The RMS survey: far-infrared photometry of young massive stars

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ABSTRACT

Context. The Red MSX Source (RMS) survey is a multi-wavelength campaign of follow-up observations of a colour-selected sample of candidate massive young stellar objects (MYSOs) in the galactic plane. This survey is returning the largest well-selected sample of MYSOs to date, while identifying other dust contaminant sources with similar mid-infrared colours including a large number of new ultra-compact (UC) H\textsc{ii} regions.

Aims. To measure the far-infrared (IR) flux, which lies near the peak of the spectral energy distribution (SED) of MYSOs and UC\textsc{ii} regions, so that, together with distance information, the luminosity of these sources can be obtained.

Methods. Less than 50\% of RMS sources are associated with IRAS point sources with detections at 60 \textmu m and 100 \textmu m, though the vast majority are visible in Spitzer MIPS\textsc{g}al or IRAS Galaxy Atlas (IGA) images. However, standard aperture photometry is not appropriate for these data due to crowding of sources and strong spatially variable far-IR background emission in the galactic plane. A new technique using a 2-dimensional fit to the background in an annulus around each source is therefore used to obtain far-IR photometry for young RMS sources.

Results. Far-IR fluxes are obtained for a total of 1113 RMS candidates identified as young sources. Of these 734 have flux measurements using IGA 60 \textmu m and 100 \textmu m images and 724 using MIPS\textsc{g}al 70 \textmu m images, with 345 having measurements in both data sets.

Key words. stars: formation – stars: pre-main sequence – H\textsc{ii} regions – infrared: general – techniques: photometric – surveys

1. Introduction

Massive (M \geq 8 M_\odot) stars dominate the regions they form in due to their strong ionising radiation and their powerful outflows and winds. However, understanding the early stages of their formation and evolution remains difficult due to their rarity, the embeddedness of the formation process and the short timescale over which it takes place. Of particular interest is the massive young stellar object (MYSO) phase, which is mid-infrared (IR) bright but radio quiet, where core hydrogen burning has probably begun but an H\textsc{ii} region has yet to form. As such, MYSOs represent one of the earliest phases of the life of a massive star where it is truly massive while major accretion may still be ongoing, as evidenced by their powerful bipolar molecular outflows (e.g. Beuther et al. 2002a,c).

A major stumbling block to studying MYSOs, and by extension massive star formation in general, is the lack of large well-selected samples of sources from which to differentiate the global properties of all such sources from individual peculiarities and variations. Until recently, the relatively small catalogues of primarily serendipitously discovered MYSOs by Wynn-Williams (1982) and Henning et al. (1984) represented the largest collection of MYSOs but, with only of order 50 sources between them, this is a factor of 10 smaller than expected (Lumsden et al. 2002). This led Lumsden et al. (2002) to identify a sample of \sim 2000 candidate MYSOs from the Midcourse Space Experiment (MSX) point source catalogue (Egan et al. 1999, 2003) using colour selection criteria.

The Red MSX Source (RMS) survey is a campaign of follow-up observations designed to identify other dusty objects which contaminate this sample of MYSO candidates (ultra compact (UC) H\textsc{ii} regions, low-mass YSOs, Evolved stars, proto-planetary nebulae (PNe) and PNe) and gain information about the sources (Hoare et al. 2005; Mottram et al. 2006; Urquhart et al. 2008b). These include: \sim 1” resolution continuum radio observations to identify regions of ionised emission (Urquhart et al. 2007a, 2009); ground-based \sim 1” resolution mid-IR imaging where Spitzer IRAC (3.6 \textmu m, 4.5 \textmu m, 5.8 \textmu m and 8.0 \textmu m) imaging from the GLIMPSE survey (Benjamin et al. 2003; Churchwell et al. 2009) is unavailable to identify multiplicity in the \sim 18” MSX beam, obtain better astrometry and ensure that the survey does not bias against YSOs near H\textsc{ii} regions (e.g. Mottram et al. 2007); 13CO molecular line observations in order to obtain kinematic distances (Urquhart et al. 2007b, 2008a); and near-IR spectroscopy (e.g. Clarke et al. 2006) to distinguish between YSOs and evolved stars.

