

Hanle effect in the CN violet system with LTE modeling^{*}

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ABSTRACT

Context. Weak entangled magnetic fields with mixed polarity occupy the main part of the quiet Sun. The Zeeman effect diagnostics fails to measure such fields because of cancellation in circular polarization. However, the Hanle effect diagnostics, accessible through the second solar spectrum, provides us with a very sensitive tool for studying the distribution of weak magnetic fields on the Sun.

Aims. Molecular lines are very strong and even dominate in some regions of the second solar spectrum. The CN $B^2\Sigma - X^2\Sigma$ system is one of the richest and most promising systems for molecular diagnostics and well suited for the application of the differential Hanle effect method. The aim is to interpret observations of the CN $B^2\Sigma - X^2\Sigma$ system using the Hanle effect and to obtain an estimation of the magnetic field strength.

Methods. We assume that the CN molecular layer is situated above the region where the continuum radiation is formed and employ the single-scattering approximation. Together with the Hanle effect theory this provides us with a model that can diagnose turbulent magnetic fields.

Results. We have succeeded in fitting modeled CN lines in several regions of the second solar spectrum to observations and obtained a magnetic field strength in the range from 10–30 G in the upper solar photosphere depending on the considered lines.

Key words. line: formation – Sun: magnetic fields – molecular processes – polarization – radiative transfer – scattering

1. Introduction

The understanding of the solar surface magnetic field plays a key role in modern solar physics since the magnetic field is connected to and even controls most phenomena of solar activity. The solar magnetic field is structured in a fractal-like pattern with a significant degree of self-similarity over many scales (Stenflo & Holzreuter 2003), while the field strength covers many orders of magnitude reaching up to a few kG. Observationally, we can resolve only a small fraction of the magnetic structures, which can be accurately measured via the Zeeman effect diagnostics. However, about 99% of the photospheric volume is filled by unresolved, entangled magnetic fields with mixed polarity, not accessible with the Zeeman effect due to cancellation of the polarimetric signal, that might contribute a substantial amount of magnetic energy to the solar photosphere (Trujillo Bueno et al. 2004).

These hidden magnetic fields can be accessed with the Hanle effect (Stenflo 1982), which modifies the properties of coherent scattering processes. The Hanle effect can thus only be employed through the spectrum formed by the coherent scattering, which is called the “second solar spectrum” (Ivanov 1991) due to its richness and significance. The Hanle effect usually leads to a reduction and rotation of line polarization as compared to the non-magnetic case. The net-influence of rotation cancels out in a turbulent magnetic field, but the depolarization remains visible for most geometries even in a turbulent magnetic field. Moreover, the Hanle effect is extremely sensitive to the weak field. Therefore, the second solar spectrum provides us with a

very sensitive tool for studying the distribution of weak magnetic fields on the Sun and gives the possibility to measure spatially unresolved mixed polarity magnetic fields.

Molecular lines are particularly valuable for magnetic field diagnostics thanks to their broad range of magnetic sensitivities within narrow spectral regions. This allows us to employ the differential Hanle effect (Stenflo et al. 1998), which greatly reduces the model dependence of deduced magnetic field strengths, and has led to the first unambiguous detection of molecular Hanle effect in the second solar spectrum (Fluri & Berdyugina 2004). Recent studies of the molecular Hanle effect have yielded field strengths of the order of 10 G for the quiet solar photosphere, assuming a single-valued, isotropic, turbulent magnetic field (Faubert & Arnaud 2003; Berdyugina & Fluri 2004; Trujillo Bueno et al. 2004), although larger field strengths are possible depending on the assumed collision rate, in particular when interpreting MgH lines (Asensio Ramos & Trujillo Bueno 2005; Bommier et al. 2006).

Diagnostics with molecular Hanle effect has concentrated so far only on two molecules, C₂ and MgH. It is very important to extend the technique to other molecules because in general different molecules form at different heights due to their large temperature sensitivity and can thus sample other layers of the solar atmosphere. Combined with atomic Hanle effect, which is well established since the work by Faubert-Scholl (1992), only such a multitude of lines will allow us to constrain the 3-dimensional structure of the unresolved solar magnetic field.

In the present paper we apply the Hanle effect to a new molecule, namely the CN violet system, for which we have recently developed a complete theory of the Hanle effect taking into account the Paschen-Back effect and interference between

^{*} Tables 2–4 are only available in electronic form at <http://www.aanda.org>

doublet components (Shapiro et al. 2007). After describing our radiative transfer method we fit our model to observations in three spectral regions of the second solar spectrum, which enables us to determine the magnetic field strength in the quiet solar photosphere.

2. Radiative transfer model

In this section we introduce our method for solving the radiative transfer problem. It is based on the modeling strategy developed by Faurobert & Arnaud (2003), who have calculated scattering polarization at the line centers of C₂ and MgH transitions. Here, we expand their method to obtain wavelength dependent spectra not restricted to line centers.

The radiative transfer model is optimized for relatively weak lines whose scattering polarization is calculated within the framework of a plane-parallel atmosphere using the single-scattering approximation. The CN lines are assumed to form in a homogeneous atmospheric layer situated above the continuum formation region. We fully account for line blending and for depolarization of the continuum polarization by line absorption and scattering. In this respect, the current model is more complex than the one employed by Berdyugina & Fluri (2004) for very weak C₂ lines. The CN bands that are modeled in this paper include stronger lines and require a consistent treatment of continuum depolarization, although we increase at the same time the model dependence of the magnetic field diagnostics, because we rely on a proper computation of the continuum polarization. Nonetheless, the model remains simple and ideally suited for differential Hanle effect diagnostics.

First we introduce the necessary notation and the source function. Then we discuss in detail the two-layer approximation, the free model parameters, and the fitting procedure for comparing the model with observations.

2.1. Source function

The total, monochromatic optical depth is given by

$$d\tau(\nu) = - \left(\alpha_c + \sum_{i=1}^n \alpha_\ell^i \varphi^i(\nu) \right) dz, \quad (1)$$

where $\alpha_c = k_c + \sigma_c$ is the continuum opacity (with k_c and σ_c the continuum absorption and scattering coefficients), $\alpha_\ell^i = k_\ell^i + \sigma_\ell^i$ is the integrated line opacity in the i th line (with k_ℓ^i and σ_ℓ^i the line absorption and scattering coefficients), and $\varphi^i(\nu)$ is the normalized Voigt profile function. The opacities in molecular lines depend on the main line parameters and on the temperature of the CN layer T_{eff}

$$\alpha_\ell^i \propto g_\ell^i f_\ell^i \exp(-E_\ell^i/kT_{\text{eff}}), \quad (2)$$

where g_ℓ^i and E_ℓ^i are the statistical weight and energy, respectively, of the lower level, and f_ℓ^i is the absorption oscillator strength, which can be calculated with the pure Hund's case (b) expressions (Schadee 1978; Whiting & Nicholls 1974) because the orbital angular momenta for both the upper and lower states are equal to zero in the CN violet system.

Now we can write the radiative transfer equation in the usual form

$$\mu \frac{\partial \mathbf{I}}{\partial \tau} = \mathbf{I} - \mathbf{S}, \quad (3)$$

where $\mathbf{S} \equiv \mathbf{S}(\tau, \nu, \mu)$ is the total source function

$$\mathbf{S} = \frac{k_c \mathbf{B}_{\text{th}} + \sum_{i=1}^n k_\ell^i \varphi^i(\nu) \mathbf{B}_{\text{th}} + \sigma_c \mathbf{S}_c + \sum_{i=1}^n \sigma_\ell^i \varphi^i(\nu) \mathbf{S}_\ell^i}{\alpha_c + \sum_{i=1}^n \alpha_\ell^i \varphi^i(\nu)}, \quad (4)$$

with \mathbf{B}_{th} the Planck function, $\mathbf{S}_c \equiv \mathbf{S}_c(\tau, \nu, \mu)$ the continuum source function, and $\mathbf{S}_\ell^i \equiv \mathbf{S}_\ell^i(\tau, \nu, \mu)$ the scattering part of the line source function of the i th line.

In our further discussion we employ the single-scattering approximation, since the CN lines for which we will apply this theory are optically thin. Then, the scattering part of the line source function takes the form

$$\mathbf{S}_\ell^i = \frac{1}{2} \int_0^1 \mathbf{P}_\ell^i(\mu, \mu', W_2^i, W_H^i) \mathbf{I}_{\text{inc}}(\mu') \exp\left(-\frac{\tau_{\text{inc}}(\nu) - \tau}{\mu'}\right) d\mu', \quad (5)$$

where $\mathbf{I}_{\text{inc}} = \mathbf{I}(\tau_{\text{inc}}, \nu, \mu)$ is the Stokes vector at the lower boundary with an optical depth τ_{inc} .

