The 2009 outburst of H 1743-322 as observed by RXTE

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ABSTRACT

Context. Six outbursts have been observed by RXTE from the source since 2003. Following “a failed outburst” with a missing high/soft state in 2008, the most recent one occurred in 2009, exhibiting all spectral transitions.

Aims. An outburst was observed by RXTE from the black hole candidate H 1743-322 in 2009. We investigate this particular outburst in the context of the source outburst history.

Methods. We analyze the RXTE observations of the 2009 outburst of H 1743-322, as well as the observations of the previous five outbursts for comparison. All available RXTE observations taken since 2003 sum up to a total exposure of ~ 1200 ks.

Results. The hardness-intensity diagram (HID) shows a complete counter-clockwise q-track for the 2009 outburst. It leaves the low/hard state but does not reach the leftmost edge of the overall HID. Near maximum, the X-ray spectra show a power law with photon index < 2.4 ; the peak luminosity during the whole outburst reached only 20% of the Eddington luminosity. The highest disk fractional luminosity is about 0.75 for the 2009 outburst, which is close to the lower limit observed from the black hole XRB sample in their state transitions from the intermediate to the high/soft state. While the lowest hardness (6–19 keV/3–6 keV) values in the HID is about 0.3–0.4 for the 2009 outburst, the timing analysis shows that a transition to the high soft state occurred. During the low/hard state of the 2009 outburst, the inner radius of the accretion disk is found to be closer to the central black hole and anti-correlated with the disk temperature implying reprocessing of the hot corona on the disk's soft X-rays. In the luminosity diagram of the corona versus the disk, the tracks of the outbursts in 2003 and 2009 cross the line which represents a roughly equal contribution to the emission from the thermal and the non-thermal components; the track of the 2008 outburst has its turn-over falling on this line. This may indicate an emission balance between the corona and the disk that prevents the state transition from going further than the low/hard state.

Key words. X-rays: binaries – X-rays: stars – X-rays: individuals: H 1743-322 – black hole physics

1. Introduction

From an observational point of view, the low-mass X-ray binaries (LMXB), which contain a the compact star such as a neutron star or a black hole candidate (BHC), have at least two states during their outbursts, the low/hard state (LHS) and the high/soft state (HSS), that can be distinguished on the basis of their spectral behaviour. Many models have been developed to understand the nature of black hole candidates (BHCs), but they basically all assume that their spectra are generated by two (thermal and non-thermal) components. The differences among these models arise from the locations in which these components are generated (for reviews, see, e.g., Homan & Belloni 2005; Remillard & McClintock 2006; and Belloni 2009). When the accretion rate increases from a quiescent state, the inner accretion disk region evaporates to form hot plasma or an advection-dominated accretion flow region (ADAF, Narayan et al. 1996), and the seed photons are Comptonized to X-ray photons. The LHS is more luminous than the quiescent state, and most binaries have been discovered when this transition occurs (Nowak 1995; Fender et al. 2004; Belloni 2009). It is understood that for BHCs after the LHS, the soft X-ray (e.g., in 2–10 keV) luminosity of most outbursts still increase to more than 10% of Eddington (McClintock & Remillard 2006), the inner disk being hotter and closer to the BHCs than in the LHS. However, circular orbits

around BHs with radii below a certain value are unstable, and the disk-radius lower limit extends only to the innermost stable circular orbit (ISCO). For a Schwarzschild BH, the ISCO is at $3GM/c^2 \approx 3.0 M_{\text{BH}}/M_{\odot}$ km. For the BHCs beyond the HSS, the luminosity of some outbursts can increase to the Eddington luminosity, placing it in a very high state (VHS, Miyamoto et al. 1991), where the contribution from the soft thermal component with typical temperature ~ 1 keV and the power-law component (with a photon index ~ 2.5) are comparable.

There are two main intermediate states between the LHS and HSS: the so-called hard intermediate state (HIMS) and the soft intermediate state (SIMS), where the X-ray spectra are described by a power law with photon index 1.6–2.5. They can be distinguished in the hardness-intensity diagrams (HIDs). In the latter, the evolution of BHCs forms a “q-track” pattern, e.g., see the outburst evolution of GX 339-4 (Fender et al. 2004; Belloni et al. 2005; Homan & Belloni 2005; Del Santo et al. 2008; Belloni 2009, for a review, see Homan & Belloni 2005; Dunn et al. 2010). Both the LHS and HSS are rather stable states, which can be observed on timescales of months or years, while the remaining states exhibit strong and rapid variations.

Besides their spectral characteristics, another important property of BHC states is their fast aperiodic variability (e.g., van der Klis 1994; Vaughan et al. 2003; McHardy et al. 2004; Uttley et al. 2005). The root mean square (rms), i.e., integrated

power spectral densities (PSDs) over a broad range of frequencies, is usually treated as an indicator of the variability amplitude. As for HIDs, the hardness-rms diagrams (HRDs), where the fractional rms is plotted versus hardness, are useful for characterizing the aperiodic variability behaviour in different states (e.g., Gleissner et al. 2004; Belloni 2009). The fractional rms in the LHS, of typically $\sim 30\%$, increases slowly with flux. For the HSS, the fractional rms is typically $<10\%$, weaker and located at the lower right corner of the HRDs. In the latter, the intermediate states of course have intermediate rms values. The evolution of rms values is smooth in the HRDs except for the SIMS, which is marked by a sharp drop (Belloni et al. 2005; Fender et al. 2006).

