

HD 97048: a closer look at the disk[★]

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ABSTRACT

Aims. Large ground-based instruments, like VISIR on the VLT, providing diffraction-limited (~ 0.3 arcsec) images in the mid-infrared where strong PAH features appear, enable us to see the flaring structure of the disks around Herbig Ae stars. Although progress has been made in modelling the disk with radiative transfer models able to reproduce the spectral energy distribution (SED) of Herbig Ae stars, the constraints brought by images have not been yet fully exploited. We investigate whether these new observational imaging constraints can be accounted for by predictions based on existing models of passive, centrally irradiated hydrostatic disks made to fit the SEDs of the Herbig Ae stars.

Methods. The images taken by VISIR in the 8.6 and 11.3 μm aromatic features reveal a large flaring disk around HD 97048 inclined to the line of sight. To analyse the spatial distribution of these data, we use a disk model that includes the most up to date understanding of disk structure and physics around Herbig Ae stars with grains in thermal equilibrium, in addition to transiently-heated PAHs.

Results. We compare the observed spatial distribution of the PAH emission feature and the adjacent continuum emission with predictions based on existing full disk models. Both SED and spatial distribution are in very good agreement with the model predictions for common disk parameters.

Conclusions. We take the general agreement between observations and predictions as strong support for the physical pictures underlying our flared disk model.

Key words. circumstellar matter – stars: formation – stars: pre-main sequence – stars: individual: HD 97048

1. Introduction

Herbig Ae/Be stars are thought to be the intermediate mass ($\sim 2\text{--}10 M_{\odot}$) counterparts of T Tauri stars (Waters & Waelkens 1998). They have a strong infrared (IR) excess coming from warm circumstellar dust. Similar to T Tauri stars, this dust is believed, through the interpretation of spectral energy distributions (SEDs), to reside in a flared/flat circumstellar disk (Meeus et al. 2001).

PAHs (Polycyclic Aromatic Hydrocarbons) undergo transient heating: they do not reach thermal equilibrium with the radiation field but absorb individual photons, experiencing a rapid increase in temperature, and slowly cool, thus re-radiating the absorbed energy in the infrared (IR). This radiation allows us to see much further in the disk since these grains can reach high temperatures far away from the star. As a result, PAH emission can be used to probe the external region of disks around HAE stars (Lagage et al. 2006). Another point of interest is that, if present, PAHs can constitute an important source of opacity and are likely to play a key role in the thermal budget and chemistry of the gas, as they do in the interstellar medium. The thermal coupling between PAHs and gas via the photoelectric effect will, for example, determine the gas temperature in the upper layer of the disk where the gas and dust temperatures are not well coupled (Jonkheid et al. 2004; Kamp & Dullemond 2004).

Furthermore, PAHs are also good tracers of the presence of very small particles in the surface layers of disks and their emission can tell if and where very small particles survive settling and coagulation processes that cause the majority of the original grain population to grow to larger sizes in the same objects.

We have started a program of imaging nearby HAE stars with VISIR, the new mid-infrared instrument attached to the third 8-m unit of ESO’s Very Large Telescopes located at Cerro Paranal, Chile. One of the first targets was HD 97048 which has a strong IR excess, characteristic of a flared disk, and strong PAH features.

In a previous work, we were able for the first time to constrain the geometry of the disk, thanks to PAH emission (Lagage et al. 2006). Lagage et al. (2006) found there is a large flaring disk around HD 97048 extending at least up to 370 AU, vertically optically thick at the observed wavelength (8.6 μm) and inclined to the line of sight by $42.8_{-2.5}^{+0.8}$ degree. Moreover, with a very simple model, we were able to measure the flaring index which was found to be $1.26_{-0.05}^{+0.05}$. In this model, the PAH-emitting region is only located at the surface of the disk, whose surface scale height H_s varies with r^q (q is the flaring index) and whose flux intensity I follows r^{-p} . It is only recently that images of the flaring structure of the disk around Herbig Ae objects became available (Perrin et al. 2006; Lagage et al. 2006). It is now primordial to compare the spatial distribution constraints with existing full disk models only based on fitting the SEDs of the Herbig Ae stars.

[★] Based on observations obtained at VLT (Paranal) with VISIR. Program number 075.C-0540(C).

Table 1. Observations of HD 97048 in the different VISIR filters in imaging BURST mode. The sensitivity was calculated for each night in the BURST mode. The filter are free of any strong atmospheric line contribution.

