

LETTER TO THE EDITOR

A *Chandra* X-ray detection of the L dwarf binary Kelu-1 Simultaneous *Chandra* and Very Large Array observations

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ABSTRACT

Context. Magnetic activity in ultracool dwarfs, as measured in X-rays and $H\alpha$, shows a steep decline after spectral type M7-M8. So far, no L dwarf has been detected in X-rays. In contrast, L dwarfs may have higher radio activity than M dwarfs.

Aims. We observe L and T dwarfs simultaneously in X-rays and radio to determine their level of magnetic activity in the context of the general decline of magnetic activity with cooler effective temperatures.

Methods. The field L dwarf binary Kelu-1 was observed simultaneously with *Chandra* and the Very Large Array.

Results. Kelu-1AB was detected in X-rays with $L_X = 2.9_{-1.3}^{+1.8} \times 10^{25}$ erg s⁻¹, while it remained undetected in the radio down to a 3σ limit of $L_R \leq 1.4 \times 10^{13}$ erg s⁻¹ Hz⁻¹. We argue that, whereas the X-ray and $H\alpha$ emissions decline in ultracool dwarfs with decreasing effective temperature, the radio luminosity stays (more or less) constant across M and early-L dwarfs. The radio surface flux or the luminosity may better trace magnetic activity in ultracool dwarfs than the ratio of the luminosity to the bolometric luminosity.

Conclusions. Deeper radio observations (and at short frequencies) are required to determine if and when the cut-off in radio activity occurs in L and T dwarfs, and what kind of emission mechanism takes place in ultracool dwarfs.

Key words. radio continuum: stars – stars: activity – stars: coronae – stars: individual: Kelu-1 – stars: low-mass, brown dwarfs – X-rays: stars

1. Introduction

There is significant evidence that magnetic activity, commonly seen in low-mass stars, survives in ultracool dwarfs (e.g., Tagliaferri et al. 1990; Fleming et al. 1993; Drake et al. 1996; Fleming et al. 2000; Rutledge et al. 2000; Martín & Bouy 2002; Schmitt & Liefke 2002; Fleming et al. 2003; Briggs & Pye 2004; Stelzer 2004; Berger et al. 2005; Hallinan et al. 2006; Osten et al. 2006; Stelzer et al. 2006a,b; Phan-Bao et al. 2007). Since the latter are fully convective, a different kind of magnetic dynamo mechanism than in the Sun and in stars with tachoclines must take place in ultracool dwarfs, possibly due to turbulent magnetic fields.

A common indicator of magnetic activity in late-type stars, X-rays have so far never been detected in stars later than spectral type M9. Stelzer & Neuhäuser (2003), Berger et al. (2005), and Stelzer et al. (2006b) reported the non-detection in X-rays of L dwarfs, down to $<6.6 \times 10^{24}$ erg s⁻¹. Studies of the X-ray emission in M dwarfs show a decline in emission after spectral type M7-M8 (e.g., Fleming et al. 2003; Stelzer et al. 2006a), in parallel with the decline in the chromospheric $H\alpha$ emission (Gizis et al. 2000; Mohanty & Basri 2003; West et al. 2004). The radio luminosity in late-type stars correlates over several decades with the X-ray luminosity (Güdel & Benz 1993;

Benz & Güdel 1994); however, Berger (2002) argued that ultracool dwarfs do not follow this correlation and suggested an increase in radio activity with cooler effective temperatures. Such an increase was further supported by Berger et al. (2005) and Burgasser & Putman (2005) for late-M and early-L dwarfs, despite very low radio detection rates (e.g., Berger 2006). The decline in $H\alpha$ and X-ray activity, and the non-detections of late-L and T dwarfs in the radio, may be related to the highly neutral atmospheres of these dwarfs, which decouples the magnetic fields from the photospheric gas (Meyer & Meyer-Hofmeister 1999; Mohanty et al. 2002) and could lead to unfavorable conditions for magnetic activity.