Full Table 1 is only available in electronic form at the CDS via anonymous ftp to cdsarc.u-strasbg.fr (130.79.128.5) or via http://cdsweb.u-strasbg.fr/cgi-bin/qcat?J/A+A/510/A89/DOI: 10.1051/A&A:2010510089 (2010)
In order to obtain the luminosities of the young massive stars identified by the RMS survey, photometry is required over as much of the spectral energy distribution (SED) as possible. In particular, the SEDs of YSOs and UCH\(\text{II}\) regions peak in the far-IR between \(\sim 60 \mu m\) and \(\sim 120 \mu m\). It is therefore important that far-IR photometry is available for those candidates identified as young, i.e. associated with star formation activity. Despite covering the majority of the sky, the IRAS point source catalogue (Beichman et al. 1988) does not contain entries for \(\sim 28\%\) of young RMS sources, and a similar number have only upper limit detections at \(60 \mu m\) and/or \(100 \mu m\). In addition, the background emission at far-IR wavelengths is strong and spatially variable within the galactic plane for IRAS data which adds complexity to background subtraction. Simple aperture photometry, which estimates the background emission using the average within an annulus surrounding the source is therefore unlikely to provide accurate measurements. However, fitting the point spread function (PSF) of the detector is also difficult due to the crowded nature of sources in the galactic plane and interaction between the extended background and the PSF.

In this paper, we begin by discussing in Sect. 2 the RMS sample of massive young stars and the data used to obtain far-IR fluxes for them: HIRES reprocessed IRAS Galaxy Atlas (IGA Cao et al. 1997) \(60 \mu m\) and \(100 \mu m\) images and MIPSGAL (Carey et al. 2009) \(70 \mu m\) images. In Sect. 3 we outline a new technique for performing aperture photometry using a 2-D fit to the background within an annulus, rather than a simple average. The results of photometric measurements using this technique are presented in Sect. 4, and discussed in Sect. 5. We then summarise and reach our conclusions in Sect. 6.

2. Data

2.1. The RMS sample of massive young stars

Using the RMS follow-up observations and data from the literature (e.g. 2MASS, GLIMPSE), all of which are available on the RMS database (http://www.ast.leeds.ac.uk/RMS/), each candidate MYSO has been assigned one of several identifications (see Mottram 2008, for more details). Young sources can be assigned “YSO”, “H\(\text{II}\) region”, “H\(\text{II}/\text{YSO}\)” and “Young/Old Star”.

Sources assigned the “YSO” identification are radio-quiet mid-IR point sources associated with \(^{13}\text{CO}\) emission which do not look isolated in near-IR images. Where near-IR spectroscopy is available, these sources usually show H\(\text{I}\) emission, though not at Case-B ratios (Baker & Menzel 1938), and can also show H\(_2\), He\(_1\), Fe\(_{\text{II}}\) and [Fe\(_{\text{II}}\)] emission, as well as CO bandhead emission and/or absorption. “H\(\text{II}\) region” sources are either extended in \(\sim 1''\) resolution mid-IR imaging and/or radio loud, with Case-B H\(\text{I}\) line ratios in the near-IR. “H\(\text{II}/\text{YSO}\)” sources either have indications that both a YSO and an H\(\text{II}\) region are present within the MSX beam (18’’), or have conflicting evidence as to their identification with insufficient information currently available to determine between the “YSO” and “H\(\text{II}\)” region designations. Finally sources identified as “Young/Old Star”’s are unlikely to be H\(\text{II}\) regions or PNe as they are radio quiet and unresolved in the mid-IR, but have insufficient evidence to determine whether the source is a YSO or an evolved star. As more information becomes available identifications are updated where needed, particularly from our near-IR spectroscopy and the luminosities derived using the data presented in this paper, which will be the subject of a forthcoming paper (Mottram et al., in prep.).

![Fig. 1. A comparison of an IRAS ISSA 60 μm image (top), an IGA 60 μm image after 20 HIRES iterations (centre) and a Spitzer MIPSGAL 70 μm (bottom) image of the same region centred on G015.00+0.00. The relevant features of IGA and MIPSGAL data are discussed in Sects. 2.2 and 2.3.](image)

2.2. The IRAS Galaxy Atlas

The IRAS sky survey scanned most regions of the sky more than once, often at different scan angles, thus additional spatial information is available within the data. Cao et al. (1997) therefore undertook HIRES reprocessing of the original IRAS Sky Survey (ISSA) images (Wheelock et al. 1994) for the galactic plane in order to produce the IGA. These data have improved resolution (1.0’’ \(\times\) 1.7’’ at 60 μm, 1.7’’ \(\times\) 2.2’’ at 100 μm) with respect to both the original ISSA, as shown in Fig. 1, and the
more recent Improved Reprocessing of the IRAS Survey (IRIS) images (Miville-Deschênes & Lagache 2005), which have resolutions between 3.5′ and 5′. Though the IRIS data are better calibrated due to Miville-Deschênes & Lagache (2005) using COBE/DIRBE data, the improved resolution of the IGA is crucial when considering regions in the galactic plane.

