

Evolution of the infrared Tully-Fisher relation up to $z = 1.4$

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ABSTRACT

Context. The Tully-Fisher relation represents a connection between fundamental galaxy parameters, such as its total mass and the mass locked in stars. Therefore, the study of the evolution of this relation in the optical and infrared bands can provide valuable information about the evolution of the individual galaxies through the changes found in each band.

Aims. This work aims to study the Tully-Fisher relation at high redshift in the B , V , R , I , and K_S -bands by comparison with the local relations derived from a large sample of galaxies in the redshift range $0.1 < z < 0.3$, processed in the same way, and with the same instrumental constraints that the high-redshift sample.

Methods. Using the large amount of photometric information available in the AEGIS database, we determined the optimal procedure for obtaining reliable k -corrections. Instrumental magnitudes were then k - and extinction corrected and the absolute magnitudes derived, using the concordance cosmological model. The rotational velocities were inferred from the widths of optical lines in DEEP2 spectra. At high redshift, this method is found to provide more accurate results than using the rotation curve, because of spatial resolution limitations. Morphology was determined by visual classification of the HST images. From the above information, the Tully-Fisher relations in B , V , R , I , and K_S -bands are derived for the local and high-redshift sample.

Results. We detect evolution in the B , V , and R -bands in the sense that galaxies were brighter in the past at the same rotation velocity. The change in luminosity is more noticeable in the bluer bands. This colour evolution, unnoticed in our previous work, is detected thanks to the more reliable k -corrections carried out in this paper, which included photometry from B to IRAC bands. The change in the $(V - K_S)$ and $(R - I)$ colours (for a fixed velocity) could be interpreted as an ageing of the stellar populations as consequence of the star formation decrease since $z = 1.25$. In addition, we conclude that spiral galaxies may have doubled their stellar masses in the past 8.6 Gyr.

Key words. galaxies: evolution – galaxies: fundamental parameters – galaxies: spiral – galaxies: kinematics and dynamics

1. Introduction

The relation between the luminosity and maximum rotational velocity of spiral galaxies (Tully & Fisher 1977) is an important distance estimator, which has been used for measuring the Hubble constant H_0 (e.g., Tully & Pierce 2000). The study of this relation at different cosmic times can probe evolution in galaxy properties such as its total mass, or the relation between dark and luminous matter (Zwaan et al. 1995). Any consistent model of galaxy formation and evolution should be able to reproduce the evolution of the Tully-Fisher relation (hereafter TFR). Its study therefore has important implications for determining fundamental cosmological parameters, the study of structure formation, and the evolution of disc galaxies (Navarro & Steinmetz 2000).

Traditionally, the internal kinematics of nearby galaxies has been measured using 21 cm line widths, but the limited sensitivity of radio telescopes prevents the effective use of this method at high redshift. Optical lines become important in evolutionary studies of TFR, and some authors have studied the relation between both velocity indicators. For example, Mathewson et al. (1992) compared the projected rotation velocity measured from $H\alpha$ rotation curves, with the velocity measured from integrated HI profiles at the 50% of the maximum flux level. They obtained a difference of 10 km s^{-1} , which they attributed to the HI widths measuring not only the rotational velocity but the

internal galaxy turbulence as well. This turbulence is important mainly in the most external regions of the galaxy, where the gravity is lower. However, the optical emission does not extend so far. The contribution of turbulence to the rotational velocity is negligible when using these lines. Vogt et al. (1996, 1997) modelled rotation curves using optical long-slit observations, and found a modest luminosity evolution in the B -band TFR ($\Delta M_B(\text{high } z\text{-local}) \leq -0.4$ at $\langle z \rangle = 0.5$). However, Rix et al. (1997) and Simard & Pritcher (1998) found a stronger evolution ($\Delta M_B \approx -1.5$ at $z = 0.4$), modelling the kinematics of disc galaxies in a similar way. Subsequent works have measured the evolution of the TFR, again finding different and even conflicting results. Ziegler et al. (2002), and Böhm et al. (2004), found a relation at high redshift ($z = 1$) shallower than that measured in local samples, and evidence of a luminosity evolution with look-back time of $\Delta M_B \approx -1$ mag at redshift $z = 1$. In contrast to previous findings, these authors claimed that less massive galaxies were brighter in the past (at fixed rotational velocity), while the most massive follow the local relation. However, other groups, such as Bamford et al. (2006), have detected luminosity evolution but no slope change. Moreover, these authors found differences in luminosity evolution ranging from -0.2 to -2 mag.

To cope with contributions to line widths other than those produced by rotational velocity, Weiner et al. (2006a) measured

kinematic line widths (σ_{1D}), and spatially resolved rotation and dispersion profiles. They combined line-of-sight rotation and dispersion ($S_{0.5}^2 = 0.5V_{\text{rot}}^2 + \sigma_{2D}^2$), and found that this combination correlates well with the integrated line width, demonstrating that σ_{1D} is a robust velocity indicator, that allows constructing scaling relations with velocity for a population of diverse kinematic properties (dispersion and rotation-dominated galaxies). In another paper, [Weiner et al. \(2006b\)](#) used σ_{1D} to study the evolution in the B and J -band TFRs, finding a slope evolution in both bands, and a larger intercept evolution in the B -band (-1.5 mag versus -0.5 mag in the J -band up to $z = 1.2$). This slope change, which implies a large evolution for the most massive galaxies, is in the opposite sense to that found by [Ziegler et al. \(2002\)](#) and [Böhm et al. \(2004\)](#). Although the B -band TFR derived by [Weiner et al. \(2006b\)](#) using resolved rotation velocities also implies that the intercept has evolved, the sample is too small and noisy to measure slope evolution. In addition, [Kassin et al. \(2007\)](#) demonstrated that the scatter in the stellar-mass TFR is lower when using $S_{0.5}$ instead of V_{rot} . In [Fernández Lorenzo et al. \(2009\)](#), we analysed the optical line widths to study the evolution of the TFR in B , R , and I -bands, demonstrating that all optical emission lines can be used to determine galaxy kinematics. We also constructed our local TFRs from data derived in a similar way to the sample of high redshift galaxies, and found evidence of luminosity evolution in all three bands for the greatest redshift range of our sample, $1.1 < z < 1.3$.

In contrast to the variety of results found in the optical bands, there seemed to be a consensus about the absence of evolution in the K -band TFR ([Conselice et al. 2005](#); [Flores et al. 2006](#)). However, [Puech et al. \(2008\)](#) find K -band evolution in the sense that galaxies had been fainter in the past, a result opposite to what is found in the optical bands (for example, [Bamford et al. 2006](#)). The disagreement between this work and that of [Flores et al. \(2006\)](#) (both use 3D spectroscopy for deriving velocities) is produced by the method for correcting the rotation velocity and, above all, the local relation used as a reference. Several works have studied the local TFR in the K -band finding different results, with a slope ranging from -6.88 ([Hammer et al. 2007](#)) to -11.3 ([Verheijen 2001](#)). Identifying the reason for these disagreements is crucial for obtaining reliable conclusions about the evolution of the infrared TFR by fixing the local relation. [Masters et al. \(2008\)](#) presented a universal calibration of the TFR in the 2MASS J , H , and K -bands. In all three bands, they showed that the relation is steeper for later-type spirals, and obtained a slope of -10.017 in the K -band correcting all galaxies to Sc type.

In this era of precision cosmology, increasingly large, deep and accurate galaxy surveys are required for the study of galaxy evolution. Within these surveys, the DEEP2 project ([Davis et al. 2003, 2007](#)) provides the galaxy spectra with a highest resolution. Furthermore, the large amount of photometry information available in the Extended Groth Strip (EGS) allows an additional step in the study of the TFR: its colour evolution. Models of disc galaxy formation predict different colour evolution depending on whether the evolution is mainly due to collapse or accretion (see, for example, [Westera et al. 2002](#)). In our previous work ([Fernández Lorenzo et al. 2009](#)), we studied for first time the $(R - I)$ TFR evolution up to $z < 1.3$, but the large scatter of this relation prevented us from distinguishing between one galaxy model and another.

In the present work, we extend our previous sample to include all the galaxies in the Extended Groth Strip (EGS) with DEEP2 spectra available, in the redshift range $1.1 < z < 1.4$, with the aim of confirming or not the prima facie evidence of evolution found previously ([Fernández Lorenzo et al. 2009](#)).

We also extend our study to the V and K_S -bands, in addition to the B , R , and I -bands, to study the most probable galaxy formation model, since the [Westera et al. \(2002\)](#) models predict larger colour differences in these bands. The local sample ($0.1 < z < 0.3$) is also extended to the whole EGS. Using this new local sample, we expect to set a more reliable TFR slope at $z = 0$ to compare with our high-redshift sample.

This paper is organised as follows. In Sect. 2, a description of the data and the sample selection criteria are provided. The study of k -corrections and derivation of the absolute magnitudes and rotation velocities are described in Sect. 3. The results are presented in Sect. 4, and our last two sections provide the discussion of the results and conclusions. Throughout this article, the concordance cosmology with $\Omega_{\Lambda 0} = 0.7$, $\Omega_{m0} = 0.3$, and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ is assumed. All magnitudes are in the AB zeropoint system.

2. Data and sample selection

The sample consists of galaxies in the Groth Strip Survey (GSS) sky region. The baseline for spectroscopy target pre-selection were the galaxies for which DEEP2 spectra (Data Release 3, DR3) in this field were available in the redshift ranges $0.1 < z < 0.3$ and $1.1 < z < 1.4$. The DEEP2 project ([Davis et al. 2003, 2007](#)) is a survey using the DEIMOS multi-object spectrograph ([Faber et al. 2003](#)) in the Keck telescopes, to study the distant Universe. The grating used was the 1200 l/mm one, covering a spectral range of 6500–9100 Å with a dispersion of 0.33 Å/px, equivalent to a resolution $R = \lambda/\Delta\lambda = 4000$. Despite both 1D and 2D spectra being available, in this work we use only the integrated spectra provided by DR3 ([Horne 1986](#)), because of the limited spatial resolution at high redshift. These spectra were extracted along the locus of constant lambda from the 2D spectra, using the routine `do_extract.pro`. Therefore, the 1D spectra provided by the DEEP2 team are corrected for the effects of tilted slits.