In this Letter, we present the results of simultaneous observations of the early-L field brown dwarf binary Kelu-1 (Ruiz et al. 1997; Liu & Leggett 2005; Gelino et al. 2006) with *Chandra* and the Very Large Array (VLA). As mentioned earlier, no L dwarf has yet been detected in X-rays; this Letter presents, therefore, the first X-ray detection of an L dwarf, while the dwarf remains undetected in the radio.

2. The Kelu-1AB dwarf binary

At a distance of 18.7 ± 0.7 pc, Kelu-1 was found thanks to its high proper motion ($\mu \sim 0.3$ /yr; Ruiz et al. 1997; Dahn et al. 2002; Scholz & Meusinger 2002; Lodieu et al. 2005). Its optical

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Table 1. Observation log for VLA and *Chandra*.

Position of Kelu-1AB (Equinox: J2000; Epoch: J2006.41)	
Based on Lodieu et al. (2005)	
Right ascension	13 ^h 05 ^m 40 ^s .00
Declination	-25°41′06″.1
<i>Chandra</i> mean position	
Right ascension	13 ^h 05 ^m 40 ^s .01 ± 0.01
Declination	-25°41′05″.9 ± 0′.10
VLA (Program S6570)	
Observation date	2006 May 31 0h36–7h36 UT
Configuration & Wavelength	BnA & 3.6 cm (X band)
Flux and phase calibrators ..	1331+305 (3C 286) / 1258–223
<i>Chandra</i> (ObsId 5421)	
Observation date	2006 May 31 0h04–7h18 UT
Instrument & Mode	ACIS-S & VFAINT
Total exposure	23.8 ks

spectrum shows weak Li I absorption and H α in emission (Ruiz et al. 1997; Kirkpatrick et al. 1999). Its age is difficult to assess, but it probably lies in the range 0.3–0.8 Gyr (Liu & Leggett 2005). Basri et al. (2000) and Mohanty & Basri (2003) measured a high rotation rate ($v \sin i \sim 60 \text{ km s}^{-1}$), supported by a 1.8 h photometric period in H α (Clarke et al. 2002, 2003), which cannot be due to the orbit of a binary (Gelino et al. 2006). Kelu-1 remained undetected in the X-rays down to $L_X < 7.3 \times 10^{27} \text{ erg s}^{-1}$ (Neuhäuser et al. 1999), and in the radio ($S_{3.6} < 28 \mu\text{Jy}$, 3σ limit; Krishnamurthi et al. 1999).

It was originally classified as an L2 dwarf; however, high spatial resolution images have recently revealed its binary nature (Liu & Leggett 2005; Gelino et al. 2006). The binary is separated by about 0′.3 with a position angle of $\sim 220^\circ$ (Liu & Leggett 2005), and there is evidence of orbital motion (Gelino et al. 2006). The latter authors estimate that Kelu-1A has spectral type L2 ± 1 , whereas Kelu-1B is slightly colder with a spectral type of L3.5 ± 1 , in line with the estimates (L1.5–L3 and L3–L4.5) of Liu & Leggett (2005) based on different methods. Gelino et al. (2006) also give masses of 0.060 ± 0.01 and $0.055 \pm 0.01 M_\odot$, bolometric luminosities $\log(L_{\text{bol}}/L_\odot) = -3.83$ and -3.99 , and effective temperatures in the range 1900–2100 K and 1700–1900 K for Kelu-1A and B, respectively.

3. Observations and data reduction

The *Chandra* observation was coordinated with the VLA. Table 1 provides the details of the observations. The X-ray data were processed with the CIAO 3.3.0.1 software together with CALDB 3.2.2. The VLA data were calibrated with the AIPS software. A typical dwell cycle spent 2 min on the phase calibrator and 9.5 min on Kelu-1. The total VLA on-source time in the 3.6 cm band was 4.4 h.

