Transiting exoplanets from the CoRoT space mission*

VII. The “hot-Jupiter”-type planet CoRoT-5b


(Affiliations can be found after the references)

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ABSTRACT

Aims. The CoRoT space mission continues to photometrically monitor about 12 000 stars in its field-of-view for a series of target fields to search for transiting extrasolar planets ever since 2007. Deep transit signals can be detected quickly in the “alarm-mode” in parallel to the ongoing target field monitoring. CoRoT’s first planets have been detected in this mode.

Methods. The CoRoT raw lightcurves are filtered for orbital residuals, outliers, and low-frequency stellar signals. The phase folded lightcurve is used to fit the transit signal and derive the main planetary parameters. Radial velocity follow-up observations were initiated to secure the detection and to derive the planet mass.

Results. We report the detection of CoRoT-5b, detected during observations of the LRa01 field, the first long-duration field in the galactic anti-center direction. CoRoT-5b is a “hot Jupiter”-type planet with a radius of 1.1948 ± 0.0046 RJup, a mass of 0.467 ± 0.047 MJup, and therefore, a mean density of 0.217 ± 0.031 g cm⁻³. The planet orbits an F9V star of 14.0 mag in 4.0378962 ± 0.0000019 days at an orbital distance of 0.04947 ± 0.00026 AU.

Key words. planets and satellites: general – techniques: photometric – techniques: radial velocities

1. Introduction

CoRoT searches for the photometric signal of transiting extrasolar planets. Radial-velocity follow-up measurements help us understand the nature of the transiting body and allow us to derive its mass.

The nominal lightcurve analysis for small transiting signals has to await the completion of an observing run and detailed signal analysis. The mission “alarm-mode” (Quentin et al. 2006; Surace et al. 2008), however, can be used to quickly trigger follow-up measurements during ongoing observations of a target field. The “alarm-mode” is used to increase the transmitted time-sampling for individual stellar lightcurves in the CoRoT exoplanet channel. The sampling is increased from 512 s to 32 s if a transit-like signal is detected during the observations. It therefore provides planetary candidates early during an observing run, which are, however, biased towards relatively large planetary candidates because of the limited data set available at this point.

CoRoT-5b is the fifth secured transiting planet detected by CoRoT. As CoRoT-1b to CoRoT-4b (Alonso et al. 2008; Barge et al. 2008; Deleuil et al. 2008; Moutou et al. 2008; Aigrain et al. 2008), it was first detected by the alarm-mode. Here, we present the photometric detection of CoRoT-5b by the satellite based on pre-processed alarm-mode data, the accompanying radial-velocity observations confirming its planetary nature, and the resulting planet parameters.

2. Observations and data reduction

CoRoT-5b was detected in the LRa01-field, the second long-run field of CoRoT. The field is located near the anti-center direction of the galaxy at RA(2000): 06°46′53″ and Dec(2000): −00°12′00″ (Michel et al. 2006). The observing sequence started on October 24, 2007 and finished after 112 days duration. CoRoT observations usually have a very high duty cycle since data gaps are mainly caused by the regular crossings of the South Atlantic Anomaly (SAA), which typically last for about 10 min. During the observations of the LRa01 field, however, two longer interruptions occurred. An intermediate interruption of about 12 hours occurred eight days after the beginning of the observing run, and a longer data gap of about 3.5 days started on January 18, 2008, after a DPU reset. Finally, a duty cycle of 93% was achieved.
The alarm-mode was triggered after 29 days of observations. When seven transit-like signals were detected, the time sampling was switched to 32 s. The alarm-mode data for CoRoT-5 are based on the analysis of “white light” lightcurves, without using the color information of the CoRoT prism. In total 219,711 data points were obtained, 214,938 of it in oversampling mode. The data pipeline flags data points taken during the SAA crossing or affected by other events decreasing the data quality. When taking only unflagged data into account, the number of data points reduced to 204,092 in total and 199,917 as highly sampled.

