

LETTER TO THE EDITOR

First detection of the methylidyne cation (CH^+) fundamental rotational line with the *Herschel*/SPIRE FTS[★]

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ABSTRACT

Aims. To follow the species chemistry arising in diverse sources of the Galaxy with *Herschel*.

Methods. SPIRE FTS sparse sampled maps of the Orion bar & compact HII regions G29.96-0.02 and G32.80+0.19 have been analyzed.

Results. Beyond the wealth of atomic and molecular lines detected in the high-resolution spectra obtained with the FTS of SPIRE in the Orion Bar, one emission line is found to lie at the position of the fundamental rotational transition of CH^+ as measured precisely in the laboratory by Pearson and Drouion. This coincidence suggests that it is the first detection of the fundamental rotational transition of CH^+ . This claim is strengthened by the observation of the lambda doublet transitions arising from its relative, CH, which are also observed in the same spectrum. The broad spectral coverage of the SPIRE FTS allows for the simultaneous measurement of these closely related chemically species, under the same observing conditions. The importance of these lines are discussed and a comparison with results obtained from models of the photon dominated region (PDR) of Orion are presented. The CH^+ line also appears in absorption in the spectra of the two galactic compact HII regions G29.96-0.02 and G32.80+0.19, which is likely due to the presence of CH^+ in the cold neutral medium of the galactic plane. These detections will shed light on the formation processes and on the existence of CH^+ , which are still outstanding questions in astrophysics.

Key words. ISM: molecules – evolution – ISM: general – space vehicles: instruments – techniques: spectroscopic

1. Introduction

Molecules and radicals with C and H as constituents are expected to be the first chemical building blocks. The methylidyne cation (CH^+) was discovered at visible wavelengths 70 years ago by Douglas & Herzberg (1941), shortly after the discovery of the methylidyne (CH) radical by Swings & Rosenfeld (1937). Since that time there has been intense debate surrounding the chemistry of this species (Bates & Spitzer 1951). CH^+ far-IR detections were reported by Cernicharo et al. (1997) for the $J = 2-1$ to $4-3$ transitions in the NGC7027 PDR using ISO spectra. Falgarone et al. (2005) report the probable detection of $^{13}\text{CH}^+(1-0)$ in absorption against a galactic source. CH $N = 2-1$ transitions have been detected by Stacey et al. (1987) and Polehampton et al. (2005) toward SgrB2 and by Liu et al. (1997) in NGC7027. CH^+ is commonly detected in the visible and found to correlate with rotationally excited H_2 (Lambert & Danks 1986), which argues for its formation in energetic regions. One of the underlying questions is whether CH^+ formation proceeds via a still undefined reaction or, following the consensus in the literature, via ISM conditions liberating sufficient energy (shock chemistry,

turbulence, UV pumping) to overcome activation barriers, thereby strongly enhancing the otherwise hindered endothermic reaction rate for $\text{C}^+ + \text{H}_2 \rightarrow \text{CH}^+ + \text{H}$, with a ~ 0.4 eV barrier. CH^+ formation has been evoked as a result of shock chemistry (Pineau des Forets et al. 1986; Draine & Katz 1986, and citations). Species produced in shocks, however, should display velocity differences from the source producing the shocks. The CH^+ absorption profiles measured to date are found to lie at the same velocity as CH (Gredel 1997), which has been taken as an argument against shock chemistry production. UV pumping, leading to vibrationally excited H_2 (Stecher & Williams 1972), is not expected to be significant in the diffuse ISM, on the grounds that the mean UV field is not strong enough to generate the observed excited H_2 . However, the laboratory measured $\text{C}^+ + \text{H}_2(v = 1) \rightarrow \text{CH}^+ + \text{H}$ reaction (Hierl et al. 1997), may support such a route to CH^+ formation in PDRs (e.g. Agundez et al. 2010). Among the propositions, turbulence models are gaining support in CH^+ chemistry (Godard et al. 2009, and references therein). In this letter we present the first detection of the fundamental rotational transition of CH^+ , along several lines-of-sight, in different and distinct regimes of the ISM and discuss the potential of SPIRE FTS observations to constrain the CH chemistry in such sources.

