A weak lensing study of the Coma cluster

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ABSTRACT

Context. Due to observational constraints, dark matter determinations in nearby clusters based on weak lensing are still extremely rare, in spite of their importance for the determination of cluster properties independent of other methods.

Aims. We present a weak lensing study of the Coma cluster (redshift ~0.024) based on deep images obtained at the CFHT.

Methods. After obtaining photometric redshifts for the galaxies in our field based on deep images in the u’ (1 × 1 deg²), and in the B, V, R and I bands (42 × 52'), allowing us to eliminate foreground galaxies, we apply weak lensing calculations on shape measurements performed in the u’ image.

Results. We derive a map of the mass distribution in Coma, as well as the radial shear profile, and the mass and concentration parameters at various radii. We obtain \( M_{200c} = 5.1^{+1.3}_{-2.1} \times 10^{14} M_{\odot} \) and \( c_{200} = 5.0^{+1.3}_{-2.1} \), in good agreement with previous measurements.

Conclusions. With deep wide field images it is now possible to analyze nearby clusters with weak lensing techniques, thus opening a broad new field of investigation.

Key words. gravitational lensing – galaxies: clusters: individual: Coma

1. Introduction

Although weak lensing studies are of great importance in cosmology, in particular concerning various aspects of the analysis of large scale structures in the universe (e.g. Mellier 1999; Bartelmann & Schneider 2001; Schneider 2006), dark matter determinations in clusters based on weak lensing have been limited to relatively distant objects, usually with redshifts larger than about 0.15. This is mainly due to the fact that the cluster lensing strength is maximized if the angular diameter distance to the foreground cluster is about half that of the background source galaxies, and to the requirement of large field images, in order to cover the full cluster in a single exposure. Very few low redshift clusters have been analyzed until now in detail through weak lensing: Abell 3667 at redshift 0.055 (Joffre et al. 2000), and clusters observed in the shallow Sloan Digital Sky Survey (Sheldon et al. 2001), including Coma (Kubo et al. 2007). We present here a weak lensing study of the Coma cluster at z_Coma = 0.024 based on deep exposures in several bands obtained with the CFHT12k and MegaPrime/Megacam large cameras at the Canada France Hawaii Telescope. The J2000 location of the brightest cluster galaxy NGC 4874 is RA = 12:59:35.67, Dec = +27:57:33.7. In the following we assume the “concordance model” cosmological background with \( H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1} \), \( \Omega_m = 0.3 \) and \( \Omega_{\Lambda} = 0.7 \), for which 1’’ on the sky corresponds to a transverse scale of 0.79 kpc. All magnitudes are in the AB system.

2. Data analysis

2.1. Cluster membership and background sources

Multiband imaging for this cluster was obtained at CFHT, both in the inner parts made of two CFH12k camera pointings in BVRI bands resulting in about 0.54 deg² coverage and with Megacam (1 × 1 deg² fov) in the u’ band (see Adami et al. 2006, 2007, 2008, for a detailed description of these data).

We were able to estimate photometric redshifts of all sources down to \( R \leq 23 \) in the area covered by the CFH12k images (Adami et al. 2008). This allowed us to separate the population of faint background sources from nearby objects with \( z_{\text{phot}} \leq 0.4 \) and to eliminate the bulk of the cluster members that could bias the weak lensing study (e.g. Medezinski et al. 2007). Adami et al. (2008) used the same technique to discriminate between \( z < 0.2 \) and \( z > 0.2 \) galaxies. Here, we push the limit to \( z_{\text{phot}} = 0.4 \) to increase the number of background sources although a more conservative cut at 0.2 would not produce noticeable changes given the small number of added objects.

This photometric redshift separation is possible only in the innermost 0.54 deg². However we found that none of our results below could be significantly biased due to the extreme...
prevalence of distant sources in the faint end of the magnitude counts (a factor \(\sim 10\)–20) despite some incompleteness for \(u' \gtrsim 24.5\). We thus apply a magnitude selection \(23.5 \leq u' \leq 25\) everywhere in the field but we complement the selection by removing the few objects in the inner regions that are known to be at \(\delta_{\text{phys}} \leq 0.04\).