The alarm-mode data were processed with a first version of the data reduction pipeline (Auvergne et al. 2009). The pipeline corrects for the CCD zero offsets and gain, the sky background intensity and the telescope jitter. In addition, “hot pixels” (Pinheiro da Silva et al. 2008) affect the lightcurves, causing sudden jumps in intensity of varying duration. The lightcurve of CoRoT-5 was, however, only moderately affected by such jumps, as can be seen in Fig. 1, which shows the full lightcurve. The oversampled part of the data set was re-binned to display the whole lightcurve with a 512 s time sampling. The measured intensity decreases during the observing run, as observed for all stars in the fields. Overall, CoRoT-5 only shows a minor level of variability, without clear periodicity.

CoRoT measures stellar intensities by aperture photometry using optimized masks (Llebaria et al. 2003) that encompass the shape of the stellar point-spread-functions (PSFs). The bi-prism introduced in the light path of the exoplanet channel (Auvergne et al. 2009) causes relatively wide PSFs of unusual shapes that vary with e.g. stellar magnitude. Contaminating eclipsing binary stars within the PSF could mimic a planetary transit-like signal. Based on the pre-launch observations of the target field included in the Exo-Dat data base (Deleuil et al. 2009), the contamination of the mask of CoRoT-Exo5 is estimated to 8.4 %. Refinement of this value will be performed in a more detailed future analysis using the dedicated windowing mask for this target star. We subtracted this flux level from the lightcurve before normalization to take low level contamination into account.

The overall intensity trend and smaller scale variability of the lightcurve were removed. To do this, we resampled the lightcurve to 512 s sampling rate first and convolved this lightcurve with a fourth order Savitzky-Golay filter (similar to the treatment for CoRoT-2b, Alonso et al. 2008). Then median averages were calculated for 24 h segments of the lightcurve (excluding the transit points and the data jumps), which was fitted by a spline-curve. The original lightcurve was then divided by the spline fit. The filtered lightcurve was used for normalization and further analysis. The out-of-eclipse scatter of CoRoT-5 was determined from the standard deviation of data points in the phase-folded lightcurve. It was found to be 0.0017 mag.

3. Photometric follow-up observation

Photometric follow-up observations with higher spatial resolution than CoRoT’s (of ≈20′′× 6′′) are used to exclude the presence of nearby contaminating eclipsing binaries (Deeg et al., this volume). Such observations of CoRoT-5 were performed at the IAC 80 cm telescope at Teide Observatory, Tenerife, on the January 12, and March 11, 2008 at a spatial resolution of about 1.5′′. These data showed only one star bright enough to cause a potential false alarm, about 8′′ southwest of the target. Observations obtained during and out of a transit (“on/off photometry”) showed, however, that this contaminating star varies by less than 0.08 mag. This is far below the variation of about 0.55 mag that is required by this star in order to explain the observed signal in the CoRoT data.

4. Radial velocity follow-up observations

In January 2008, after the identification of a transit signal by the alarm-mode, CoRoT-5 was observed with the SOPHIE spectrograph installed on the 193 cm telescope at the Haute Provence Observatory. Two radial velocity measurements were taken at opposite quadrature phases of the radial velocity variation expected from the transit ephemerides assuming a circular orbit. At this time the data were found to be compatible with a radial velocity amplitude suggesting a Jupiter mass planet. Additional measurements were obtained later in the season to confirm the reality of the signal but not enough to obtain a precise measurement of the orbit eccentricity. One year later, a new series of measurements was obtained with the HARPS spectrograph installed on the 3.6 m ESO telescope at La Silla in Chile.

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**Table 1. Radial velocity measurements of the star CoRoT-5 obtained by SOPHIE and HARPS spectrographs from December 2007 to December 2008.**