The photometric data used here are part of the AEGIS survey ([Davis et al. 2007](#)), for which B , R , and I -band photometry were taken with the CFH12K mosaic camera ([Cuillandre et al. 2001](#)), installed on the 3.6-m Canada-France-Hawaii Telescope (CFHT). These magnitudes are included in the DEEP2 photometric catalogue (Data Release 1, DR1, [Coil et al. 2004](#)) and the magnitude errors from sky noise and redshift are also available. The data in the V -band ($F606W$) were taken from the HST catalogue. As IR photometry, we used the K_S -band, which was taken with Palomar WIRC ([Bundy et al. 2006](#)). Finally, for high- z galaxies, we used data taken with the Infrared Array Camera (IRAC), on the Spitzer Space Telescope ([Barmby et al. 2008](#)) in the IRAC1 (3.6 μm) and IRAC2 (4.5 μm) bands. In Table 1, we present the limited magnitude and PSF (Point Spread Function) of each measurement. B , R , and I -band magnitudes already being corrected for Galactic reddening based on [Schlegel et al. \(1998\)](#) dust maps. The V -band was corrected following the same work. The Galactic reddening for the K_S and IRAC-bands is lower (< 0.005 mag) than the error in these magnitudes, so was not been considered.

Since in the present work, we study the TFR in V and K_S -bands, we limited the sample to galaxies with these photometric bands, further selected by restricting the sample to galaxies with emission lines in their spectra, necessary for obtaining the rotation velocity. In addition, for the high-redshift sample, we only selected galaxies for which with IRAC1 and IRAC2-bands photometry available, since they roughly

Table 1. Depth and average image quality of each measurement.

Band	Limiting magnitude	PSF (")
<i>B</i>	24.50 (8σ)	1.0
<i>V</i>	28.75 (5σ)	0.1
<i>R</i>	24.20 (8σ)	1.0
<i>I</i>	23.50 (8σ)	1.0
K_S	22.5 (5σ)	1.0
IRAC1	24.00 (5σ)	1.8
IRAC2	24.00 (5σ)	2.0

correspond to NIR photometry in the rest-frame and can provide a more reliable k -correction in the K_S -band.

The inclination angle (i) must also be included as a selection criteria, since no correlation between magnitude and rotation velocity can be observed for inclinations lower than 25° (Fernández Lorenzo et al. 2009). The inclination angle was calculated from the major to minor axis ratio as found in the HST catalogue, while the inclination errors were obtained by comparing with the inclination derived using SExtractor (Bertin & Arnouts 1996) in the combined V+I HST images. In this way, we estimated the mean error to be $\pm 2.5^\circ$ for the local sample, and $\pm 6^\circ$ for the high-redshift sample. Galaxies almost edge-on are more affected by extinction, but in our sample all galaxies have inclinations lower than 80° . Therefore, the inclination of the final sample ranges between 25° and 80° .

The second selection criteria was the morphology. Full details of the morphology classification were given in Fernández Lorenzo et al. (2009) to which the interested reader is referred. To select spiral galaxies, we performed a visual classification of every galaxy using HST images. The objects were divided into five groups: elliptical/S0 (1%), spirals (66%), irregulars (5%), interacting (5%), and unknown (23%). To classify our visually unknown objects, GIM2D (Simard 1998) was used. The objects with a Sérsic index lower than 2.5 were considered as spirals. The objects visually classified as unknown that could not be fitted with GIM2D (5% of “unknown” objects), were discarded. Finally, after applying all these criteria, we were left with 128 galaxies in the local sample, and 113 in the high redshift ($1.1 < z < 1.4$) sample.

3. Data analysis

The luminosity and rotation velocity of disc galaxies are the parameters involved in the TFR. At high redshift, several, not obvious, corrections are necessary to reliably obtain these parameters. Moreover, some corrections can dramatically change the results emulating evolution. Apart from Galactic extinction, which affects the observed magnitudes of local and high redshift galaxies in an identical way, the corrections required to obtain the absolute magnitudes in each band are the k -correction and the intrinsic extinction. We now try to determine the most effective way of calculating the absolute magnitudes and kinematics of spiral galaxies.

3.1. Rest-frame magnitudes

To calculate the k -correction in the B, V, R, I , and K_S -bands, we need to know the spectral energy distribution (SED) of the galaxy. Since the SED is generally unknown, or at least not known with the required photometric accuracy, it is necessary to use an appropriate set of templates to reproduce the SED of each galaxy. In addition, the errors in the photometric information

used to fit the template and its spectral coverage can strongly affect the best-fit template and thus the k -corrections. In the present work, we use the photometry included in the AEGIS catalogues. Since different instruments were used to obtain the magnitudes at different wavelengths, the same aperture might not enclose a similar fraction of the light in each band, due to the PSF, the seeing or the pixel scale, that act more or less spreading the object image. Then, we choose the photometry as close to the total magnitudes as possible, to fit the SEDs. For the B, R , and I -bands, Coil et al. (2004) measured the total R magnitudes from an aperture that systematically contains the whole galaxy flux, whereas the B and I magnitudes were corrected to total ones using the colours $(B - R)_{1''}$ and $(R - I)_{1''}$, respectively. For the V -band, we chose the MAG_BEST in the catalogue, which provides a nearby to total magnitude of the object. For the K_S -band, we used MAG_AUTO because it is available for a larger number of objects than the aperture magnitudes, and we confirmed that there is no significant difference between aperture and MAG_AUTO photometry because of the larger errors in this band. Finally, for the IRAC-bands, we used the MAG_APER photometry in an aperture of $1.5''$ (only for the high redshift sample) because aperture corrections derived from average mosaic PSFs were applied to the aperture magnitudes but not to MAG_AUTO, so we did not use MAG_AUTO in this case. We checked that there is a good correlation between the photometry in $1.5''$ and $2.14''$ apertures, and that the whole galaxy could be enclosed inside the aperture of $1.5''$.

Since the photometric bands available in this work do not match the rest-frame optical magnitudes, and the K_S -band photometry is very noisy, we performed a careful and systematic study to determine the most reliable way of performing a k -correction on our data using various sets of templates and methods (see Appendix A). We found that the best result is obtained by performing a nonnegative linear combination of five templates based on the Bruzual & Charlot (2003) stellar evolution synthesis codes using kcorrect. However, when a noisier band is present in the data, the k -corrections calculated by this code are unsuitable, because the photometric error in the observed quantity propagates to the rest-frame magnitude (Appendix A). In this case, we found that the optimal k -correction is achieved using photometric information that roughly matches the rest-frame band for which we wish to calculate the k -correction. The rest-frame magnitude in the noisy band is then more accurately calculated directly from the best-fit template instead of k -correcting the observed data. This is the case for our K_S -band data that is roughly matched at rest-frame by the IRAC2-band. On the other hand, the k -correction in our optical bands can only be determined by interpolation, a procedure that requires available information at larger and shorter wavelengths than the band that we wish to correct. Finally, we calculated the rest-frame magnitudes used in this work by convolving the rest-frame best-fit kcorrect template with the filter response in each band.

3.2. Intrinsic extinction

The absolute magnitudes were inferred from the luminosity distance corresponding to the measured redshifts, by assuming a concordance cosmology. Finally, the absolute magnitudes were corrected for intrinsic extinction. This correction is basically based on inclination. In this work, we adopted the Tully et al. (1998) procedure, which is valid for local galaxies. Since the dust content of the galaxies may have evolved with redshift, our extinction correction may be either under or overestimated, so

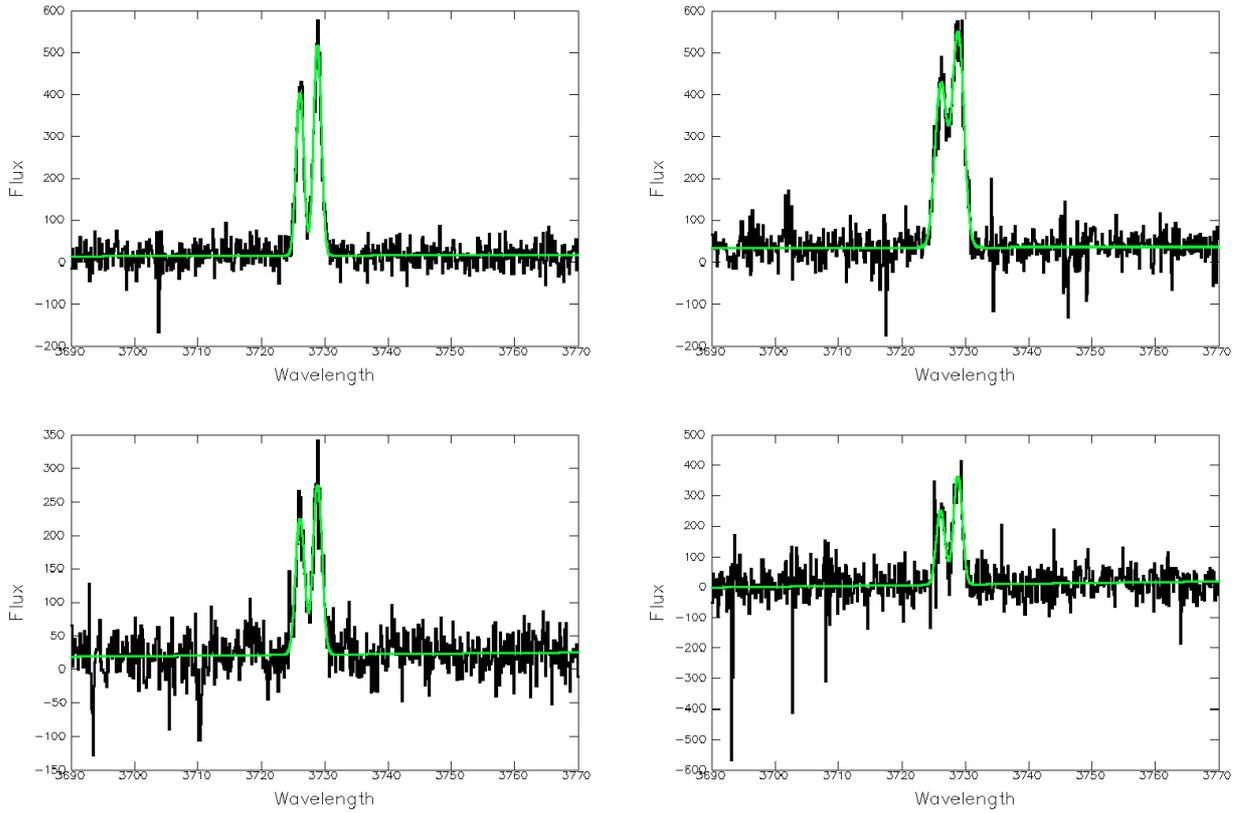


Fig. 1. Example spectra (not flux-calibrated) of the high- z galaxies in our sample. The x -axis represents the rest-frame wavelengths in \AA . The green line is the fit to the $[\text{OII}]\lambda\lambda 3727 \text{ \AA}$ double line obtained with `dipso`, imposing that both lines had the same width.

we consider both possibilities below. According to this method, the extinction A_λ as a function of inclination i in the λ -band, is defined to be

$$A_\lambda^{i-0} = \gamma_\lambda \log(a/b), \quad (1)$$

where

$$\gamma_B = -0.35(15.31 + M_B) \quad (2)$$

$$\gamma_R = -0.24(15.91 + M_R) \quad (3)$$

$$\gamma_I = -0.20(16.61 + M_I) \quad (4)$$

$$\gamma_K = -0.045(18.01 + M_K), \quad (5)$$

and a/b is the galaxy major-to-minor axis ratio. These equations are valid for the magnitudes in the Vega zeropoint system, so our magnitudes were converted into the Vega-system to perform the extinction correction. For the V -band extinction, we used the Calzetti's law (Calzetti et al. 2000) obtaining $A_V^{i-0} = 0.8 A_B^{i-0}$.