After the initial LHS, most sources have tracks extending to the left branch in the HID that step into the HSS. However, a few sources, such as XTE J1118+480 (Remillard et al. 2000; Wilson & McCollough 2000) and H 1743-322 (Capitanio et al. 2009), sometimes have tracks that do not reach the HSS with peak luminosities around $\sim 0.1 L_{\text{Edd}}$: they leave the LHS and return to quiescence without canonical “q” patterns in HIDs, forming the so-called “failed outbursts”. The outburst can move yet further from the LHS to form a SIMS. The case in which the SIMS is not followed by a HSS is referred to as a “failed state transition”. Although rather rare, failed state transitions have been occasionally observed from the persistent bright black hole binary Cyg X-1 (Belloni et al. 1996). Cyg X-1 spends most of its time in the LHS, with occasional transitions to the HSS or a failed state transition. In a failed state transition, the hard component does not become as soft as a full-fledged HSS but only reaches the SIMS with the X-ray spectra being described by a power law with a photon index ~ 2.15 and inner disk temperature of about 0.36 keV. Spectrally, the hard component does not become too soft and does not reach the leftmost part of the HIDs as in the case of other transient BHC. The rms does not go down to 10% and has no sharp QPO (for a review, see Belloni 2009).

The black hole candidate and X-ray binary H 1743-322 (also named 1H 1741-322, XTE 1746-322, and IGR J17464-3213) was proposed to be a BHC based on its “ultrasoft X-ray spectra” during outburst (White & Marshall 1983). In the RXTE era before 2009, five outbursts (referred to in what follows as #1–5, in time sequence) were detected from H 1743-322: on March 2003 (MJD 52 550), September 2004 (MJD 53 250), September 2005 (MJD 53 600), January 2008 (54 450), and October 2008 (MJD 54 750). In all transient sources for which the initial rise is observed, the start of the outburst is in the LHS. For the first and brightest outburst in 2003, in which the source moved to the VHS, McClintock et al. (2009) found that the source had many similarities with the BH binary XTE J1550-564 in the radio and X-ray bands. In this outburst, Homan et al. (2005) detected high-frequency QPOs at 240 and 160 Hz, making this source the fourth BHC with high-frequency QPOs (the others being GRO J1655-40, GRS 1915+105, and XTE J1550-564). The outbursts #2, 3, and 4 (see Swank 2004; Rupen et al. 2005; Capitanio et al. 2006; Kalemci et al. 2008) have been investigated less thoroughly because of limited observations. The fifth outburst was a failed outburst, with the X-ray spectra being described by a power law with the softest photon index ($\sim 2.1^{+0.03}_{-0.03}$) and the lowest luminosity ($\sim 0.1 L_{\text{Edd}}$). This one is the faintest outburst in the set; it is understood to lack the high/soft state because of its low (peak) accretion rate at that time (Capitanio et al. 2009). More recently, H 1743-322 experienced a new outburst initially detected by Swift/BAT on 2009 May 26 (Krimm et al. 2009). Radio activity was also detected by VLA in H 1743-322 on 2009 May 27 (Miller-Jones et al. 2009). H 1743-322 returned to a low/hard state on 2009 July 7, with

power law photon index dropping to 1.84 (Kalemci et al. 2009). Renewed activity from H 1743-322 was detected by MAXI/GSC on 2009 December 27 with a power-law photon index of 1.7, and the source was suggested to be in the LHS (Yamaoka et al. 2009).

In this paper, we carry out a systematic analysis of the most recent outburst (#6), for which only some preliminary results were reported (see e.g. Krimm et al. 2009, and references therein)¹, based on the most updated RXTE observations.

We interpret our results in the context of all the other outbursts detected from H 1743-322 (e.g., Parmar et al. 2003; McClintock et al. 2009; Capitanio et al. 2009; and Dunn et al. 2010).

2. Observations and data analysis

Public data from RXTE (Gruber et al. 1996) on H 1743-322 covers the period between March 2003 and July 2009, and include 410 RXTE/PCA pointed observations, with the identifier (OBSID) of proposal number (PN) 80137, 80138, 80144, 90058, 90115, 90421, 91050, 91428, 93427, and 94413 in the data archive of the High Energy Astrophysics Science Archive Research Center (HEASARC). These observations sum up to ~ 1200 ks of exposure time on the source, and are scattered over the six outbursts.

The analysis of PCA data is here performed by using HEASoft v. 6.6. We filter the data using the standard RXTE/PCA criteria. Only the PCU2 (in the 0–4 numbering scheme) has been used for the analysis, because this PCU was 100% on during the observations, we selected the time intervals with the constraints of elevation angle $>10^\circ$ and pointing offset $<0.02^\circ$, where elevation is the angle above the limb of the Earth and the pointing offset is the angular distance between the pointing and the source. The background file used to analyzed the PCA data consistent of the most recent bright source available at the HEASARC website²; the detector breakdowns were removed³. Data from cluster 1 of the HEXTE system are used to produce the higher energy spectra. An additional 1% systematic error is added to the spectra because of calibration uncertainties, if not otherwise specified. The X-ray spectra are fitted with XSPEC v12.5.0 and the model parameters are estimated with 90% confidence level.

3. Results

3.1. Lightcurves

The lightcurves of the six outbursts (#1–6, in time sequence) in the energy band 1.5–12 keV are shown in Fig. 1, using 6400 s time bins for every observation identifier (OBSID). These lightcurves use PCA standard 2 mode data for each energy channel from which the background has been removed. For comparison, we present together the RXTE/ASM (1.5–12 keV)⁴ and Swift/BAT (15–50 keV)⁵ daily-averaged lightcurves. The

¹ We notice a contemporary work by Motta et al. (2010) based on the same RXTE data. Their results partially overlap ours.