Due to the shape of the IRAS detector pixels and the detector layout used, the IRAS sky survey had different resolution parallel and perpendicular to the scanning direction of the satellite. This caused the point source function (PSF) of raw IRAS images to be elliptical, rather than circular—a feature which is retained in the IGA (see Fig. 1), although IRIS images are smoothed to a circular PSF with resolution closer to the larger dimension (i.e. cross-scan, see Miville-Deschênes & Lagache 2005). The angle between the major axis of the PSF and the x-axis of the image also varies across the IGA images due to the changing scan direction of the IRAS satellite across the sky as it orbited the earth and the earth orbited the sun.

As a part of HIRES processing, beam sample map images were produced which have the PSF placed at the same positions on each tile on a smoothed version of the 20 iteration image in order to provide a reasonable estimate of the local background. These were used to measure the local beam angle for each source. In addition, HIRES processing results in a negative ringing artifact around the central maximum of the PSF (see Fig. 1) which must be avoided during aperture photometry, caused by interaction between the PSF and the non-zero background (see Cao et al. 1997, Sect. 5.2 for a more detailed discussion).

Despite the improved resolution of the IGA over the original IRAS data, the resolution is not as good as that of MSX (18.3”) and source contamination may still be an issue (see Fig. 2). Nearby sources in the IRAS point source catalogue (PSC) were identified in order to remove contaminating sources from the sky annulus, and flag sources where the flux measurement is potentially contaminated by other nearby sources. However, the PSC is not necessarily complete within the galactic plane and may not include cases where several sources cluster to give the appearance of a resolved source. Therefore, all sources were subsequently visually inspected and cases where the RMS target was dominated by other nearby sources were flagged as unusable.

2.3. MIPSGAL

For sources in the inner galactic plane (5° < l < 63° and 298° < l < 355°) to |b| ≤ 1°, higher resolution data are available from the Spitzer Space Telescope (SST) MIPS Galactic plane (MIPSGAL) survey (Carey et al. 2009), which has a resolution of 18” at 70 μm. Though fully reduced and calibrated mosaic images have not yet been released for this waveband, observational data with basic reduction is available publicly for the whole survey region. This was obtained and mosaiced in galactic
coordinates with a pixel scale of 4.8″ pixel$^{-1}$ using the program MONTAGE (Berriman et al. 2006). Note that we do not consider the MIPSGAL 24 μm data as the majority of RMS sources are saturated, we already have MSX 21 μm data and this waveband is not crucial to the determination of the total IR flux and therefore luminosity for our sources.

The MIPS 70 μm detector suffers from a non-linear pixel response for surface brightnesses above ~66 MJy sr$^{-1}$ (Dale et al. 2007) in wide-field mode or 1.5 mJy pixel$^{-1}$, corresponding to a point source flux of ~1 Jy. Though this is not a particularly large effect in extragalactic studies, the brightness of the galactic plane at these wavelengths results in background emission levels that are often of order 1000 MJy sr$^{-1}$. We therefore derive a pixel correction using the following method, which is based on that used by Dale et al. (2007), with the additional inclusion of errors on the data in our linear regression. For each standard star measured using aperture photometry in wide-field observations by Gordon et al. (2007, their Table 4), the measured flux in mJy can be calculated from the measured calibration factor, the predicted flux (derived from Engelbracht et al. 2007) and the mean calibration factor obtained from PSF fitting photometry of the same observations. The measured and predicted fluxes are then both converted to surface brightnesses by dividing by 1.5 mJy pixel$^{-1}$, after which a linear bisector least-squares fit to those sources with fluxes ≥1 Jy (see Fig. 3) is used to obtain the pixel correction of the form:

$$F_{\text{true}} = a(F_{\text{meas}})^b.$$  (1)

We find $a = 0.558 \pm 0.051$ and $b = 1.11 \pm 0.02$, which compares well with the results of Dale et al. (2007). The differences between our correction and that of Dale et al. (2007) come from the use of a linear bisector rather than simply a least squares linear fit. This correction was applied to all target images before comparison of corrected and uncorrected MIPSGAL fluxes for all 707 detections. The solid lines show linear bisector least-squares fit to the data which the dashed line again indicates the line of concordance.