We used the proper motion properties of Kelu-1 as given by Lodieu et al. (2005) to determine the position of the binary at the epoch of observations (see Table 1). We extracted X-ray events inside a circle centered at the position of Kelu-1 and with a radius of 1″ (Fig. 1). We used the energy range of 0.2–6.0 keV to reduce the background contamination. We obtained 4 events of energy 0.63, 0.86, 1.19 and 1.38 keV. There is no strong clustering of the event arrival times (the 3rd and 4th events are separated by 91.1 s, much longer than the frame time, 3.1 s), suggesting that the events do not arise from flares. We estimated the background level by extracting events in a concentric annulus of 3″ and 30″ radii that avoided nearby sources. We extracted 102 background events, which corresponds to a

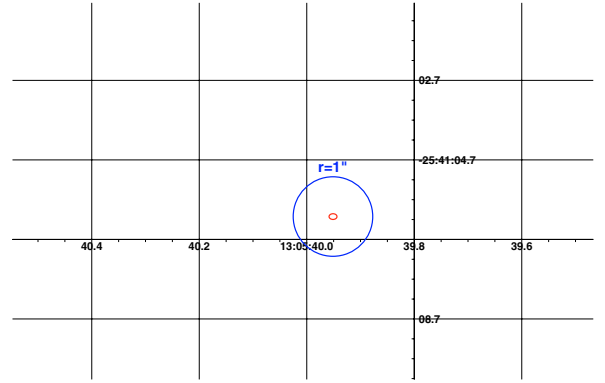


Fig. 1. *Chandra* 0.2–6.0 keV image centered on Kelu-1AB. The small red ellipse corresponds to the expected position (with errors) of Kelu-1. The extraction region is shown as a blue circle of 1″ radius. The nearby background is very low, with sky pixels containing either 0 or 1 event.

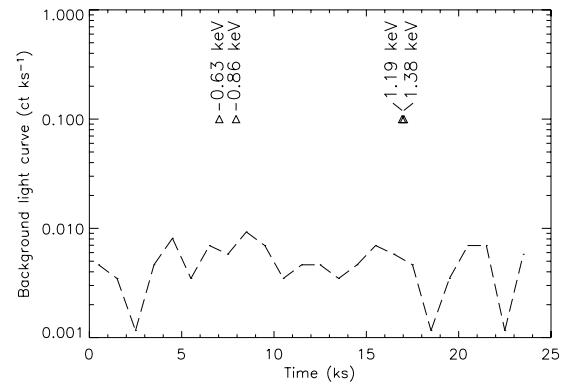


Fig. 2. *Chandra* 0.2–6.0 keV light curve of the background around Kelu-1AB, scaled down to the extraction area of the source. The 4 events at the position of Kelu-1AB are placed arbitrarily at a y -value of 0.1; the event energies are also labeled.

background level in the extraction region for Kelu-1 of less than 1 event (0.11 event). The average background count rate was, thus, 0.005 ct ks^{-1} (Fig. 2). Using the Kraft et al. (1991) approach, the 68% Bayesian confidence range for the number of source events for Kelu-1 is 2.18–6.27. With the approach of Ayres (2004), we determined that the X-ray detection of Kelu-1 was a 4.3σ detection for a one-sided Gaussian distribution.

In a similar fashion as for ϵ Ind Bab (Audard et al. 2005), we simulated a single temperature plasma model (APEC 1.3.1; Smith et al. 2001) with solar abundances in XSPEC to determine the conversion factor from count rate to X-ray luminosity. Using the distance of Kelu-1, we estimate a 0.1–10 keV X-ray luminosity of $L_X = 2.9 \times 10^{25} \text{ erg s}^{-1}$, while the 68% Bayesian confidence range based on the Kraft et al. (1991) approach corresponds to $(1.6\text{--}4.7) \times 10^{25} \text{ erg s}^{-1}$. Note that the above estimates work for plasma temperatures of 0.4 to 1 keV, which could be expected from old brown dwarfs. As mentioned in Audard et al. (2005), the values can be higher by factors of 1.2–2.0 and even 3.25 in the extreme case of a plasma temperature of 0.1 keV.

