

LETTER TO THE EDITOR

LoCuSS: A *Herschel* view of obscured star formation in Abell 1835[★]

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ABSTRACT

We present *Herschel*/PACS, MMT/Hectospec and XMM-Newton observations of Abell 1835, one of the brightest X-ray clusters on the sky, and the host of a strong cool core. Even though Abell 1835 has a prototypically “relaxed” X-ray morphology and no signs of ongoing merger activity in strong- and weak-lensing mass maps, it has a complex velocity distribution, suggesting that it is still accreting significant amounts of mass in the form of smaller satellite systems. Indeed, we find strong dynamical segregation of star-forming dusty galaxies from the optically selected cluster population. Most *Herschel* sources are found close to the virial radius of the cluster, and almost a third appear to be embedded within a filament feeding the cluster from the SW. We find that the most luminous infrared galaxies are likely involved in galaxy-galaxy interactions that may have triggered the current phase of star formation.

Key words. galaxies: clusters: individual: A1835 – galaxies: evolution – galaxies: star formation – infrared: galaxies

1. Introduction

Galaxy clusters have long been thought of as particularly active sites of galaxy evolution (Dressler 1980). Early optical studies suggested that clusters affect stellar populations in infalling substructure mainly by turning off star formation as galaxies encounter the dense cluster environment (Lewis et al. 2002). However, when these same objects were observed in the infrared by IRAS, ISO and *Spitzer*, a significant fraction of cluster galaxies was revealed to contain heavily obscured sites of star formation (see e.g. Metcalfe et al. 2005, and references therein). When this obscured activity is accounted for, clusters emerge as environments where star formation is both quenched and induced, perhaps even at different stages of the same physical process. Some studies have claimed to observe triggered star formation in strongly disrupted, merging clusters (e.g. Geach et al. 2006) but the intrinsic scatter in this correlation is large, with some seemingly relaxed clusters presenting much higher star formation activity than their actively merging counterparts (Haines et al. 2009).

Many studies in recent years have attempted to separate cluster samples into two distinct populations: merging clusters vs.

clusters that are dynamically relaxed. The global properties of the latter are expected to be tightly correlated with mass, and therefore directly relatable to statistical predictions from cosmological models.

While there is evidence that some X-ray cluster properties may in fact be bimodal, e.g. with respect to the presence or absence of a cool core (Sanderson et al. 2009), cosmological simulations have shown that even the most relaxed cool core clusters can still be accreting significant amounts of mass (Poole et al. 2008), at a mean rate which is approximately linearly proportional to the mass of the system (McBride et al. 2009). Generally, only the most massive of these accretion events (mass ratios <10:1) are expected to disturb the dense cluster core. Massive cool core clusters can thus quietly accrete a significant fraction of their total mass without large disturbances to their central ICM, which makes them arguably “cleaner” targets for studies of cluster galaxy evolution.

The Local Cluster Substructure Survey (LoCuSS¹) *Herschel* key programme (Smith et al. 2010) was designed to probe dusty star-forming galaxies in the infall regions of a sample of 30 galaxy clusters ($z \sim 0.2$) spanning a wide range of mass and merging histories.

Abell 1835 is the most X-ray luminous cluster in this sample ($L_{X,\text{bol}} = 5.32 \pm 0.15 \times 10^{45}$ erg s⁻¹, Zhang et al. 2008) and the

[★] *Herschel* is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA.

¹ <http://www.sr.bham.ac.uk/locuss>

host of a strong cool core (Peterson et al. 2001). Its strong- and weak-lensing derived mass distributions (Richard et al. 2010; Okabe et al. 2010), and X-ray morphology (Smith et al. 2005) are typical of undisturbed (i.e. non-merging) cool core clusters. For example, the low substructure mass fraction in its core ($f_{\text{sub}} = 0.13 \pm 0.01$) suggests that it has grown in mass by $<10\%$ in the previous few Gyr (Smith & Taylor 2008).

In this Letter, we focus on a sample of 100 μm -selected spectroscopically confirmed cluster galaxies in Abell 1835. We use the Photodetector Array Camera and Spectrometer (PACS) (Poglitsch et al. 2010) on the *Herschel* satellite (Pilbratt et al. 2010) to measure far infrared fluxes and from these we estimate total infrared luminosities for these objects. We then use these data to examine the spatial and velocity distribution of obscured star formation sites in Abell 1835. We assume $H_0 = 70 \text{ km s}^{-1}$, $\Omega_M = 0.3$, $\Omega_\Lambda = 0.7$. All virial masses and radii are stated with respect to the virial overdensity, $\Delta_{\text{virial}}(z = 0.25) = 121$ (Bryan & Norman 1998).