A clear advantage of lensing by nearby deflectors is that the geometric distance factor \(D_{\text{LS}}/D_S\) (e.g. Bartelmann & Schneider 2001), which is the ratio of angular distances between deflector and source \(D_{\text{LS}}\) and to the source \(D_S\), is almost constant and close to unity for most of the faint background sources. Given the low redshift of Coma, even a source at \(z \sim 0.2\) will experience lensing effects with \(D_{\text{LS}}/D_S \approx 0.87\). In order to check this small redshift dependance, we used the CFHTLS-deep photometric redshifts (Ilbert et al. 2006) measured in four independent 1 deg\(^2\) fields to infer the redshift distribution of sources having \(23.5 \leq u' \leq 25\). The mean and median photometric redshifts are 0.87 and 0.92 respectively and the sample contains a negligible amount of low redshift \((z < 0.04)\) galaxies. We estimate the mean \(\bar{\mu} = \left(D_{\text{LS}}/D_S\right)\) averaged over the whole population of sources to be \(\bar{\mu} = 0.962\). This is the value we will consider below for further weak lensing mass calibrations.

### 2.2. Shear measurement

In Adami et al. (2009), we presented measurements of the systematic orientation of faint Coma cluster galaxies. The small size of these galaxies required a careful deconvolution of the Point Spread Function (PSF) smearing effect which was done with the same procedure we use here. It is based on the KSB technique (Kaiser et al. 1995) and the pipeline has been applied to ground-based imaging from CFH12k and Megacam cameras (e.g. Gavazzi et al. 2004; Gavazzi & Soucail 2007). We used the \(u'\) channel to measure shapes because it is the one with largest spatial coverage, and also the most homogeneous since the images in the other bands are built from two adjacent pointings. The median \(u'\) band seeing FWHM is 1.02 arcsec. Despite this relatively poor value, we were able to successfully correct for PSF smearing.

The quality of PSF anisotropy corrections can be seen in the left panels of Fig. 1 in Adami et al. (2009). The good overall image quality is reflected in the raw and corrected ellipticities of stars. The low scatter in the corrected ellipticities, 0.0028 rms, is also a good assessment of the controlled systematics.

Based on mock images designed to simulate Megacam data for the CFHTLS survey and kindly made available to us by Rowe, we were able to assess the reliability of our shear measurements, especially for the isotropic part of the PSF smearing effect. These simulations are similar to STEP2 simulations (Massey et al. 2007) as they capture the complex shape of sources by the use of shapelet models. Following notations of the STEP shear measurement project (Heymans et al. 2006; Massey et al. 2007), residual additive calibration errors are found to be \(c \leq 0.003\) and multiplicative \(m \leq 0.05\), which is significantly smaller than the statistical errors present in our data (see below).

Because redder bands are usually used for weak lensing studies, we checked that PSF-corrected ellipticities of central objects measured in both Megacam-\(u'\) and CFH12k-I bands are consistent. This is found to be the case with great accuracy as shown in Fig. 1. We checked that the ellipticity difference \(e_{r'} - e_U\) (for both components 1 and 2 taken together or considered individually) does not correlate with either \(e_U\), \(u'\) magnitude, I magnitude or \(r_{h,0'}\) the \(u'\)-band half light radius.

### 2.3. Convergence map

From the source catalogue, we can infer the shear field \(\gamma(\theta)\) and deduce the associated convergence field \(\kappa(\theta)\) which is the projected surface mass density expressed in units of the critical density \(\Sigma_{\text{cr}} = c^2(4\pi G)^{-1}D_S/(D_D D_S)\). Here \(D_D/D_S\) is replaced by the source population average inverse 1/\(\bar{\mu}\) which leads to a critical density \(\Sigma_{\text{cr}} = 1.69 \times 10^{-10}\) \(h_7^2\) \(M_\odot/\text{kpc}^2\).