Filter	Central wavelength (μm)	half band width (μm)	Sensitivity (mJy/10 σ /1h)	Elementary exposure time (ms)	Total time on source (s)	Date	Seeing (arcsec)	Airmass	Standard star
SIV	10.49	0.16	5.6	25	600	17/06/2005	0.80	1.74	HD 91056
PAH2	11.26	0.59	4.5	50	800	25/01/2005	0.55	1.66	HD 85503
NeII	12.27	0.18	7.7	16	320	17/06/2005	0.80	1.74	HD 91056

In this paper, we are interested in particular to check whether these new observational spatial constraints can be accounted for by predictions based on existing models of passive centrally irradiated hydrostatic disks made to fit the SEDs of the Herbig Ae stars. We explore whether a more sophisticated disk model could explain the SED of HD 97048 and at the same time, the spatial distribution of the circumstellar material with few free parameters. To the best of our knowledge, few tests like that (Doucet et al. 2006) have been performed yet for disks around pre-main sequence stars of intermediate mass. To do that, we used the model of Dullemond et al. (2001) and Dullemond & Dominik (2004a) which can account for the global shape of the SED of a quite large number of Herbig Ae stars. The dust model takes into account grains in thermal equilibrium and stochastically-heated PAH. The structure of the disk is calculated at hydrostatic equilibrium and a radiative transfer in two dimensions is used to calculate the emission of the different grain populations.

The paper is organized as follows. In Sect. 2, we present the knowledge on HD 97048 and its circumstellar material. In Sect. 3, we describe the observations and the data results. In Sect. 4, we describe the disk model and in Sect. 5, we compare the observed SED and spatial distribution of the circumstellar material to the model predictions.

2. HD 97048

HD 97048 is a nearby HAe star of spectral type Be9.5/A0 located in the Chamaeleon I dark cloud, at a distance of 180 pc (van den Ancker et al. 1997; Whittet et al. 1997). It is surrounded by a large amount of circumstellar material left from the star formation process, which produces a large infrared (IR) excess over the stellar emission (large IR ~ 0.35 – $0.40 L_*$, Acke & van den Ancker 2004; Van Kerckhoven et al. 2002). HD 97048 has been classified as a HAe star of group I with evidence of flared disk (Meeus et al. 2001), since its SED rises in the IR (Acke & van den Ancker 2004). Spectroscopic observations of the IR excess have also revealed the presence of strong PAH features at 3.3, 6.2, 7.7, 8.6, 11.3 microns and nano-diamond features in the 3.4–3.5 μm region (Siebenmorgen et al. 2000; Van Kerckhoven et al. 2002). No silicate emission band at 10 μm appears in the spectra of HD 97048. Recent mid-IR long slit spectroscopic observations with TIMMI2 show that the aromatic emission features at long wavelengths (i.e., 8.6 and 11.3 μm) are extended and come mostly from a region of 200–300 AU, likely a disk (van Boekel et al. 2004). These results have been confirmed recently by imaging data taken in the PAH band at 8.6 μm (Lagage et al. 2006). The large extended part of the mid-IR emission seen on scales of 5 to 10 arcsec by Prusti et al. (1994) and Siebenmorgen et al. (2000) is most likely due to an extended envelope of transiently heated very small grains and PAHs surrounding the star and the disk system. Using Adaptive Optics high angular resolution (~ 0.1 arcsec) spectroscopic observations, Habart et al. (2004, 2005) were also able to

spatially resolve the emission in the aromatic and diamond features around 3 μm and found that the emission must be within 30 AU, closely related to the star-disk system. Weintraub et al. (2005) have reported near-IR molecular hydrogen emission for this object.

3. VISIR observations

3.1. Observations

The observations were performed using the ESO mid-infrared instrument VISIR installed on the VLT (Paranal, Chili), equipped with a DRS (former Boeing) 256 \times 256 pixels BIB detector. The object was observed on 2005, January 25th and June 17th. It was observed in the PAH bands at 8.6 and 11.3 μm and in the adjacent continuum. The data at 8.6 μm were used in Lagage et al. (2006). In this paper, we focus on the data obtained at 11.3 μm . A summary of the observations is given in Table 1. These data were taken with an imaging mode of VISIR that allows diffraction-limited image in the N band: the BURST mode.