<table>
<thead>
<tr>
<th>BJD</th>
<th>RV</th>
<th>Error</th>
</tr>
</thead>
<tbody>
<tr>
<td>≈−2 400 000</td>
<td>km s⁻¹</td>
<td>km s⁻¹</td>
</tr>
<tr>
<td>SOPHIE</td>
<td></td>
<td></td>
</tr>
<tr>
<td>54463.4939000</td>
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</tr>
<tr>
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<tr>
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<tr>
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</tr>
<tr>
<td>54544.3463300</td>
<td>48.925</td>
<td>0.026</td>
</tr>
<tr>
<td>HARPS</td>
<td></td>
<td></td>
</tr>
<tr>
<td>54548.583775</td>
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<td>0.014</td>
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<td>54550.577783</td>
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<tr>
<td>54771.850953</td>
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<td>0.010</td>
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<td>54773.847921</td>
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<td>0.008</td>
</tr>
<tr>
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<td>48.929</td>
<td>0.009</td>
</tr>
<tr>
<td>54805.748602</td>
<td>48.852</td>
<td>0.012</td>
</tr>
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</table>
including the Rossiter effect. Red: SOPHIE, green: HARPS.

**Fig. 2.** Radial velocity measurements and Keplerian fit to the data including the Rossiter effect. Red: SOPHIE, green: HARPS.

H. Rauer et al.: CoRoT-5b 283

(2003). Both sets of data (SOPHIE and HARPS) have been processed as in Bouchy et al. (2008). Radial velocities (RV) were computed by weighted cross-correlation (Baranne et al. 1996; Pepe et al. 2005) with a numerical G2-spectral template excluding spectral orders below 4200 Å. Radial velocity values are listed in Table 1 and plotted in Fig. 2.

We analyzed the cross-correlation function computed from the HARPS spectra using the line-bisector technique according to the description in Queloz et al. (2001) to detect possible spectral distortions caused by a faint background eclipsing binary mimicking a small RV amplitude signal. No correlation between the RV data and the bisector span was found at the level of the uncertainty on the data (Fig. 3).

The stability of the bisector, combined both with the amplitude of the radial velocity and the accuracy of transit of the lightcurve, is enough to discard an alternate background eclipsing binary scenario. In the case of a hypothetical background eclipsing binary, obtaining a sine-shaped radial-velocity signal would require a superimposed spectrum moving with the same systemic velocity as the brightest component, and on an RV range corresponding to the sum of the width of both CCF line profiles. This prerequisite constrains both on the mass of the potential eclipsing component and its companion. The example of HD41004 provides us with an interesting benchmark (Santos et al. 2002). This system was detected with a similar radial velocity amplitude but with a strong bisector correlation, and could be explained by a superimposed spectrum with 3% flux of the bright star. If one scales down this result to CoRoT-5, which has no bisector correlation, one finds that the contrast ratio between the brightest star and the hypothetical eclipsing binary is such that the eclipse must be very deep and the radius of the eclipsing stars much smaller than CoRoT-5. Considering the quality of the CoRoT lightcurve such a binary scenario does not match the transit ingress and egress timing and the detailed shape of the curve.

**5. Properties of the central star**

We determined the fundamental parameters of the host star carrying out a spectral analysis of the set of HARPS spectra acquired for radial velocity measurements. The individual spectra were reduced with the HARPS standard pipeline. The extracted spectra were corrected for cosmics impacts, for the Earth and the stars velocity, and then corrected for the blaze function and normalized, order by order, to increase the signal-to-noise (S/N). The S/N level in the continuum is around 40 in the range 5000–6500 Å and it decreases to 15 towards the blue at 4000 Å.

Spectroscopic observations of the central star have also been performed in January 2008 with the AAOmega multi-object facility at the Anglo-Australian Observatory. By comparing the low-resolution (R = 1300) AAOmega spectrum of the target with a grid of stellar templates, as described in Frasca et al. (2003) and Gandolfi et al. (2008), we derived the spectral type and luminosity class of the star (F9 V).