Shear and convergence are related by:

\[
\kappa(\theta) = \frac{2}{\bar{\mu}} \int H \left[ \gamma(\theta - \theta_c) - \gamma(\theta) \right] d^2\theta. \quad (1)
\]

\(K(\theta) = \frac{1}{\bar{\mu}} \Sigma_{\text{cr}} \exp(\theta / \theta_s)\) is a complex convolution kernel (Kaiser & Squires 1993). The shear field is smoothed with a Gaussian \(G(\theta) \propto \exp(-\theta^2 / \theta_s^2)\) with \(\theta_s = 170'' \approx 80 h_7^{-1}\) kpc. The convergence field is smoothed by the same filter. The convergence map presents correlated noise properties (van Waerbeke 2000) characterised by \(\eta_{\text{mg}} = 11\) arcmin\(^2\), the number density of background sources. By bootstraping the orientation of background sources, we estimated this noise level to be \(\sigma_N = 0.0134\) with little variation over the field (except near edges and masked bright stars).

The top panel of Fig. 2 shows the reconstructed convergence map 1 deg\(^2\) field of view with contours showing a \(> 5\sigma\) significance central peak along with various substructures. The high physical resolution of \(\sim 80 h_7^{-1}\) kpc that is made possible by the vicinity of Coma somewhat balances the low associated lensing efficiency. The position of the brightest cluster galaxies with \(r \leq 14\) is shown as red crosses. For comparison we also show in the lower panel of Fig. 2 the luminosity-weighted \(r\)-band luminosity distribution coming from bright \(r \leq 19\) cluster members that lie on the Red Sequence defined by the relation \(0.72 \leq (g - r) + (r - 14) \times 0.035 \leq 0.87\). The convergence contours are overlaid. For this latter study we used the more extended SDSS DR7 photometry, which is better suited for investigation of the bright end of the luminosity function\(^1\). This panel also shows contours of X-ray emissivity in excess of a smooth \(\beta\)-model profile (Neumann et al. 2003).

\(^1\) http://cas.sdss.org/astro/en/tools/
3. Radial mass profile

We now investigate the azimuthally-averaged tangential shear profile $\gamma_t(R)$ which is simply related to the azimuthally-averaged projected surface mass density profile $\Sigma(R)$ by the relation

$$\Sigma_{\text{crit}} \gamma_t(R) \equiv \Delta \Sigma(R) = M(< R)/(\pi R^2) - \Sigma(R),$$

(2)

where we have defined the frequently used rescaled shear $\Delta \Sigma$ and $M(< R)$ is the projected mass enclosed by radius $R$.

We measure the radial shear profile average in circular annuli centered on the peak of the convergence map ($\alpha = 12^h:59^m:39.007$, $\delta = +27^\circ:57^\prime:55.93$) which is only 40" east of NGC 4874, i.e. in the direction of NGC 4889, which are respectively the brightest and second brightest member galaxies. Figure 3 shows the radial shear profile $\gamma_t(R)$ out to the edge of the Megacam field of view (~1 Mpc). The bottom panel of this figure shows the same profile once galaxies are rotated by $45^\circ$, which is the curl or B-mode component of the ellipticity field. In the absence of systematic PSF correction residuals, this rotated shear profile should be consistent with zero at all scales. This is what we observe.

We attempted to fit a radial shear profile as predicted by the NFW mass density profile of the form

$$\rho(r) = \rho_s (r/r_s)^{-1} (1 + r/r_s)^{-2}$$

(3)

coming from cosmological simulations (Navarro et al. 1997). The corresponding lensing quantities were derived by Bartelmann (1996). The two relevant quantities that we fit are the mass $M_{200c}$ enclosed in the radius $r_{200c}$ in which the mean density is 200 times the critical density $\rho_{\text{crit}}$ and the concentration parameter $c_{200c} = r_{200c}/r_s$. This implies that $\rho_s = \rho_{\text{crit}} 200^3/[\ln(1 + c) - c/(1 + c)]$. Figure 4 shows the constraints we obtain on these two parameters. Marginal distributions yield the following constraints:

$$M_{200c} = 5.1_{-1.3}^{+1.3} \times 10^{14} \, h^{-1}_7 M_\odot$$

and

$$c_{200c} = 1.9_{-0.3}^{+0.8}$$

(2) Assuming a flat uniform prior in log $M_{200c}$ between 13 and 16 and on $c_{200c}$ between 0.01 and 20 to infer respectively the marginal PDFs $p(c)$ and $p(M_{200c})$. 

