Resolving the L/T transition binary SDSS J2052-1609 AB[★] (Research Note)

M. B. Stumpf¹, K. Geißler², H. Bouy³, W. Brandner¹, B. Goldman¹, and Th. Henning¹

¹ Max-Planck-Institut für Astronomie, Königstuhl 17, 69117 Heidelberg, Germany
 e-mail: stumpf@mpia.de

² Department of Physics and Astronomy, State University of New York, Stony Brook, NY 11794-3800, USA
 e-mail: geissler@astro.sunysb.edu

³ Centro de Astrobiología, INTA-CSIC, PO Box 78, 28691, Villanueva de la Cañada, Madrid, Spain

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ABSTRACT

Context. Binaries provide empirically derived key constraints for star formation theories, including the overall binary fraction, mass ratio distribution, and the separation distribution. They play crucial roles in calibrating the output of theoretical models, such as absolute magnitudes, colors, and effective temperature depending on mass, metallicity, and age.

Aims. We present first results of our ongoing high-resolution imaging survey of late type brown dwarfs. The survey aims at resolving tight brown dwarf binary systems to better constrain the T dwarf binary fraction. We intend to follow up on the individual binaries to determine orbital parameters.

Methods. Using NACO at the VLT, we performed AO-assisted near-infrared observations of SDSS J2052-1609. High-spatial resolution images of the T1 dwarf were obtained in *H* and *K_S* filters.

Results. We resolved SDSS J2052-1609 into a binary system with a separation of $0.1009 \pm 0.001''$. Archival data from HST/NICMOS taken one year before our observations proves the components are co-moving. Using the flux ratio between the components we infer *J*, *H*, and *K_S* magnitudes for the resolved system. From the near-IR colors, we estimate spectral types of $T1_{-4}^{+1}$ for component A and $T2.5 \pm 1$ for component B. A first estimate of the total system mass yields $M_{\text{tot}} \geq 78 M_{\text{Jup}}$, when assuming a circular orbit.

Key words. stars: low-mass – brown dwarfs – stars: individual: SDSS J2052-1609 AB

1. Introduction

Observational studies of binaries provide the key to assessing fundamental astrophysical parameters such as mass, radius, or density of compact objects for a wide range of astronomical objects ranging from asteroids to brown dwarfs and stars. For brown dwarfs, the continuous cooling with increasing age results in a degeneracy of the mass-luminosity relation; therefore, for a brown dwarf of unknown age, dynamical mass estimates based on the determination of orbital parameters of a binary provide the only precise means to derive masses.

After the discovery of the first binary brown dwarfs (Martín et al. 1999a,b), systematic surveys have led to the discovery of more than 70 brown dwarfs in binary or multiple systems (e.g., Bouy et al. 2003; Burgasser et al. 2003; Gizis et al. 2003). Astrometric follow-up observations of some of the shorter period binaries provided first system-mass estimates (e.g. Bouy et al. 2004; Konopacky et al. 2010) and revealed that mass-luminosity predictions based on theoretical evolutionary models for brown dwarfs still needed to be calibrated against observations.

Binaries have also been invoked to explain spectral and flux peculiarities observed in brown dwarfs at the transition between L and T dwarfs. Burgasser et al. (2006b) find a higher incident of binaries among L/T transition objects compared to dwarfs of earlier or later spectral types, and they suggest that spectral features of several L/T transition brown dwarfs could be explained by a superposition of spectra of a mid to late L and a T dwarf. Meanwhile, Goldman et al. (2008) conclude that the larger binary fraction among L/T brown dwarfs is not statistically significant and that a much larger observation sample is required.

Liu et al. (2006) reported the first detection of a flux ratio reversal between primary and secondary in a T dwarf binary in the *J*-band. Several more flux reversal L/T transition binaries have been identified in the meantime, supporting the idea of this so-called “crypto-binarity” (see e.g., Looper et al. 2008, and references therein; Stumpf et al. 2010). As a result, binaries composed of a late L dwarf and an early T dwarf could nicely explain both the observed near infrared spectral energy distribution and the spectral characteristics of L/T transition dwarfs.

As part of a high-resolution near-IR imaging survey, which had the aim of finding tight T dwarf binaries, we observed SDSS J205235.31-160929.8 (hereafter SDSS J2052-1609) and resolved it into a new close L/T transition binary system. Originally, SDSS J2052-1609 was identified as an apparent single object in an SDSS survey by Chiu et al. (2006) and classified as a $T1 \pm 1$ dwarf. Furthermore, SDSS J2052-1609 has a measured proper motion of $\mu = 0.483 \pm 0.022$ arc-sec/yr (Faherty 2010, in prep.; see Burgasser et al. 2010).