3. Two dimensional background fitting photometry

In order to perform background subtraction when the background itself is spatially variable (e.g. see Fig. 1), a 2-dimensional fit to the background is the next logical progress from a simple average. As in standard aperture photometry, an annulus around the source of interest is selected, though in this case the surface in the annulus is fitted with a 2-D second order polynomial using a least-squares multiple linear regression method (see Bevington 1969). While a higher dimension fit could be used, it is questionable whether an annulus around a source would uniquely constrain such a fit, or whether the variation in diffuse background truly changes quickly on small spatial scales. Ideally, the background should vary on large scales compared with the size of a point source so that it is fully resolved.

It is possible, however, that other nearby sources lie partially or wholly within the sky annulus (e.g. see Fig. 2). In such cases the coordinates of these sources can be used to remove pixels up to a given radius around such sources from the area used for the fit, so that they do not affect it. It is not desirable to have too low an area from which to derive the fit, so a cut-off percentage is used below which no more sources are considered for removal.

Having derived the fit to the pixels within the annulus, the fit is then removed from the original image and the total counts within the aperture calculated. The total error in the aperture counts is the combination in quadrature of the noise in the sky within the aperture, the error in the fit to the sky within the

![Figure 3](image-url)
annulus and the errors due to the detector (e.g. read noise, gain etc.). With modern detectors it is often difficult to recover this last error explicitly from the released mosaiced images, but it can be considered a limiting minimum error $\sigma_{\text{lim}}$, and can therefore be obtained through measurement of regions free from sources of emission. The total error is therefore given by:

$$\sigma_{\text{Tot}} = \left( N_{\text{ap}} \frac{\Sigma_{\text{ap}} (c_o - c_f)^2}{N_{\text{an}}} + N_{\text{ap}}^2 \frac{\Sigma_{\text{ap}} (c_o - c_f)^2}{N_{\text{an}}} + \sigma_{\text{lim}}^2 \right)^{1/2}$$ (2)

where $c_o$ and $c_f$ are the pixel counts of the included pixels within the sky annulus of the original image and 2-D fit respectively while $N_{\text{ap}}$ and $N_{\text{an}}$ are the number of pixels included in the aperture and annulus.

In order to test this method, the results of aperture photometry and 2-D background fitting photometry were compared using TIMMI2 10.4 $\mu$m images from Mottram et al. (2007) of sources in their “clean” subsample (see Fig. 4). The images for these sources have approximately flat uniform backgrounds and isolated unresolved targets, so the results should be the same using either method. The same aperture and inner annulus radii were used as for the standard photometry for TIMMI2 observations (i.e. 12 and 15 pixels respectively) for both methods. However it was found that the 2-D aperture fitting photometry requires a larger sky annulus in order to constrain the fit properly. An outer radius of 30 pixels was therefore used for the aperture fitting photometry, as opposed to the outer radius of 20 pixels used for the standard aperture photometry measurements. The flux was also measured using the 2-D aperture fitting photometry at four locations on each image which were free of source emission and the mean calculated in order to obtain a limiting background flux for each source. The mean of this limiting background flux over all sources was then calculated to obtain the limiting error $\sigma_{\text{lim}}$.

The linear bisector least-squares fit to the photometric results for the subsample of TIMMI2 sources (see Fig. 4) has a slope of 1.0 within the errors, and the correlation of aperture photometry to 2-D background fitting photometry fluxes is also 1.0. The two methods are therefore consistent with each other within the calculated errors. In addition to comparisons between the two methods for the same image, comparisons were also performed between the original TIMMI2 fluxes and 2-D background fitting photometry obtained from images which had a diagonal ramp in the background counts and a large bright 2-D gaussian source near the target added to them. These returned very similar results to the first test, so the new method works well even in challenging circumstances.

In order to obtain accurate photometry for IGA data, which has an elliptical beam profile (see Fig. 2), both standard aperture photometry and aperture fitting photometry routines were also developed to use elliptical instead of circular apertures and annuli.

4. Results

4.1. IGA results

Aperture fitting photometry was undertaken on all 1336 RMS candidate MYSOs identified as young sources, of which 602 were too badly contaminated in IGA data for reliable results to be obtained, leaving 734 sources with acceptable measurements. The aperture and inner and outer sky annulus major axis radii of 11, 28 and 46 pixels were used for the 60 $\mu$m data (as shown in Fig. 2) and 13, 30 and 50 pixels for the 100 $\mu$m data for all sources, after careful examination of isolated sources. The ratios of the major to minor axis radii used were 1.9 for 60 $\mu$m and 1.3 for 100 $\mu$m. The limiting flux error $\sigma_{\text{lim}}$ was measured using background regions in 13 images for both 60 $\mu$m and 100 $\mu$m data, resulting in errors of 0.51 Jy and 5.97 Jy respectively. The much larger limiting error for the 100 $\mu$m images is probably due to the brightness of diffuse galactic dust emission at this wavelength.