² pca_bkgd_cmbrightrvle_eMv20051128.mdl

³ For more information see: http://heasarc.gsfc.nasa.gov/docs/xte/recipes/pca_breakdown.html

⁴ See the lightcurves provided by the All-Sky Monitor (ASM) team at the Kavli Institute for Astrophysics and Space Research at the Massachusetts Institute of Technology, <http://xte.mit.edu/>

⁵ See the Swift/BAT transient monitor results provided by the Swift Team at <http://swift.gsfc.nasa.gov/docs/swift/results/transients>

Fig. 1. The ASM(1.5–12 keV), PCA(1.5–12 keV), BAT(15–50 keV) lightcurves and PCA hardness ratio (6–19 keV/3–6 keV) of H 1743-322 from March 2003 to July 2009. The bin size for the lightcurves is one day, and for the hardness 6400 s. The PCA lightcurves prior to outbursts #6 are also available from McClintock et al. (2009) and Capitanio et al. (2009).

ASM lightcurves (Fig. 1) enclose all the six outbursts from H 1743-322, including the brightest one in 2003 and the subsequent five, fainter ones, of which the last four were caught by Swift/BAT at 15–50 keV. The lightcurves for the outbursts #1–5 were also analyzed by Parmar et al. (2003), McClintock et al. (2009), Capitanio et al. (2009), and references therein. We thus compare these with the new outburst that occurred in 2009. The latter was well covered by RXTE/PCA for almost its entire evolution from the LHS to the quiescent state. The ASM lightcurve has an outburst at soft X-rays with a profile similar to that of the BAT lightcurve at hard X-rays, with the flux peak at soft X-rays lagging that at hard X-rays by roughly 10 days. For the 2009 outburst, the Swift/BAT flux increased since MJD 54977 and reached a peak on MJD 54987, followed by a decay of more than 60 days. The hardness ratio, defined as the PCU2 flux ratio (6–19 keV)/(3–6 keV), was computed to illustrate at first glance (see also Fig. 1) the spectral changes during the outburst evolution.

3.2. Hardness-intensity diagrams

We investigate the six outbursts by deriving the evolution of the hardness (6–19 keV)/(3–6 keV) against intensity in the 3–21 keV band (Fig. 2). The brightest outburst, #1, which occurred in April 2003, has a rather complete evolution track in the HID: it shows the LHS state, the intermediate states, the HSS, and the quiescent state. For outburst #2 (occurred in July 2004, around MJD 53200), a partial HID track including the HSS and the transition to quiescence were recorded by the PCA. For outburst #3 (occurred in August 2005, around MJD 53600), the initial LHS state is missed, and it has an intermediate hardness at the flux peak. Only the decay part of a typical track was observed by the PCA for outburst #4, whereas the whole track is clearly evident for outburst #5. The latter is the faintest outburst and was reported to have failed the transition beyond the LHS. The most recent outburst, #6, is the 3rd one with an almost full HID track to be recorded by the PCA. In June 2009, the outburst #6 reached the flux peak but with a hardness located in between the outbursts #1 and #5.

3.3. Spectra

The spectra of the outburst #6 are extracted for each OBSID of the combined PCU 2 data in the 3–20 keV band and the HEXTE cluster 1 data in the 20–100 keV band. We take the DISKBB model plus a power-law/cutoff power-law to fit the data, which is similar to the modeling of BHC sources, assuming a fixed column density of 1.6×10^{22} atoms/cm² (given by Capitanio et al. 2009) and an iron line fixed at 6.4 keV. The reduced χ^2 of the fit for each OBSID are generally ~ 0.8 –1.5 with 62(63) d.o.f.s, showing mostly good fits⁶. The derived DISKBB temperature (T_{in}) and the corresponding observed inner radius R_{in}^* of the emission area are shown in Table 1 and Fig. 3. Here, $R_{\text{in}}^* = R_{\text{in}} \times (\cos \theta)^{1/2}$, where $R_{\text{in}} = D_{10} \times (N_{\text{disk}} / \cos \theta)^{1/2}$ is the inner radius of the disk, N_{disk} is the normalization on the disk blackbody component via the DISKBB models assuming $\theta = 70^\circ$ and a distance of 10 kpc, following McClintock et al. (2009). The total flux in the energy band 1.5–100 keV is also given in Table 1. The value of L_{Edd} is $\sim 10^{39}$ erg s^{−1} based on the assumption of a black hole mass of $10 M_\odot$ (McClintock et al. 2009).

At the initial part of the 2009 outburst, the flux increased in the first 6 OBSIDs (region I, where the hardness decreases). The spectral fitting in this period shows that the temperature of the inner disk drops from 1.07 to 0.75 keV, whereas the radius of the inner disk increases from 5 km to 11 km, and the index of the power law spectrum steepens from 1.33 to 2.06, the spectra being almost cut off at energies around 60–80 keV (Table 1 and Fig. 3). As the outburst decayed from the flux maximum (region II, where the hardness remains at the lowest level, and the peak luminosity is $\sim 0.2 L_{\text{Edd}}$), the disk temperature decreases slowly from 0.85 keV to 0.66 keV, whereas the radius of the inner disk remains at ~ 30 km and the power-law index at ~ 2.3 . When the outburst went into quiescence (region III, where the hardness starts to increase again), the disk temperature and the inner disk radius drop to 0.6 keV and 7 km, respectively.