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Fig. 1. Timeline for the observations of SDSS J2052-1609 AB. The orientation is the same for all images. The 2MASS image shows the unresolved system, and the circle indicates the 4.8'' region around the target in which no other object is visible. The object in the northwest corner of the image is the PSF reference object 2MASS J205234-160913. The HST/NICMOS and VLT/NACO images show the resolved system where SDSS J2052-1609 B is the component northeast of SDSS J2052-1609 A.

Just recently, Burgasser et al. (2010) has identified SDSS J2052-1609 as a binary candidate using a spectral template fitting technique. Subsequently it was ranked as a strong binary candidate, since the best-fit composite spectrum returned a much better fit than any single dwarf template. SDSS J2052-1609 was best fitted by a composite spectrum of the L7 dwarf SDSS J115553.86+055957.5 (hereafter SDSS J1155-0559, Knapp et al. 2004) and the T2 dwarf 2MASS J11220826-3512363 (hereafter 2MASS J1122-3512, Tinney et al. 2005), and Burgasser et al. (2010) estimated average component types of $L7.5 \pm 0.6$ and $T2 \pm 0.2$ for the SDSS J2052-1609 AB system.

Here we present the observations that resolved SDSS J2052-1609 AB using NACO at the Very Large Telescope (VLT).

2. Observation and data reduction

2.1. NACO

The near-infrared imaging observations of the T1 dwarf SDSS J2052-1609 were obtained with NACO at the Very Large Telescope (VLT) on Cerro Paranal. The NACO system provides high-resolution AO-assisted imaging in the near-infrared at an 8-m class telescope. The observations were carried out in service mode as a part of our T dwarf high-resolution imaging campaign. SDSS J2052-1609 was observed on June 19, 2009 in H ($1.65 \mu\text{m}$) and K_S ($2.15 \mu\text{m}$) broad-band filters using the CONICA S27 camera. The S27 camera provides a field of view (FoV) of $28'' \times 28''$ and a pixel scale of $0.0271''$. For AO correction, the wavefront sensing was performed on a reference source chosen from the GSC-II (V2.2.01). The star S3313312149 has $R = 12.85$ mag and is $37.3''$ away from SDSS J2052-1609. During the observations the target was at a mean airmass of 1.148 in H and 1.104 in K_S , with average seeing conditions between $0.8''$ in H and $0.6''$ in K_S . The observations were executed using an 8-point dither pattern to allow for sky subtraction and bad pixel correction, which resulted in total integration times of 1200 s in H -band and 960 s in K_S -band. Standard data reduction steps such as flat-fielding, sky subtraction, and bad pixel corrections were performed before averaging the eight dither positions with the Eclipse jitter (Devillard 1997) software package. The final reduced NACO images of SDSS J2052-1609 AB are shown in Fig. 1 and represent the detection images of the resolved system.

For determining the relative photometry (flux ratio) and astrometry (separation and position angle (PA)) of the new

system, we used the IDL-based simultaneous PSF-fitting algorithm from Bouy et al. (2003), adapted to the VLT/NACO data. As PSF reference star, we used the nearby point source 2MASS J205234-160913, which was observed in the same FoV as our target, hence under the same observing conditions. To determine the statistical error of this fit, the algorithm was also applied to each individually reduced image. Finally the results per filter were averaged, and the uncertainties were calculated from the standard deviation. To account for the systematic errors an uncertainty of 1% for the pixel scale and 0.4° for the orientation (e.g. Köhler 2008; Eggenberger et al. 2007) was added at the end.

2.2. HST and 2MASS

After resolving SDSS J2052-1609 as a candidate binary, we checked the 2MASS and HST data archives for previous epoch data. We downloaded the available 2MASS J , H , and K_S imaging data taken on June 1, 1999. The Two Micron All-Sky Survey provides near-infrared imaging data with a pixel scale of $1''$. Since SDSS J2052-1609 has a measured proper motion of 0.483 arcsec/yr , the dwarf should have moved by approximately $4.8''$ during the last 10 years, corresponding to ~ 4.8 pixel in the 2MASS image. The 2MASS images show no object around SDSS J2052-1609 within a radius of $\sim 5''$.