During their analysis, Cao et al. (1997) found that the IGA after 20 processing iterations overestimates the fluxes of isolated sources with respect to the IRAS PSC. The final flux results therefore include division by the IGA(20)/IGA(1) factors measured by Cao et al. (1997) of 1.02 and 1.10 for IGA 60 $\mu$m and 100 $\mu$m data respectively. This correction results in a systematic error in the fluxes due to the error in measurements of Cao et al. (1997) of 7% and 14% in the measured 60 $\mu$m and 100 $\mu$m fluxes respectively. This systematic error in the fluxes is in addition to the error derived during the 2-D background fitting photometry, which is a random error. Therefore these two errors are kept separate and the systematic error is not included in the results presented in Table 1. The correction factors were measured by Cao et al. (1997) using relatively isolated sources with well defined backgrounds, so may well be larger for RMS candidates which are often in complex star formation regions with high levels of background emission.

An example of these results is presented in Table 1, while the full version of this table is available at the CDS. The MSX PSC name for each target is given in Col. 1 and the IGA measurements and errors are presented in Cols. 2 and 3 respectively. The coordinates of the nearest IRAS PSC entry within 0.7$^\circ$ (the selection of which will be discussed in Sect. 5.1) of the MSX coordinates of the target are given for comparison in Cols. 5 and 6. The offset between the MSX and IRAS PSC coordinates is given in Col. 7 while the IRAS PSC fluxes with errors are given in Cols. 8 and 9. The errors in the IGA fluxes were calculated using Eq. (2), and are sometimes quite small. However, Cao et al. (1997) estimate that the overall photometric accuracy of the IRAS Galaxy Atlas is $\sim$25%, so the dominant error in these results is probably from the image data rather than the measurements.

Of the 734 targets with IGA fluxes, 682 have an IRAS PSC source within 0.7$^\circ$ but 172 of these PSC sources are upper limit detections at either 60 $\mu$m, 100 $\mu$m or both. Therefore,
Table 1. Example of measured IGA 60 μm, IGA 100 μm and MIPSGAL 70 μm photometry, see text for details. A full version of this table is available online at the CDS.