Under good seeing (≤ 0.5 arcsec in the visible), the images in the mid-IR are diffraction-limited even on an 8 m class telescope. Unfortunately, the median seeing experienced at Paranal is close to 0.8 arcsec, which degrades the angular resolution. Indeed, for a seeing of 0.8 arcsec in the visible, assuming that the wavelength dependence of the seeing follows a $\lambda^{-1/5}$ law, the seeing value at 10 μm is 0.4 arcsec, which is larger than the diffraction limit of 0.3 arcsec. This represents a 5 pixel movement on the detector with the smallest field of view of VISIR (0.075''/pixel). In order to reach the best spatial resolution with VISIR, we implemented a new imaging mode on bright objects, the *BURST mode*. The principle is to take short enough exposure images (≤ 50 ms) to freeze the turbulence; the coherence time of the atmosphere at 10 μm is around 300 ms at Paranal for a good seeing. In order to correct for the turbulence by offline processing, the data are stored every 1000 elementary images in one nodding position for a chopping frequency of 0.25 Hz in north/south direction. The nodding direction is perpendicular to the chopping direction with an amplitude of 8''. After classical data reduction in mid-IR, a cube of 500¹ images chopped and noded (4 beams/image) is obtained. Because of the turbulence, each source on an image moves independently and as a result, we have to individually extract the 4 sources in each image (4 quarters) of the cube and shift and add the image with the ones corresponding to the same quarter. Finally, we shift and add the four final images of the four quarters.

3.2. Results

In this section, we present the images of HD 97048 obtained in the PAH emission filter, as well as in the adjacent continuum

¹ 1000 divided by 2 because of the 2 chopper positions.

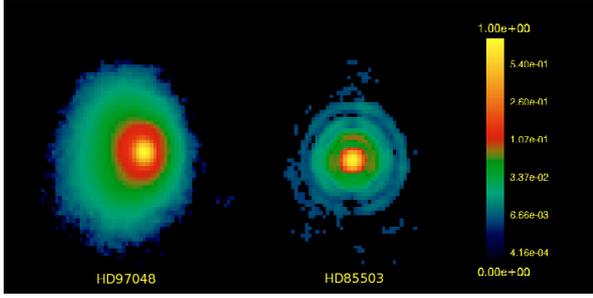


Fig. 1. *Left*, HD 97048 in the PAH2 filter (centered at $11.3 \mu\text{m}$) with VISIR camera. *Right*, PSF reference (HD 85503) with the same filter. The two objects are normalized to see the extension. The object is more extended than the reference with an east/west asymmetry in the emission.

emission, and compare them with the standard star in the same filters.

Figure 1 shows the image of HD 97048 in the PAH filter centered at $11.3 \mu\text{m}$ and compared with the standard star HD 85503. HD 97048 is quite extended – up to 2 arcsec – compared to the reference star. The full-width at half-maximum (FWHM) of HD 97048 is 1.2 times that of the standard star, close to the diffraction-limit, for almost all filters (Table 2). In addition, comparing the emission in the PAH band and in the adjacent continuum (Fig. 2), the emission in PAH is much more extended than in the continuum. Figure 3 shows the brightness spatial distribution along a cut in the north/south and east/west direction in the PAH filter and in the adjacent continuum (SIV) compared to the standard star in the SIV filter. In the PAH filter, the disk is extended in the direction north/south up to 320 AU (surface brightness of $53 \text{ mJy}/''^2$) whereas the extension in the continuum at $10.5 \mu\text{m}$ goes to 135 AU (surface brightness of $110 \text{ mJy}/''^2$). These results are in agreement with those already found in the previous study at $8.6 \mu\text{m}$ (Lagage et al. 2006). The continuum is extended, when comparing the object and the reference star in the SIV filter. By comparing the VISIR and ISO (Infrared Space Observatory) fluxes, we found that the PAH emission peak to the continuum ratio is stronger in the ISO spectrum. This is certainly due to the surrounding nebula, also included in the larger beam of ISO, as already suggested by van Boekel et al. (2004). Based on the fluxes measured in the wavelength bands at 8.6 and $11.3 \mu\text{m}$, we estimate that the nebula contamination is about 40%, in agreement with that found by van Boekel et al. (2004).

4. Disk model

In order to analyse the observations, we use the disk model described in Dullemond et al. (2001), Dominik et al. (2003) and Dullemond & Dominik (2004a), including the most up to date understanding of disk structure and physics around Herbig Ae/Be stars. In the following, we describe the disk structure and the radiation transfer used and give the details of the adopted dust model.