[★] *Herschel* is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA.

Table 1. Observations summary.

Source	Obs. date	Ra	Dec	Int(s)
HD37041	09/09/13	05 ^h 35 ^m 22.83 ^s	-0 ^d 24'57.67"	266.4
G29.96-0.02	09/09/13	18 ^h 46 ^m 04.07 ^s	-02 ^d 39'21.88"	266.4
G32.80+0.19	09/09/21	18 ^h 50 ^m 30.87 ^s	00 ^d 02'00.85"	266.4

2. Observations with the SPIRE FTS

The SPIRE FTS measures, simultaneously, the source spectrum across two wavebands: SLW, covering 14.9–33.0 cm⁻¹ (303–671 μm) and SSW covering 32.0–51.5 cm⁻¹ (194–313 μm). Each band is imaged with a hexagonal bolometer array with pixel spacing of approximately twice the beam-width. The *FWHM* beam-widths of the SLW and SSW arrays vary between 29–42" and 17–21" respectively. The source spectrum, including the continuum, is obtained by taking the inverse transform of the observed interferogram. For more details on the ESA *Herschel* Space Observatory, the SPIRE FTS, and the FTS calibration and data reduction procedures, the reader is referred to the articles by Pilbratt et al. (2010), Griffin et al. (2010), and Swinyard et al. (2010), respectively, in this volume. Our observations are part of the "Evolution of Interstellar Dust" key program of the SPIRE consortium (Abergel et al. 2010). The source data presented here include the Orion Bar as well as the G29.96–0.02 and G32.80+0.19 compact HII regions. A short observation log is given in Table 1. For example, G29.96-0.02 was observed with the high-resolution mode of the SPIRE FTS on the 13th of September, 2009 at 19:49 (*Herschel* observation ID, 1342183824). Two scan repetitions were observed giving an on-source integration time of 266.4 seconds. The unapodized spectral resolution was 0.04 cm⁻¹ (1.2 GHz). After apodization (using extended Norton-Beer function 1.5; Naylor & Tahic 2007) the *FWHM* of the resulting instrument line shape is 0.0724 cm⁻¹ (2.17 GHz). While unapodized FTS spectra provide the highest spectral resolution, the instrument line shape, which in the case of an ideal FTS is the classical sinc function, is characterized by relatively large secondary oscillations having negative lobes. An iterative spectral line fitting routine has been developed to extract line parameters from unapodized FTS spectra. This algorithm fits a continuum (either a low order polynomial or a blackbody variant) and a series of lines using the Levenberg-Marquardt least squares method. The fitting procedure weights the spectral intensity at a given frequency of an averaged spectrum by the statistical uncertainty at that frequency. The fitting routine returns the line centres, intensities, and widths together with their associated errors. The transitions of methylidyne and its cation accessible to the FTS are summarized in Fig. 1 and Table 2.

The Orion Bar

The recorded spectra for ¹²CH⁺ are overlaid on published maps by Lis et al. (1998) in Fig. 2. The position of the O6 θ¹ Ori C exciting star is also shown on the map. The expected geometry along the line of sight can be found in Fig. 3 from Pellegrini et al. (2009) and in references therein. A cut through the expected best exposed projected line of sight is displayed (Fig. 2, left panel) and the integrated fluxes along this cut are shown in the right panel. Estimates of the column density corresponding to the positions shown in Fig. 2, which have been derived assuming thermal equilibrium with a mean Orion Bar temperature (~85 K), determined from ground based observations (Hogerheijde et al. 1995), are given in Table 3. This table also presents abundance values for temperatures of 50 and 200 K, which probe different