3.3. Kinematic line widths

Most local velocities have been derived historically from radio measurements, usually from 21 cm line widths at 50 per cent of the peak intensity (for example, Giovanelli et al. 1997). As already known (Bosma 1981), the rotation curves are not perfectly flat at large radii. Since the observed HI and $\text{H}\alpha$ gas emission do not span the same radii, the velocities measured from both lines can provide different results. In addition, turbulent motions broaden the HI profile and affect both the optical versus radio velocity width determinations (Mathewson et al. 1992). Moreover,

the existence of three types of rotation curves, depending on the relation between the maximum velocity and the velocity of the flat region (Verheijen 2001), can complicate the comparison. Mathewson et al. (1992) compared the rotation velocity measured from $\text{H}\alpha$ rotation curves with the velocity measured from integrated HI profiles at the 50% level. They measured a ratio of $2V_c(\text{H}\alpha)/W_{50}^i(\text{HI}) \sim 0.94$, which they attributed to the contribution of the turbulence to the measured HI velocity, thus overestimating the rotational velocity obtained using this line, as already mentioned. This contribution is more important in the most external regions of the galaxy where the gravity is lower, while the optical emission does not spread to this region. Applying diverse corrections, Rix et al. (1997) found a factor of 0.86 ± 0.04 between $2V_c(\text{H}\alpha)$ and $W_{20}^i(\text{HI})$, which must be considered when comparing our results with those obtained for other local samples using 21 cm.

In the DEEP2 survey, the long axis of every individual slitlet is orientated along the major axis of each galaxy, with sufficient accuracy to derive rotation velocities. Masks in each field are generally oriented at 2 different mask PAs, which are -49° and 41° in EGS. In a given mask, each slitlet can be oriented across the range $\pm 30^\circ$ with respect to the long axis of the slitmask. By combining the mask and slit PAs, the maximum slit misalignments (δ) are either 11° or 19° . We checked this for a subsample of galaxies, in the ACS images, and found good agreement between the slit and the major axis, differences always being lower than 20° . This slit misalignment causes a negligible error in the velocity as already pointed out by Giovanelli et al. (1997) for PA offsets smaller than 15° . Weiner et al. (2006a) concluded that the spatially resolved kinematics (2D) depends strongly on slit position angle alignment with galaxy major axis, but that integrated

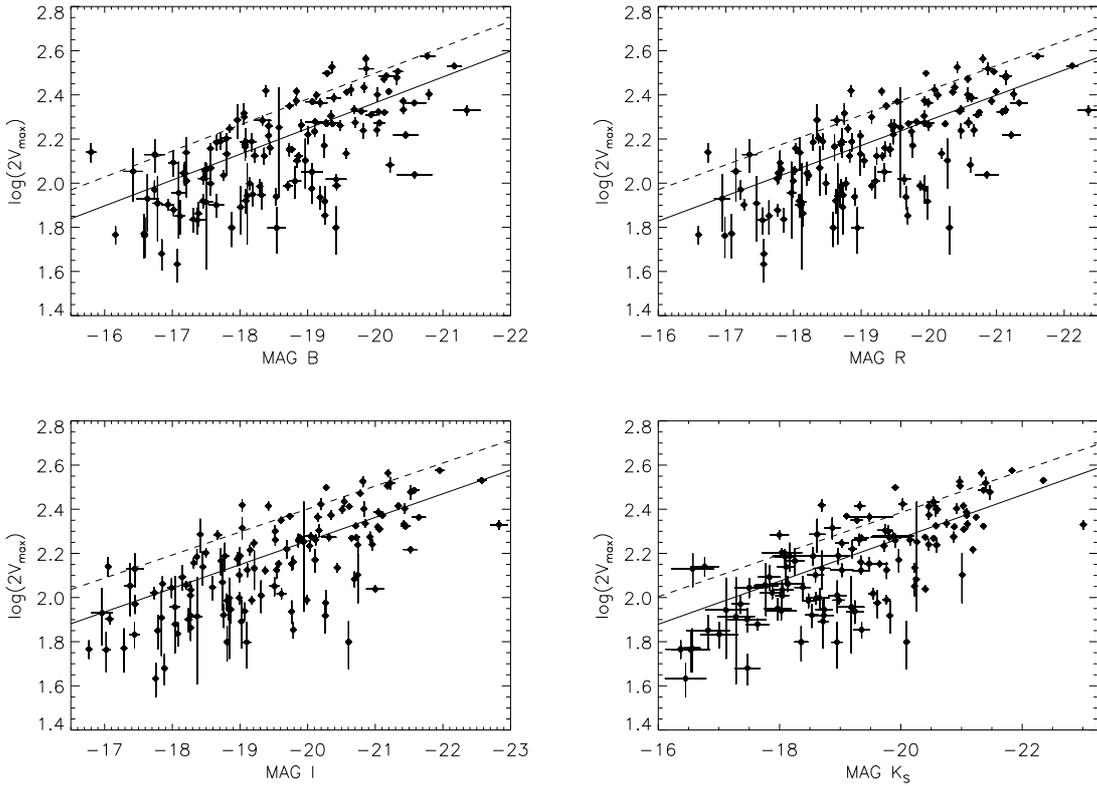


Fig. 2. Comparison between local TFRs obtained from DEEP2 data ($0.1 < z < 0.3$) in $B, R, I,$ and K_S -bands, and Verheijen (2001) local TFRs (dashed lines). The solid line represents the linear fit to the DEEP2 data. The parameters of the fit are shown in Table 2.

line width does not. For example, applying the standard $\cos^{-1}(\delta)$ correction, the difference in $\log(V_{\max})$ for a $\delta = 20^\circ$ would be ~ 0.012 for a galaxy similar to the Milky Way at $z = 0.2$, which is lower than the error in $\log(V_{\max})$ for our local sample. We therefore chose not to apply the slit misalignment correction.

The optical line used in this work to derive the rotation velocity is not the same for each sample of galaxies because of the limited DEEP2 spectral range. We measured the $H\alpha$ line for the local sample, and the $[\text{OII}]\lambda\lambda 3727 \text{ \AA}$ double line for the high-redshift sample. In Fernández Lorenzo et al. (2009), we demonstrated that both optical emission lines can be used to determine disc-rotation velocities with the aim of comparing different redshift samples. To calculate the rotation velocity, we used

$$2V_{\max} = \frac{\Delta\lambda c}{\lambda_0 \sin(i) (1+z)}, \quad (6)$$

where i is the inclination angle, λ_0 the line central wavelength at $z = 0$, and $\Delta\lambda$ the line width at 20% of peak intensity. Equation (6) is derived from the Doppler effect, applying the corrections related to cosmology and inclination. We determined $\Delta\lambda$ with the Gaussian fitting routines in the Starlink package *dipso* (Howarth et al. 1996), which calculate the full-width half-maximum (FWHM). The spectral range and resolution of the DEEP2 spectra were designed to resolve the $[\text{OII}]\lambda\lambda 3727 \text{ \AA}$ doublet, and all our spectra fulfill this aim. In Fig. 1, we represent four cases with different S/N and line blending. To measure the $[\text{OII}]\lambda\lambda 3727 \text{ \AA}$ double line, we assumed that both lines have the same width. We corrected for the instrumental width as in Fernández Lorenzo et al. (2009). Where the profile is Gaussian, the line width at 20% of the peak intensity can be compared with

other velocity width measurements such as the *FWHM* and velocity dispersion, σ , using

$$\sigma = \frac{FWHM}{2.35} = \frac{W_{20}}{3.62}. \quad (7)$$

Since our objects have been classified visually as spiral galaxies (see Sect. 2), we assume that all galaxies rotate. However, we cannot verify the accuracy of this assumption with the existing data for the high redshift sample.

4. Results

4.1. Local relations

In Fernández Lorenzo et al. (2009), we derived the local TFRs by fitting the DEEP2 data points in the redshift range $0.1 < z < 0.3$. In this work, we extend the local sample to the entire Groth field with ACS data available, to determine a more reliable TFR slope at $z = 0$ for comparison with our high-redshift sample. Although the bands that we wish to study are V and K_S , we derived the local TFRs in $B, R,$ and I -bands as well, to compare with our previous results.

We performed an error-weighted least squares fitting to the DEEP2 data points in the redshift range $0.1 < z < 0.3$, to estimate the slope a , and the intercept b of the local TFR in each band. A more detailed explanation of the fitting procedure can be found in Fernández Lorenzo et al. (2009). We adopted $\log(2V_{\max})$ as the dependent variable in the fit. This is the so-called inverse TFR, which is less sensitive to luminosity incompleteness bias (Willick 1994; Schechter 1980). The results obtained for each band are given in Table 2. In addition, we show

Table 2. Parameters of the Tully-Fisher relations obtained by fitting DEEP2 data in the redshift range $0.1 < z < 0.3$ used as local reference.