2. Observational data

PACS data were obtained in scan map mode on December 24th 2009, at both 100 and 160 μm . The details of the observing mode, data reduction procedure and source extraction and photometry can be found in (Smith et al. 2010). The 90% completeness flux limits are 21 mJy at 100 μm (86 sources) and 28 mJys at 160 μm (91 sources).

We conducted a wide-field spectroscopic survey of this field with Hectospec at MMT in February 2009, confirming cluster membership of 439 galaxies (Hardegree-Ullman et al., in prep.). In addition, Czoske (2004) obtained redshifts for 630 cluster members with VIMOS/VLT, bringing the total number of unique redshifts for this cluster to 851. This extensive spectroscopic coverage means that, out of 91 *Herschel* detections with $f_{160 \mu\text{m}} > 28 \text{ mJy}$, 64 have redshifts, of which 33 lie within $\delta v < 6000 \text{ km s}^{-1}$ of the cluster redshift. The total number of spectroscopically confirmed cluster galaxies with $>3\sigma$ detections in both 100 μm and 160 μm bands is 45, and we plot their location in Fig. 1.

X-ray maps can be used both to separate the AGNs from the starburst galaxies and to identify galaxy groups around clusters, and we have analysed the *XMM* observation ID 0147330201 to characterize the X-ray appearance of A1835. We decompose the 0.5–2 keV combined *XMM* image into unresolved and extended emission, following the wavelet technique of Finoguenov et al. (2009), before extracting point source fluxes at the locations of *Herschel* sources. Additionally, we make use of previously published wide-field data from *Spitzer*/MIPS 24 μm and UKIRT/WFCAM *J, K* bands (Haines et al. 2009), and Subaru/Suprime-CAM (Okabe et al. 2010).

3. Results

3.1. Infrared luminosities, temperatures and SFRs

At the redshift of the cluster, $z = 0.25$, the PACS 100 and 160 μm bands are expected to straddle the peak of the spectral energy distribution for sources with dust temperatures of $T \sim 30 \text{ K}$. The 160 $\mu\text{m}/100 \mu\text{m}$ color should thus be sensitive to dust temperature differences, with redder objects having cooler dust components. We plot observed 160 $\mu\text{m}/100 \mu\text{m}$ vs. 100 $\mu\text{m}/24 \mu\text{m}$ colors for the sources in our cluster in Fig. 2, and overplot the synthetic colors from commonly used templates from the literature: Rieke et al. 2009 (R09), Dale & Helou 2002 (DH02),

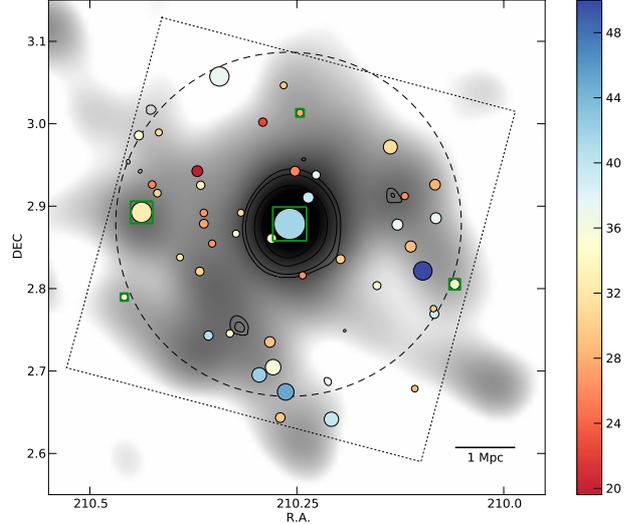


Fig. 1. Location of *Herschel* detected, spectroscopically confirmed, cluster members on the sky, overlaid on a *K*-band luminosity density image of the cluster produced from all spectroscopically confirmed members and corrected for the spatial incompleteness of our spectroscopic survey. The black solid contours correspond to 4σ wavelet detections of X-ray extended sources in the *XMM* data. The dashed circle marks $r_{\text{vir}} = 2.93 \text{ Mpc}$, the dotted square the limits of the PACS field. The sizes of the points scale linearly with L_{TIR} , and their color scales with the temperature of the best fit modified blackbody SED, ranging from 20 K for the coldest galaxies to $\sim 60 \text{ K}$ for the hottest. Galaxies identified as AGN through emission lines or X-ray emission are marked with a green square.