<table>
<thead>
<tr>
<th>Object</th>
<th>IGA 60 μm flux (Jy)</th>
<th>IGA 100 μm flux (Jy)</th>
<th>MIPSGAL 70 μm flux (Jy)</th>
<th>Nearest IRAS PSC Entry</th>
<th>IRAS 60 μm flux (Jy)</th>
<th>IRAS 100 μm flux (Jy)</th>
</tr>
</thead>
<tbody>
<tr>
<td>G00.0732+00.2295</td>
<td>134.48 ± 2.14</td>
<td>132.32 ± 12.52</td>
<td>59.35 ± 3.20</td>
<td>16:19:48.7</td>
<td>19:27:42</td>
<td>0.210</td>
</tr>
<tr>
<td>G00.1109+00.4011</td>
<td>1145.32 ± 1.86</td>
<td>2323.30 ± 10.73</td>
<td>885.59 ± 3.56</td>
<td>18:11:33.2</td>
<td>1810.03</td>
<td>219</td>
</tr>
<tr>
<td>G00.1723-00.0656</td>
<td>181.62 ± 2.85</td>
<td>299.14 ± 10.16</td>
<td>114.20 ± 3.16</td>
<td>18:10:25.2</td>
<td>19:17:51</td>
<td>0.078</td>
</tr>
<tr>
<td>G00.9454-00.0373</td>
<td>1423.15 ± 3.95</td>
<td>1848.77 ± 14.14</td>
<td>1119.95 ± 3.27</td>
<td>18:11:53.0</td>
<td>18:36:20</td>
<td>0.116</td>
</tr>
<tr>
<td>G01.5932-00.5708</td>
<td>94.24 ± 1.75</td>
<td>229.58 ± 9.25</td>
<td>17.77 ± 4.01</td>
<td>18:15:10.0</td>
<td>18:17:30</td>
<td>0.105</td>
</tr>
<tr>
<td>G02.9090+00.2607</td>
<td>3212.46 ± 15.41</td>
<td>1811.66 ± 27.53</td>
<td>2916.63 ± 5.33</td>
<td>18:14:39.0</td>
<td>17:52:03</td>
<td>0.183</td>
</tr>
<tr>
<td>G03.2097+01.4136</td>
<td>1214.14 ± 4.21</td>
<td>2999.06 ± 14.72</td>
<td>1116.61 ± 3.59</td>
<td>18:14:47.8</td>
<td>17:32:48</td>
<td>0.511</td>
</tr>
<tr>
<td>G03.6562-00.5997</td>
<td>403.97 ± 2.27</td>
<td>989.68 ± 9.72</td>
<td>403.22 ± 3.12</td>
<td>18:17:24.4</td>
<td>17:22:13</td>
<td>0.027</td>
</tr>
<tr>
<td>G04.3885+00.4760</td>
<td>661.08 ± 4.96</td>
<td>1681.81 ± 26.79</td>
<td>169.60 ± 3.57</td>
<td>18:17:24.7</td>
<td>17:06:26</td>
<td>0.023</td>
</tr>
<tr>
<td>G04.0329-00.5155</td>
<td>265.45 ± 16.31</td>
<td>1195.19 ± 8.10</td>
<td>70.79 ± 3.64</td>
<td>18:17:50.7</td>
<td>16:59:55</td>
<td>0.040</td>
</tr>
<tr>
<td>G04.2071-00.1105</td>
<td>346.53 ± 4.85</td>
<td>566.17 ± 14.71</td>
<td>201.11 ± 3.48</td>
<td>&lt;145.00</td>
<td>923.00</td>
<td>&lt;25%</td>
</tr>
<tr>
<td>G04.2166-00.6344</td>
<td>70.15 ± 4.50</td>
<td>398.38 ± 13.86</td>
<td>57.22 ± 3.21</td>
<td>18:18:38.5</td>
<td>16:53:33</td>
<td>0.088</td>
</tr>
<tr>
<td>G04.6703+00.4795</td>
<td>138.96 ± 1.63</td>
<td>594.71 ± 35.98</td>
<td>40.19 ± 3.47</td>
<td>&lt;145.00</td>
<td>145.00</td>
<td>&lt;82%</td>
</tr>
<tr>
<td>G05.0755+01.1212</td>
<td>249.51 ± 4.06</td>
<td>575.18 ± 14.22</td>
<td>211.48 ± 3.29</td>
<td>18:18:27.6</td>
<td>15:53:36</td>
<td>0.086</td>
</tr>
<tr>
<td>G05.0939+00.1913</td>
<td>247.00 ± 2.36</td>
<td>480.57 ± 8.71</td>
<td>75.13 ± 3.14</td>
<td>&lt;145.00</td>
<td>145.00</td>
<td>&lt;34%</td>
</tr>
<tr>
<td>G05.4805+00.3900</td>
<td>175.60 ± 1.20</td>
<td>712.90 ± 6.77</td>
<td>23% ± 6.77</td>
<td>&lt;145.00</td>
<td>145.00</td>
<td>&lt;38%</td>
</tr>
<tr>
<td>G06.1438+00.0074</td>
<td>44.23 ± 4.26</td>
<td>161.02 ± 14.41</td>
<td>85.67 ± 3.17</td>
<td>18:20:03.7</td>
<td>14:53:31</td>
<td>0.260</td>
</tr>
</tbody>
</table>

Notes. (1) No IRAS PSC entry within 0.7° of MSX target coordinates.

Fig. 5. A comparison of measured IGA 60 μm fluxes with IRAS 60 μm (left) and IGA 100 μm fluxes with IRAS 100 μm (right) PSC fluxes for 510 sources with PSC detections. The solid line shows the linear-bisector least-squares fit to the data while the dashed line indicates the line of concordance.

IGA fluxes have been obtained for 224 sources which have no far-IR detection in the IRAS PSC at these wavelengths. Figure 5 shows the comparison between IGA and IRAS PSC fluxes for the 510 sources which have IGA fluxes and true IRAS PSC detections at 60 μm and 100 μm. The IGA fluxes are, on average, slightly larger than the IRAS PSC fluxes by a factor of 1.12 ± 0.05 at 60 μm and 1.35 ± 0.09 at 100 μm, however the fluxes in the IRAS PSC did not include a correction for hysteresis at 60 μm, which when included tends to increase these fluxes and would account for the increase seen (Cao et al. 1997, give IGA(1)/PSC = 1.12). The increase in the 100 μm data may be due to flux increases during HIRES processing as our sources are often in regions of considerable background emission which can cause interactions between this and the point source. However, the difference may also be due to more accurate background subtraction.

