measurements if there are several per satellite orbit and to the individual Hipparcos abscissa error if there is only one measurement.}
\end{figure}
a slower increase up to 60 \textit{M}_J. Surprisingly, there is no companion left of minimum mass higher than 60 \textit{M}_J.

In summary, the emerging cumulative distribution does not support a uniform distribution function of companion masses. The possibility of all companions in the 13–25 \textit{M}_J range being stellar companions seen at low inclination, is both statistically unlikely and also inplausible because their astrometric signature would be large and probably detectable. Instead, these companions supposedly have \textit{M}_2 \approx \textit{M}_2 \sin i and may represent the high-mass tail of the planetary distribution.

We note that the distribution found here, which is based on radial velocities, is compatible with the first brown-dwarf companions found in transiting systems, such as CoRoT-3b.
(\(M_2 = 21.7 \, M_J\), Deleuil et al. 2008) and CoRoT-15b \((M_2 = 63.3 \, M_J\), Bouchy et al. 2010).

5.3. Metallicities of the brown-dwarf host stars

We showed that the companion of HD 137510 is a brown dwarf. In addition, the companion of HD 190228 can be considered to be a brown dwarf at 95% confidence. The host stars have metallicities of \([\text{Fe/H}] = 0.373 \, \text{dex} \) (Butler et al. 2006) and \([\text{Fe/H}] = -0.24 \, \text{dex} \) (Perrier et al. 2003), respectively. Compared to the mean metallicity \([\langle \text{Fe/H} \rangle = -0.09 \, \text{dex} \) with root-mean-square \(\sigma_{\langle \text{Fe/H} \rangle} = 0.24 \, \text{dex} \) derived for the uniform sample of 451 nearby Sun-like stars by Sousa et al. (2008), the two brown-dwarf hosts are therefore metal-rich in one case and unremarkable in the other. The question of whether the hosts of brown dwarf companions have a particular metallicity distribution, as the planet hosts do (Santos et al. 2001), cannot be answered with this small sample, thus remains open.

6. Conclusions

We have presented the radial-velocity solutions of 15 stars with potential brown dwarf companions, of which 9 are new discoveries. At this point, the CORALIE and HARPS surveys have increased the number of known potential brown-dwarf companions of Sun-like stars by 20%.

Eight companions in our search sample were identified as being of stellar nature with masses of above \(80\, M_J\). These include one companion with a mass of \(90\, M_J\) and thus very close to the brown-dwarf mass boundary. The upper and lower mass limits of three companions were derived, but enclose the star-brown dwarf mass boundary and can therefore not be used to identify the object's nature. Yet, the companion of HD 190228 is a brown dwarf at 95% confidence. Upper mass limits were set for the companions of nine stars and revealed a brown dwarf orbiting HD 137510. The masses of 12 companions could not be constrained.

Our findings are in accordance with the presence of the brown-dwarf desert, because only three brown dwarfs are identified in our search sample of 33 companions: HD 190228b (at 95% confidence), HD 137510b, and HD 38529b (Benedict et al. 2010). In contrast, eight brown-dwarf candidates were identified as Ms dwarfs. On the basis of the CORALIE planet-search sample, we obtain an upper limit of 0.6% for the frequency of brown-dwarf companions around Sun-like stars and confirm their presence in the brown-dwarf desert, because only three brown dwarfs are identified by the Hipparcos astrometric precision and will make it possible to determine the masses of many radial-velocity companions.

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