In the radio regime, Kelu-1 remained undetected during the observation down to a rms flux density of $14 \mu\text{Jy}$, i.e., a 3σ upper limit for the radio luminosity at 3.6 cm of $L_R < 1.76 \times 10^{13} \text{ erg s}^{-1} \text{ Hz}^{-1}$. After combining with the archival VLA data from Krishnamurthi et al. (1999), Kelu-1 still remains undetected down to an rms flux density of $11 \mu\text{Jy}$, i.e., $L_R < 1.38 \times 10^{13} \text{ erg s}^{-1} \text{ Hz}^{-1}$ (3σ).

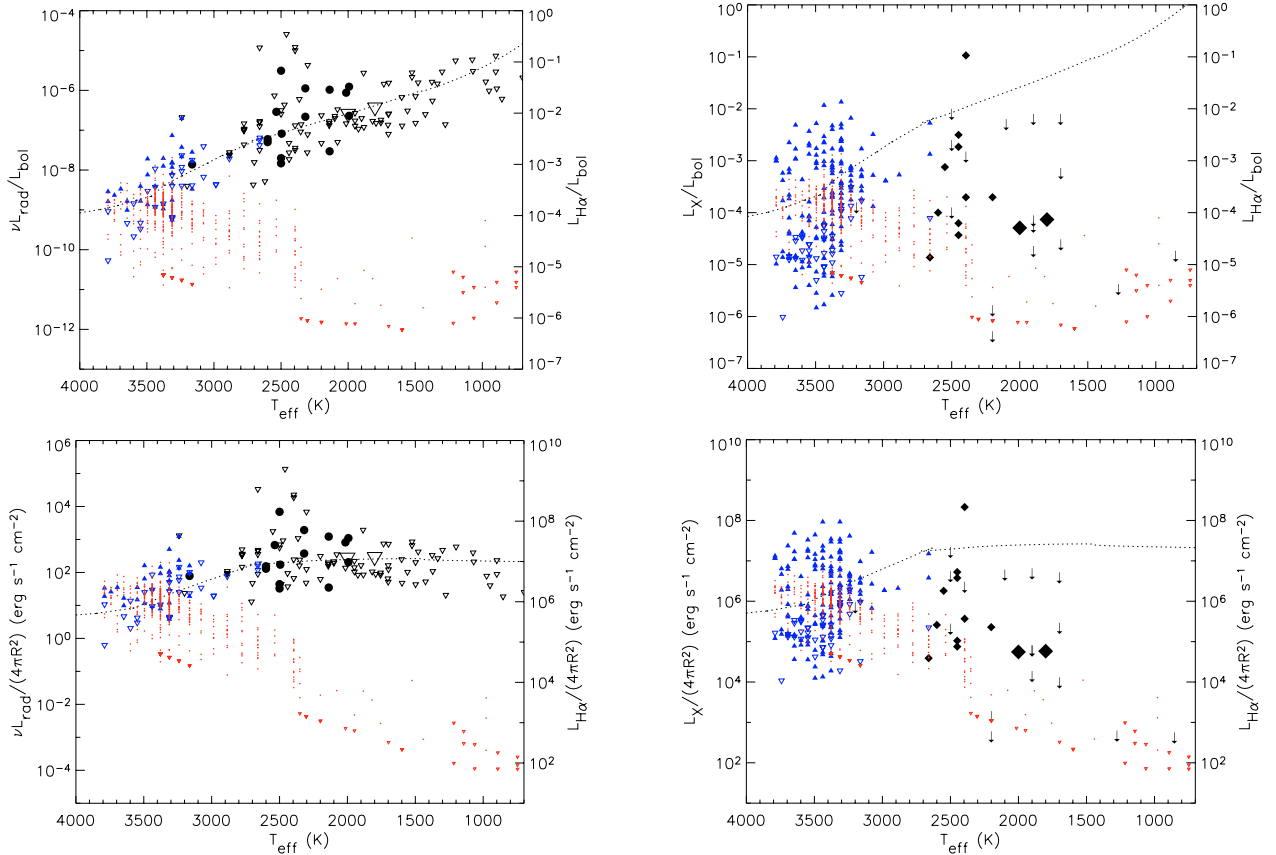


Fig. 3. *Top:* luminosity to bolometric luminosity ratio in the radio (*left panel*) and in the X-rays (*right panel*). The ratio for H α is shown in both panels with its scale on the right y -axis and as small symbols in red. *Bottom:* similar as above, for surface fluxes. Ultracool dwarfs are shown in black (filled circles in the radio and diamonds in X-rays for detections, empty downward-pointing triangles for radio upper limits and downward-pointing arrows for X-ray upper limits), while M dwarfs are shown in blue (filled upward-pointing triangles for detections, empty downward-pointing triangles for upper limits). The Kelu-1AB points (we assumed that the X-ray flux and radio upper limit were either due to Kelu-1A or to Kelu-1B) are shown as larger symbols. Note that flaring levels (if known) were purposely discarded. The dotted lines show the variations of L/L_{bol} ratios or the surface fluxes vs. T_{eff} for an arbitrary constant luminosity (in radio, X-rays, or H α).

4. Discussion

The top panels of Fig. 3 show the $\nu L_{\text{R}}/L_{\text{bol}}$ and $L_{\text{X}}/L_{\text{bol}}$ ratios for ultracool dwarfs and for M dwarfs, and the $L_{\text{H}\alpha}/L_{\text{bol}}$ ratio for comparison, as a function of effective temperature¹. As noted in previous studies, the X-ray ratio significantly decreases with decreasing effective temperature, in line with the H α ratio. On the