The HST archive proved to be more fruitful. NICMOS imaging observation of SDSS J2052-1609 AB have been obtained under program GO 11136 (PI: M. Liu). The observation was carried out on June 24, 2008 and used the NICMOS1 (NIC1) camera with an FoV of $11'' \times 11''$ and a pixel scale of $0.043''$. Observations of SDSS J2052-1609 AB are available in several filters. Here we used the $F110W$ and $F170M$ filters, with total integration times of 12 s and 60 s, respectively. The telescope orientation angle during all observations was -27.9068° (E of N).

The data analysis is based on the pipeline-reduced frames provided by the HST archive. To derive the magnitudes from the HST/NIC1 images, we applied aperture photometry to measure the total flux of the system, since the two components of the binary were too tight to derive their absolute photometry separately. With IDL-based routines, the count rate within an aperture of 11.5 pixel was measured, corrected to a nominal infinite aperture, and converted into the “Arizona” Vega photometric

Table 1. Photometric properties of the resolved SDSS2052-1609 AB system.

Filter	Flux ratio	Δ mag [mag]	Comp. A [mag]	Comp. B [mag]	A + B [mag]
F110W ^a	0.787 ± 0.012	0.26 ± 0.02	17.77 ± 0.03	18.02 ± 0.03	17.14 ± 0.02
F170M ^a	0.576 ± 0.005	0.60 ± 0.01	15.86 ± 0.02	16.46 ± 0.02	15.37 ± 0.02
<i>J</i> ^b	0.908 ± 0.149	0.10 ± 0.18	17.00 ± 0.13	17.10 ± 0.13	16.30 ± 0.12
<i>H</i> ^c	0.591 ± 0.007	0.57 ± 0.01	15.92 ± 0.12	16.49 ± 0.12	15.41 ± 0.12
<i>K_S</i> ^c	0.513 ± 0.005	0.72 ± 0.01	15.57 ± 0.15	16.30 ± 0.15	15.12 ± 0.15
<i>J</i> – <i>H</i>			1.08 ± 0.05	0.61 ± 0.05	0.88 ± 0.04
<i>H</i> – <i>K_S</i>			0.35 ± 0.19	0.19 ± 0.19	0.29 ± 0.19
<i>J</i> – <i>K_S</i>			1.43 ± 0.20	0.81 ± 0.20	1.17 ± 0.19

Notes. ^(a) The NICMOS photometry is on the Arizona Vega System ($M_{\text{Vega}} = 0.02$); ^(b) due to the lack of direct observations, the *J*-band magnitudes are based on the NICMOS photometry which were transformed from the Vega system into the 2MASS system as described in the text. ^(c) based on 2MASS *H*, *K_S* photometry for the unresolved sources.

system (Table 1), using the photometric keywords and flux zero points provided on the STScI webpage¹.

Similar to the VLT/NACO data, we applied the PSF-fitting algorithm from Bouy et al. (2003) to derive the astrometric parameters of the binary, as well as the flux ratio of the two components. This time we built and used a library of 6 reference PSFs: one TinyTim synthetic PSF and five natural PSFs from other HST/NICMOS programs that targeted brown dwarfs of similar spectral types (L7 – T1.5) in the corresponding filters (GO 10143 and 10879, PI: I. Reid). For a precise measurement of the separation, PA, and flux ratio, the algorithm was used on each of the 4 images per filter. The final results are the average per filter of the individual results, and the error was calculated from the standard deviation.

3. Results

3.1. Photometry

Table 1 lists the flux ratio and the corresponding magnitude difference obtained from the binary fitting procedure. For the *H* and *K_S* filter, the 2MASS magnitudes of the combined system are given in column “A + B”, while the aperture photometry result is listed for the NICMOS filters. Individual magnitudes for the two components were then calculated from these system magnitudes using the flux ratio determined during the binary fitting procedure. To estimate the individual *J*-band magnitudes of the components in the 2MASS system, we applied the relation defined by Burgasser et al. (2006b). With it, the NICMOS *F110W* and *F170M* magnitudes are transformed into MKO *J*- and *H*-band magnitudes. Afterwards, the MKO *J*–*H* colors were transformed into the 2MASS system² to calculate the 2MASS *J*-band magnitudes (given in bold in Table 1). For comparison, the derived *J*-band magnitude of the unresolved system ($J = 16.30 \pm 0.12$ mag) is, within its errors, in very good agreement to the published 2MASS magnitude ($J = 16.33 \pm 0.12$ mag). Finally, the given *J*-band flux ratio was derived from the separately transformed component magnitudes.