Band	This work			Verheijen (2001)			
	a	b	σ_{total}	A	B	A	B
B	0.030 ± 0.045	-0.117 ± 0.002	0.129	0.25 ± 0.38	-8.55 ± 0.15	1.24 ± 0.82	-8.5 ± 0.4
V	-0.026 ± 0.046	-0.116 ± 0.002	0.126	-0.22 ± 0.40	-8.62 ± 0.15		
R	0.004 ± 0.042	-0.114 ± 0.002	0.121	0.04 ± 0.37	-8.77 ± 0.15	1.54 ± 0.82	-8.9 ± 0.4
I	0.114 ± 0.039	-0.107 ± 0.002	0.119	1.07 ± 0.36	-9.34 ± 0.17	3.05 ± 0.83	-9.6 ± 0.4
K_S	0.316 ± 0.040	-0.098 ± 0.002	0.119	3.23 ± 0.41	-10.20 ± 0.21	5.04 ± 0.92	-10.5 ± 0.4

Notes. In addition, we show the values calculated by Verheijen (2001) for B, R, I and K_S -bands.

the values calculated by Verheijen (2001) for the local ‘‘RC/FD sample (without NGC 3992)’’ case, by fitting V_{max} (the same sample used for comparison in our previous work). In Fig. 2, we represented both relations in the B, R, I , and K_S -bands. As in our previous work, the slope of the TFR in each band is very similar to that found locally by Verheijen (2001), even for the K_S -band, and consistent within his errors. Moreover, our K_S -band TFR slope is consistent with that derived by Masters et al. (2008) after correcting for incompleteness, morphology, and luminosity dependence. Unlike the slope, we obtained a zeropoint of the relations lower than those of Verheijen (2001). However, this difference is very similar in all bands, so it may be related to the true rotation velocity versus observed line width relation. Rix et al. (1997) simulated observations of disc velocity fields with true circular velocity V_c , and found that $\langle \sigma \rangle / V_c = 0.6$, where the observed $\langle \sigma \rangle$ represents an average over all inclination angles. This result implies that $2V_c = 0.92 W_{20}$. In our case, a ratio of $\langle \sigma \rangle / V_c \sim 0.4$, which implies that $2V_c = 1.38 W_{20}$, would be necessary to explain the difference found in our local TFRs. Rix et al. (1997) found that factors as line profile asymmetries, fiber size, or inclination effects work towards reducing $\langle \sigma \rangle / V_c$. The spectra that we use were not acquired using fibers, but may be affected by e.g. slit misalignment, which also reduces $\langle \sigma \rangle / V_c$. Since we wish to compare our high redshift sample with a sample of local galaxies observed with the same instrumentation and analysed following the same procedures, both samples are expected to be affected in a similar way by these factors. Consequently, the possible differences found in the TFR should be meaningful.

4.2. High-redshift sample

In Fig. 3, we present B, V, R, I , and K_S -band TFRs, for the redshift range $1.1 < z < 1.4$, to test the evidence of evolution found in our previous work but with a larger sample of galaxies. To examine the intercept evolution, we fixed the slope to have the same value as for the local TFR in each band. The large scatter in the high- z relation is noticeably smaller when we limit the sample to galaxies with inclination $i > 45^\circ$ (see Fig. 3), which is the lowest value used in some works (for example, Nakamura et al. 2006). When we fitted the TFR separately for objects with $i < 45^\circ$, we found a smaller amount of evolution than for all objects together. For galaxies with $i > 45^\circ$, the opposite result was found, i.e. we measured greater evolution in the zeropoint. For galaxies with low inclination, an error in the inclination angle yields larger velocity errors than for more inclined galaxies. Nevertheless, the error in the extinction will be more noticeable for galaxies with high inclination. If the error in the extinction were the reason for the differences found for both inclination ranges, then for the K_S -band there would be no difference in the evolution found for galaxies with low and high inclination, because the extinction is negligible in this

Table 3. Magnitude evolution found in each band assuming a constant slope.

Band	ΔM	$\sigma_{\Delta M}$
B	-1.125	0.535
V	-0.708	0.547
R	-0.646	0.531
I	-0.298	0.545
K_S	-0.213	0.610

band. However, when both inclination ranges were fitted independently for the K_S -band, we found that the difference was identical to those for the other bands. Therefore, a velocity error caused by an inclination error seems to be the reason for the disagreement between both inclination ranges. Although in terms of the velocity, the results are more reliable for galaxies with $i > 45^\circ$, the evolution in the zeropoint of the TFRs does not change significantly if the whole sample is used. Therefore, we fitted all galaxies together to increase, into the bargain, the statistical significance of the result. Measuring the scatter in the TFR and its evolution are important to constraining models of galaxy formation. Masters et al. (2008) observed that the scatter in the local K -band TFR increases with decreasing rotation velocity. In our local sample, the scatter is 1.21 mag in the K_S -band, but reduces to 0.9 mag when we consider objects with $\log(2V_{\text{max}}) > 2.2$. For $\log(2V_{\text{max}}) = 2.2$, the scatter in the K -band TFR obtained from Masters et al. (2008) is ~ 0.9 mag. However, the scatter in our high-redshift sample (1.23 mag) is larger than the scatter of ~ 0.5 mag obtained in Masters et al. (2008) for $\log(2V_{\text{max}}) = 2.5$, that is the average velocity of our high redshift sample. Although the scatter does not differ significantly from those of our local relations (in agreement with the result found by Conselice et al. 2005, for their samples at $z < 0.7$ and $z > 0.7$), it is two to three times higher than those of the local relations of Verheijen (2001), Masters et al. (2008), or Hammer et al. (2007). This is probably due to the measurement errors being larger than for $z = 0$ galaxies.

As we assumed that the slope does not change with the redshift, we were able to infer the evolution in the magnitude for each band, from the difference between the high-redshift and local TFRs for a fixed velocity. The results are shown in Table 3. For the B -band, we found a change in magnitude of $\Delta M_B = -1.13$ at $z = 1.25$, which is larger than 2σ , and consistent with our previous work (Fernández Lorenzo et al. 2009), and the results of other authors (e.g., Böhm et al. 2004; Bamford et al. 2006). However, for both R and I -band, we found weaker evidence of evolution, which can probably be attributed to the more reliable k -corrections carried out in this work. The most interesting result is a smaller difference in the intercept as the band wavelength increases, the change in the zeropoint for I and K_S -bands being smaller than 1σ . The difference in magnitude

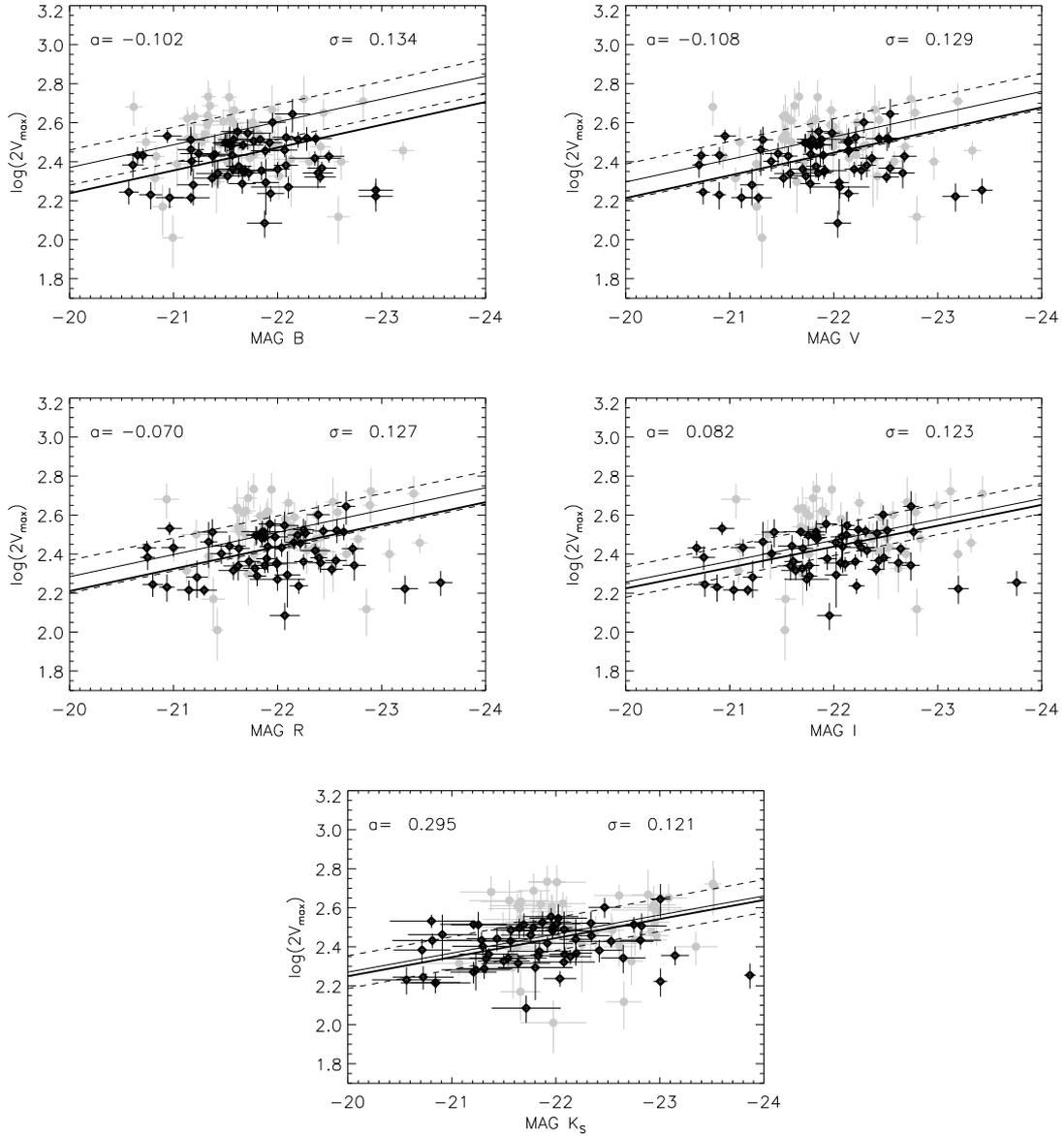


Fig. 3. Tully-Fisher relation in B , V , R , I , and K_S -bands for the redshift range $1.1 < z < 1.4$. Black diamonds are galaxies with inclination $45^\circ < i < 80^\circ$, and grey points are galaxies with $25^\circ < i < 45^\circ$. Thin and dashed lines represent the local TFR and its 2σ uncertainty in the offset. The thick line is the weighted least squares fit to all points. The fit was performed by fixing the slope to the local one in each case.

found for the K_S -band is compatible with no evolution, and agrees with the results of [Conselice et al. \(2005\)](#) and [Flores et al. \(2006\)](#). Nevertheless, [Puech et al. \(2008\)](#) found evolution in the K_S -band TFR, in the sense that galaxies were fainter in the past. They attributed their disagreement with [Flores et al. \(2006\)](#) to both the local relation slope used as reference and the more accurate measurement of the rotation velocities of [Puech et al. \(2008\)](#). However, if the [Puech et al. \(2008\)](#) data are compared with the [Verheijen \(2001\)](#) local relation (the same used by [Flores et al. 2006](#)), no evolution is found. Therefore, the local TFR slope plays a fundamental role in the evolution of the K_S -band TFR, and is studied separately in the next section.