Within the weak-field approximation and under the assumption of an isotropic turbulent magnetic field the phase matrix \mathbf{P}_ℓ^i is parametrized as

$$\mathbf{P}_\ell^i = \mathbf{E}_{11} + \frac{3}{4} W_2^i W_H^i \mathbf{P}^{(2)}. \quad (6)$$

The matrix \mathbf{E}_{11} represents isotropic, unpolarized scattering, while $\mathbf{P}^{(2)}$ describes linearly polarized coherent scattering (e.g. Stenflo 1994). The line polarizability W_2^i , calculated as in Berdyugina et al. (2002), corresponds to the branching ratio of polarized and unpolarized scattering. The physics of the Hanle effect is contained in the Hanle depolarization factor

$$W_H = 1 - 0.4 \left(\frac{\gamma_H^2}{1 + \gamma_H^2} + \frac{4\gamma_H^2}{1 + 4\gamma_H^2} \right), \quad (7)$$

with

$$\gamma_H = 0.88 \frac{g_L B}{\Gamma_R + D^{(2)}}. \quad (8)$$

The variables Γ_R and $D^{(2)}$ are the radiative damping and depolarizing collision rates, respectively, and B is the strength of the single-valued, isotropic, microturbulent magnetic field. For the analysis in this paper we neglect depolarizing collisions. In the CN violet system we are dealing with ${}^2\Sigma - {}^2\Sigma$ transitions, so that the effective Landé factors g_L of the upper state of the two doublet components can be expressed through the quantum number of the total angular momentum J or, alternatively, of the angular momentum N , i.e. the total angular momentum without spin,

$$\begin{aligned} g_{L1} &= 1/(N + \frac{1}{2}) = 1/J, \\ g_{L2} &= -1/(N + \frac{1}{2}) = -1/(J + 1), \end{aligned} \quad (9)$$

where the indices 1 and 2 refer to the first and second doublet components with $J = N - 1/2$ and with $J = N + 1/2$, respectively.

With our choice of the phase matrix we implicitly neglect interferences between different fine structure components and perturbations by the Paschen-Back effect. The latter becomes relevant in the CN violet system only for field strengths exceeding 100 G (Shapiro et al. 2007), while the fields detected in the present study are weaker than 30 G.

Let us rewrite Eq. (4) in the form

$$\mathbf{S} = \frac{\sigma_c + \sum_{i=1}^n \sigma_\ell^i \varphi^i(\nu)}{\alpha_c + \sum_{i=1}^n \alpha_\ell^i \varphi^i(\nu)} \left(\sum_{i=1}^n \chi_i \mathbf{S}_\ell^i + \chi_c \mathbf{S}_c \right) + \frac{k_c + \sum_{i=1}^n k_\ell^i \varphi^i(\nu)}{\alpha_c + \sum_{i=1}^n \alpha_\ell^i \varphi^i(\nu)} \mathbf{B}_{\text{th}}, \quad (10)$$

where the coefficients χ_i and χ_c correspond to the conditional probabilities that scattering occurs in the i th line or in the continuum, respectively. In general they depend on frequency and are defined as

$$\chi_i = \frac{\sigma_\ell^i \varphi^i(\nu)}{\sigma_c + \sum_{j=1}^n \sigma_\ell^j \varphi^j(\nu)}, \quad \chi_c = \frac{\sigma_c}{\sigma_c + \sum_{j=1}^n \sigma_\ell^j \varphi^j(\nu)}. \quad (11)$$

The emergent Stokes vector is given by the formal solution of the radiative transfer equation

$$\mathbf{I}_{\text{em}}(\nu, \mu) = \int_0^{\tau_{\text{inc}}(\nu)} \mathbf{S}(\tau, \nu, \mu) e^{-\tau/\mu} \frac{d\tau}{\mu} + \mathbf{I}_{\text{inc}}(\nu, \mu) e^{-\tau_{\text{inc}}(\nu)/\mu}. \quad (12)$$

2.2. Two-layer approximation

We assume that the CN molecular layer is situated above the region where the continuum radiation is formed. The CN layer is isothermal with a homogeneous CN density, and its lower boundary is irradiated by the solar continuum radiation. While passing through the molecular layer this initial incident radiation can be scattered (but only once, as we adopt the single scattering approximation) or absorbed by CN molecules.

The ‘‘continuum’’ layer corresponds to the lower photosphere of the average quiet Sun, defines the lower boundary condition of the CN layer, and is thus crucial for the formation of the CN lines. In particular, the continuum layer determines the anisotropy of the continuum radiation, without which scattering polarization would cancel out, and it sets the degree of continuum polarization, which is subsequently depolarized by the line opacity.

In our modeling strategy we take the continuum intensity and polarization from previous work, rather than recalculating these values. The continuum intensity is obtained from the analytical expressions for limb-darkening derived from observations of the quiet Sun (Neckel 1996). We assume the radiation field to be azimuthally symmetric. The values of the continuum polarization correspond to the calculations by Stenflo (2005) for $\mu = 0.1$. The continuum polarization for arbitrary μ -values are computed by scaling the value at $\mu = 0.1$ with the analytical expressions introduced by Fluri & Stenflo (1999), which describe the μ -dependence of the continuum polarization and are based on theoretical modeling of the quiet solar photosphere.

The CN layer is described by the source function and the radiative transfer equation given in Sect. 2.1. In Eqs. (5) and (12) we insert the Stokes vector \mathbf{I}_{inc} , which is incident into the CN layer at the lower boundary and given by the frequency and angle dependent continuum intensity and polarization formed in

the continuum layer. The same equations require also the total optical thickness $\tau_{\text{inc}} = \mathcal{T}_{\text{CN}}$ of the CN layer. Assuming $\alpha_c = 0$, since the continuum is formed below the CN layer, the total optical thickness of the CN layer \mathcal{T}_{CN} can be calculated with the expression

$$\mathcal{T}_{\text{CN}}(\nu) = \mathcal{T}_{\text{norm}} \cdot \sum_{i=1}^n \alpha_\ell^i \varphi^i(\nu), \quad (13)$$

where the frequency independent coefficient $\mathcal{T}_{\text{norm}}$ is a free parameter of the model. It depends on the CN density and the geometrical thickness of the CN layer.

Within the two-layer approximation the total source function, Eq. (10), is further simplified since the opacity in the CN layer is only due to CN molecules, i.e. $k_c = \sigma_c = 0$. We assume that the ratio of the scattering and absorption probabilities $c_{\text{sc}} = \sigma_\ell^i / \alpha_\ell^i$ is the same for all molecular lines. This is basically equivalent to the assumption that the inelastic collision rate is independent of the involved transitions because the life time of the excited CN state $B^2\Sigma$ is almost independent of its quantum numbers (with the exception of some states with very small J numbers). Furthermore we neglect in our model the elastic collisions. Equation (10) is then reduced to the form

$$\mathbf{S} = c_{\text{sc}} \left(\sum_{i=1}^n \chi_i \mathbf{S}_\ell^i \right) + (1 - c_{\text{sc}}) \mathbf{B}_{\text{th}}. \quad (14)$$

Here the first term with the summation represents the scattered diffuse part of the radiation field, restricted within the CN layer to line contributions, while the second term corresponds to the thermal part.

The emergent Stokes vector is calculated employing the formal solution of the radiative transfer equation. To improve the accuracy of our results we divide the CN layer into k sublayers with equal optical thickness. The source function is assumed to be constant within each sublayer and evaluated at the center of each layer, giving us k different values $\mathbf{S}_1 \dots \mathbf{S}_k$. Introducing the above assumptions into Eq. (12), we find for the emergent radiation field at a given frequency

$$\mathbf{I}(\mu) = \mathbf{I}_{\text{inc}}(\mu) e^{-\mathcal{T}_{\text{CN}}/\mu} + \sum_{j=1}^k \mathbf{S}_j(\mu) \left(e^{-\mathcal{T}_{\text{CN}}(j-1)/k\mu} - e^{-\mathcal{T}_{\text{CN}}j/k\mu} \right). \quad (15)$$

The first term describes the radiation transmitted through the CN layer while the summation in the second term adds up the contributions from the different CN sublayers to the emergent spectrum. Note that the depth dependence of the source function within the CN layer is introduced by the exponential factor in Eq. (5), which describes absorption of the incident radiation and thus the height variation of the radiation field even within the isothermal slab. The scattering coefficient, on the other hand, depends only on the probability of the collisional de-excitation of the upper CN states, which under our assumptions remains constant for the whole molecular layer.