4.2. MIPSGAL results

2-D background fitting photometry was attempted for 843 Red MSX Source survey candidate MYSOs identified as young sources which lie within the MIPSGAL survey region. However, 85 were either heavily saturated or have a much brighter source nearby so that no usable information could be obtained from the image (e.g. see the left hand plot of Fig. 6) and 34 are visible but saturated. Of the 724 sources where flux measurements were possible, 693 are detected and measurable in MIPSGAL images, though nearby sources had to be masked from the sky annulus for some of these sources (e.g. see the right hand plot of Fig. 6). In addition, 12 sources are not detected at 70 μm and therefore probably not MYSOs or HII regions and 19 sources are non-detections due to there being bright sources or complex nearby, and so have large upper limits on the flux associated with them. For an isolated unresolved source, aperture and sky radii of 13, 16 and 32 pixels were used, while these radii were set by inspection of the images for more complex sources.

The point source function (PSF) of Spitzer MIPS observations is circular but slightly non-gaussian and includes a low-level extended contribution. An aperture correction was therefore calculated for the MIPSGAL images by performing photometry on the image of an unresolved source with a relatively flat background using various aperture radii and inner and
outer sky annulus radii, then dividing the results of photometry with a particularly large aperture which measures as close to the total source flux as possible. This ratio then provides an aperture calibration which is multiplied by the measured flux in order to arrive at the total flux for a given source. Calculation of the aperture correction was undertaken after pixel non-linearity correction of the images (see Sect. 2.3). A sample of correction values are shown in Table 2, and compare well to those obtained by Gordon et al. (2007) using model PSFs for blackbodies of 10,000 K, 60 K and 10 K, particularly those for the 10 K blackbody. The errors in the calculated corrections are generally below 10%, and are primarily due to the limiting error associated with the measurements used to derive the correction factors. The limiting error $\sigma_{\text{lim}}$ was found to be 2.68 Jy, measured as discussed previously, using 17 images with relatively flat backgrounds.

For non-detections, four measurements of the background flux were taken using the standard aperture and annulus radii and the mean of the absolute values of the measurements calculated. A 3σ upper limit on the flux for the source is then given by three times this mean. For those sources which are non-detections due to bright nearby emission, the flux was calculated in the same way as for a standard source, and this flux used as an upper limit to the true MIPSagal 70 $\mu$m flux for the source. Though the saturation limits of MIPSagal 70 $\mu$m images have yet to be fully characterised, it is certainly a problem for bright RMS sources or sources in or near bright regions of the galactic plane.

Example MIPSagal 70 $\mu$m 2-D background fitting photometry results are shown in Col. 4 of Table 1, while the full version of this table is available at the CDS. All sources which either have a flux measurement or are non-detections for whatever reason are included in the full table.

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shows the mean IRAS 60 $\mu$m to MIPSagal 70 $\mu$m flux ratio as a function of offset between the IRAS and MSX coordinates. From this figure, as well as visual inspection of MSX and MIPSagal images, the radius out to which the nearest IRAS PSC candidate is considered a probable counterpart to the RMS candidate was set at $0.7\arcmin$. This is more than twice the resolution of MSX ($18.3\arcsec$) and almost three times the quoted IRAS PSC cross-scan positional error ($15.6\arcsec$). This limit was also used to identify probable IRAS counterparts for sources with IGA observations (see Sect. 4.1).