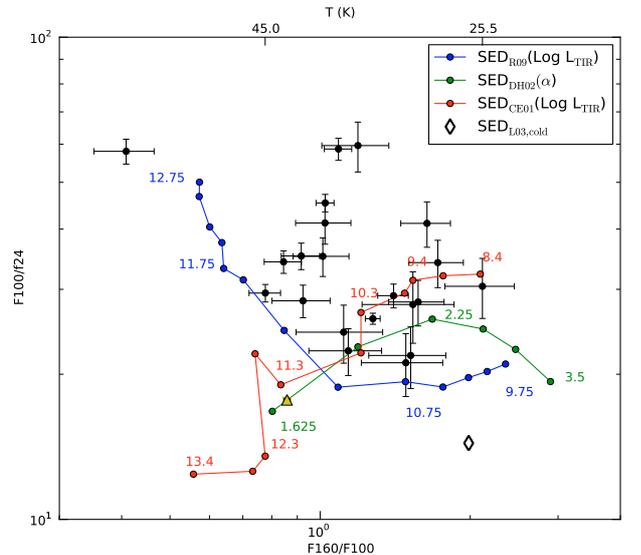


Fig. 2. 160 $\mu\text{m}/100 \mu\text{m}$ vs. 100 $\mu\text{m}/24 \mu\text{m}$ colors, compared with synthetic colors derived from SED templates in the literature. The 100 $\mu\text{m}/24 \mu\text{m}$ colors of the cluster galaxies appear significantly reddened compared to the templates. The BCG (yellow triangle) has the lowest 100 $\mu\text{m}/24 \mu\text{m}$ color of the sample, and is well fit by a Dale & Helou starburst SED with $\alpha \approx 1.65$. Error bars are based on statistical noise, and do not include the overall PACS calibration error. The stated PACS calibration uncertainty of $<10\%$ at 100 μm is nevertheless not large enough to reconcile the data with the template colors.

Lagache et al. 2003 (L03), Chary & Elbaz 2001 (CE01). All the SED templates are poor representations of these galaxies, which appear too red in their 100 $\mu\text{m}/24 \mu\text{m}$ colors. The CE01 templates appear to fit the colors of $\sim 1/2$ of the sample, but with

templates corresponding to total infrared luminosities that are one to two orders of magnitude fainter than the actual measured luminosities – see Smith et al. (2010) and Rawle et al. (2010) for a more extended discussion of this 100 μm excess.

We estimate total infrared luminosities, L_{TIR} , for these galaxies in two ways: first, we fit the CE01 templates (redshifted to the cluster rest-frame) to the 24, 100 and 160 μm emission and integrate the best-fit spectrum, normalized to the observed fluxes, from rest-frame 3–1100 μm . In order to obtain estimates for the dust temperatures in these systems, we also fit a modified blackbody spectrum to the 100 μm and 160 μm emission: $f_\nu \propto [1 - \exp(-\tau_\nu)]B_\nu$, where B_ν is the Planck function, and τ_ν is the frequency-dependent optical depth, ν/ν_0^β . We assume $\beta = 1.5$ as the emissivity power law index, and $\nu_0 = 3$ THz as the frequency at which the emission becomes optically thick, but note that the integrated luminosity depends very weakly on the values assumed for $\beta \in [0.7-2]$ and ν_0 , while the best fit temperature varies at most by 20%. The L_{TIR} derived from the blackbody fits are systematically lower than the ones obtained from the best-fit templates by a factor of $\sim 0.6-0.8$, since they underestimate the mid-IR contribution to the total IR flux. The cluster members range in L_{TIR} from $4-79 \times 10^{10} L_\odot$, and all but one have dust temperatures in the range 20–45 K.

The BCG is much brighter than the rest of the galaxies in the cluster, with fluxes of 255.6 ± 3.9 mJy (160 μm) and 297.8 ± 3.1 mJy (100 μm) and an estimated L_{TIR} of 7.9×10^{11} , consistent with previous *Spitzer*/*MIPS* observations of Abell 1835 (Egami et al. 2006). We derive star formation rates from the combined $\text{H}\alpha$ (from our spectroscopic measurements) and L_{TIR} emission using the following relation from Kennicutt et al. (2009): $\text{SFR}[\text{Kroupa}](M_\odot \text{yr}^{-1}) = 5.5 \times 10^{-42} [L(\text{H}\alpha)_{\text{obs}} + 0.002 L_{\text{TIR}}] (\text{erg s}^{-1})$. This yields SFRs in the range: 2–20 $M_\odot \text{yr}^{-1}$ for the general cluster population. The BCG’s high luminosity and special location at the center of a massive cool core separate it from the rest of the cluster population and we therefore exclude it from the subsequent analysis.