other hand, the radio ratio increases (Berger 2002; Burgasser & Putman 2005). However, it should be noted that most radio observations of L and T dwarfs only provide upper limits (e.g., Berger 2006). The top panels also show as a dotted line the L/L_{bol} ratio (in radio, X-rays, or H α) vs. T_{eff} for an arbitrary constant radio/X-ray/H α luminosity. For ultracool dwarfs, the ratios increase because their bolometric luminosity depends essentially

¹ Data taken from the literature; radio: Krishnamurthi et al. (1999); Burgasser & Putman (2005); Berger (2006) and this paper for ultracool dwarfs, White et al. (1989); Güdel et al. (1993); Leto et al. (2000) for M dwarfs; X-rays: Rutledge et al. (2000); Martín & Bouy (2002); Fleming et al. (2003); Stelzer & Neuhäuser (2003); Audard et al. (2005); Berger et al. (2005); Burgasser & Putman (2005); Stelzer et al. (2006a,b) for “old” ultracool dwarfs ≥ 100 Myr, Doyle (1989); White et al. (1989); Güdel et al. (1993); Delfosse et al. (1998); Leto et al. (2000) for M dwarfs; H α : Hawley et al. (1996); Delfosse et al. (1998); Gizis et al. (2000); Burgasser et al. (2003); Mohanty & Basri (2003). If unavailable, L_{bol} , T_{eff} , and R/R_{\odot} were calculated from polynomial fits based on some of the above data, or from Golimowski et al. (2004), and Stefan-Boltzmann’s law.

only on the effective temperature, i.e., $\log L_{\text{bol}} \propto \log T_{\text{eff}}$. Indeed, radii in ultracool dwarfs vary little ($R \sim 0.09 R_{\odot}$). The L/L_{bol} ratios, therefore, increase with decreasing T_{eff} for a constant luminosity (in radio, X-rays, or H α). The radio upper limits in L and T dwarfs are also consistent with $\nu L_{\text{R}}/L_{\text{bol}}$ for a constant radio luminosity, suggesting that the current radio data do not go deep enough to detect any cut-off, if present. In contrast, the X-ray and H α observations are deep enough to detect it.