The subdivision of the molecular layer, expressed by Eq. (15), is a further sophistication of the model introduced by Faurobert & Arnaud (2003), who have assumed one single value of the source function within the molecular layer. The sublayers take care of the depth dependence of the source function within the CN layer, which results from the optical depth τ in the exponent of Eq. (5). Note that this does not contradict to our assumption of an isothermal CN layer with constant density. Our

test calculations have shown that the emergent Stokes vector becomes independent of the number of sublayers for k greater than five.

2.3. Free parameters

The model described above contains several free parameters, which have to be determined by fitting the theoretical spectrum to observations. These are the temperature of the isothermal CN layer T_{eff} , the scattering probability c_{sc} , the optical depth normalization coefficient $\mathcal{T}_{\text{norm}}$, and the magnetic field strength B .

The temperature T_{eff} affects the emergent spectrum in two ways. First of all, it enters via thermal term B_{th} of the source function (cf. Eq. (14)). This strongly influences the intensity, but not the other Stokes parameters, because of the strong temperature dependence of the Planck function and because the contribution from the thermal part is relatively large in the case of CN lines. Secondly, the line opacities depend on temperature via the Boltzmann factor (cf. Eq. 2), which affects the whole Stokes vector, although in a much weaker way than the thermal term.

The scattering probability c_{sc} defines the balance between the thermal and nonthermal parts of the total source function and influences the whole Stokes vector. The contribution to the Stokes vector from the CN layer linearly depends on c_{sc} .

The optical depth normalization coefficient $\mathcal{T}_{\text{norm}}$ describes the relative contributions to the Stokes vector from the CN layer and the continuum layer. It effectively determines the optical thickness of the CN layer, thus scaling both the molecular line depths in the intensity spectrum and the degree of line polarization.

Our main goal is the determination of the magnetic field strength B assuming a turbulent, isotropic field distribution within the CN layer. The magnetic field enters via Eq. (8) into the line source function. Since we consider weak fields and neglect Zeeman broadening, the magnetic field strength has virtually no effect on Stokes I and only causes differential Hanle effect in the second solar spectrum (see Sect. 3.1).

Usually, the influence of B on the Q/I_c spectrum can be easily distinguished from the effects of c_{sc} and $\mathcal{T}_{\text{norm}}$. The turbulent magnetic field causes a depolarization of the CN lines. The same can in principle also be obtained with a smaller scattering probability c_{sc} . Even an increased $\mathcal{T}_{\text{norm}}$ parameter can mimic the same behavior because it leads to stronger continuum depolarization. The crucial point, however, is that c_{sc} and $\mathcal{T}_{\text{norm}}$ affect all molecular lines with comparable strength, while the magnetic field leads to differential depolarization, because the Hanle effect efficiency strongly depends on the J quantum number (Fig. 1). Thus, the relative strength of CN lines in the second solar spectrum mainly depends on the magnetic field, but not so much on the other free parameters. Therefore, it is possible to diagnose the magnetic field strength when fitting several lines with different Hanle effect sensitivities.

2.4. Fitting procedure

We simultaneously fit the modeled polarization spectrum to observed Q/I_c data and the modeled intensity spectrum to I/I_c observations. During the fitting procedure we vary the magnetic field strength B , the scattering probability c_{sc} , the effective temperature T_{eff} and the optical depth normalization coefficient $\mathcal{T}_{\text{norm}}$. The most suitable values for these free parameters are

determined with a least-square method. More precisely, we minimize the quantity

$$\chi^2 = \frac{\zeta}{n} \sum_{i=1}^n \left(\frac{(Q/I_c)_{\text{obs}} - (Q/I_c)_{\text{th}}}{\sigma_{Q/I_c}} \right)^2 + \frac{(1-\zeta)}{n} \sum_{i=1}^n \left(\frac{(I/I_c)_{\text{obs}} - (I/I_c)_{\text{th}}}{\sigma_{I/I_c}} \right)^2, \quad (16)$$

where n is the number of data points and σ_{Q/I_c} and σ_{I/I_c} are the absolute errors of the observed polarization and intensity, which we set to 0.008% and 0.005%, respectively. The coefficient ζ , with possible values in the range from 0 to 1, allows us to give larger weight either to the intensity or polarization fit. As the signal in polarization is smaller and as only polarization is affected by the magnetic field we overweight the fit to the polarization fit by choosing $\zeta = 7/8$ (but the result of the fitting procedure is almost independent on the assumed ζ value).

The χ^2 value in Eq. (16) is a measure of the discrepancy between observations and the theoretical model, and describes the accuracy of the fit. Wavelength points that are heavily affected by atomic blends are excluded from the summation and the fitting procedure since we neglect atomic line opacity in our model.

3. Diagnostics with the CN violet system

In this section we apply the previously developed theory to the interpretation of observations of the CN $B^2\Sigma - X^2\Sigma$ system, the so-called violet system, in several spectral regions, with the goal to determine the turbulent magnetic field strength in the quiet Sun. From the discussion in Sect. 2.3 it is clear that suitable wavelength regions should contain molecular lines with different Hanle effect sensitivities. Therefore, we first investigate the differential Hanle effect in the CN violet system. This will guide us to choose three ideal spectral windows in different molecular bands, which are then analyzed and discussed.

3.1. Differential Hanle effect

Figure 1 illustrates the different Hanle effect sensitivities of synthetic line profiles for the example of the CN violet (0, 0) band (other CN violet bands give the same qualitative results). The relative degree of line polarization in the non-magnetic case is defined by three factors: the oscillator strength (which increases with the J number), the scattering polarizability W_2 (which decreases with J and asymptotically approaches a value of 0.1), and the Boltzmann factor (which also decreases with J). However, the final polarization obtained within the lines not only depends on the intrinsic line polarization. Rather it results from a combination of different contributions, namely the continuum polarization, continuum depolarization due to the line opacity, and line polarization, where only the latter is subject to the Hanle effect. It is clearly visible in Fig. 1 that stronger lines lead to additional continuum depolarization (cf. Fluri & Stenflo 2003). Furthermore, we can identify the expected diverse behavior of the lines in a magnetic field. While the intensity profiles remain unaffected, the line polarization is depolarized with increasing field strength, reflecting the J dependence of W_H . A magnetic field of 10 G already leads to a nearly saturated Hanle effect in the lines with $N = 1$ and 5, while the polarization in the $N = 70$ and 80 lines remains almost unchanged compared to the non-magnetic case. The latter lines require a magnetic field strength in the order of 100 G for a considerable depolarization.