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**Fig. 2.** Top panel: convergence map for the Coma cluster. Green contours represent signal-to-noise ratios of 1, 2, ..., 5, corresponding to $\kappa = 0.01, 0.02, ..., 0.05$. Red crosses represent the bright cluster members lying on the red sequence. Bottom panel: gray-scale view of the luminosity distribution of cluster red sequence members with the overlaid contours in green. The blue dotted contours show the excess X-ray emission over a smooth $\beta$-model X-ray emissivity map (Neumann et al. 2003).

**Fig. 3.** Top panel: radial shear profile measured in Coma as a function of projected distance. The best-fit NFW profile is overlaid. Bottom panel: B-mode shear profile showing the negligible residuals in the rotated shear component.

**Fig. 4.** $M_{200c}$ and $c_{200c}$ parameters constraints of the NFW density profile. The $M-c$ relation predicted by simulation and WMAP5 cosmology (with its 1$\sigma$ scatter) are shown in red.
Based on a deep mass-concentration relation ($M_{200c}/10^{14} h^{-1}_{70} M_\odot$) calculated from the virial mass $M_{200c}$ and the corresponding concentration $c_{vir}$, which is better suited for the assumed cosmology. 

The south-west extension of the convergence map is more puzzling. It does not match the light distribution. The groups G8 and G9 from Adami et al. (2005) explain part of this extension, but are not luminous enough to fully account for this mass concentration. Another explanation would be massive low redshift structures (in order to still efficiently induce a weak lensing signal and to be spread largely enough over the sky) but that would not belong to Coma. The SDSS data used by Gott et al. (2005) exhibit such a z $\sim$ 0.08 “Great Wall”, probably not luminous enough though to explain the south west mass concentration. 

Moreover, the magnitude limit of the SDSS spectroscopy is not very deep, leaving room for other possible large scale structures not appearing in the SDSS sample. This will have to be confirmed by a more detailed knowledge of the immediate background of the Coma cluster (Adami et al., in preparation).

4. Discussion and conclusion

We have performed a weak lensing analysis of the Coma cluster, based on a deep exposure and on images in other bands which allowed to derive photometric redshifts in order to remove as many foreground galaxies as possible from the background population. We thus obtain a map of the mass distribution in Coma, the radial shear profile, as well as the mass and concentration parameter.

The peak of the convergence map is nearly coincident with the X-ray center (Neumann et al. 2003) or the brightest galaxy NGC 4874. The difference is $\sim 27 h^{-1}_{70}$ kpc below the spatial resolution of the convergence map. We can also notice the correspondence between the ($\alpha \sim 194.4, \delta \sim 27.82$) 2$\sigma$ detection and the G12/G14 groups of Adami et al. (2005) and between the ($\alpha \sim 194.5, \delta \sim 28.03$) 2$\sigma$ detection and the northern part of the west X-ray substructure (Neumann et al. 2003). Finally, the convergence map shows a 4$\alpha$ ($\alpha \sim 195, \delta \sim 28$) extension overlapping the position of NGC 4889, confirming the existence of a mass concentration around this galaxy. Other substructures visible in X-rays and in the optical (e.g. the NGC 4911 group) do not show up in the convergence map, indicating that we are dealing with low mass systems.

Table 1. Summary of mass estimates with and without priors on the mass-concentration relation (Macciò et al. 2008).

| Parameter | No prior | Prior $p(c|\mathcal{M})$ |
|-----------|----------|--------------------------|
| $M_{200c}/10^{14} h^{-1}_{70} M_\odot$ | 5.1$\pm$4.3 | 9.7$^{+4.1}_{-3.5}$ |
| $c_{200c}$ | 5.0$^{+3.2}_{-2.5}$ | 3.5$^{+1.1}_{-0.9}$ |
| $r_{200c}/h_{70}$ Mpc | 1.8$^{+0.6}_{-0.3}$ | 2.2$^{+0.3}_{-0.2}$ |
| $M_{200c}/10^{14} h^{-1}_{70} M_\odot$ | 6.1$^{+1.1}_{-0.3}$ | 11.1$^{+6.7}_{-4.1}$ |
| $c_{vir}$ | 6.7$^{+4.1}_{-3.3}$ | 4.9$^{+1.7}_{-1.4}$ |
| $r_{vir}/h_{70}$ Mpc | 2.5$^{+0.8}_{-0.5}$ | 3.6$^{+1.1}_{-0.7}$ |

References


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