The hardness at the left end of the HID is usually a diagnostic of the softest spectral state that the outburst reached. In general, the LHS has the greatest hardness, with the SIMS and

⁶ We also checked the model COMPTT, but we generally derived a larger χ^2 , around 1.7 or beyond with 62(63) d.o.f.s.

Fig. 2. The hardness-intensity diagrams (HIDs) of the six outbursts observed by RXTE/PCA. The hardness is defined as the ratio of count rates in the 6–19 keV band to the 3–6 keV band, as extracted from the lightcurves. The intensity is in units of count rates in the 3–21 keV band. The gray points represent all of the six outbursts together, each outburst being plotted in colour against this gray background. For comparison, we mosaic the six HIDs in the *bottom panel*. The time bin-size for each data point is 6400 s. Note that the HIDs were reported in Dunn et al. (2010) for the outbursts #1–5 and in Capitanio et al. (2009) for outbursts #4–5.

the HSS following it at decreasingly lower hardness levels. The hardness at the left end of HID usually corresponds to the softest spectral state that the outburst reached. McClintock et al. (2009) suggested a borderline where the disk fraction is beyond 75% to discriminate the HSS from the intermediate state. To investigate this for the six outbursts of H 1743-322, we selected data at the left end of the HID for each outburst, and performed spectral fits. The results are shown in Table 2. We find that only outbursts #1 and #2 have a disk fraction well beyond 0.75 (0.95 ± 0.01 and 0.94 ± 0.02 , respectively), #3 and #5 have it well below 0.75 (0.62 ± 0.01 and 0.29 ± 0.02 , respectively), and the others (#4 and #6) have it around 0.75 (0.73 ± 0.02 and 0.75 ± 0.02 , respectively). These results indicate that only the outbursts #1 and #2 clearly exhibited a HSS.

In the HSS of the BHC, R_{in} is assumed to remain almost constant (e.g., see Tanaka & Lewin 1995; and McClintock et al. 2009). The evidence of a constant inner disk radius in the HSS is that the bolometric luminosity of the disk, L_{disk} , is approximately proportional to the temperature of the inner disk T_{in} . Assuming that the observed disk flux is proportional to $\cos \theta$, we use $R_{\text{in}} = R_{\text{in}}^* / (\cos \theta)^{1/2}$ to calculate L_{disk} . In Fig. 4, we show the fit of the relationship between the bolometric luminosity of the disk and the temperature of the inner disk T_{in} in the HSS, with $L_{\text{disk}} \propto T_{\text{in}}^4$. The reduced χ^2 is smaller

in outburst 6 (1.3/27 d.o.f.) than in outburst 1 (40/99 d.o.f.), with both showing deviations from the T^4 relationship. This deviation has commonly been observed in BH XRBs, for instance, in GRO J1655-40, XTE 1550-564, and GRS 1915+105 (see the review by Trümper & Hasinger 2008). The possible explanation of this phenomenon might be the flux dependence of the spectral hardening factor due to the Compton scattering of the soft photons in the corona region (Ohsuga et al. 2002, 2005; Kubota & Makishima 2004; McClintock et al. 2009). The HSS data of the outbursts #1 and #6 at the HSS are the primary concerns of this plot. For comparison, we show the LHS data of outbursts #1 and #6, and spectral data of the outburst #5 as well. We note that for the outburst #1, the data points were taken from Table 1 of McClintock et al. (2009). For outburst #5, we extracted the spectra from every OBSIDs, since they were not available in Capitanio et al. (2009). For the HSS of both outbursts, #1 and #6, Fig. 4 shows that the fit results are consistent with each other, which means that the inner disk R_{in} remained at a similar distance during the HSS in these two outbursts, despite the large differences in the outburst evolution of both flux and spectrum.

The energy-division diagrams (see Remillard & McClintock 2006), which were first used by Muno et al. (1999) in describing the behaviour of GRS 1915+105, represent the contributions of the thermal and the non-thermal components. Since the

Table 1. The spectral fit results of the outburst in 2009.