5. Discussion

Several authors have studied the infrared TFR in both the K and K_S -bands. To compare with these works, we need to know

the relation between both bands. From [Grocholski & Sarajedini \(2002\)](#), a Bessell K -band is related to the 2MASS K_S -band as $K = K_S - (-0.044 \pm 0.003) - (0.000 \pm 0.005)(J - K)$. Assuming that the 2MASS and Palomar K_S -bands are equivalent, the difference between the K and K_S -bands is within the errors in our K_S -band absolute magnitudes, and therefore we can compare our results with those obtained in other works.

5.1. The slope in the local K_S -band TFR

To study the discrepancies found by the authors in the slope of the local K_S -band TFR, we used the Millenium simulation database ([Springel et al. 2005](#)) produced by the Virgo consortium. From this simulation, the catalogue that we used, contained in the millimil database, is the table DeLucia2006a, from which we took the data corresponding to maximum rotational velocity (v_{\max}), Johnson B , V , R , I , and K absolute rest-frame

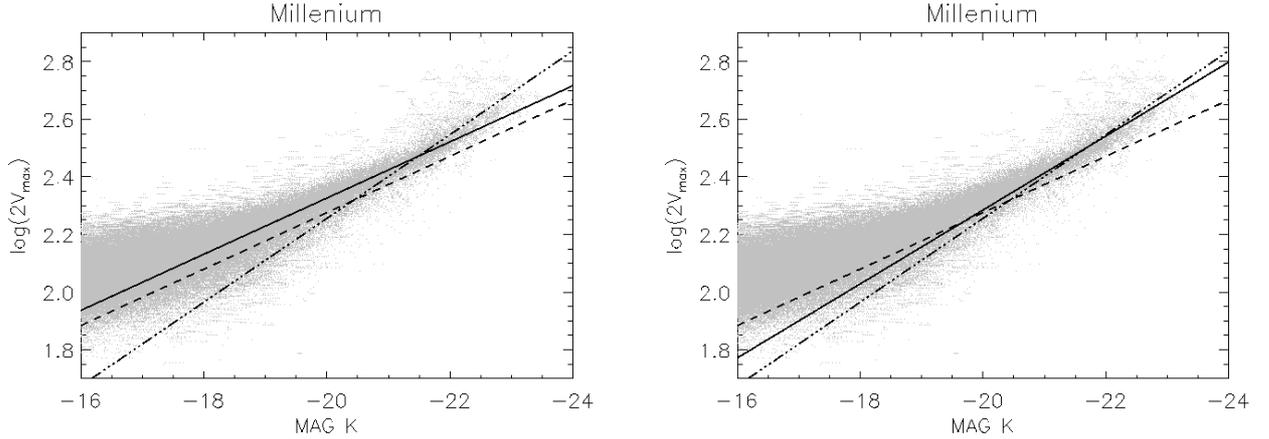


Fig. 4. Local Tully-Fisher relation in the K_S -band for the Millenium mock galaxies. The grey points are the objects with $M_{\text{bulge}} < 0.15 M_{\text{star}}$ and $M_K < -16$. The dashed line is the relation obtained with our local sample and the dot-dash line is the local relation obtained by Hammer et al. (2007) for a sample of SDSS galaxies with $\log(V_{\text{flat}}) \geq 2.2$. The solid line is the fit to the data for $M_K < -18$ (left) and $M_K < -21$ (right).

magnitudes (dust extinction included), redshift, mass of bulge (M_{bulge}), and the total mass in stars (M_{star}) (De Lucia & Blaizot 2007; Croton et al. 2006). Using the upper value of bulge-to-total (B/T) mass ratio for late-type galaxies from Laurikainen et al. (2007), we selected the objects with $M_{\text{bulge}} < 0.15 M_{\text{star}}$ as being representative of spiral galaxies. In Fig. 4, we present the data for $M_K < -16$. The K -band magnitudes were converted into the AB system using $M_K(\text{AB}) = M_K(\text{Vega}) + 1.9$ (Hewett et al. 2006). The slope of the local TFR obtained using simulated galaxies seems to depend on the magnitude range considered. For instance, when we fit galaxies with $M_K < -18$, we obtain a slope very close to our sample of local galaxies. However, for $M_K < -21$, the slope is closer to the K -band TFR found by Hammer et al. (2007), who restricted their sample to SDSS galaxies with $\log(V_{\text{flat}}) \geq 2.2$ (the local TFR used by Puech et al. 2008). For the other bands, we found the same variations in the slope as found for the K -band.

If we assume that the slopes of the TFRs depend on the magnitude range considered, then our local relations cannot be compared with our high redshift sample, which consists of galaxies brighter than $M_K = -20.5$. To be able to compare our high redshift data with a local relation derived for the same magnitude range, we used the Millenium data. We calculated the difference in the zeropoint between our local sample and Millenium data by setting the same slope across the same magnitude range as for our local one. As happened when comparing with Verheijen (2001) local relations, we obtained a similar difference in the intercept of the relation for all bands, whose value corresponds to the ratio of rotation velocity to line width. Nevertheless, the difference in the zeropoint was found to be lower, the zeropoint of the Verheijen (2001) relations being greater than the Millenium local zeropoints. We then calculated the Millenium local TFR relations for the brightest galaxies and we applied the shift to the zeropoint required to compare with our high redshift sample. Finally, we fit the data of the high redshift sample by setting the slope to the new local one determined using Millenium data. We found a greater change in luminosity for all bands, but for the I -band, the difference in the zeropoint of the high redshift relation with respect to the Millenium local one was still found to be within 2σ (where σ is the uncertainty in the local offset), and for the K_S -band the difference is within $\sim 1\sigma$. Therefore, we found smaller luminosity evolution as passing from blue to

redder bands, consistent with the results obtained in the present work. This result supports our measurement of no evolution of the TFR for the K_S -band.

5.2. The colour evolution of disc galaxies

The colours of galaxies provide information about their stellar content and, by using evolutionary models, the history of star formation. Studying the colour evolution versus rotational velocity thus allows us to study the change in the stellar content of a galaxy relative to its total mass. In Fig. 5, we present the colours ($V - K_S$) and ($R - I$) versus rotational velocity for local and high-redshift galaxies. For ($V - K_S$), we have a larger dispersion, and the evolution of ~ 0.35 mag found by setting the slope to the theoretical local ($V - K_S$) relation, is within the dispersion (smaller than 1σ). However, there is a clear change of ~ 0.33 mag with redshift, in the ($R - I$) colour at fixed velocity ($>2\sigma$). The galaxies are therefore redder today, which may represent an ageing of the stellar population caused by a star formation decrease in the past 8.6 Gyr. As in the case of the evolution in R and I -band TFRs, we found a different result from that found in Fernández Lorenzo et al. (2009), because of the more reliable k -corrections carried out in the present work, where we used additional photometric information.

The colour evolution of simulated disc galaxies with redshift has been investigated by other studies. In Fernández Lorenzo et al. (2009), we compared the results with Westera et al. (2002), who studied the colour evolution of disc galaxies in two models of galaxy formation called the accretion and collapse models. Using these models, the change in colours observed here can be more accurately reproduced by the accretion model, but the predicted evolution in the ($V - K$) colour would be twice that found here, and the change in magnitude is in the opposite sense to that predicted by the simulations. In a subsequent paper, Westera et al. (2007) investigated the effect of the initial mass function (IMF) on the colour evolution of disc galaxies by considering two different IMFs: Salpeter and Kroupa. The evolution in the ($g - K$) colour derived from their work (for the Kroupa IMF without absorption) would again be larger than the result found here, assuming $(g - K) \approx (V - K_S)$, mainly because of the change in the K -band of ~ 1 mag predicted by the simulations. However, this change in K would lead, according

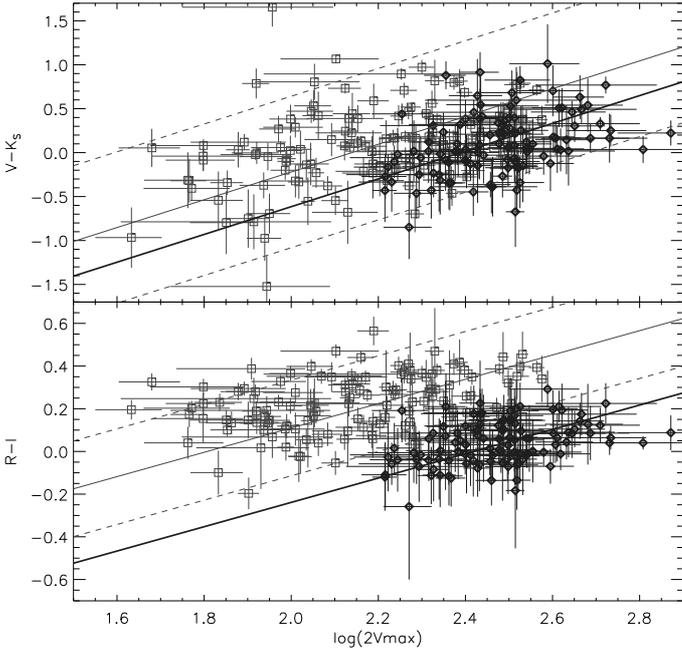


Fig. 5. Colour ($V - K_S$) vs. rotation velocity (up) and colour ($R - I$) vs. rotation velocity ($down$). Black diamonds represent the galaxies in the redshift range $1.1 < z < 1.4$, while open squares are the galaxies in the redshift range $0.1 < z < 0.3$. Difference between the local ($V - K_S$) TFRs (up) and ($R - I$) TFRs ($down$) (thin lines) and its 2σ uncertainty (dashed lines) are shown. Thick lines are the fit to the local redshift data when setting the slope to the local relation in each case.

to their inferred evolution of the M_*/L_K , to a growth of a factor ~ 6 in the stellar mass since $z = 1.25$. Nevertheless, according to the M_*/L_r and M_*/L_i determined in the same work, this change in stellar mass would produce an evolution of $\Delta M_r = 0.36$ and $\Delta M_i = 0.26$ from $z = 1.25$ to $z = 0$, which would imply that galaxies were fainter in the past. This result is the opposite of that found here, and that found by van Starkenburg et al. (2006) for the R -band. Our results are therefore inconsistent with the simulations of Westera et al. (2002, 2007).