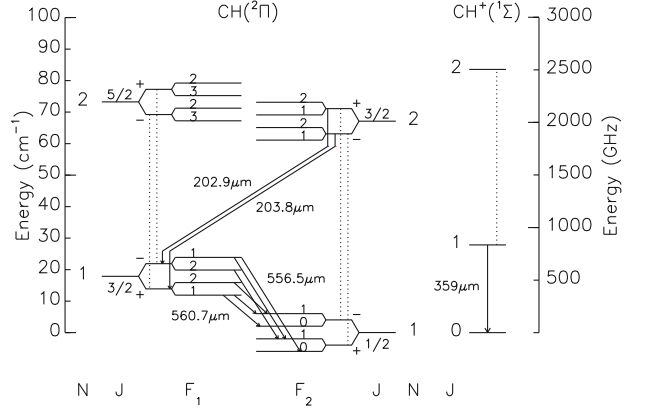

Fig. 1. Energy diagram for CH and CH⁺ in the SPIRE FTS range. Expected transitions are shown with arrows. Dotted lines are transitions lying at higher frequencies accessible by other *Herschel* instruments.

Table 2. CH⁺ and CH transitions in the SPIRE/FTS range.

Freq (GHz)	Transition	A(s ⁻¹) ^a	Ref.
CH ⁺	<i>J</i> = 1 ← 0		
835.079043(1.0)	1 – 0	5.96 × 10 ⁻³	c
CH	<i>N</i> = 1 <i>J</i> = 3/2 ← 1/2		
532.721333 (100)	F1 = 1+ ← F2 = 1-	2.1 × 10 ⁻⁴	b
532.723926 (40)	F1 = 2+ ← F2 = 1-	6.2 × 10 ⁻⁴	b
532.793309 (50)	F1 = 1+ ← F2 = 0-	4.1 × 10 ⁻⁴	b
536.761145 (50)	F1 = 2- ← F2 = 1+	6.4 × 10 ⁻⁴	b
536.781945 (50)	F1 = 1- ← F2 = 1+	2.1 × 10 ⁻⁴	b
536.795678 (50)	F1 = 1- ← F2 = 0+	4.2 × 10 ⁻⁴	b
CH	<i>N</i> = 2 ← 1 <i>J</i> = 3/2 ← 3/2		
1470.6861(30)	F2 = 1- ← F1 = 2+	8.8 × 10 ⁻⁴	d
1470.6890(30)	F2 = 1- ← F1 = 1+	4.4 × 10 ⁻³	d
1470.7356(30)	F2 = 2- ← F1 = 2+	4.8 × 10 ⁻³	d
1470.7385(30)	F2 = 2- ← F1 = 1+	5.3 × 10 ⁻⁴	d
1477.2901(30)	F2 = 1+ ← F1 = 1-	4.4 × 10 ⁻³	d
1477.3104(30)	F2 = 1+ ← F1 = 2-	8.8 × 10 ⁻⁴	d
1477.3630(30)	F2 = 2+ ← F1 = 1-	5.3 × 10 ⁻⁴	d
1477.3832(30)	F2 = 2+ ← F1 = 2-	4.7 × 10 ⁻³	d

Notes. ^(a) Calculated using CDMS data parameters; ^(b) Amano (2000); ^(c) Pearson & Drouin (2006); ^(d) Brown & Evenson (1983).

optical depths in the PDR. As noted above, CH⁺ is highly reactive, easily destroyed by collisions and is not expected to be found in thermal equilibrium in the presence of strong UV pumping. Detailed analysis and interpretation of the measured fluxes is ongoing, but is beyond the scope of this letter.

One of the striking features when comparing these plots is the prominence of CH⁺ with respect to the ¹³CO(8-7) transition fluxes at the A and B positions, away from the Bar, probably corresponding to the surface of the cavity perpendicular to the line of sight. This tendency is reversed on the Bar (C position). From a preliminary analysis of ¹³CO SLW data along the A and B line of sight, an excitation temperature of 85 K₋₄₅⁺¹⁵ is derived from the rotational diagram, corresponding to *N*_B(¹³CO) = 4.5–16.9 × 10¹⁵ cm⁻². We chose to focus on the B line of sight to avoid the effects of vignetting which are known to occur on the outer ring of bolometers. Such column densities would correspond to *A_V* in the range 1–6 at most, and fractional abundances in the range 2–7 × 10⁻¹⁰ for CH⁺. We compare these estimates with a Meudon PDR code (Le Petit et al. 2006) model applied to the Orion Bar environment, described in more detail in Habart et al. (2010). CH⁺ abundances in A and B are only marginally compatible with a classical Orion PDR model (Fig. 3, red curves), whereas taking into account the