No	OBSID	Time	T_{in}	R_{in}^*	Γ_1	N	E_{cut}	Flux	χ^2_{red} (d.o.f.)
		(MJD)	keV	km			keV	10^{-9}	
1	94413-01-02-00	54 980.40	$1.07^{+0.22}_{-0.14}$	5^{+3}_{-1}	$1.33^{+0.07}_{-0.09}$	$0.26^{+0.04}_{-0.04}$	$62.00^{+16.00}_{-12.00}$	7.02 ± 0.15	0.90(62)
2	94413-01-02-02	54 980.85	$1.09^{+0.19}_{-0.12}$	5^{+2}_{-1}	$1.30^{+0.06}_{-0.07}$	$0.24^{+0.03}_{-0.03}$	$59.00^{+11.00}_{-8.00}$	7.33 ± 0.12	0.96(62)
3	94413-01-02-01	54 981.96	$1.05^{+0.20}_{-0.14}$	6^{+3}_{-2}	$1.40^{+0.08}_{-0.09}$	$0.32^{+0.05}_{-0.05}$	$67.00^{+22.00}_{-14.00}$	7.32 ± 0.18	0.88(62)
4	94413-01-02-05	54 982.28	$0.93^{+0.11}_{-0.10}$	8^{+3}_{-2}	$1.42^{+0.08}_{-0.08}$	$0.32^{+0.05}_{-0.05}$	$69.00^{+18.00}_{-13.00}$	7.30 ± 0.13	0.94(62)
5	94413-01-02-04	54 983.33	$0.85^{+0.09}_{-0.08}$	11^{+4}_{-2}	$1.63^{+0.07}_{-0.07}$	$0.49^{+0.22}_{-0.24}$	$81.00^{+31.00}_{-18.00}$	6.88 ± 0.13	1.20(62)
6	94413-01-02-03	54 984.38	$0.75^{+0.06}_{-0.06}$	18^{+6}_{-4}	$2.06^{+0.03}_{-0.03}$	$1.14^{+0.09}_{-0.09}$		7.47 ± 0.12	1.21(63)
7	94413-01-03-00	54 987.26	$0.86^{+0.03}_{-0.03}$	29^{+3}_{-2}	$2.24^{+0.05}_{-0.05}$	$2.15^{+0.29}_{-0.28}$		13.40 ± 0.19	1.29(63)
8	94413-01-03-01	54 988.24	$0.83^{+0.02}_{-0.02}$	32^{+2}_{-2}	$2.24^{+0.06}_{-0.06}$	$1.22^{+0.20}_{-0.19}$		9.97 ± 0.16	0.99(63)
9	94413-01-03-07	54 988.63	$0.87^{+0.03}_{-0.02}$	28^{+2}_{-2}	$2.27^{+0.03}_{-0.04}$	$1.87^{+0.20}_{-0.20}$		11.60 ± 0.18	0.96(63)
10	94413-01-03-05	54 989.16	$0.79^{+0.02}_{-0.02}$	34^{+3}_{-2}	$2.31^{+0.07}_{-0.08}$	$1.47^{+0.23}_{-0.21}$		8.65 ± 0.17	0.85(63)
11	94413-01-03-06	54 989.22	$0.80^{+0.02}_{-0.02}$	34^{+2}_{-2}	$2.22^{+0.06}_{-0.07}$	$0.98^{+0.17}_{-0.16}$		8.84 ± 0.16	0.76(63)
12	94413-01-03-02	54 990.27	$0.88^{+0.03}_{-0.03}$	26^{+3}_{-2}	$2.21^{+0.05}_{-0.06}$	$1.80^{+0.26}_{-0.25}$		12.00 ± 0.18	1.40(63)
13	94413-01-03-03	54 990.33	$0.86^{+0.03}_{-0.03}$	27^{+3}_{-2}	$2.28^{+0.04}_{-0.04}$	$2.26^{+0.25}_{-0.25}$		12.20 ± 0.19	0.97(63)
14	94413-01-03-04	54 990.39	$0.85^{+0.03}_{-0.03}$	28^{+3}_{-2}	$2.23^{+0.05}_{-0.05}$	$1.89^{+0.24}_{-0.23}$		11.97 ± 0.19	1.17(63)
15	94413-01-03-08	54 991.77	$0.81^{+0.02}_{-0.02}$	32^{+2}_{-2}	$2.21^{+0.05}_{-0.05}$	$0.97^{+0.13}_{-0.12}$		8.60 ± 0.15	1.12(63)
16	94413-01-03-09	54 992.43	$0.83^{+0.02}_{-0.02}$	30^{+2}_{-2}	$2.26^{+0.06}_{-0.06}$	$1.41^{+0.21}_{-0.19}$		9.74 ± 0.16	1.10(63)
17	94413-01-03-10	54 993.86	$0.80^{+0.02}_{-0.02}$	31^{+2}_{-2}	$2.14^{+0.05}_{-0.06}$	$0.70^{+0.11}_{-0.10}$		7.52 ± 0.13	1.07(63)
18	94413-01-04-00	54 994.25	$0.79^{+0.02}_{-0.02}$	32^{+2}_{-2}	$2.15^{+0.08}_{-0.09}$	$0.69^{+0.17}_{-0.15}$		7.49 ± 0.14	1.44(63)
19	94413-01-04-01	54 995.24	$0.78^{+0.02}_{-0.02}$	32^{+3}_{-2}	$2.18^{+0.08}_{-0.08}$	$0.75^{+0.16}_{-0.14}$		7.26 ± 0.14	0.90(63)
20	94413-01-04-03	54 996.28	$0.77^{+0.02}_{-0.02}$	36^{+2}_{-2}	$2.12^{+0.01}_{-0.11}$	$0.42^{+0.14}_{-0.11}$		6.71 ± 0.14	1.18(63)
21	94413-01-04-02	54 996.48	$0.77^{+0.02}_{-0.02}$	35^{+2}_{-2}	$2.16^{+0.06}_{-0.07}$	$0.58^{+0.10}_{-0.09}$		6.86 ± 0.13	1.39(63)
22	94413-01-04-08	54 997.27	$0.77^{+0.02}_{-0.02}$	35^{+3}_{-2}	$2.17^{+0.10}_{-0.11}$	$0.56^{+0.17}_{-0.15}$		6.75 ± 0.14	1.74(63)
23	94413-01-04-05	54 997.40	$0.77^{+0.02}_{-0.02}$	34^{+2}_{-2}	$2.15^{+0.08}_{-0.08}$	$0.50^{+0.10}_{-0.10}$		6.35 ± 0.12	1.20(63)
24	94413-01-04-06	54 999.75	$0.74^{+0.02}_{-0.02}$	36^{+2}_{-2}	$2.14^{+0.10}_{-0.11}$	$0.38^{+0.12}_{-0.10}$		5.74 ± 0.12	1.29(63)
25	94413-01-04-04	55 000.40	$0.74^{+0.02}_{-0.02}$	32^{+3}_{-2}	$2.17^{+0.06}_{-0.06}$	$0.63^{+0.10}_{-0.09}$		5.91 ± 0.12	0.72(63)
26	94413-01-05-00	55 001.32	$0.75^{+0.02}_{-0.02}$	31^{+3}_{-2}	$2.18^{+0.08}_{-0.09}$	$0.66^{+0.15}_{-0.13}$		5.72 ± 0.12	1.05(63)
27	94413-01-05-01	55 002.36	$0.72^{+0.01}_{-0.01}$	37^{+2}_{-2}	$1.96^{+0.11}_{-0.12}$	$0.17^{+0.05}_{-0.04}$		4.68 ± 0.10	1.06(63)
28	94413-01-05-02	55 003.41	$0.70^{+0.01}_{-0.01}$	38^{+2}_{-2}	$2.05^{+0.06}_{-0.09}$	$0.22^{+0.05}_{-0.05}$		4.53 ± 0.10	1.06(63)
29	94413-01-05-03	55 004.39	$0.73^{+0.02}_{-0.02}$	29^{+3}_{-2}	$2.22^{+0.05}_{-0.03}$	$0.74^{+0.11}_{-0.10}$		5.28 ± 0.11	1.10(63)
30	94413-01-05-04	55 005.37	$0.73^{+0.02}_{-0.02}$	28^{+3}_{-2}	$2.13^{+0.07}_{-0.07}$	$0.52^{+0.10}_{-0.09}$		4.56 ± 0.09	0.86(63)
31	94413-01-05-05	55 006.36	$0.71^{+0.02}_{-0.02}$	21^{+4}_{-3}	$2.14^{+0.07}_{-0.07}$	$0.51^{+0.09}_{-0.09}$		4.38 ± 0.09	1.01(63)
32	94413-01-05-06	55 007.40	$0.71^{+0.03}_{-0.03}$	26^{+3}_{-2}	$2.15^{+0.05}_{-0.06}$	$0.55^{+0.08}_{-0.07}$		4.12 ± 0.08	1.27(63)
33	94413-01-06-00	55 008.65	$0.66^{+0.02}_{-0.02}$	33^{+4}_{-3}	$2.22^{+0.06}_{-0.07}$	$0.54^{+0.09}_{-0.08}$		3.82 ± 0.10	1.00(63)
34	94413-01-06-01	55 013.09	$0.70^{+0.04}_{-0.04}$	24^{+5}_{-3}	$2.02^{+0.12}_{-0.13}$	$0.29^{+0.10}_{-0.08}$		3.00 ± 0.09	1.05(63)
35	after 94413-01-07-00	55 029.00	$0.70^{+0.08}_{-0.08}$	7^{+4}_{-2}	$1.87^{+0.04}_{-0.04}$	$0.11^{+0.01}_{-0.01}$		1.04 ± 0.02	1.18(63)