5.3. Comparison with other works

As we discuss in Fernández Lorenzo et al. (2009), the evolution found in the optical bands is more likely due to a change in luminosity than in velocity. Conselice et al. (2005) and Flores et al. (2006) obtained the same result as this work, i.e. no evolution in the K -band TFR, whereas Puech et al. (2008) detected a change in the K -band TFR zeropoint, which they attributed to an average brightening of galaxies since $z \sim 0.6$ by 0.66 ± 0.14 mag. They explain this brightening as a growth in stellar mass by a factor of ~ 2.5 since $z = 0.6$, as evaluated from the evolution in $\log(M_{\text{stellar}}/L_K)$ found by Drory et al. (2004). They claim that the result is consistent with the gaseous O/H phase abundance of $z \sim 0.6$ emission-line galaxies, which is half of that found in present-day spirals (Rodrigues et al. 2008).

However, assuming a change in $\log(M_{\text{stellar}}/L_K)$ with redshift in the way described by Drory et al. (2004), but no evolution in the K -band TFR as found in the present work, we conclude that the spiral intermediate-mass galaxies have doubled their stellar masses since $z = 1.25$, which agrees rather well with expectations from stellar mass density studies (Drory et al. 2004). From the mass-to-light ratios in u , r , i , and K -bands of

Westera et al. (2007, who found a similar result for the K -band to that of Drory et al. 2004, in the past 9 Gyr), but considering our growth in stellar mass by a factor of 2, the evolution in the optical bands would be $\Delta M_u = -1.3$; $\Delta M_r = -0.79$; $\Delta M_i = -0.67$ (for a Salpeter IMF without absorption), which are consistent with the result found in the present work. Assuming at least the same evolution in the K -band magnitude at $z = 1.25$ as that found by Puech et al. (2008) at $z = 0.6$, the stellar mass should have grown at least by a factor of 3.5 when applying the evolution in $\log(M_{\text{stellar}}/L_K)$ found by Drory et al. (2004). However, the stellar mass should have grown by a factor of 2.6 since $z = 1.25$, if the change in the gaseous O/H abundance were attributed to all gas being transformed into stars (Rodrigues et al. 2008). Given the short lives of the most massive stars responsible for the enrichment of the ISM (interstellar medium), and the time required for cooling the gas to form new stars, the change in O/H in the past 8.6 Gyr will not only be due to a rise in O, but also a drop in H. The change in stellar mass inferred from the metallicity evolution is therefore probably an upper limit.

The stellar mass can also increase by means of galaxy merger processes or the accretion of small galaxies. In the first case, the galaxy disc would be destroyed if the mass of the merger were larger than $\sim 15\%$ of the parent galaxy mass. The accretion of small galaxies, although a probable contribution to increasing the stellar mass, is not enough to explain a growth of a factor 3.5. Hence, our result of a doubling in stellar mass since $z = 1.25$ seems more plausible and consistent. The increase in stellar mass is probably caused by gas being transformed into stars, which would be consistent with there being a higher SFR density at $z \sim 1-2$ than today (Bouwens et al. 2009). This gas could be present either in the disc, the halo, or both. Tacconi et al. (2010) found empirical evidence that the higher star formation rates at $z \sim 1-2$ are the consequence of large molecular gas reservoirs (both due to continuous gas inflow from the halos, and minor but not major mergers) and not of a higher star formation efficiency than that at $z \sim 0$. Nevertheless, doubling the amount of gas at $z = 1.25$, to account for a doubling of the stellar mass at $z = 0$, would result in higher extinctions in the optical bands than those derived from the Tully et al. (1998) equations used in this work, because they have been obtained using local galaxy samples. Were this the case, the evolution found in the optical bands would be larger than the evolution observed in this work.

6. Conclusions

We have studied the evolution of the TFR in B , V , R , I , and K_S -bands, using two sample of galaxies selected from the GSS, one in the redshift range $0.1 < z < 0.3$, which is representative of local galaxies, and a high redshift sample in the redshift range $1.1 < z < 1.4$. The rotation velocity of all 241 spiral galaxies were measured from optical lines widths, using DEEP2 spectra. Morphology was determined from HST images, and the absolute magnitudes were derived from a carefully study of the k -correction made with a large set of photometric information. The results of this study can be summarized as follows.

We analysed four sets of templates and concluded that the SED of a spiral galaxy is more accurately reproduced by the nonnegative linear combination of five templates based on the Bruzual & Charlot (2003) stellar evolution synthesis codes obtained by kcorrect. However, when the data include a noisy band, the k -corrections calculated by this code are unsuitable. We also found that the most reliable k -correction is obtained from information in a observed band that roughly match the rest-frame band. In this case, the rest-frame magnitude in the noisy

band is more accurate when calculated directly from the best-fit template.

When the observed photometry does not match the rest-frame bands, it is possible to calculate the rest-frame magnitudes via an interpolation method, as long as enough information at longer and shorter wavelengths is available. However, an extrapolation method is not reliable for galaxies at $z \sim 1.25$ since the results can vary by more than one magnitude mainly because of the poorer fit of the template.

New local TFRs were constructed by increasing the sample of our previous work (Fernández Lorenzo et al. 2009) in the redshift range $0.1 < z < 0.3$. We inferred a similar difference in the zeropoint of the TFR for all bands when we compared with the Verheijen (2001) local relations, which we assumed to equal the ratio of rotation velocity to line width.

We derived the high redshift $1.1 < z < 1.4$ TFRs by fixing the slope to the local one in each band, to study the evolution of the zeropoints. We confirmed the evolution in the B -band TFR found in our previous work (Fernández Lorenzo et al. 2009), in the sense that galaxies were brighter in the past, and no evolution in the K_S -band TFR. Furthermore, for the other bands we detected a gradual evolution, with a change in luminosity becoming more noticeable as the band became bluer.

We studied the slope of the local K_S -band TFR by comparing with Millenium simulated data. We found that the slope varies depending on the luminosity range used in the fit. Considering this effect, we calculated new local K_S -band TFR from the Millenium data in the same luminosity range as our high redshift sample, confirming again that there is no evolution in the K_S -band TFR.

Assuming the change in the stellar mass-to-light ratio M/L_K found by Westera et al. (2007), which is consistent with that of Drory et al. (2004), but no evolution in the K_S -band, we have found that galaxies have probably doubled their stellar mass in the past 8.6 Gyr. This growth is probably mainly due to gas being transformed into stars, as supported by the change in the gaseous O/H abundance found by Rodrigues et al. (2008). From the mass-to-light ratios derived by Westera et al. (2007) for the optical bands (IMF of Salpeter), and assuming this change in stellar mass, we inferred an evolution in the optical luminosity similar to that found in the present work. However, the luminosity and colour evolution simulated by Westera et al. (2007) is not consistent with that found here and would imply a growth of ~ 6 in the stellar mass from $z = 1.25$ to $z = 0$.

Therefore, the galaxies, despite having doubled their stellar mass in the past 8.6 Gyr, are nowadays fainter and redder in the optical bands, which may reflect an ageing of the stellar populations caused by the star formation density decrease observed in the past 8.6 Gyr. This is supported by the change found in the $(V - K_S)$ and $(R - I)$ colours (for a fixed velocity), in the sense that galaxies were bluer in the past.

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part of the activities of the German Astrophysical Virtual Observatory. We thank the SAO/NASA Astrophysics Data System (ADS) that is always so useful.

Appendix A: Deriving the k -corrections

We present a systematic study using various sets of templates and methods to establish the most reliable k -correction for our data.

We performed the following tests:

- (i) We compared the results obtained by fitting different sets of spiral templates to the data, to establish the set of templates that more closely fits our high redshift sample (see Sect. A.1).
- (ii) We estimated the accuracy of the rest-frame photometry reproduced by the templates, since our photometry does not match the rest-frame optical magnitudes, and the K_S -band errors are large. We calculated the best-fit template of $z = 1.25$ galaxies with a known SED (i.e.: templates of local galaxies redshifted to $z = 1.25$) from the photometric information in the same bands and using the same weights that we evaluated for our real sample. We then determined the differences between the rest-frame optical magnitudes calculated from the original rest-frame SED, and those from the best-fit template (see Sect. A.2).
- (iii) We compared the accuracy of the k -corrections derived using this interpolation method with those obtained using an observed band that roughly match the rest-frame magnitude (see Sect. A.3).

A.1. Optimal set of templates

To evaluate the templates that more closely describe each of our galaxies, we used the code `bpz`, which implements the Bayesian photometric redshift method described in Benítez (2000). We restricted the code redshift range to 0.95–1.4, where all the galaxies of our sample are located. This code was applied using four different sets of spiral galaxy templates:

- (i) The templates of SWIRE, formed by 7 spirals ranging from early to late types (S0–Sdm), and generated with the GRASIL code (Silva et al. 1998).
- (ii) The templates of Poggianti (1997), that consist of models of Sa and Sc spirals computed for various redshifts from 0 to 3.
- (iii) The Bruzual & Charlot (2003, hereafter BC03) collection of galaxy templates used by Tremonti et al. (2003) in the analysis of SDSS (Sloan Digital Sky Survey) galaxy spectra. The library includes 39 templates, obtained from 13 models and three metallicities $Z = 0.008, 0.02, \text{ and } 0.05$ (see Table A.1 for the description of the 13 models).
- (iv) The templates produced by the routine `kcorrect` (Blanton & Roweis 2007), which are a nonnegative linear combination of five templates based on the BC03 stellar evolution synthesis code.