We find that, of all galaxies with far-infrared emission detected by *Herschel*, only 2 have X-ray counterparts down to $L_X \approx 4 \times 10^{41} \text{ erg s}^{-1}$, which is below the classical $10^{42} \text{ erg s}^{-1}$ threshold for AGN definition. Only four galaxies (including the couple of X-ray point sources) in the sample are spectroscopically classified AGN according to their position in the classic BPT diagram (Baldwin et al. 1981), defined by the emission line flux ratios of $[\text{NII}] 6587/\text{H}\alpha$ versus $[\text{OIII}] 5007/\text{H}\beta$. We thus conclude that the bulk of the infrared luminous population is being heated by star formation and not an active nucleus, and we exclude the 4 likely AGN from our analysis.

3.2. Probing the environments of obscured star formation

It is apparent from Fig. 1 that, apart from the BCG, there are few luminous infrared galaxies in the central regions of the cluster. In fact, the radial surface density profiles in Fig. 3 show that most of the infrared luminosity is concentrated at large ($r \sim 2.5$ Mpc) radial distances, which is not the case for the *K*-band luminosity density which, as an approximate tracer of stellar mass, declines smoothly with radius.

Upon closer inspection, however, it is clear that this overdensity at the virial radius does not occur isotropically around the cluster, but rather peaks along specific directions. The velocity distribution of *K*-band selected cluster galaxies (black line in Fig. 4) is distinctly non-Gaussian, and shows signs of significant substructure. The most luminous infrared sources, located toward the south of the cluster, are dynamically segregated from

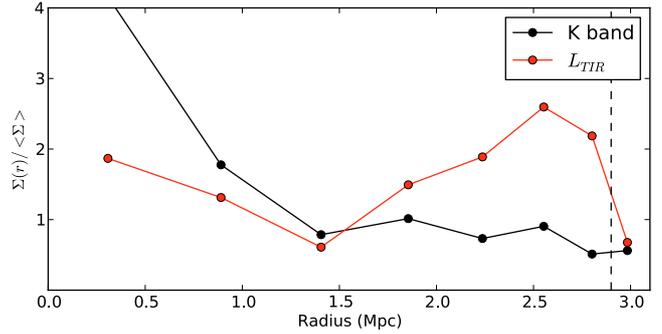


Fig. 3. Radial profile of *K*-band (black) and total infrared (red) luminosity surface density relative to the average surface density in concentric annuli within $r < 3$ Mpc (excluding the BCG). The dashed vertical line marks r_{vir} . The *K*-band luminosity density is an approximate tracer of stellar mass, and declines continuously with radius, whereas the infrared luminosity peaks close to the virial radius of the cluster.

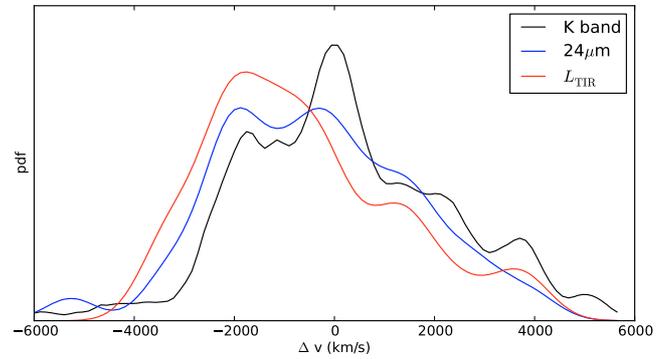


Fig. 4. Redshift probability density function for galaxy populations observed at different wavelengths. The density function is weighted by flux/luminosity, and corrected for the incompleteness in our spectroscopic survey. *K*-band: black, *MIPS* 24 μm : blue, L_{TIR} : red.

the rest of the cluster population, clustering around an elongated structure in the foreground ($\delta v \sim 2000 \text{ km s}^{-1}$) of the cluster, which is also traced out in the *K*-band luminosity density maps (see Fig. 1). A gradient is observed in the velocities projected along the axis of elongation, where the relative velocity with respect to the cluster center progressively decreases as the inner region of the cluster is approached.

Figure 5 shows Subaru *i'*-band cutouts of the most luminous infrared galaxies, in decreasing order of L_{TIR} . Many galaxies appear to have undergone recent interactions – \sim half of them have close companions, and half present distorted morphologies indicative of recent merging. We can attempt to quantify the importance of mergers in our small sample by comparing the fraction, f_c , of *Herschel* sources with close ($\Delta r < 50$ kpc and $\Delta v < 300 \text{ km s}^{-1}$) companions to that of a population of non-IR luminous cluster galaxies that has been matched, both in *K*-band luminosity and radial distance from the cluster center, to the IR luminous sample (excluding in both samples all identified AGNs). We find that, while the total sample of IR luminous galaxies is indistinguishable from the matched sample of IR faint galaxies ($f_c \sim 12_{-5}^{+8}\%$ vs. $10_{-2}^{+2}\%$, respectively), the IR bright ($\log_{10} L_{\text{TIR}} > 10.8$) galaxies in Abell 1835 are marginally more likely to have a close neighbour ($f_c \sim 30_{-16}^{+21}\%$ vs. $9_{-2}^{+3}\%$ in the respective matched sample).