3.2. Spectral type and distance

The derived colors of the binary components are given in Table 1. A comparison to mean *JHK_S* colors of L and T dwarfs,

¹ http://www.stsci.edu/hst/nicmos/performance/photometry/postnics_keywords.html

² <http://www.astro.caltech.edu/~jmc/2mass/v3/transformations/>

Fig. 2. Location of SDSS J2052-1609 A and B in the *J* – *K_S* vs. *J* – *H* plane relative to known dwarfs from the *DwarfArchives*. The numbers indicate the individual spectral types from 7 = L7 to 14.5 = T4.5. The two green numbers indicate the location of SDSS J1155+0559 (L7.5) and SDSS J1122-3512 (T2), the two best-fit spectral composite components derived by Burgasser et al. (2010).

calculated from the dwarfs listed at the *DwarfArchives*³ (provided in the 2MASS photometric system) shows good agreement with spectral types of T0 – T1 for component A and spectral types of T2 – T3 for the B component (see also Fig. 2 for illustration). An additional comparison with the average *J* – *K_S* colors vs. SpT in Table 5 of Faherty et al. (2009) yields the same result. In fact, the colors of SDSS J2052-1609 A are almost identical to the colors of SDSSp J083717.22-000018.3 (hereafter SDSS J0837-0000: $J - H = 1.11 \pm 0.28$ mag, $H - K_S = 0.31 \pm 0.23$ mag, $J - K_S = 1.43 \pm 0.26$ mag), which has an assigned optical SpT of T0 (Kirkpatrick et al. 2008) and T1 in the near-IR (Burgasser et al. 2006a). A different approach to estimating the spectral types was done using the SpT – NICMOS color relation given in Burgasser et al. (2006b). For the B component, the calculation yields a spectral type of T2, while the spectral type is roughly T1 for component A. Thus, these estimated spectral types are in good agreement with the spectral types yielded by the *JHK_S* color comparison.

In contrast, the spectral binary fitting of the unresolved SDSS J2052-1609 AB by Burgasser et al. (2010) returned spectral

³ <http://www.DwarfArchives.org>

Fig. 3. Comparison of the flux distribution in the NICMOS/NIC1 *F170M* filter bandpass between a blue L7, a T0, and the T2 spectral standard. All spectra are normalized to the peak flux in the 1.0–1.3 μm range.

Table 2. Estimated distances of SDSS J2052-1609 AB in parsec.

	A (L7.5)	A (T1)	B (T2)	A + B (T1)
<i>J</i>				22.3 ± 3.2
<i>H</i>	31.8 ± 6.4	27.2 ± 4.0	34.6 ± 5.0	21.6 ± 3.1
<i>K_S</i>	37.2 ± 8.9	28.6 ± 4.8	36.7 ± 6.1	23.2 ± 3.7

types of $L7.5 \pm 0.6$ and $T2 \pm 0.2$. While the spectral type of the B component agrees with our estimation, the spectral type of the A component differs by 3–4 subclasses. The A component of their best-fit composite template, SDSS J1155+0559, is an L7 dwarf with unusually blue near-IR colors ($J - H = 0.96 \pm 0.11$ mag, $H - K_S = 0.59 \pm 0.10$ mag, $J - K_S = 1.54 \pm 0.11$ mag). Figure 3 compares the spectra of SDSS J1155+0559 (L7, Burgasser et al. 2010), SDSS J0837-0000 (T0, Burgasser et al. 2006a), and the T2 spectral standard SDSS J125453.90-012247.4 (SDSS J1254+0122, Burgasser et al. 2004), which were derived from the SpeX prism spectral libraries⁴. It illustrates their flux emitted in the *F170M* filter, which is centered on the methane absorption band. The assumption that SDSS J2052-1609 A is an unusually blue late type L dwarf with no methane absorption could explain the large difference in the *F170M* magnitudes of components A and B, which is not expected if the spectral types are as similar as T1 and T2. With only the near-IR colors to rely on, we assign preliminary spectral types of $T1^{+1}_{-4}$ and $T2.5 \pm 1$ for SDSS J2052-1609 A and B, respectively, in which the larger uncertainty of the A component accounts for the still unexplained large $\Delta F170M$ mag of the components. Only resolved spectroscopy of the system components can unambiguously determine their spectral types.