For each galaxy, we calculated the redshifted best-fit template to be

$$\lambda = \lambda_0 (1 + z) \quad (\text{A.1})$$

$$f_\lambda = \frac{f_{\lambda_0}}{(1 + z)}, \quad (\text{A.2})$$

where λ_0 is the wavelength and f_{λ_0} is the flux density in $\text{erg s}^{-1} \text{cm}^{-2} \text{Å}^{-1}$ of the template at $z = 0$. We then normalized

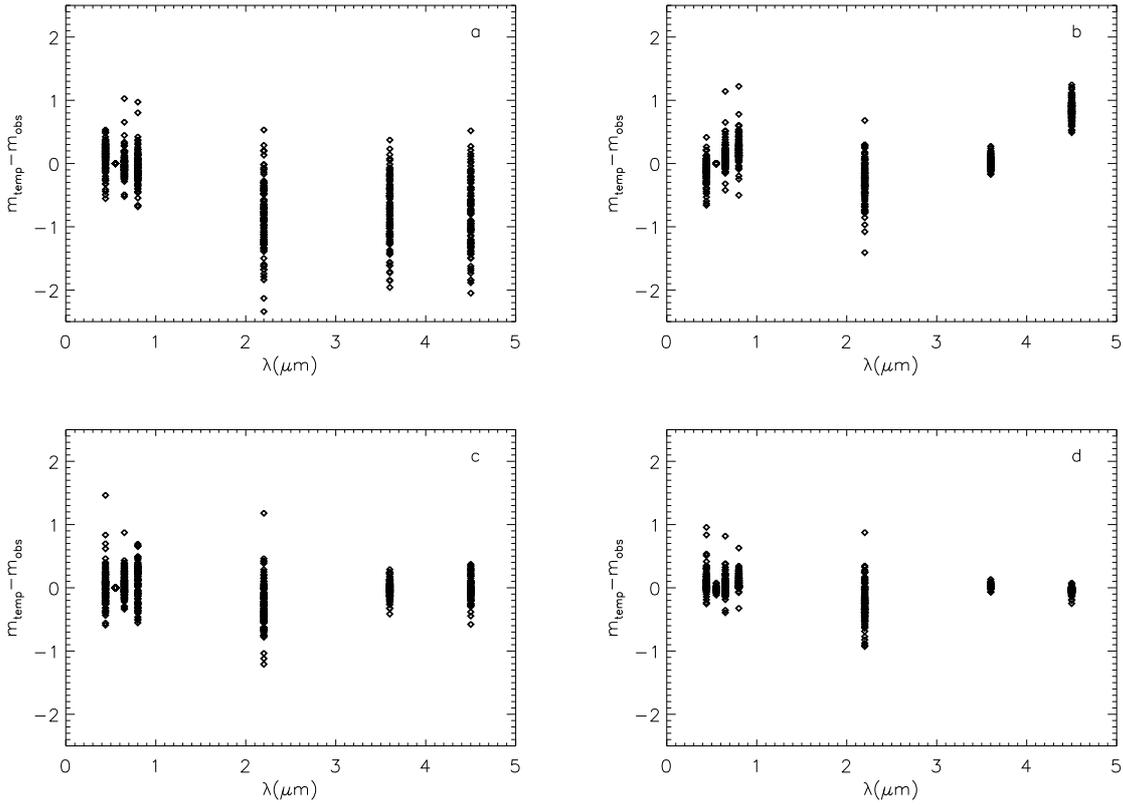


Fig. A.1. Differences between the B , V , R , I , K_S , IRAC1, and IRAC2-band magnitudes derived from the best-fit template normalized to the observed V -band magnitude (m_{temp}), and the corresponding observed magnitudes (m_{obs}), for our high-redshift sample. **a)** Templates of SWIRE. **b)** Templates of Poggianti. **c)** Templates of BC03. **d)** Templates provided by `kcorrect`.

to the V -band to obtain the SED of each galaxy. As this template is the closest representation of the photometry observed, we can infer which set of templates reproduces our data the most accurately by calculating the difference between the observed magnitude m_{obs} and the magnitude m_{temp} obtained from the redshifted template, in each band. To calculate m_{temp} , we projected the template into the filter Q , using the expression

$$m_Q = -2.5 \log \left(\frac{\sum_l (\lambda_{l+1} - \lambda_l) \lambda_l f_{\lambda_l} T_{Q\lambda_l}}{\sum_l (\lambda_{l+1} - \lambda_l) \frac{c}{\lambda_l} f_{\text{AB}}(\nu) T_{Q\lambda_l}} \right), \quad (\text{A.3})$$

where $T_{Q\lambda_l}$ is the filter transmission in the Q -band at the wavelength λ_l , $f_{\text{AB}}(\nu) = 3.631 \times 10^{-20} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ Hz}^{-1}$ is the flux density of the AB standard source, and c is the speed of light.

In Fig. A.1, we present the result for all galaxies in each set of templates. For our objects, with a median redshift of $\langle z \rangle = 1.25$, B , V , R , and I -bands correspond to UV rest-frame wavelengths, while K_S , IRAC1, and IRAC2-bands correspond to NIR rest-frame wavelengths. For all sets of templates, the larger difference found in the K_S -band is caused by the photometric errors in this magnitude (the fits of `bpz` and `kcorrect` are weighted by the photometric errors). However, there is a clear difference between UV and NIR in the comparison with the SWIRE templates (Fig. A.1a). For this set, the galaxies are bluer than the templates (Fig. A.1a). For the templates of Poggianti (1997) (Fig. A.1b), the opposite happens, i.e. the galaxies are redder than the templates. However, with the set of BC03 (Fig. A.1c), the agreement

Table A.1. Description of the templates from BC03 and number of galaxies (N_G) in our sample that were fitted to each one by `bpz`.

Templates	Model description	Age	N_G
1–3	Constant star formation	6 Gyr	36
4–6	Instantaneous-burst	5 Myr	0
7–9	Instantaneous-burst	25 Myr	1
10–12	Instantaneous-burst	100 Myr	21
13–15	Instantaneous-burst	290 Myr	32
16–18	Instantaneous-burst	640 Myr	7
19–21	Instantaneous-burst	900 Myr	5
22–24	Instantaneous-burst	1.4 Gyr	3
25–27	Instantaneous-burst	2.5 Gyr	0
28–30	Instantaneous-burst	5 Gyr	1
31–33	Instantaneous-burst	11 Gyr	0
34–36	Exp. declining SFR $\tau_{\text{SFR}} = 5 \text{ Gyr}$	12 Gyr	8
37–39	Exp. declining SFR $\tau_{\text{SFR}} = 9 \text{ Gyr}$	12 Gyr	28

Notes. Each model was derived for 3 metallicities: $Z = 0.008, 0.02, 0.05$.

between observed-frame bands and the magnitudes obtained from the templates, is better. The combination of the templates of BC03 achieved using `kcorrect` (Fig. A.1d) provides a similar result, but with smaller dispersions. In this case, the template supplied by the code does not need to be normalized to V -band. In addition to the best-fit template, `kcorrect` evaluates the k -correction for each band used in the fit. This k -correction is

the difference between the redshifted and rest-frame band, both calculated from the best-fit template. Our rest-frame magnitudes would then be obtained by subtracting these k -corrections from the observed magnitudes. However, the photometric errors in the observed quantities propagate into the rest-frame magnitudes obtained in this way. This has a strong effect on colour evolution studies if one of the observed bands have substantially larger errors than the others. This is the case for our K_S -band data, so we used another approach to derive the rest-frame K_S magnitudes. Since the observed IRAC2-band roughly matched the rest-frame K_S -band, and the fit of the template is quite good in the wavelength range corresponding to IRAC2 (see Fig. A.1d), we directly used the magnitudes obtained in this band from the rest-frame best-fit template instead of k -correcting the data. For the optical bands, the rest-frame magnitudes calculated from the best-fit template or using the k -corrections of `kcorrect` are equivalent in terms of the TFR, because the average difference between both methods is almost null (see Fig. A.1d). As in K_S -band, the best-fit templates determined by `kcorrect` were then used, but not the k -corrections provided by this software.

A.2. The reliability of the photometry

Since our photometry does not match the rest-frame optical magnitudes, we need to know how reliably `kcorrect` can reproduce the optical SED of a galaxy at $1.1 < z < 1.4$ from the photometric information available. To determine this, we redshifted the templates from BC03 to our median redshift, $z = 1.25$, and we calculated the magnitudes in $B, V, R, I, K_S, IRAC1$, and $IRAC2$ -bands. We then introduced this photometry in `kcorrect` to obtain the best-fit template. We assumed the average magnitude error in our data to be the standard deviation used by the code in every band. In this way, the bands were weighted as the real data. We then calculated the difference in magnitude between the non-redshifted templates from BC03 and the template produced by `kcorrect` for the rest-frame B, V, R, I , and K_S -bands. These differences are represented in Fig. A.2 and provide a first order estimation of the accuracy of the rest-frame magnitudes obtained in the previous section. Because we did not include photometric information for the optical rest-frame wavelength range, the `kcorrect` template provides interpolated magnitudes for these bands. The general trend found for the B -band is that `kcorrect` provides a fainter rest-frame magnitude that the template from BC03 does, while the K_S -band rest-frame magnitudes are brighter. However, the average difference for V, R , and I -bands is ~ 0 when considering the full set of templates. For the constant star formation model (first 3 templates) and the two exponential models (last 6 templates), the error in the SED reconstructed by `kcorrect` is ~ 0.1 , whereas the poorest fit corresponds to a simple stellar population. In general, the fitting of more recent bursts infers optical rest-frame magnitudes that are brighter than the real magnitudes of the templates from BC03, while the optical information about the oldest burst reconstructed by `kcorrect` tends to imply that it is fainter than reality. For the simple stellar populations that reproduce the data for some of our galaxies, the errors in some cases are ~ 0.4 . A recalculation of the k -corrections in our previous work (Fernández Lorenzo et al. 2009), by adding the K_S -band photometry to the B, R , and I -bands used in this work, shows that our previous rest-frame magnitudes were overestimated, especially in the I -band. Moreover, the best-fit template of BC03 provided by `bpz` is different if we use only the information in the B, R and I -bands, or if we introduce also the infrared information.