As for the previous planet host stars, we used different methods to derive Corot-5 atmospheric parameters: line profile fitting with the SME (Valenti & Piskunov 1996) and the VWA packages (Brunt et al. 2002, 2008). We find general agreement and here we quote the results from VWA. The star has a very low projected rotational velocity, \( v \sin i = 1 \pm 1 \text{ km s}^{-1} \). More than 600 mostly non-blended lines were selected for analysis in the wavelength range 3990–6810 Å. VWA uses atmosphere models from the grid by Heiter et al. (2002) and atomic parameters from the VALD database (Kupka et al. 1999). The abundance determined for each line is computed relative to the result for the same line in the solar spectrum from Hinkle et al. (2000), following the approach of Bruntt et al. (2008). The results for CoRoT-5 are shown in Table 2. Using these parameters for the atmospheric model, we determined the abundances of 21 individual elements. The uncertainty on the abundances includes a contribution of 0.04 dex due to the uncertainty on the fundamental parameters. The abundance pattern is shown in Fig. 4. The overall metallicity is found as the mean abundance of the elements with at least 20 lines (Si, Ca, Ti, Cr, Fe, Ni) giving \( [M/H] = -0.25 \pm 0.04 \). We did not include Mn, as this has a significantly lower abundance. The metallicity and the 1-σ error bar is indicated by the horizontal bar in Fig. 4. There is no evidence of the host star being chemically peculiar, except Mn.

The fundamental parameters of the parent star, its mass and radius were subsequently derived using stellar evolutionary tracks as presented in Deleuil et al. (2008) plotted in a \( M^{1/3}/R - T_{\text{eff}} \) HR diagram. The stellar density parameter was...
derived from the lightcurve fitting (see Sect. 7). We determined the mass and radius of the star to: $M_{\text{star}} = 1.00 \pm 0.02\,M_{\odot}$ and $R_{\text{star}} = 1.186 \pm 0.04\,R_{\odot}$. As a final check, we calculated the corresponding surface gravity $\log g = 4.311 \pm 0.033$ while the spectroscopic value is $4.19 \pm 0.03$. These two values of $\log g$ are comparable with each other at the $3\sigma$ level. Based on our photometric analysis, we estimate the age of the star to be $5.5–8.3$ Gyr. The spectra show no sign of Ca II emission or of a strong Li I absorption line, which is consistent with a relatively evolved star.

Table 2. Parameters of the parent star CoRoT-5.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
<th>Source</th>
</tr>
</thead>
<tbody>
<tr>
<td>RA</td>
<td>06h45m07s</td>
<td>Exo-Dat</td>
</tr>
<tr>
<td>Dec</td>
<td>00°48'55''</td>
<td>Exo-Dat</td>
</tr>
<tr>
<td>Epoch</td>
<td>2000.0</td>
<td></td>
</tr>
<tr>
<td>Type</td>
<td>F9V</td>
<td>Exo-Dat, AAOmega</td>
</tr>
<tr>
<td>V</td>
<td>14.0</td>
<td>Exo-Dat</td>
</tr>
<tr>
<td>GSC2.3 ID</td>
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<td></td>
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<tr>
<td>2MASS ID</td>
<td>06450653+0048548</td>
<td></td>
</tr>
<tr>
<td>$v\sin i$ [km s$^{-1}$]</td>
<td>$1 \pm 1$</td>
<td>VWA</td>
</tr>
<tr>
<td>$\xi_t$ [km s$^{-1}$]</td>
<td>$0.91 \pm 0.09$</td>
<td>VWA</td>
</tr>
<tr>
<td>Teff [K]</td>
<td>6100 $\pm$ 65</td>
<td>VWA</td>
</tr>
<tr>
<td>log g</td>
<td>4.19 $\pm$ 0.03</td>
<td>VWA</td>
</tr>
<tr>
<td>$[M/H]$</td>
<td>$-0.25 \pm 0.06$</td>
<td>VWA</td>
</tr>
<tr>
<td>$M_{\text{star}} [M_{\odot}]$</td>
<td>$1.00 \pm 0.02$</td>
<td>Evolut. tracks</td>
</tr>
<tr>
<td>$R_{\text{star}} [R_{\odot}]$</td>
<td>$1.186 \pm 0.04$</td>
<td>Evolut. tracks</td>
</tr>
<tr>
<td>$M^{(1/3)}<em>{\text{star}} / R^{(1/3)}</em>{\text{star}}$</td>
<td>$0.843 \pm 0.024$</td>
<td>lightcurve</td>
</tr>
<tr>
<td>age [Gyr]</td>
<td>$5.5–8.3$</td>
<td>+Evolut. tracks</td>
</tr>
</tbody>
</table>

Fig. 4. Stellar abundances of CoRoT-5. Abundances found from neutral lines are marked by circles, for ionized lines box symbols are used.