Notes. RXTE/PCA and RXTE/HEXTE data are used for the spectral fit. The OBSID, time, the inner temperature of disk, the inner radius of the disk, power law index, break for the energy, and unabsorbed flux in the energy band 1.5–100 keV (in units of 10^{-9} erg cm $^{-2}$ s $^{-1}$) are given. PCA spectra are derived for each OBSID (~ 3000 s) and from the combining OBSIDs for the last 11 data segments. Here $R_{\text{in}}^* = R_{\text{in}} \times (\cos \theta)^{1/2}$ and $R_{\text{in}} = D_{10} \times (N_{\text{disk}} / \cos \theta)^{1/2}$ is the inner radius of the disk, θ is the inclination angle, N_{disk} is the normalization on the disk blackbody component, and D_{10} is the distance of the source in units of 10 kpc.

energy-division diagrams are more effective than the HIDs in disentangling intrinsic differences between the states of BH binaries, we show the Eddington ratios of the disk bolometric luminosity versus the power-law luminosity at 1–100 keV (Fig. 5). Here we consider only the outbursts #1 (excluding the VHS, defined by McClintock et al. 2009), #5, and #6, for which whole HID tracks were observed by the PCA. The dashed line in Fig. 5 marks the spots where the contribution from the power-law and disk fractions are equal. We find that any spectral state beyond the LHS moves across this line, as in the case of the outbursts

#1 and #6. An additional support of this argument comes from the energy-division diagram of the failed outburst #5: the peak at the right end of the track falls on the equal-division line.

3.4. Hardness-rms diagrams

To investigate the aperiodic variability of outburst #6, we derived the lightcurves in time bins of 1 ms (1/1024 s) using the PCA event mode data. We adopted the method of Gleissner et al. (2004) to calculate the rms by integrating the normalization

Fig. 3. The time evolution of the ASM (1.5–12 keV), PCA (1.5–12 keV), and BAT (15–50 keV) lightcurves, the PCA hardness ratio (defined using count rates corresponding to 6–19 keV and 3–6 keV), the inner temperature of disk T_{in} , the observed inner radius of the disk R_{in}^* , the power-law index Γ , the Fe luminosities L_{Fe} , and the total luminosities ($L_{\text{disk}} + L_{\text{pow}}$) for the recent outburst #6 of H 1743-322. The time bin-size for both ASM and BAT is one day, and for the hardness curve is 6400 s. The spectral bin-size is the same as in Table 1.