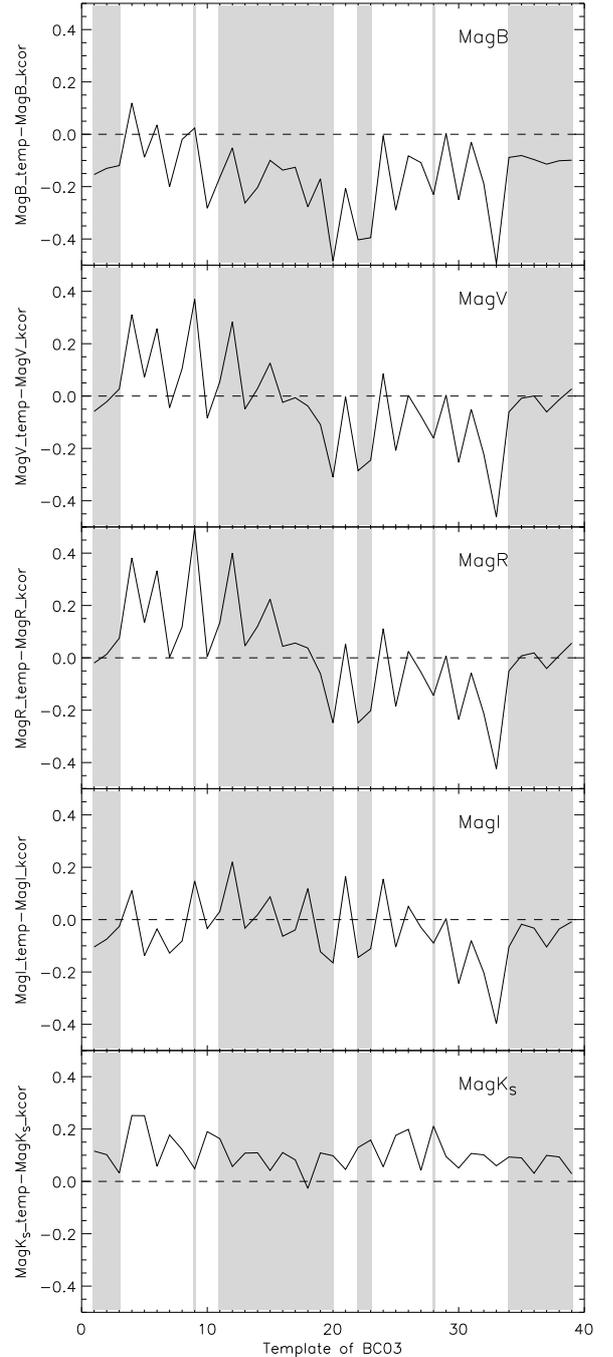


Fig. A.2. Difference between the magnitudes calculated from each rest-frame template of BC03 and from the best-fit template (at $z = 0$) calculated by `kcorrect` from the photometry of the redshifted template at $z = 1.25$. See Table A.1 for the description of the templates from BC03. The grey zones represent the templates that were fitted by `bpz` for some of our galaxies.

We repeated the procedure used in the Fig. A.2, but introducing only the optical bands (B, V, R , and I) in the `kcorrect` fit, so that the rest-frame magnitudes are extrapolated rather than interpolated. The error in the k -corrected magnitudes was found to be between 0.5 and 1.5 mag for most of the templates. The extrapolation method is therefore unsuitable for high redshift galaxies.

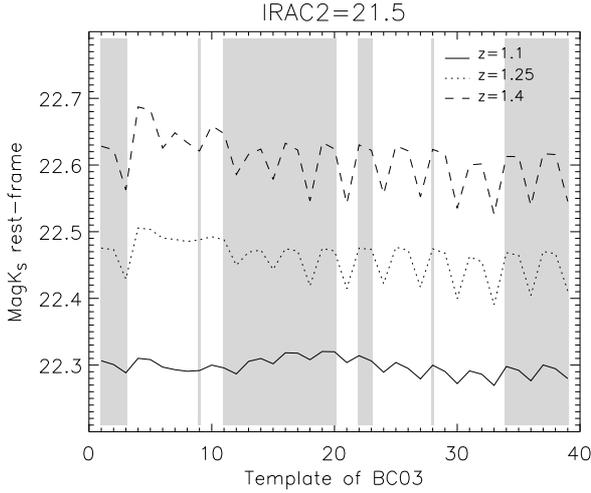


Fig. A.3. Rest-frame K_S -band magnitude calculated from each of the templates from BC03 normalized to an IRAC2 magnitude of 21.5 in the redshifts $z = 1.1$, 1.25, and 1.4. The description of the templates is the same as in Fig. A.2.

A.3. Comparison of methods

We studied the k -corrections obtained using an observed band that roughly matches the rest-frame magnitude. In our redshift range, IRAC2 provides information about the SED region corresponding to rest-frame K_S -band. Hence, if we normalize the redshifted best-fit template to the IRAC2-band, and calculate the K_S magnitude from this template at $z = 0$, we can achieve a higher accuracy in the rest-frame K_S -band magnitudes. The main difficulty in using this procedure is the error in the choice of the template. To quantify this error, we compared the k -corrected K_S -band magnitude obtained from all templates of BC03 redshifted to $z = 1.1$, 1.25, and 1.4, and normalized to the same IRAC2 magnitude. In Fig. A.3, we present our results. The greatest difference between two templates is ~ 0.1 mag at $z = 1.4$. We repeated the same procedure with the SWIRE templates, obtaining a similar result. Using photometry that roughly matches the rest-frame K_S -band, the k -correction is quite independent of the template, mainly due to the smaller difference in the infrared SEDs between early and late-types spirals (see, for example, Mannucci et al. 2001). However, for optical B , V , R , and I -bands, this does not hold. Some of the galaxies have J -band photometry from Palomar data, which corresponds to the rest-frame V -band magnitude at $z = 1.25$. Following the same procedure as used with the K_S -band data (Fig. A.3), but normalizing the templates to $m_J = 22.5$, we found that the maximum difference between the rest-frame V -band of two different templates was ~ 0.3 . Since the errors in the J -band are larger than the IRAC2 errors, this method is not as effective for the V -band as for the K_S -band k -correction.

Finally, we applied both methods to our sample of galaxies instead of redshifted templates. In Fig. A.4, we present the rest-frame magnitudes calculated from the best-fit template of BC03 normalized to the IRAC2-band, versus the rest-frame magnitudes calculated from the best-fit template of `kcorrect`. In addition, we present the result for the V -band, but the template from BC03 being normalized to the observed-frame V -band rather than the J -band. For the K_S -band, the difference between both methods is negligible, but we found a larger scatter between the values of rest-frame V -band. Even so, both values provide us

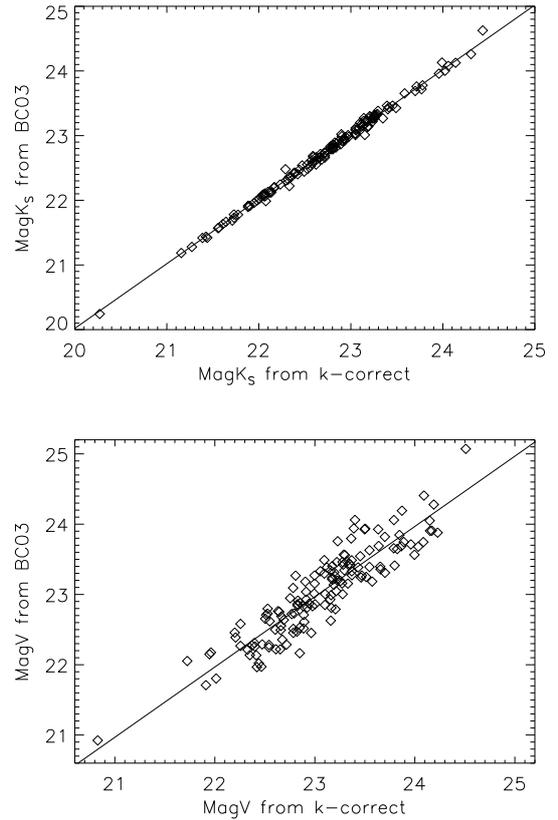


Fig. A.4. Comparison between the k -corrected magnitude obtained from `kcorrect` best-fit template and from the best-fit template of BC03. For K_S -band (up), the normalization of the template from BC03 was done to the observed IRAC2-band. For V -band (down), the template of BC03 was normalized to the observed V -band. The solid line shows the least-square fits with slope 1.

the same mean result because the averaged differences are almost zero.

A.4. Aperture effects

All the photometry used in this work is part of the AEGIS project and has already been corrected for aperture and PSF effects. However, the magnitudes derived at the same aperture for all bands are not available. We therefore decided to compare the results obtained in this work with those obtained using other optical photometry. We used the aperture photometry in the g , r , i , and z -bands, which are provided with the CFHTLS data and are also part of AEGIS. The aperture used by this team is $1.5''$, enough to cover the whole galaxies in the high redshift sample. We then calculated the best-fit template using the aperture photometry in the g , r , i , z , K_S , IRAC1, and IRAC2 bands, with an aperture of $1.5''$ (for the noisy K_S -band we used `mag_AUTO`). In Fig. A.5, we present the rest-frame B and K_S -bands calculated this way, using the total magnitudes in the B , V , R , I , K_S , IRAC1, and IRAC2-bands determined as described in Sect. 3.1 of this paper. The dispersion obtained in the rest-frame B -band is mainly due to the interpolation, as we found previously (Fig. A.2). For the rest-frame K_S -band, we obtain very good agreement between the magnitudes derived from aperture and total photometry, since these k -corrections are basically dependent on the IRAC2-band, which roughly corresponds

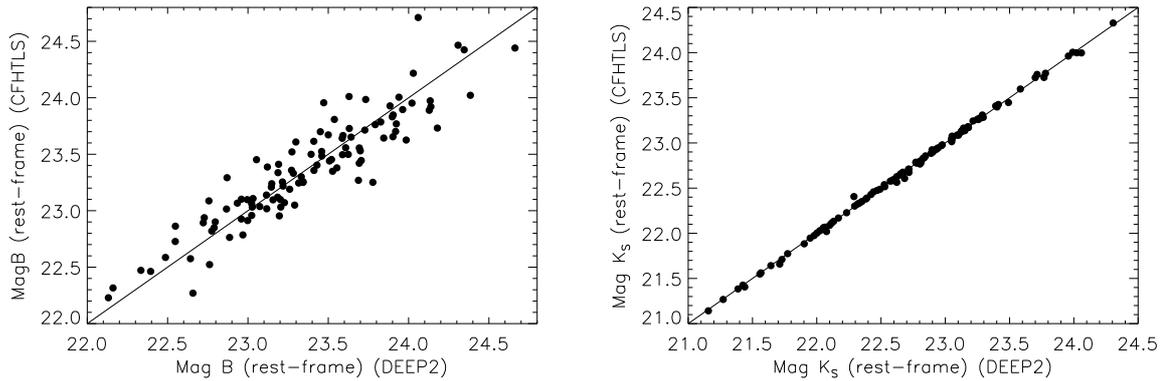


Fig. A.5. Comparison between the rest-frame B -band (*left*) and K_S -band (*right*) magnitudes calculated from the best-fit template obtained using the $g, r, i, z, K_S, IRAC1,$ and $IRAC2$ -bands aperture magnitudes in $1.5''$ aperture, and those obtained using the $B, V, R, I, K_S, IRAC1,$ and $IRAC2$ total magnitudes. The solid lines represent the one to one relations.

to the rest-frame K_S -band. In this case, the k -correction is not too dependent on the template, as we explained in Sect. A.3 (and Fig. A.3).

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