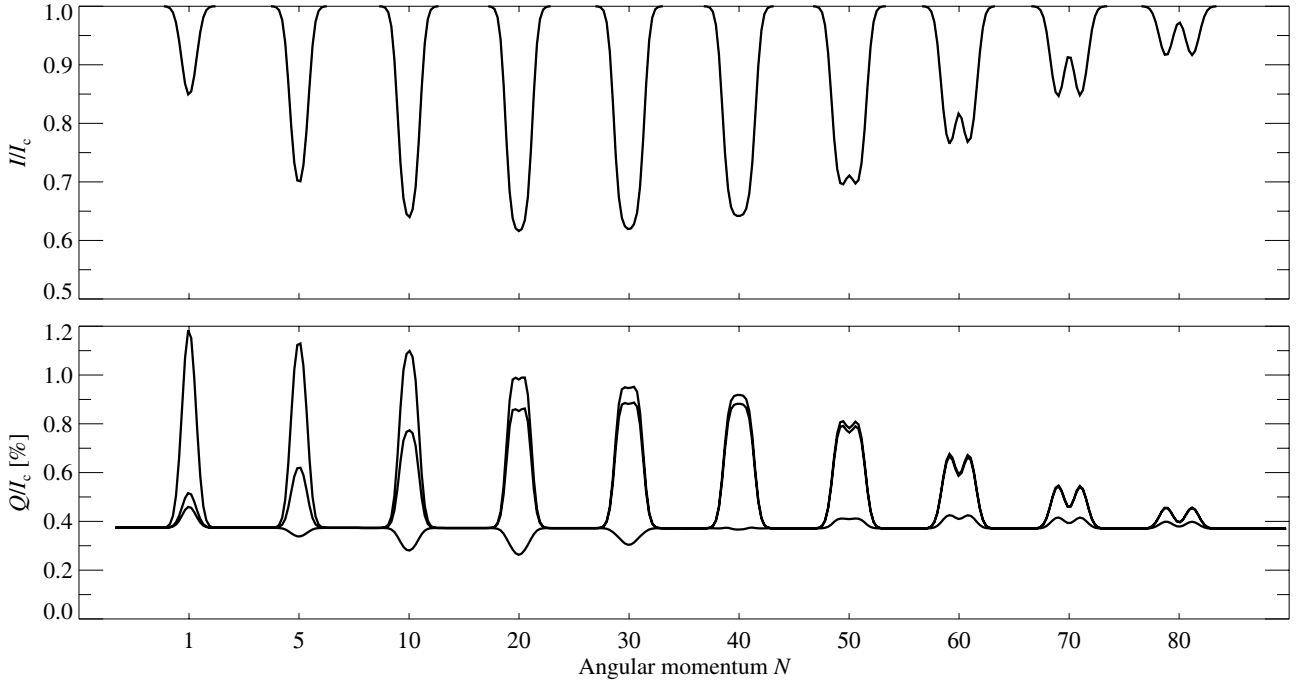


Fig. 1. Differential Hanle effect in synthetic Stokes profiles of R-branch lines in the CN violet (0,0) band. For selected angular momenta N of the lower state we show the complete doublet, which causes the double peak structure for large N . Three values for the turbulent magnetic field strength were chosen, giving identical Stokes I profiles but large modifications in the Stokes Q profiles: 0 G (upper curve), 10 G (intermediate curve), and 100 G (lower curve). In this and all other figures we assume an elastic collision rate equal to zero.

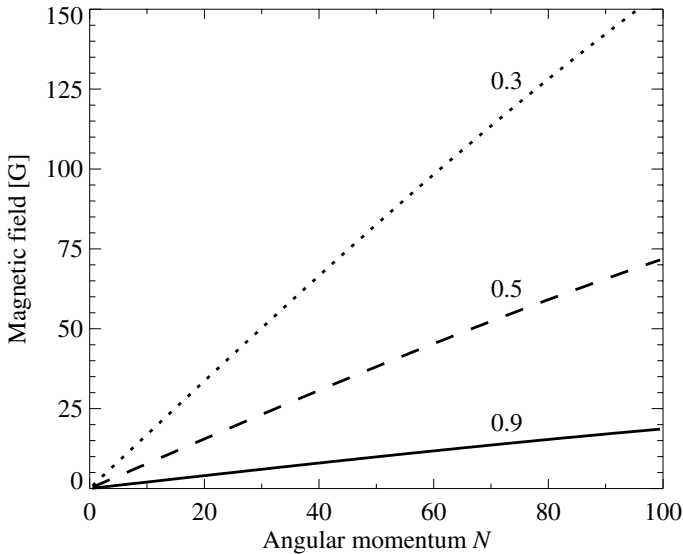


Fig. 2. The dependence of the Hanle depolarization factor W_H , given in Eq. (7), on the magnetic field strength and the N quantum number for the CN violet (0,0) band. The three curves correspond to the depolarization factors W_H equal to 0.9 (solid), 0.5 (dashed), and 0.3 (dotted), the last of which is already close to the asymptotic saturation limit of 0.2.

The magnetic field necessary for sizable depolarization by the Hanle effect is shown in Fig. 2 for different N quantum numbers, again for the CN violet (0,0) band. The different curves correspond to three values of W_H and identify the magnetic field range in which a specific line is sensitive to the Hanle effect. For example, lines with $N = 40$ react to the Hanle effect mainly in the field range from about 10 G to 60 G. For stronger fields these lines are maximally depolarized. On the other hand, their Hanle effect depolarization is negligible for weaker fields, which

makes them ideal reference lines below 10 G when compared to lines with much smaller N .

The ideal choice of spectral windows thus depends on the actual magnetic field strength. Based on previous results we selected narrow wavelength regions containing lines with N greater than 40 and others with N around 10. An additional selection criterion is a small number of atomic or molecular blends since they would significantly alter the intensity and polarization spectra.

3.2. Observational data

We compare our model calculations with the data from the second solar spectrum atlas compiled by Gandorfer (2005). All spectra were recorded with the UV version of the Zurich Imaging Polarimeter ZIMPOL II (Gandorfer et al. 2004) at the 1.5 m McMath-Pierce facility of the National Solar Observatory on Kitt Peak (Arizona, USA). The spectrograph slit was set parallel to the solar limb at the distance corresponding to the cosine of the heliocentric angle $\mu = 0.1 \pm 0.02$. Due to the telescope guiding and seeing, the value of μ was not completely stable, resulting in a small uncertainty of the exact limb distance. Because the scattering polarization is very sensitive to small variations in μ very close to the solar limb, μ behaves like a semi-free parameter for our analysis within strict limits defined by the error bars.

One of the main problems for our fitting procedure is the uncertainty of the zero point of the polarization scale in the observations. To solve this problem, we selected several wavelengths as continuum reference points, which are far away from very strong lines and least affected by blends. Based on the theoretical continuum polarization (Stenflo 2005), we shifted the polarization scale of the observations appropriately. The same wavelength points were used to renormalize the I/I_c data.

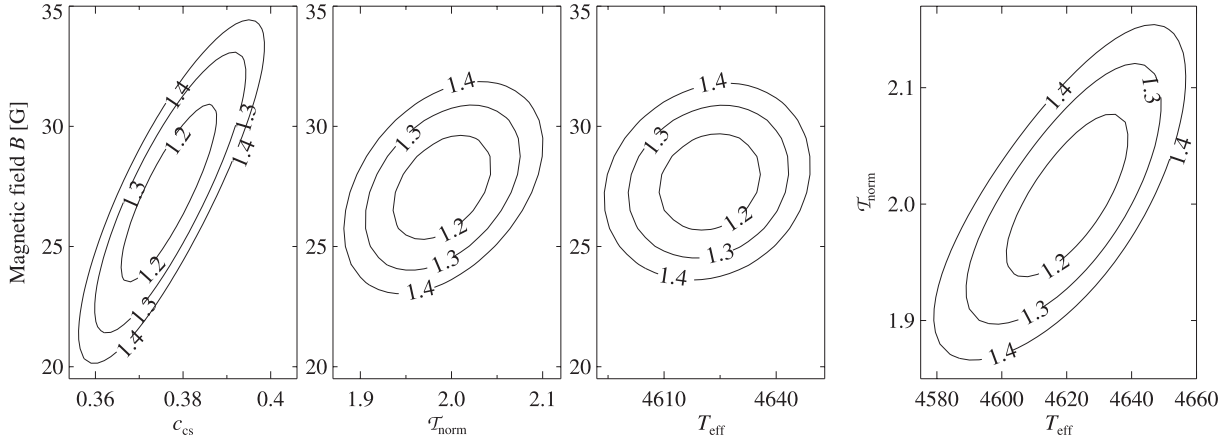


Fig. 3. χ^2 -contours of the fitting procedure in the first wavelength region (3865.9–3866.8 Å). The panels show 2-dimensional cuts in the parameter space through the minimum $\chi^2_{\min} = 1.13$. Numbers on the contours indicate the values of χ^2 as defined in Eq. (16).

We have chosen three wavelength regions for our calculations. These regions were not observed simultaneously and correspond, in principle, to different areas of the solar surface. Therefore measured parameters can vary from region to region.

The line parameters for all three regions are given in Tables 2–4 (available only in electronic form), which provide also weaker CN lines considered in the fit.

3.3. Region 1: diagnostics near 3866 Å

The first wavelength region 3865.9–3866.8 Å contains four CN violet system doublets with $N = 6, 7$ and 52 from the (1, 1) vibrational band and with $N = 69$ from the (0, 0) vibrational band. Because of the relatively large Landé factors ($g_N \sim 1/N$), the lines with $N = 6$ and 7 are very sensitive to the magnetic field and reach the saturated Hanle effect already at 5–10 G, while the lines with $N = 52$ and 69 are much less affected by the magnetic field (see Fig. 2). Therefore, this region is ideal for applying the differential Hanle effect technique, in particular since it contains lines with small and large N quantum numbers from the same vibrational band, as the relative strength of the lines from different bands can be influenced by other factors. Furthermore, all lines can be measured simultaneously in a single observation because they are spread over less than 1 Å.