Using the approximate spectral types, we can estimate the distance by employing the relation between absolute magnitude and spectral type fromLooper et al. (2008). While the unresolved system had a photometric distance of 22 ± 1.7 pc, the resolved system will be farther away. For each of the individual components, we calculated the distance based on the *H* and *K_S* magnitudes (see Table 2), and the uncertainties were derived from the photometric, as well as the absolute magnitude vs. SpT relation uncertainties. Assuming spectral types T1 and T2 for SDSS J2052-1609 A and B, we derived an average distance of

Table 3. Astrometric parameters for the SDSS J2052-1609 AB system.

Parameter		HST/NIC1 <i>F110W</i> (24/06/2008) ^a	VLT/NACO <i>K_S</i> (19/06/2009)
Separation	ρ [mas]	102.7 ± 1.1	100.9 ± 1.0
Position Angle	θ [deg]	50.4 ± 0.2	67.3 ± 0.5
estimated distance	d [pc]	31.8 ± 2.3	
Semi-major axis	a [AU]	$\geq 3.2 \pm 0.5$	
Orbital period	P [years]	21 ± 1	
System mass	$M_{\text{tot}} [M_{\text{Jup}}]$	≥ 78	

Notes. ^(a) *HST* observations by M. Liu that are present in the public archive.

31.8 ± 2.5 pc for the binary system. Given the large uncertainty in the spectral type of component A, the average distance may increase up to 35.8 ± 3.4 pc if we assume a spectral type of L7.5.

3.3. System parameters

To derive the most reliable astrometric parameters from the 2008 HST/NICMOS and 2009 VLT/NACO observations, we only used the *F110W* and *K_S*-band data to determine the first orbital parameters, since those two filters provide results that are notably better than the diffraction limit in these filters. In addition, the NACO observations in the *K_S*-band are obtained at a much better Strehl ratio (with an average value of 69% compared to 54% in *H*-band). The derived separations and position angles are listed in Table 3. The observations are separated by a little bit less than one year. During this time the separation between the two components changed from 102.7 mas to 100.9 mas. With this total decrease of only 1.8 ± 1.5 mas, the separation between components A and B remained almost constant over the period of one year. This leads to the conclusion that the two components must be co-moving, since SDSS J2052-1609 AB has a reported proper motion of $\mu = 0.483 \pm 0.022''/\text{yr}$. In the same time interval, the position angle changed by $16.9^\circ \pm 0.5^\circ$.

The insignificant change in the separation allowed for a first assumption of a circular orbit seen face-on and an estimate of the orbital period of the system, therefore, the orbital movement of 16.9° per year translates into an orbital period of 21 ± 1 yr. With an averaged separation of 101.8 ± 1.5 mas and the new determined photometric distance of 31.8 pc, the semi-major axis of the system corresponds to a projected separation of 3.2 ± 0.5 AU. Using Kepler's law, this yields a first estimate of the total system mass of $\approx 78 M_{\text{Jup}}$. Depending on the inclination (if $i > 0$) and eccentricity (e) of the system, the true semi-major axis is on average larger than the observed separation (Fischer & Marcy 1992), resulting in an increase in the total system mass. Only in the case of an eccentric orbit seen face-on and observed close to the periastron would the estimated total system mass be an upper limit. An ongoing astrometric monitoring program will help to better determine these orbital parameters in the future.

4. Summary

VLT/NACO observations of SDSS J2052-1609 resolved the brown dwarf into a close binary system with a separation of only ~ 101.8 mas. Previous, unpublished observations conducted with HST/NICMOS confirmed the binary nature of the system and allowed for a first estimation of the total system mass. Assuming a

⁴ <http://www.browndwarfs.org/spexprism>

circular orbit with an orbital period of 21 ± 1 years, we estimated a system mass of $\geq 78 M_{\text{Jup}}$. The near-IR colors of the individual components suggest spectral types of $T1^{+1}_{-4}$ and $T2.5 \pm 1$.

A divergent estimate of the spectral types comes from spectral binary fitting (Burgasser et al. 2010), suggesting spectral types of L7.5 for component A and T2 for component B. Their best-fit composite template consists of two late types with unusually blue near-IR colors. Our obtained photometry agrees with both scenarios: a “normal” T + T dwarf binary or an L/T transition system consisting of a blue L and an early T dwarf.

Upcoming resolved spectroscopy with SINFONI at the VLT will finally determine the spectral types of the system.

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