Fig. 5. The O–C diagram of the CoRoT-5b system. No clear period variation can be seen.

6. Period determination and transit timing variations

In total, 27 individual transit events are clearly seen, separated by an orbital period of about 4.03 days. One event was lost in a data gap.

First, we estimated the mid-times of each transit by applying the so-called Kwee-van Woerden method (Kwee & van Woerden 1956). This method mirrors the lightcurve around a pre-selected time-point, $T$, computes the differences of original and mirrored lightcurves and then searches for an optimum $T$. The O–C diagram of the system was constructed, based on the resulting transit times and an initial guess of the period. A linear fit of this diagram yielded an improved estimate of the period. This period value was then refined with the following procedure. The lightcurve was phase-folded using this previously determined period and then averaged. The size of the bin used was 0.001 in phase (or to 5.81 min, using the final period). Then, this lightcurve was fitted (see the next section) by a theoretical transit lightcurve. The transit mid-times were then determined again by cross-correlating the observed and the theoretical lightcurve. This resulted in more precise mid-times of the transit and a new O–C curve. Another linear fit to this O–C diagram yielded a better period value, and the whole procedure was repeated. The final O–C diagram can be seen in Fig. 5. The resulting ephemeris is given in Table 4.

There is no obvious period variation present in the O–C diagram. The first part of the lightcurve was obtained with the 512 s sampling rate, so the first seven minima typically consist of only 20 data points. Thus, they have larger scatter and uncertainties. The next twenty minima were obtained with the high sampling rate (32 s) and typically consist of a few hundred data points, leading to much higher accuracy. If one takes only these high-resolution minima into account, the constancy of the period is clearer. However, we cannot exclude that small period variations are present in the system. The upper limit of such a period variation was estimated by a quadratic fit to the data, which showed that it should be less than 0.42 s/cycle.

7. Analysis of parameters of CoRoT-5b

The final phase-folded lightcurve of the transit event is seen in Fig. 6. The transit signal shows a depth of about 1.4% and lasts for about 2.7 h. We derived the planetary parameters by fitting simultaneously the lightcurve of CoRoT-5 with the SOPHIE and HARPS radial velocities. A planetary model on a Keplerian orbit in the formalism of Giménez (2006a) and Giménez (2006b) was fitted to the data using a Markov Chain Monte-Carlo (MCMC)
Table 4. The derived planet parameters.

<table>
<thead>
<tr>
<th>Derived physical parameters</th>
<th>Value</th>
<th>Units</th>
</tr>
</thead>
<tbody>
<tr>
<td>Transit epoch $T_0$</td>
<td>2454400.19885 ± 0.0002</td>
<td>HJD</td>
</tr>
<tr>
<td>Orbital period $P$</td>
<td>4.0378962 ± 0.0000019</td>
<td>days</td>
</tr>
<tr>
<td>Orbital semi-major axis $a$</td>
<td>0.049474 ± 0.00026</td>
<td>AU</td>
</tr>
<tr>
<td>Orbital inclination $i$</td>
<td>85.83 ± 0.99</td>
<td>degrees</td>
</tr>
<tr>
<td>Orbital eccentricity $e$</td>
<td>0.096</td>
<td></td>
</tr>
<tr>
<td>Argument of periastron $\omega$</td>
<td>−2.24 ± 0.04</td>
<td>rad</td>
</tr>
<tr>
<td>Planet radius $R_p$</td>
<td>1.388 ± 0.046</td>
<td>$R_J$</td>
</tr>
<tr>
<td>Planet mass $M_p$</td>
<td>0.467 ± 0.024</td>
<td>$M_J$</td>
</tr>
<tr>
<td>Mean planet density $\rho_p$</td>
<td>0.217 ± 0.031</td>
<td>g cm$^{-3}$</td>
</tr>
<tr>
<td>Planetary surface gravity $\log g_p$</td>
<td>7.77 ± 0.14</td>
<td>cgs</td>
</tr>
<tr>
<td>Zero albedo equilibrium temperature $T_{eq}$</td>
<td>1438 ± 39</td>
<td>K</td>
</tr>
</tbody>
</table>

Fig. 6. Top: phase-folded lightcurve of CoRoT-5b. Bottom: residuals of fitted transit curve.