Figure 3 shows the accuracy of the fitted spectrum illustrated by χ^2 -contours within the parameter space. All four free parameters are well constrained by the data, and their best values are listed in Table 1. In particular we find a turbulent magnetic field strength of 27 ± 5 G. The magnetic field B and the scattering probability c_{cs} exhibit a somewhat linear dependence (Fig. 3, leftmost panel), which is easily understood. A greater scattering probability leads to additional line polarization that has to be compensated with a stronger Hanle effect and thus a larger magnetic field. Nonetheless, both B and c_{cs} remain restricted to well defined limits because of the differential sensitivity to the Hanle effect of the involved CN lines and the roughly line independent scattering probability. Similarly, the parameters T_{norm} and T_{eff} also show a tendency towards a linear dependence, because a greater effective temperature reduces the line depth in the intensity spectrum that can be balanced by additional line opacity.

The modeled spectra are compared with observations in Fig. 4. For the calculated curves we have assumed the best values of the parameters T_{eff} , c_{cs} , and T_{norm} as listed in Table 1 while varying the magnetic field strength. The differential Hanle effect is clearly visible and constrains the magnetic field strength to

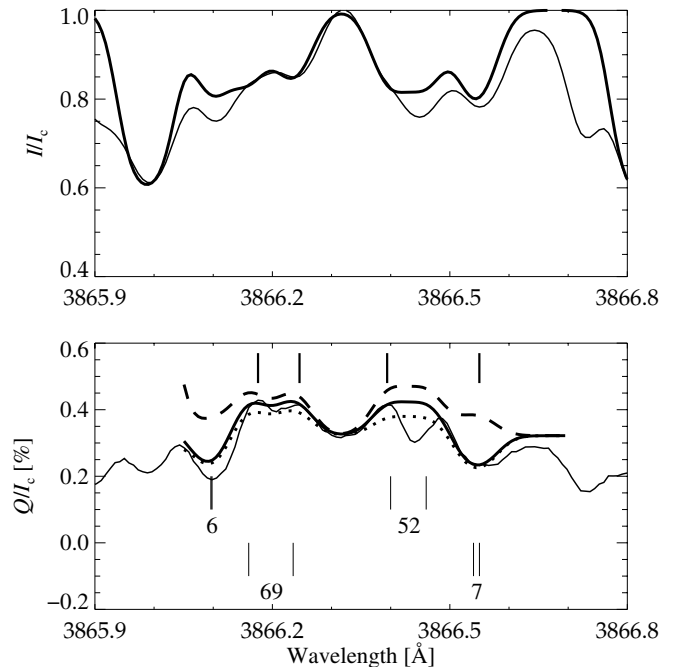


Fig. 4. Fits to the observations (thin solid lines) of Stokes I/I_c and Q/I_c in the first wavelength region (3865.9–3866.8 Å). The assumed magnetic field strengths in the calculated spectra are 0 G (no depolarization, thick dashed line), 27 G (best fit, thick solid line), and 45 G (too much depolarization, dotted line). The labels at the bottom of the lower panel indicate the position of the lines in the CN doublets with the corresponding N quantum number of the lower states. The labels at the top of the lower panel indicate the wavelength points which were chosen for the fitting procedure.

well defined limits. It has to be stronger than 15 G to saturate the highly sensitive $N = 7$ doublet. The upper limit for the field strength is defined by the only modest observed depolarizations of the lines with $N = 52$ and 69. The observed doubling of the $N = 52$ doublet probably appears due to a strong Ti I blend neglected in our model. We obtain the best fit for the magnetic field strength of 27 G.

3.4. Region 2: diagnostics in the (1, 1) bandhead

Molecular bandheads have a high potential for successful differential Hanle effect diagnostics. They contain a mixture of lines with different J -numbers and correspondingly different

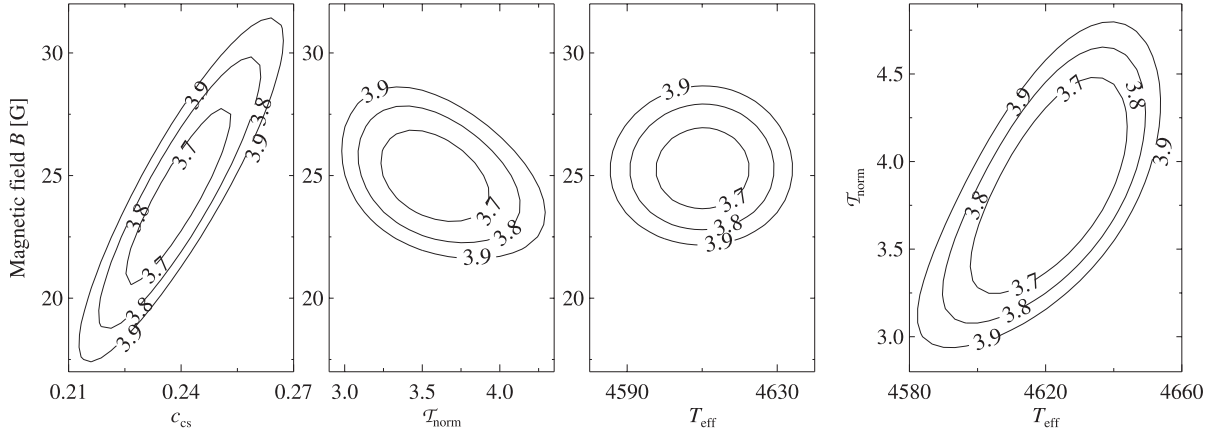


Fig. 5. χ^2 -contours as in Fig. 4 but for the second wavelength region close to the (1, 1) bandhead. We obtain a minimum value $\chi^2_{\min} = 3.62$.

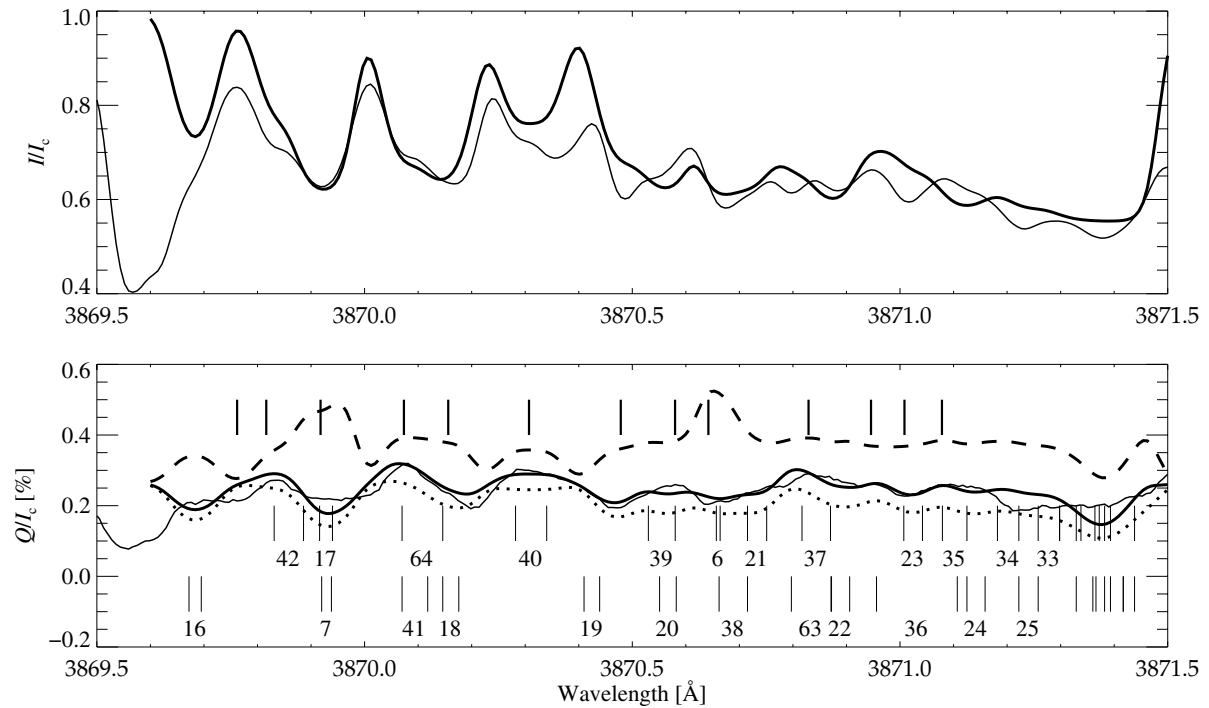


Fig. 6. Fits to the observations (thin solid lines) of Stokes I/I_c and Q/I_c in the second wavelength region close to the (1, 1) bandhead. The assumed magnetic field strengths in the theoretical spectra are 0 G (no depolarization, dashed line), 25 G (best fit, thick solid line) and 45 G (too much depolarization, dotted line). Labels at the bottom of the lower panel indicate the N quantum numbers of the lower states and the line positions within each doublet (close to the bandhead N is omitted due to overlap). The labels at the top of the lower panel indicate the wavelength points which were selected for the fitting procedure.

magnetic field sensitivities. On the other hand, all these lines belong to the same rotational branch and vibrational band so that many of the relevant line formation parameters nearly coincide. Here we present the results of our fitting procedure in the (1, 1) bandhead of the CN violet system.