While the L/L_{bol} ratio defines the amount of power radiated in a wavelength regime compared to the bolometric luminosity, and while it is considered a good measure of magnetic activity in late-type stars, the ratio might be less adequate in ultracool dwarfs. Perhaps the surface flux (i.e., the ratio of the luminosity and the dwarf’s surface, $L/(4\pi R^2)$; Schmitt 1997) may better trace magnetic activity at the bottom of the main sequence². The bottom panels of Fig. 3 show the surface fluxes in radio, X-ray, and H α . If the luminosity were constant (in radio, X-rays, or H α), we would observe no significant decrease. This is approximately the case for the radio (despite the lack of detections below 2000 K), in stark contrast with the decrease in H α and X-rays. Note that we again plotted a dotted line of the surface flux for an

² Since R/R_{\odot} is almost constant for ultracool dwarfs, we could replace the surface flux by the luminosity as an activity indicator. Nevertheless, to allow comparison with early-M dwarfs, which have larger radii, we use the surface flux in our discussion.

arbitrary constant luminosity. The slight increase in radio surface flux observed in late-M/early-L dwarfs compared to early-M dwarfs is only due to the decrease in radius with decreasing T_{eff} for M dwarfs, but it is consistent with a constant luminosity, suggesting that the radio-emitting mechanism is similar in M dwarfs and in detected ultracool dwarfs, and that it does not lose its emission strength with decreasing T_{eff} . In contrast, the X-ray surface flux and luminosity in ultracool dwarfs declines in a similar fashion as $H\alpha$.

It appears that the radio emission reaches a plateau in luminosity across M dwarfs and ultracool dwarfs (at least above 2000 K), while magnetic activity in the chromosphere ($H\alpha$) and in the hot coronal loops (X-rays) declines with decreasing T_{eff} . This result may point toward a different kind of mechanism of radio and X-ray/ $H\alpha$ in ultracool dwarfs and in late-type stars, which could explain the observed deviations of ultracool dwarfs from the Güdel-Benz L_R-L_X relation. Cyclotron maser emission is possibly the dominant radio emission mechanism, as claimed by Hallinan et al. (2006, 2007). Such a mechanism was indeed proposed for M flare stars using a dipole magnetic trap model (Bingham et al. 2001; Kellett et al. 2002). The non-detection of Kelu-1 could be due to a lack of sensitivity, or simply the inclination of Kelu-1's rotation axis does not allow the beam of coherent radio emission to cross our line of sight. Or perhaps we observed at too high frequencies (since $\nu_c = 2.8B$, with ν_c in GHz and B in kG, $\nu_c = 8.4$ GHz requires $B \sim 3$ kG).

5. Conclusions

We have presented the first X-ray detection of an L binary dwarf, Kelu-1AB, while the binary remains undetected in the radio. The suggested increase in L_R/L_{bol} in ultracool dwarfs may be an artifact of the dependence of L_{bol} almost solely on T_{eff} in ultracool dwarfs. We suggest that the radio luminosity stays constant, at least down to $T_{\text{eff}} \sim 2000$ K, and may drop for cooler temperatures, although current radio surveys of ultracool dwarfs lack sensitivity. In contrast, magnetic activity as measured in X-rays and $H\alpha$ slowly declines with decreasing T_{eff} . The main dominant radio emission mechanism in ultracool dwarfs may not be gyrosynchrotron as in earlier-type main-sequence stars but coherent emission by electron cyclotron maser. The slower spin-down of late-M and L dwarfs than G-K dwarfs may also lead to a different behavior in the radio and in X-rays for ultracool dwarfs. There is a need to go deeper in the radio regime (and shorter frequencies if electron-cyclotron maser is the main radio emission mechanism) to determine if and when the radio emission declines in the increasingly neutral atmospheres of ultracool L and T dwarfs.

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