PSDs in the 0.1–64 Hz range. The resulting track for the evolution of the rms versus the hardness for outburst #6 is plotted in Fig. 6. As shown, the HRD has a typical profile of a BHC. During the LHS and HIMS (region I), the rms is always above 20%, and is positively correlated with the hardness. After the transition to the high state (region II), the rms dropped to ~5% and then moved to another LHS (region III, the LHS after the high state) with the rms values around 10%.

The PSD for all the OBSIDs in region I show the band-limited noise and type-C QPOs at ~1 Hz, which are typical of both LHS and HIMS. No QPO is detected in the other OBSIDs in the 0.01–500 Hz band, except for two OBSIDs (94413-01-03-02 & 94413-01-03-03), which are indicated by red points in Fig. 6, with type-B QPOs being found at ~4 Hz in region II. As shown already in the outbursts of other sources (e.g., Belloni 2009), this resembles the SIMS, which is usually located in the HRD within a small region beneath the transition from the HIMS to the HSS along the overall evolution. We note that the following OBSID 94413-01-03-04 (MJD 54990.39) was carried out half an hour later, with no QPO detected. If the SIMS is correlated with a jet as suggested in Belloni et al. (2005) and Fender et al. (2006), this may represent short-lived jitters on

timescales of around 1 h. All the other OBSIDs with the lower hardness in region II of the HRD have rms values of only a few per cent, no QPO, and a relatively larger disk fraction (at most around 75%). These results, together with the constant inner disk radius, indicate that the outburst #6 most likely reached the HSS, although not at the leftmost region of the HID as in the case of the 2003 outburst, and with a disk fraction almost equal to the borderline of the canonical HSS.

4. Discussion and summary

We have analyzed the 2009 outburst from H 1743-322 and compared our results with those of the other five outbursts observed by RXTE since 2003. We have found that the 2009 outburst of H 1743-322 differs from the other five outbursts observed to date. The HID shows the third complete counter-clockwise q-track for this source during the 2009 outburst. Interestingly, this track is located between the huge one in 2003 – which has a complete state transition to HSS – and the prior outburst from the source in 2008, which is a failed outburst. In the HID, the 2009 outburst leaves the LHS but does not reach the leftmost edge attained by the 2003 outburst. The softest point in HID has

Table 2. The spectral fit results of the softest OBSID during the six outbursts.

No	OBSID	Time	T_{in}	R_{in}^*	Γ_1	N	Flux	χ^2_{red} (d.o.f.)	$F_{\text{disk}}/\text{Flux}$
		(MJD)	keV	km			10^{-9}		
1	80146-01-15-10	52900.87	$0.91^{+0.02}_{-0.01}$	32^{+1}_{-1}	$2.35^{+0.28}_{-0.41}$	$0.14^{+0.15}_{-0.09}$	8.84 ± 0.13	0.92(36)	0.95 ± 0.01
2	90058-10-01-00	53197.32	$0.85^{+0.01}_{-0.01}$	32^{+1}_{-1}	$2.58^{+0.24}_{-0.24}$	$0.21^{+0.17}_{-0.10}$	6.38 ± 0.13	0.48(36)	0.94 ± 0.02
3	91050-06-05-00	53604.50	$0.83^{+0.02}_{-0.02}$	29^{+1}_{-1}	$2.28^{+0.07}_{-0.04}$	$0.83^{+0.14}_{-0.12}$	7.40 ± 0.16	0.79(63)	0.62 ± 0.01
4	93427-01-01-00	54481.98	$0.74^{+0.02}_{-0.02}$	34^{+2}_{-2}	$1.93^{+0.06}_{-0.05}$	$0.16^{+0.06}_{-0.05}$	4.47 ± 0.11	0.97(63)	0.73 ± 0.02
5	93427-01-12-00	54764.28	$0.78^{+0.04}_{-0.04}$	20^{+3}_{-2}	$2.14^{+0.06}_{-0.07}$	$0.86^{+0.16}_{-0.14}$	5.64 ± 0.11	0.91(63)	0.29 ± 0.02
6	94413-01-05-01	55002.36	$0.72^{+0.01}_{-0.01}$	37^{+2}_{-2}	$1.96^{+0.11}_{-0.12}$	$0.17^{+0.05}_{-0.04}$	4.68 ± 0.10	1.06(63)	0.75 ± 0.02

Notes. RXTE/PCA and RXTE/HEXTE data are used except for the first 2 outbursts. The OBSID, time, the inner temperature of disk, the inner radius of the disk, power law index, break for the energy, unabsorbed flux in the energy band 1.5–100 keV in units of $10^{-9} \text{ erg cm}^{-2} \text{ s}^{-1}$, and the fraction of DISKBB component are given. We notice that the spectral results of OBSID 80146-01-15-10 and of OBSID 93427-01-12-00 were also reported by McClintock et al. (2009) and Capitanio et al. (2009), respectively.