The fitting procedure is illustrated in Fig. 5 by the χ^2 -contours. Due to numerous neglected blends from atoms and molecules other than CN the quality of the fit is worse than in the first wavelength region (Sect. 3.3). Furthermore, the noise level in this region is also a bit higher. Nonetheless, we determine a magnetic field strength of 25 ± 5 G in good agreement with the first region.

The modeled spectra are compared to observations in Fig. 6. The bandhead of the (1, 1) band is clearly visible near 3871.5 Å. The selected wavelength region also contains a few lines from the (0, 0) band. Contrary to the first region, lines with

different Hanle effect sensitivities can not be easily distinguished as all lines are blended with each other. Even when neglecting the magnetic field it is possible to fit the overall average degree of polarization. However, the exact shape of the Q/I_c spectrum can only be reproduced when accounting for a magnetic field strength within well defined limits.

In this wavelength region we had to adopt a cosine of the heliocentric angle $\mu = 0.12$, which still lies within the error limits of the observed data. In fact, we failed to obtain good fits in this region with $\mu = 0.10$, because the continuum polarization in the incident radiation field was too large.

3.5. Region 3: diagnostics close to the (0, 0) bandhead

The CN lines near the (0, 0) bandhead are very strong and heavily blended. This results in a significant optical thickness larger than 1 at the (0, 0) bandhead so that the single scattering

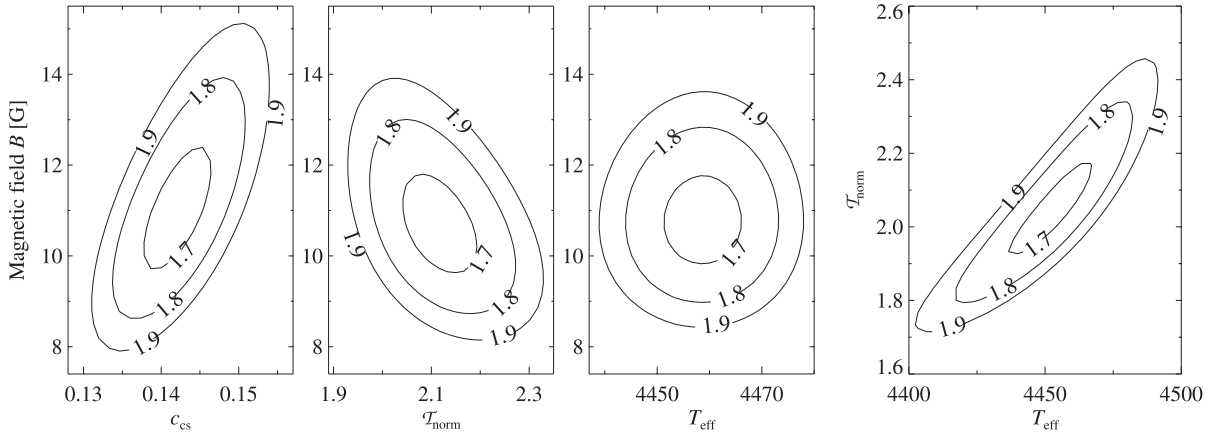


Fig. 7. χ^2 -contours as in Fig. 4 but for the third wavelength region near the (0, 0) bandhead. We obtain a minimum value $\chi^2_{\min} = 1.66$.

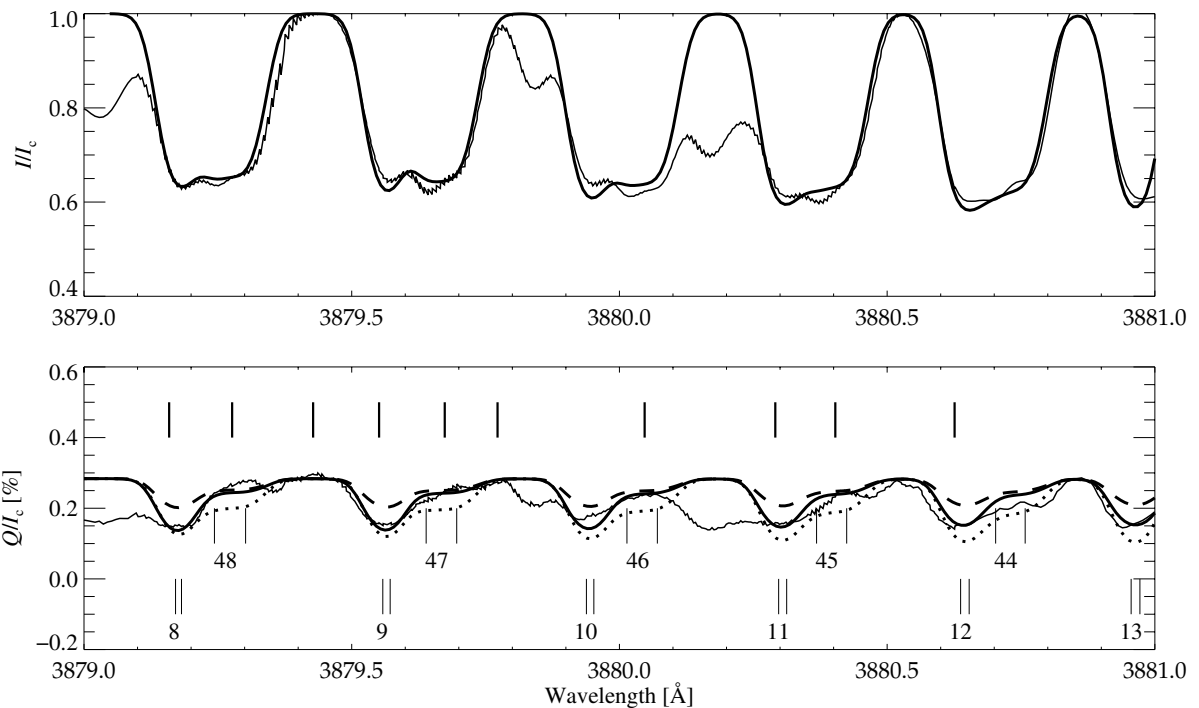


Fig. 8. Fits to the observations (thin solid lines) of the CN violet system in the third wavelength region close to the (0, 0) bandhead. The magnetic field strengths in the modeled spectra were set to 0 G (no depolarization, dashed line), 11 G (best fit, thick solid line), and 25 G (too much depolarization, dotted line). Labels at the bottom of the lower panel indicate the CN doublets with the N quantum numbers of the lower states. The labels at the top of the lower panel indicate the wavelength points chosen for the fitting procedure.

approximation loses its validity. Therefore, we selected a region which is 4 Å away from the bandhead.

This third wavelength region contains a mixture of lines with $N \sim 10$ and $N \sim 40$ from the P branch. The main features relevant for diagnosing the magnetic field in this spectral window are groups of neighboring lines containing a doublet with $N \sim 10$ and one with $N \sim 40$, which are particularly well seen in the intensity spectrum. The doublets within one group can however be distinguished due to a consistent wavelength shift and different polarization characteristics. The lines with $N \sim 10$ have a larger effective polarizability $W_2 \sim 0.5$ (compared to W_2 close to the asymptotic limit of 0.1 for lines with $N \sim 40$), which in principle causes a larger line polarization in a non-magnetic medium. The Hanle effect introduces a competing contribution to the relative line polarization because the lines with small N are much more sensitive to the magnetic field. From Fig. 8 we see that the consistent modeling leads to smaller polarization and the expected

stronger Hanle effect in the lines with small N . Therefore, we can indeed infer the magnetic field strength from the Q/I_c spectrum.

The results of the fitting procedure are presented in Figs. 7 and 8. Again all the free parameters are well constrained and we reach a reasonable quality of the fit, except in three narrow regions due to missing blend lines in our model at 3879.0 Å, 3879.85 Å, and 3880.2 Å, which is clearly seen in the intensity spectrum. The best fit is obtained for a magnetic field strength of 11 ± 2 G, which is significantly smaller than in the first two analyzed spectral windows.