The derived planet parameters are summarized in Table 4. The major uncertainties on the planet are, as usual, introduced mainly from the uncertainty of the stellar parameters.

code described in Triaud et al. (in prep.) but using $e$, $\cos \omega$ and $e \sin \omega$ instead of $e$ and $\omega$ as free parameters for better error estimation. In the fit a quadratic limb-darkening law was assumed at $u_1 = 0.616$ and $u_2 = 0$. In the initial burn-in phase of the MCMC adjustment, 15,000 steps were chosen to allow the fit to converge. A further 50,000 steps were used to derive the best parameters and their errors. In the fit, there are eight fitted parameters plus two $\gamma$ velocities and a normalization factor, totalling 11 free parameters. In addition, the fit assumed the presence of a Rossiter-McLaughlin effect with the two fixed parameters $v \sin i = 1.0$ km s$^{-1}$ and $\lambda = 0$ ($\lambda$: angle between stellar rotation axis and normal vector of the orbital plane). A Bayesian penalty is added to the $\chi^2$ creating a prior for $M_* = 0.99 ± 0.02$. The fit to the rv measurements is shown in Fig. 2, and the derived fitting parameters are shown in Table 3.

In addition, a model transit curve (Mandel & Agol 2002) was fitted to the photometric phase folded transit curve separately. The parameters fitted are the center of transit, the planet radius expressed in stellar radii, the semi-major axis in stellar radii and the orbital inclination. In this fit the limb-darkening coefficients ($u_1$ and $u_2$) were free parameters, assuming a quadratic limb-darkening law. The fitting method follows a Metropolis-Hastings algorithm, which is a kind of Markov Chain Monte-Carlo procedure. The fitting procedure was performed ten times with different starting values to find the global minimum in $\chi^2$. The errors of the fit were estimated from the standard deviations of the points in the chain. In addition to the transit curve, a third light component is included as a free parameter in the fit. In this way, we could check whether another contaminant is present, which remained unresolved in the photometric follow-up. However, no such additional source of light was found. The transformation between contamination factor $c$ and the third light $l_3$ is $c = l_3/(1 - l_3)$. We had $c = 0.005 ± 0.024$. Since we already removed the known contaminant factor from the lightcurve (see Sect. 2), we could therefore conclude that no further observable contaminant is present in the lightcurve of CoRoT-5. The planet parameters derived from this fit agree with the simultaneous fitting within the error bars, so we do not report them again here.

The resulting planetary parameters based on the MCMC approach with fixed limb-darkening coefficients and without any third light are summarized in Table 4. The major uncertainties on the planet are, as usual, introduced mainly from the uncertainty of the stellar parameters.

8. Summary

We report the discovery of a “hot-Jupiter-type” planet, CoRoT-5b, orbiting a type F9V star of 14.0 mag. The planet mass and radius were derived to $0.467^{+0.024}_{-0.024} M_J$ and $1.388^{+0.036}_{-0.037} R_J$, respectively. It orbits its central star at $0.0002$ AU orbital distance. The determined eccentricity is low (see Table 4), but further radial velocity measurements would be needed for a more accurate determination.

CoRoT-5b has a density of $0.217^{+0.031}_{-0.025}$ g cm$^{-3}$, similar to the planets WASP-12b and WASP-15b (Hebb et al. 2009; West et al. 2009), implying that it belongs to the planets with the lowest mean density found so far. As such, it is found to be larger by 20% than standard evolution models (Guillot et al. 2006) would predict. Standard recipes that account for missing physics (kinetic energy transport or increased opacities) can explain this large size, and predict that the planet is mostly made of hydrogen-helium, with at most 28% of heavy elements (maximum value obtained in the kinetic energy model, assuming 0.5% of the incoming energy is dissipated at the planet center). Thus, CoRoT-5b supports the proposed link between the metallicity of planets and of their host star.

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