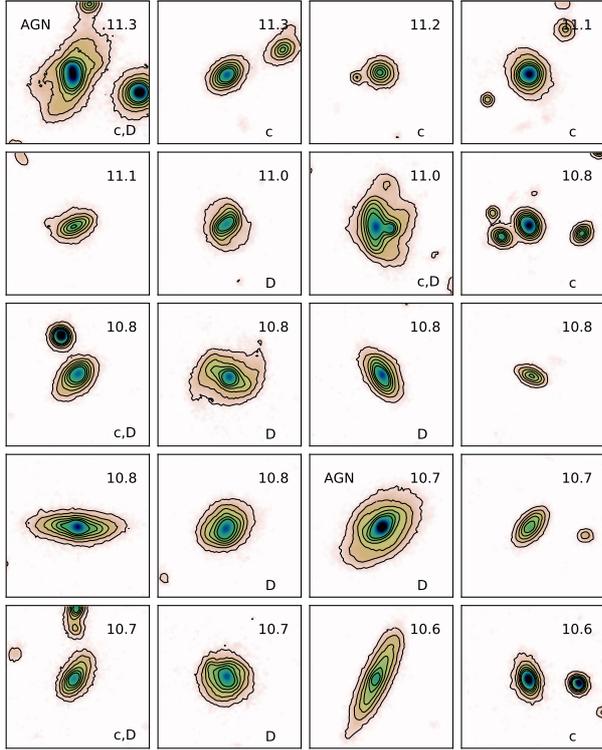


Fig. 5. Subaru i' band $\sim 60 \times 60$ kpc image cutouts centered on the IR brightest cluster galaxies, in decreasing order of $\log_{10} L_{\text{TIR}}$ (upper right corner). A significant fraction has close companions (c), and half show low surface brightness tidal features, or asymmetric morphologies indicative of a recent interaction (D).

4. Summary and discussion

Star formation can be triggered by several different physical processes prevalent in clusters and their infall regions; disentangling these effects is one of the main aims of the LoCuSS *Herschel* key programme. The most promising candidates, ICM ram pressure, galaxy mergers, and galaxy harassment can all perturb the ISM to produce a new generation of stars, but the relative importance of these processes is a strong function of clustercentric radius (cf. Fig. 1 in Smith et al. 2010). The location of the peak of IR luminosity density just inside the virial radius of Abell 1835 suggests that merging may be driving the bulk of the star formation in these galaxies. This interpretation is corroborated by the large fraction of disturbed morphologies in the infrared luminous sample, and also by our preliminary analysis of their local environment: when compared to a matched sample of galaxies with no detectable infrared emission, infrared bright galaxies are more likely to have close neighbours, although the larger sample of 30 LoCuSS clusters should be used to confirm this correlation.

The L_{TIR} density profile for Abell 1835 is not spherically symmetric: the brightest galaxies are clustered together in an elongated structure towards the south of the cluster that is

dynamically segregated from the rest of the cluster. The structure's K -band morphology and velocity gradient toward the cluster core suggests that these galaxies may be embedded within a filament that is still currently feeding the cluster with new, actively star forming galaxies.

XMM imaging of this field reveals the presence of two groups in the outskirts of Abell 1835 (see contours in Fig. 1), one of which appears to be embedded in the filament that is also feeding the bright LIRGs into the cluster. The group is dominated by a large ($L_K \approx 2 \times 10^{12} L_{\odot}$, $z = 0.245$) elliptical galaxy with no detectable *Herschel* emission.

The two groups have X-ray luminosities of 4.6 ± 0.6 and $7.5 \pm 0.8 \times 10^{42} \text{ erg s}^{-1}$ (0.1–2.4 keV), from which we estimate masses, using the weak lensing calibration of the $L_X - M$ relation of Leauthaud et al. (2010), of $M_{200} = 4.3$ and $5.8 \times 10^{13} M_{\odot}$ respectively. Together these amount to $\geq 1/10$ th of the mass of the central cluster, reinforcing the view that Abell 1835 continues to grow by feeding on smaller systems from the surrounding large scale structure.

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