Figure 8 shows the comparison between the measured MIPSagal 70 $\mu$m photometry and IRAS PSC 60 $\mu$m (left) and 100 $\mu$m (right) fluxes for all 243 sources with measurable MIPSagal fluxes and detections at both 60 $\mu$m and 100 $\mu$m in the IRAS PSC within 0.7$\arcmin$. Sources which have upper limit IRAS PSC fluxes (i.e. $\sigma = 100\%$) at either 60 $\mu$m or 100 $\mu$m are not included. Linear bisector least-squares fits to the data give relationships between the MIPSagal 70 $\mu$m and IRAS fluxes of the form shown in Eq. (1), with powers equal to one within the errors and flux ratios of $F_{\text{IRAS 60}}/F_{\text{MIPS 70}} \approx 1.32 \pm 0.21$ and $F_{\text{IRAS 100}}/F_{\text{MIPS 70}} \approx 2.82 \pm 0.65$. The expected flux ratios assuming a blackbody of temperature 50 K (or 40 K) are $F_{\text{IRAS 60}}/F_{\text{MIPS 70}} = 1.12$ (or 0.95) and $F_{\text{IRAS 100}}/F_{\text{MIPS 70}} = 1.88$ (or 2.43), so the observed ratios may be in part due to the differences in filter profiles between IRAS and MIPSagal. However, the mean ratios derived from fits to the SED models of RMS sources using the model fitter of Robitaille et al. (2007) are 0.93 and 1.07 for $F_{\text{IRAS 60}}/F_{\text{MIPS 70}}$ and $F_{\text{IRAS 100}}/F_{\text{MIPS 70}}$ respectively (Mottram et al., in prep.). The model fitter convolves the models with the observed filters using the standard colour correction, so the model fluxes and flux ratios are directly comparable to the observations. Though the IRAS filters are much wider than the MIPSagal filter, this effect is included in the ratios produced by the model fitter as well, so the large ratio between IRAS PSC and MIPSagal fluxes is probably due to contamination in the larger IRAS beam, particularly in the 100 $\mu$m band. The smaller model ratio of $F_{\text{IRAS 60}}/F_{\text{IRAS 100}}$ than would be expected for a blackbody of relevant temperature is probably the result of the distribution of dust temperatures within the SED models.

The general spread of sources in Fig. 8 is due to several factors. The SED of all young sources will be similar, but not identical, leading to some fluctuation of the flux ratios. As mentioned above, sources in more crowded regions are likely to have more confused IRAS fluxes, leading to both larger flux ratios and larger variation. For sources closer to the galactic centre, the line of sight of observations passes through more of the galactic plane, so there is more chance of confusion.

### 5.2. Comparison of MSX and MIPSagal colours

Figure 9 shows a colour–colour plot using MSX and MIPSagal 70 $\mu$m fluxes. Sources with RMS identifications of “YSO” are shown with “+” signs, “HII/YSO” with squares, “H II region” with circles “Young/Old star” with triangles. The vertical solid line indicates the RMS $F_{\text{MSX 21}}/F_{\text{MSX 8}}$ selection cut, while downward arrows indicate sources with non-detections at 70 $\mu$m. Note that sources with fluxes above this cut within the 1$\sigma$ errors are still included in the sample.
in GLIMPSE images and so are most likely evolved stars. For comparison the SEDs of G09.4332+00.1542 and G30.7569+00.0247, a typical YSO, are shown in Fig. 10. It can be seen that inclusion of far-IR flux data is important for discriminating between young and evolved sources. A further two (G30.8625−00.0629 and G33.7202+0.1762) appear similar to the previous sources in GLIMPSE images but are also associated with unresolved radio detections, so are most likely PNe. All four sources show 13CO detections, probably due to diffuse clouds either behind or in front of them. The last (G316.7754−00.0447) is certainly in a star forming region as both MSX and GLIMPSE images show multiple sources and diffuse emission nearby. However no point sources are evident in any of the Spitzer IRAC bands, so this source is either a more evolved H II region or a photo-dissociation region (PDR) which has bright mid-IR emission. RMS sources already identified as evolved are not included in this paper, and have not had their MIPS fluxes measured.

In general the H II regions have larger $F_{\text{MSX} 21}/F_{\text{MSX} 8}$ and $F_{\text{MIPS} 70}/F_{\text{MSX} 21}$ ratios than the YSOs, which is consistent with the H II regions being more embedded than the YSOs. However, there is a population of YSOs which have similar colours to the H II regions.

6. Summary and conclusions

In this paper, we have presented a new technique for background subtraction in aperture photometry using a 2-D fit to the background within the sky annulus. This was then used to obtain photometry from MIPS, MSX, IRAS, SCUBA sub-millimetre and MAMBO 1.2 mm fluxes (Skrutskie et al. 2006; Egan et al. 2003; Beichman et al. 1988; Di Francesco et al. 2008; Beuther et al. 2002b) in addition to those presented in this paper. These and the SEDs of all young RMS sources will be discussed further in Mottram et al., in prep.

These far-IR flux measurements, along with other photometric data, provide information about the spectral energy distribution (SED) of young RMS sources, and thus allow the determination of the luminosities of this sample of sources. This will allow the systematic study of the properties of a large, well-selected sample of MYSOs as a function of luminosity, and will be the subject of forthcoming papers. Information on all sources discussed in this paper is available via the RMS database, which can be found at www.ast.leeds.ac.uk/RMS/.

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