Fig. 4. Bolometric disk luminosities L_{disk} versus the temperature of the inner disk T_{in} for outbursts #1 (light blue), #5 (violet), and #6 (black); filled symbols indicate the thermal state OBSIDs, and open symbols for the LHS OBSIDs. The fit $L_{\text{disk}} \propto T_{\text{in}}^4$ was made to the thermal states of the outbursts. The time-bin size is one OBSID. The $T_{\text{in}}-L_{\text{disk}}$ relation for outburst #1 was also reported by McClintock et al. (2009).

only intermediate hardness values of 0.2–0.3, with X-ray spectra described by a power law with photon index < 2.4 , and a peak luminosity at $\sim 0.2 L_{\text{Edd}}$. This behaviour is similar to that of the “failed state transition” seen in the persistent bright black hole binary Cyg X-1 (Belloni 2009). Cyg X-1 is not a transient black hole system: it is found most of the time in the LHS with occasional transitions to the HSS, and a series of “failed state transitions” (see e.g., Fig. 8 in Belloni 2009). The “failed state transitions” only reach the SIMS with the X-ray spectra described by a power law with a photon index of ~ 2.15 and an inner disk temperature of $\sim 0.36 \text{ keV}$. For H 1743-322, the leftmost edge in HID of the 2009 outburst has a disk fraction of 75%, which is close to the lower limit observed for the XRB sample that complete the state transitions to the HSS (McClintock et al. 2009). The timing analysis shows, however, the typical features of reaching a HSS. Thus this outburst may be regarded as an interesting event, one in which the “failed state transitions” can be surpassed to reach the HSS. Therefore, H1743-322 could have a failed outburst at low flux and hardness level (as indicated in the 2008 outburst),

and high soft states with a rather wide range in the flux and the hardness values (as indicated in the 2003 and 2009 outbursts).

Advection-dominated accretion flow (ADAF) models (e.g. Esin et al. 1997) predict that the inner radius of the accretion disk should be radially truncated at hundreds of gravitational radii in the LHS and the disk should extend to the ISCO in the HSS. Indeed, GX 339-4 (Motta et al. 2009) and some other BHC follow this phenomenology. However, we note that in the latest outburst of H 1743-322, the accretion disk is closer to the innermost stable circular orbit in LHS rather than in a softer state. Similar conclusions based on the apparent change in the normalization of the DISKBB model were presented by Belloni et al. (1997) for the black hole binary GRS J1915+105 in the LHS using RXTE/PCA data, and Miller et al. (2006) for Swift J1753.5-0127 using XMM-Newton data. Reis et al. (2010) investigated 8 stellar mass black holes in LHS in the 0.5–100 keV broad band energy range using observations from XMM-Newton, Suzaku, Chandra, Swift, and RXTE. Their results also suggested that the disk in the LHS can remain at the ISCO, rather than being

Fig. 5. Left panels: plot of the luminosity contribution of the power law versus the disk component in units of the Eddington luminosity, L_{Edd} . The dashed lines in the left panels show where the two components are equal. Right panels: time evolution of the ratio of the former contributions.

truncated at hundreds of gravitational radii. The similarity of the outbursts from H 1743-322 with those seen in a typical black hole transient such as GX 339-4 is that the strongest outburst was followed by a sequence of weaker ones. A cool truncated accretion disk can be left over at around ISCO from the strongest outburst, and shows up in the low hard state of the following weaker ones (Meyer-Hofmeister et al. 2009).

For the 2009 outburst, we have found that the inner radius of the accretion disk has an anti-correlation with disk temperature. This trend was also present in the 2003 outburst (McClintock et al. 2009). One possible origin of this feature is reprocessing in the hot corona of the soft X-rays from the accretion disk, which can lead to an underestimate of the inner radius of the accretion disk. By assuming that the Comptonized component illuminates the disk, the irradiated disk can have a higher temperature and a smaller, inner disk radius than a non-irradiated disk (e.g. Yao et al. 2001; Gierliński et al. 2008). Gierliński et al. (2008) suggest that the disk spectra are largely unaffected in the HSS, in which the power-law flux is less than that of the disk, and the unaffected disk does not hold for LHS with $L_{\text{pow}}/L_{\text{disk}} > 1$. The underestimation factor of the inner radius depends on the strength of $L_{\text{pow}}/L_{\text{disk}}$ (see Fig. 6 in Gierliński et al. 2008). For example, in the case $L_{\text{pow}}/L_{\text{disk}} = 5$, the DISKBB will model the inner disk radius with a factor ~ 4 lower. In the LHS of the 2009 outburst from H 1743-322, we have $L_{\text{pow}}/L_{\text{disk}} > 10$. Therefore the inner disk could be underestimated by at least a factor > 5 . By taking this into account, the inner radius of the accretion at LHS should be comparable to or larger than that in other spectral states.

The transition from the low hard state to other states is usually accompanied by a balance between the thermal and the non-thermal emissions. We have found for the first time that this balance might account for the 2008 failed outburst of

Fig. 6. The hardness-rms diagrams (HRDs) of outburst #6. The arrows show the evolution in time, and the three regions denote the LHS/HIMS (region I), the HSS (region II) and the LHS/HIMS (region III) during the decay of the outburst. The two red points (almost on top of each other) most likely correspond to a short SIMS. Each data point is derived from one OBSID.

H 1743-322. As shown in Fig. 6, the tracks of the 2003 and 2009 outbursts go well across the so-called equal-division line, which represents equal thermal and non-thermal contributions to the overall observed emission, to leave LHS for more luminous states. Instead, the track of the 2008 outburst has a turn-over on top of the equal-division line. This may be understood if the shower of soft X-ray emission from the accretion disk cools the corona. The equal-division line in Fig. 6 may be indicative of emission balance between the corona and the inner accretion disk, beyond which the spectral state will transit to more luminous states through the LHS.

In summary, we have investigated the spectral and temporal evolution of the 2009 outburst, presenting the third complete q-track from H 1743-322. We have found that the 2009 outburst can reach the HSS by just overcoming a “failed state transition”. This finding, in addition to the other two complete q-tracks from the source, one showing a HSS at rather low hardness in 2003, and a failed outburst in 2008, may constitute an interesting testbed for further investigation of the evolution of the spectral states in BH XRBs.

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