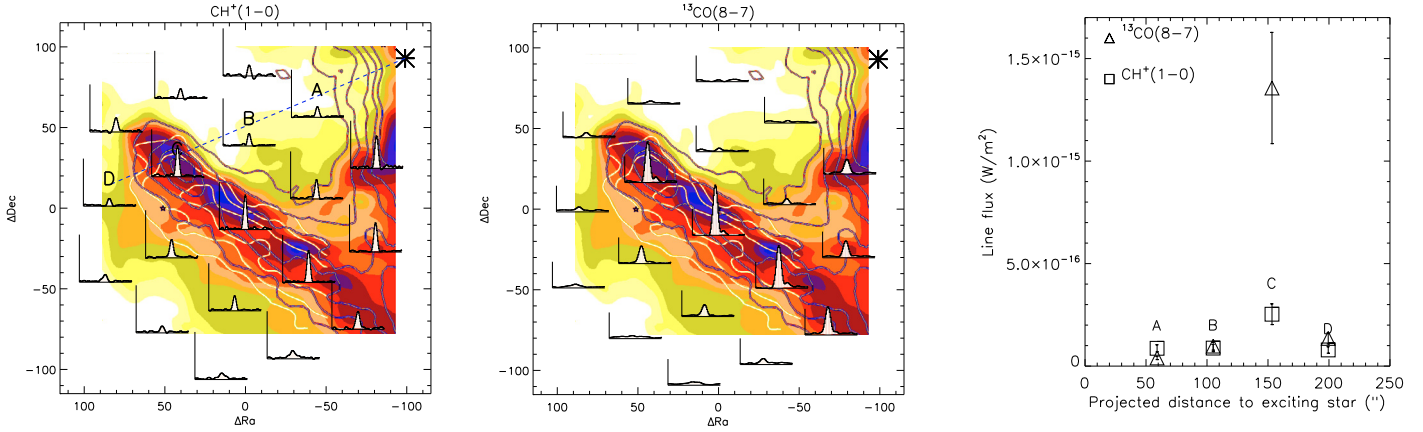


Fig. 2. *Left panel:* $^{12}\text{CH}^+(1-0)$ continuum subtracted FTS spectra in the Orion bar region, overlaid on previous observations by Lis et al. (1998), which combine dust and gas emission. The color map is the $^{12}\text{CO}(6-5)$ emission, the contours extending over the whole map are the continuum emission at $350\ \mu\text{m}$ and the contours displayed only for ΔRa above -50 correspond to $^{13}\text{CO}(6-5)$ emission. The Orion PDR exciting star is also shown and a cut passing through the bar discussed in the text. The stratification of the PDR appears clearly on Lis et al. (1998) data, where the ^{12}CO color map tracing hot gas is displaced toward the star with respect to ^{13}CO contours tracing the dense molecular regions. *Middle panel:* same as CH^+ map for $^{13}\text{CO}(8-7)$. *Right panel:* line fluxes along the cut presented in left panel for ^{13}CO and $\text{CH}^+(1-0)$