3.6. Discussion

The above comparison of the modeled spectra to the observed data was performed independently in the three selected wavelength regions. The obtained sets of free parameters for

Table 1. Least squares fits to the magnetic field strength B , the optical depth normalization coefficient $\mathcal{T}_{\text{norm}}$, the scattering probability c_{sc} , the temperature of the isothermal CN layer T_{eff} and cosine of the heliocentric angle μ in three considered regions of spectrum

Region	B , [G]	$\mathcal{T}_{\text{norm}}$	c_{sc}	T_{eff} , K	μ
1	27 ± 5	2.0	0.375	4620	0.1
2	25 ± 5	3.5	0.240	4620	0.12
3	11 ± 2	2.3	0.142	4460	0.11

each region are collected in Table 1. They agree quite well in spite of the simple modeling, but there exist some differences that should be addressed.

The magnetic field strength found in the third region lies significantly below the value determined for the first two spectral regions. This is however not necessarily a contradiction, because the analyzed data were not observed simultaneously and sampled in principle different spatial positions on the Sun.

Furthermore, the line strength and accordingly the optical thickness changes from the first region to the third one. This influences the formation heights of the CN lines, which leads to different physical conditions in the corresponding atmospheric layers. Lines in the third region are stronger than in the first two and therefore form higher in the atmosphere where the temperature is smaller. A comparison of the contribution functions calculated for the typical lines in each spectral window shows that the difference in temperatures is around 100 K, consistent with the difference of 160 K obtained with our least-squares method (see Table 1). Our results suggest that the magnetic field strength may decrease with height, but a detailed interpretation of these small differences lies outside the scope of the current model and will be clarified by future work.

In addition, the continuum normalization and polarization in the third region remains an issue due to broad wings of nearby strong atomic lines including the H8 line from the Balmer series. The intensity normalization has little effect on the inferred magnetic field strength because it affects both the theoretical and observed Q/I_c signal in the same way. However, the additional opacity from those nearby atomic lines would probably reduce the continuum polarization and slightly increase the magnetic field strength by a few Gauss.

Overall, the physical parameters determined with the fit lie in a reasonable range. The small scattering probability is consistent with expectations and with a relatively small influence of non-LTE effects. The total optical thickness of the CN layer is close to but still smaller than 1 (maximal optical thicknesses in the considered regions are 0.26, 0.5 and 0.32 for the first, second and third regions correspondingly) in accordance with our assumptions. Finally, the temperature T_{eff} of the CN layer around 4500–4600 K that is required for the best fits, corresponds well to the temperature in the upper solar photosphere where the CN lines are formed.

An interesting extension of our method is the interpretation of the center-to-limb variations of the continuum and line radiation field, as it was done for Sr I 4607 Å line by Stenflo et al. (1997) and for the C₂ and MgH molecular lines by Faurobert & Arnaud (2003). Such observation for the CN violet system are not available yet, but could be very useful for constraining the atmospheric models.

4. Conclusions

We have employed the differential Hanle effect for the interpretation of observations of the CN violet system, which allowed us to determine the turbulent magnetic field strength in the upper solar photosphere. This expands the number of molecules used for Hanle effect diagnostics, which was so far only based on C₂ and MgH. The main conclusions can be summarized as follows:

- We have developed a new model that can diagnose the turbulent magnetic field via the differential Hanle effect. It is applicable for optically thin lines, relies on simple and few assumptions, and consistently accounts for continuum polarization, depolarization of the continuum by line opacity, intrinsic line polarization by coherent scattering, and the Hanle effect.
- Lines from the CN violet system are well suited as a diagnostic tool based on the differential Hanle effect. They are very prominent in the second solar spectrum, and there exist several narrow spectral regions containing simultaneously lines with small and large J quantum numbers. These lines exhibit different sensitivities to the Hanle effect, since the effective Landé factor is inversely proportional to J in the CN violet system.
- We have identified three spectral windows that are ideal for Hanle effect diagnostics and contain a small number of atomic and molecular blends.
- In all three spectral regions we obtained a good quality of the fit and determined turbulent magnetic field strengths of 27 ± 5 G, 25 ± 5 G, and 11 ± 2 G, respectively.

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Table 2. Line parameters used for the fitting procedure in the first spectral region (3865.9–3866.8 Å). The columns list the wavelength (λ), the line identification (ID), the polarizability (W_2), the lifetime of the upper level (τ_R), the turbulent magnetic field strength at which the line polarization is reduced to 30% of its non-magnetic value neglecting elastic collision (B_{30}), the vibrational number which coincides for the upper and lower states (v), and the oscillator strength (f_{vJS}).

λ , Å	ID	W_2	τ_R , 10^{-8} s	B_{30}	v	f_{vJS}	λ , Å	ID	W_2	τ_R , 10^{-8} s	B_{30}	v	f_{vJS}
3865.978	$R_1(12)$	0.1239	7.2719	22	0	0.391	3865.998	$R_2(12)$	0.1260	7.2736	22	0	0.362
3866.013	$Q_{12}(12)$	0.3982	7.2736	22	0	0.001	3866.096	$P_1(6)$	0.0604	7.4397	9	1	0.154
3866.090	$Q_{21}(6)$	0.3916	7.4397	9	1	0.002	3866.098	$P_2(6)$	0.0545	7.4635	9	1	0.130
3866.160	$P_1(69)$	0.0957	7.0063	111	0	2.016	3866.086	$Q_{21}(69)$	0.3999	7.0063	111	0	0.000
3866.235	$P_2(69)$	0.0957	7.0069	111	0	1.987	3866.400	$P_1(52)$	0.0944	7.2450	86	1	1.254
3866.344	$Q_{21}(52)$	0.3999	7.2450	86	1	0.000	3866.460	$P_2(52)$	0.0943	7.2456	86	1	1.230
3866.540	$P_1(7)$	0.0650	7.4251	11	1	0.178	3866.533	$Q_{21}(7)$	0.3938	7.4251	11	1	0.002

Table 3. Line parameters used for the fitting procedure in the second spectral region (in the (1, 1) bandhead), with the same parameters as in Table 2.