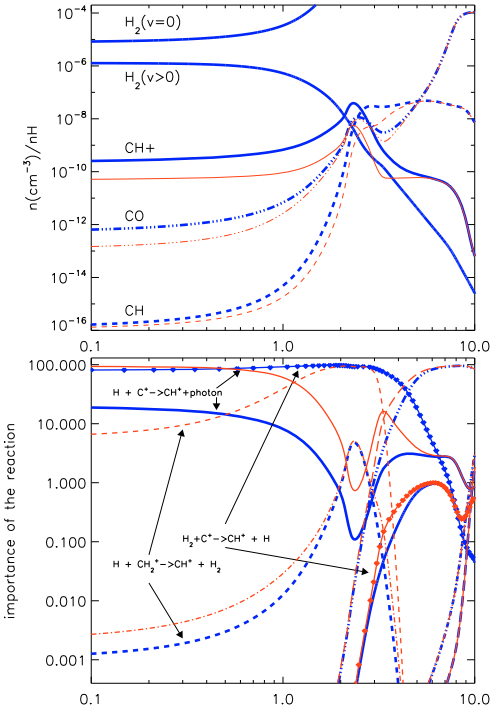


Fig. 3. Orion Bar isochore PDR model, using the Meudon code, as a function of the visual extinction normal to the front. The adopted PDR parameters are an incident flux of $\chi = 10^4$ and a density of $n_{\text{H}} = 10^5\ \text{cm}^{-3}$. The upper panel shows model results in which the state by state vibrational excitation of H_2 in $v \geq 1$ (thick blue line) is considered as an internal energy involved in overcoming potential energy barriers, notably for the $\text{H}_2 + \text{C}^+$ reaction. The thin red line shows the results for the classical PDR model. The lower panel is a fractional representation of the relative importance of the methylidyne cation production reactions as a function of A_V .

vibrationally excited H_2 states (blue curves) where about 6% of H_2 is in an excited vibrational state, enhances the formation rate by a factor of about 30 at low A_V . The relative importance of the CH^+ production reactions, for both models, is shown in the lower panel of Fig. 3. At low to moderate A_V , corresponding to the front and surface of the PDR, in the model using only thermal rates, the radiative association reaction of $\text{H} + \text{C}^+$ dominates

Table 3. Column densities along the cut.

Position	^{13}CO ($10^{15}\ \text{cm}^{-2}$)	CH^+ ($10^{12}\ \text{cm}^{-2}$)	CH ($10^{12}\ \text{cm}^{-2}$)
A	1.2(85 K) 1.3(200 K) 3.6(50 K)	2.2(85 K) 3.8(200 K) 1.9(50 K)	
B	4.8(85 K) 3.1(200 K) 13.8(50 K)	2.2(85 K) 3.8(200 K) 1.9(50 K)	
C	67.4(85 K) 43.5(200 K) 191.7(50 K)	6.4(85 K) 5.5(50 K) 11(200 K)	$\leq 4.5(85\ \text{K})$ $8.4(200\ \text{K})$ $3.6(50\ \text{K})$
D	6.7(85 K) 4.3(200 K) 19.1(50 K)	2.0(85 K) 1.7(50 K)	

Notes. The angular size assumes a filled Gaussian beam of $36''$ at CO, CH^+ transition frequencies and $33.9''$ at CH ones.

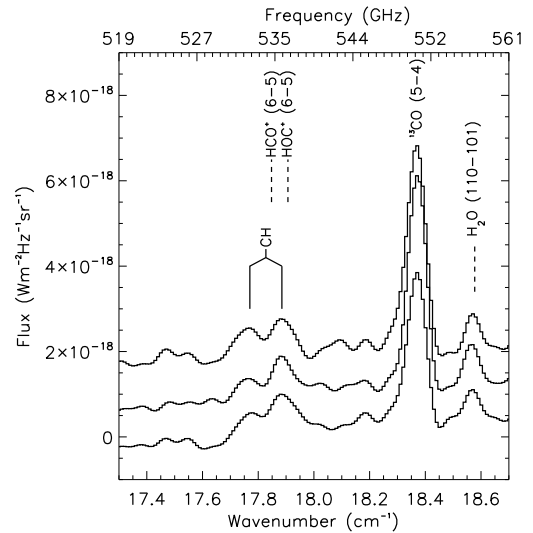


Fig. 4. Continuum subtracted spectra of the 3 positions on the Orion bar, displaying a doublet as expected for CH.

(blue curves). This reaction is, however, superseded by the $\text{H}_2 + \text{C}^+$ ion molecule reaction when the activation barrier can be overcome by the vibrational excitation of H_2 (red curves).