λ , Å	ID	W_2	τ_R , 10^{-8} s	B_{30}	v	f_{vJS}	λ , Å	ID	W_2	τ_R , 10^{-8} s	B_{30}	v	f_{vJS}
3869.672	$P_1(16)$	0.0829	7.3832	26	1	0.394	3869.655	$Q_{21}(16)$	0.3988	7.3832	26	1	0.001
3869.695	$P_2(16)$	0.0819	7.3845	26	1	0.370	3869.831	$P_1(42)$	0.0931	7.2954	70	1	1.015
3869.786	$Q_{21}(42)$	0.3998	7.2954	70	1	0.000	3869.886	$P_2(42)$	0.0929	7.2959	70	1	0.991
3869.920	$R_1(7)$	0.1397	7.2917	14	0	0.246	3869.938	$R_2(7)$	0.1457	7.2980	14	0	0.217
3869.947	$Q_{12}(7)$	0.3953	7.2980	14	0	0.002	3869.916	$P_1(17)$	0.0838	7.3806	28	1	0.418
3869.898	$Q_{21}(17)$	0.3990	7.3806	28	1	0.001	3869.940	$P_2(17)$	0.0829	7.3817	28	1	0.394
3870.070	$P_1(64)$	0.0954	7.0428	103	0	1.871	3870.001	$Q_{21}(64)$	0.3999	7.0428	103	0	0.000
3870.146	$P_2(64)$	0.0953	7.0433	103	0	1.842	3870.070	$P_1(41)$	0.0929	7.2999	68	1	0.991
3870.026	$Q_{21}(41)$	0.3998	7.2999	68	1	0.000	3870.118	$P_2(41)$	0.0928	7.3004	68	1	0.967
3870.146	$P_1(18)$	0.0846	7.3780	29	1	0.442	3870.127	$Q_{21}(18)$	0.3991	7.3780	29	1	0.001
3870.176	$P_2(18)$	0.0838	7.3790	29	1	0.418	3870.282	$P_1(40)$	0.0928	7.3043	66	1	0.967
3870.239	$Q_{21}(40)$	0.3998	7.3043	66	1	0.000	3870.340	$P_2(40)$	0.0926	7.3048	66	1	0.944
3870.460	$P_1(19)$	0.0854	7.3754	31	1	0.466	3870.440	$Q_{21}(19)$	0.3992	7.3754	31	1	0.001
3870.489	$P_2(19)$	0.0846	7.3763	31	1	0.442	3870.530	$P_1(39)$	0.0926	7.3086	65	1	0.944
3870.488	$Q_{21}(39)$	0.3998	7.3086	65	1	0.000	3870.580	$P_2(39)$	0.0924	7.3091	65	1	0.920
3870.551	$P_1(20)$	0.0861	7.3728	33	1	0.489	3870.530	$Q_{21}(20)$	0.3992	7.3728	33	1	0.001
3870.582	$P_2(20)$	0.0854	7.3736	33	1	0.466	3870.657	$R_1(6)$	0.1457	7.2987	12	0	0.217
3870.664	$R_2(6)$	0.1538	7.3078	12	0	0.187	3870.672	$Q_{12}(6)$	0.3938	7.3078	12	0	0.002
3870.715	$P_1(21)$	0.0867	7.3701	35	1	0.513	3870.693	$Q_{21}(21)$	0.3993	7.3701	35	1	0.001
3870.751	$P_2(21)$	0.0861	7.3708	35	1	0.489	3870.662	$P_1(38)$	0.0924	7.3128	63	1	0.920
3870.621	$Q_{21}(38)$	0.3998	7.3128	63	1	0.000	3870.715	$P_2(38)$	0.0922	7.3133	63	1	0.896
3870.797	$P_1(63)$	0.0953	7.0498	102	0	1.842	3870.729	$Q_{21}(63)$	0.3999	7.0498	102	0	0.000
3870.871	$P_2(63)$	0.0953	7.0503	102	0	1.813	3870.872	$P_1(22)$	0.0872	7.3674	36	1	0.537
3870.849	$Q_{21}(22)$	0.3994	7.3674	36	1	0.001	3870.906	$P_2(22)$	0.0867	7.3681	36	1	0.513
3870.817	$P_1(37)$	0.0922	7.3170	61	1	0.896	3870.777	$Q_{21}(37)$	0.3998	7.3170	61	1	0.000
3870.870	$P_2(37)$	0.0920	7.3174	61	1	0.872	3871.007	$P_1(23)$	0.0878	7.3647	38	1	0.561
3870.983	$Q_{21}(23)$	0.3994	7.3647	38	1	0.001	3871.042	$P_2(23)$	0.0872	7.3653	38	1	0.537
3870.956	$P_1(36)$	0.0920	7.3210	60	1	0.872	3870.917	$Q_{21}(36)$	0.3998	7.3210	60	1	0.000
3871.107	$P_2(36)$	0.0918	7.3214	60	1	0.848	3871.125	$P_1(24)$	0.0882	7.3618	40	1	0.585
3871.099	$Q_{21}(24)$	0.3995	7.3618	40	1	0.000	3871.159	$P_2(24)$	0.0878	7.3624	40	1	0.561
3871.222	$P_1(25)$	0.0887	7.3589	41	1	0.609	3871.195	$Q_{21}(25)$	0.3995	7.3589	41	1	0.000
3871.079	$P_1(35)$	0.0918	7.3249	58	1	0.848	3871.042	$Q_{21}(35)$	0.3998	7.3249	58	1	0.000
3871.258	$P_2(25)$	0.0882	7.3594	41	1	0.585	3871.298	$P_1(26)$	0.0891	7.3559	43	1	0.633
3871.270	$Q_{21}(26)$	0.3996	7.3559	43	1	0.000	3871.125	$P_2(35)$	0.0916	7.3254	58	1	0.824
3871.382	$P_2(27)$	0.0891	7.3533	45	1	0.633	3871.417	$P_2(28)$	0.0895	7.3501	46	1	0.657
3871.338	$P_1(27)$	0.0895	7.3528	45	1	0.657	3871.309	$Q_{21}(27)$	0.3996	7.3528	45	1	0.000
3871.182	$P_1(34)$	0.0916	7.3287	56	1	0.824	3871.145	$Q_{21}(34)$	0.3997	7.3287	56	1	0.000
3871.329	$P_2(26)$	0.0887	7.3564	43	1	0.609	3871.364	$R_1(5)$	0.1538	7.3083	11	0	0.187
3871.372	$R_2(5)$	0.1655	7.3224	11	0	0.158	3871.379	$Q_{12}(5)$	0.3916	7.3224	11	0	0.002
3871.222	$P_2(34)$	0.0913	7.3292	56	1	0.800	3871.258	$P_1(33)$	0.0913	7.3324	55	1	0.800
3871.223	$Q_{21}(33)$	0.3997	7.3324	55	1	0.000	3871.298	$P_2(33)$	0.0910	7.3329	55	1	0.776
3871.393	$P_1(30)$	0.0905	7.3430	50	1	0.728	3871.361	$Q_{21}(30)$	0.3997	7.3430	50	1	0.000
3871.438	$P_2(30)$	0.0902	7.3435	50	1	0.705	3871.393	$P_1(29)$	0.0902	7.3464	48	1	0.705
3871.362	$Q_{21}(29)$	0.3996	7.3464	48	1	0.000	3871.438	$P_2(29)$	0.0898	7.3469	48	1	0.681
3871.382	$P_1(28)$	0.0898	7.3496	46	1	0.681	3871.352	$Q_{21}(28)$	0.3996	7.3496	46	1	0.000
3871.366	$P_1(31)$	0.0908	7.3396	51	1	0.752	3871.333	$Q_{21}(31)$	0.3997	7.3396	51	1	0.000
3871.417	$P_2(31)$	0.0905	7.3401	51	1	0.728	3871.329	$P_1(32)$	0.0910	7.3361	53	1	0.776
3871.295	$Q_{21}(32)$	0.3997	7.3361	53	1	0.000	3871.360	$P_2(32)$	0.0908	7.3365	53	1	0.752

Table 4. Line parameters used for the fitting procedure in the third spectral region (close to the (0, 0) bandhead), with the same parameters as in Table 2.

$\lambda, \text{\AA}$	ID	W_2	$\tau_R, 10^{-8} \text{ s}$	B_{30}	v	f_{vJS}	$\lambda, \text{\AA}$	ID	W_2	$\tau_R, 10^{-8} \text{ s}$	B_{30}	v	f_{vJS}
3879.174	$P_2(8)$	0.0650	7.3078	12	0	0.217	3879.185	$P_1(8)$	0.0686	7.2987	12	0	0.246
3879.177	$Q_{21}(8)$	0.3953	7.2987	12	0	0.002	3879.247	$P_1(48)$	0.0939	7.1431	78	0	1.407
3879.195	$Q_{21}(48)$	0.3999	7.1431	78	0	0.000	3879.306	$P_2(48)$	0.0938	7.1435	78	0	1.378
3879.566	$P_1(9)$	0.0716	7.2917	14	0	0.275	3879.557	$Q_{21}(9)$	0.3963	7.2917	14	0	0.002
3879.580	$P_2(9)$	0.0686	7.2980	14	0	0.246	3879.649	$P_1(47)$	0.0938	7.1485	77	0	1.378
3879.598	$Q_{21}(47)$	0.3999	7.1485	77	0	0.000	3879.707	$P_2(47)$	0.0937	7.1489	77	0	1.349
3879.952	$P_1(10)$	0.0740	7.2863	16	0	0.304	3879.942	$Q_{21}(10)$	0.3970	7.2863	16	0	0.001
3879.967	$P_2(10)$	0.0716	7.2909	16	0	0.275	3880.029	$P_1(46)$	0.0937	7.1538	75	0	1.349
3879.980	$Q_{21}(46)$	0.3999	7.1538	75	0	0.000	3880.087	$P_2(46)$	0.0935	7.1542	75	0	1.320
3880.317	$P_1(11)$	0.0761	7.2819	17	0	0.333	3880.305	$Q_{21}(11)$	0.3975	7.2819	17	0	0.001
3880.332	$P_2(11)$	0.0740	7.2854	17	0	0.304	3880.389	$P_1(45)$	0.0935	7.1590	73	0	1.320
3880.340	$Q_{21}(45)$	0.3999	7.1590	73	0	0.000	3880.446	$P_2(45)$	0.0934	7.1594	73	0	1.291
3880.662	$P_{11}(12)$	0.0778	7.2782	19	0	0.362	3880.649	$Q_{21}(12)$	0.3979	7.2782	19	0	0.001
3880.678	$P_2(12)$	0.0761	7.2809	19	0	0.333	3880.728	$P_1(44)$	0.0934	7.1641	72	0	1.291
3880.681	$Q_{21}(44)$	0.3998	7.1641	72	0	0.000	3880.784	$P_2(44)$	0.0933	7.1645	72	0	1.262
3880.985	$P_1(13)$	0.0794	7.2749	21	0	0.391	3880.971	$Q_{21}(13)$	0.3982	7.2749	21	0	0.001