The fact that the CH^+ ion is seen extended all around the bar in the lower A_V regions and in high abundance seems to support the importance of excited H_2 in the CH^+ formation route. At lower frequencies, a pair of lines is observed at the position of the lambda-doublet CH transitions as shown in Fig. 4. In estimating the abundance of CH and CH^+ it has been assumed that the

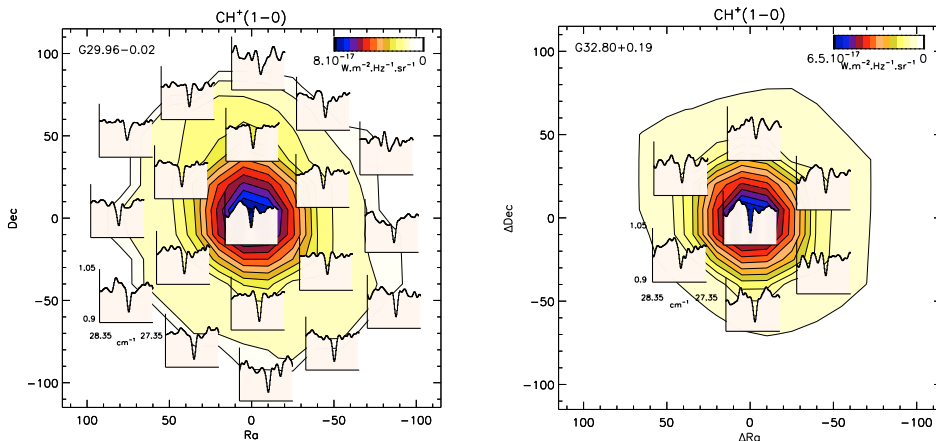


Fig. 5. *Left panel:* G29.96-0.02 CH⁺(1-0) continuum divided transmittance spectra, overlaid on the continuum intensity map at the same frequencies, extending away from the ultra compact HII source. *Right panel:* same for the G32.80+0.19 HII region. The outer detectors are vignettted so have been omitted.

opacity is negligible with the species thermalised at the temperatures noted. Only the spectra on the Bar are shown. The signal-to-noise ratio of these two lines will improve as the SPIRE FTS data processing is improved and with deeper observations toward Orion. Higher signal-to-noise is required to ensure that potential contributions from HCO⁺ and HOC⁺ line emission to this spectral region are considered, particularly for the higher frequency component.

G29.96–0.02 and G32.80+0.19

G29.96-0.02 and G32.80+0.19 are two ultra compact HII regions (e.g. Wood & Churchwell 1989), located near the Galactic plane. A global analysis of G29.96–0.02 SPIRE/FTS data can be found in this issue (Kirk et al. 2010). Both HII regions display a CH⁺ absorption line at a few percent of the continuum level. From a study of the line-to-continuum ratio maps shown in Fig. 5, it appears that for each continuum level across the maps, the amount of absorption remains essentially the same, within the signal-to-noise limits for the outer pixels. This suggests that the methylidyne ion is, as expected from its electronic visible absorption transitions, present in the foreground diffuse interstellar medium in relative high abundances. The apodized resolution of the spectra is 0.0724 cm⁻¹, corresponding to about 780 km s⁻¹ for the considered CH⁺ transition. Although the absorptions appear shallow, at the SPIRE/FTS resolution, a 5% absorption corresponds to a saturated line about 38 km s⁻¹ wide. This is consistent with the velocity dispersion of the foreground ISM as seen in HI absorption (e.g. Fish et al. 2003, Fig. 1). This precludes us from deriving stringent column densities at this level of analysis.

In subsequent observations with the SPIRE FTS the CH⁺ line has also been measured in the planetary nebula NGC7027 (Wesson et al. 2010) and in the ultra luminous infrared galaxy MRK231 (Van der Werf et al. 2010). The SPIRE FTS promises to yield a harvest of methylidyne data from a wide variety of sources.

3. Conclusions

We have presented the first detection of the lowest rotational transition of the methylidyne cation CH⁺ in an astronomical source. We have also shown strong evidence for the detection of the lambda doublet emission from the CH molecule. The detection of CH⁺ in a variety of environments shows that spectral mapping of this and other species will provide a powerful probe of the physics of the interstellar medium. These results have only

been made possible by the sensitivity and broad spectral coverage of the SPIRE FTS.

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