4.1. Disk structure

We consider disks heated by irradiation from the central star, in hydrostatic equilibrium in the vertical direction, with dust and gas well mixed (flared disks, Chiang & Goldreich 1997). In this model, the stellar flux impinging with flaring angle α upon the disk is absorbed in the upper layers of the disk, which will

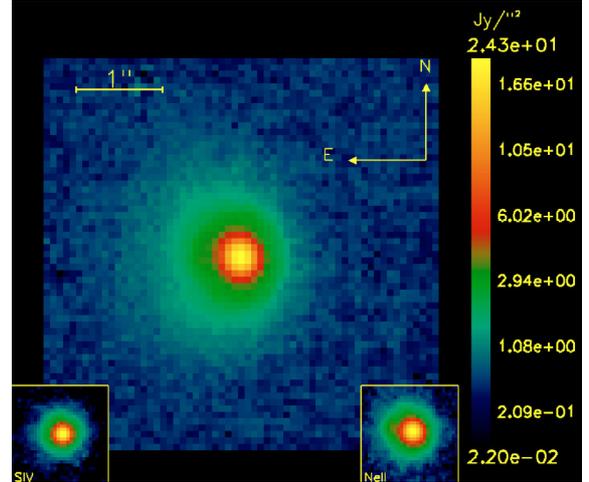


Fig. 2. HD 97048 in the PAH2 filter (*middle*), and in the adjacent continuum (SIV *left*, NeII *right*). The images have the same signal-to-noise so that it is possible to compare the extension.

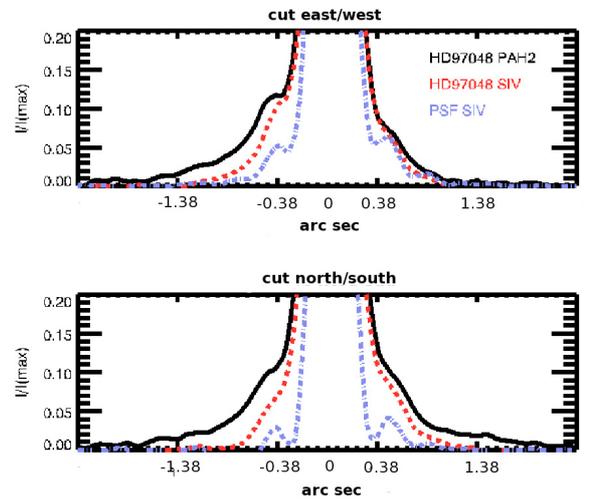


Fig. 3. Normalized intensity profiles along a cut through the VISIR images of HD 97048 in the PAH band (full line) and in the continuum SIV (dashed line) superimposed on a reference star (dashed-dot line) in the SIV band. East/West is in the upper panel and above, represents the direction South-North. Comparing the PAH band and the adjacent continuum (for the same signal-to-noise), HD 97048 is much more extended in the PAH band. Comparing the PSF and the object in SIV, we see that HD 97048 is also extended in the continuum.

reradiate half of the flux away from the disk and half down into its deeper layers. The inner boundary (rim) is directly exposed to the stellar flux and is puffed up since it is hotter than the rest of the disk. The model computes how much the inner rim puffs up, and how much of the disk behind it will be shadowed by this puffed-up rim (Dullemond et al. 2001). Once the star is given, the disk structure (i.e. pressure scale height and flaring angle) is completely defined after specifying the inner and outer radii, the surface density distribution ($\Sigma = \Sigma_0(R/R_0)^{-p}$, with R_0 a fiducial radius) and the dust model. The models are appropriate for disks that are optically thick to the stellar radiation. It is the case for disks around pre-main sequence stars, up to very large radii (e.g., 5000 AU for a disk mass $M_d \sim 0.2 M_\odot$, $p = 1$).

Table 2. Comparison of the *FWHM* for the object and the reference star (PSF) in the different filters. We give the theoretical value of the *FWHM* to show that the diffraction limit is obtained with the BURST mode. For the spatial extension, we note in the third and fourth columns the distance from the star at 10σ above the noise in the east/west direction.

Filter	Central wavelength (μm)	HD 97048					PSF		
		Flux measured (Jy)	East (10σ) (arcsec)	West (10σ) (arcsec)	<i>FWHM</i> (east/west) (mas)	<i>FWHM</i> (north/south) (mas)	<i>FWHM</i> (east/west) (mas)	<i>FWHM</i> (north/south) (mas)	<i>FWHM</i> (diffraction) (mas)
SIV	10.49	4.2 ± 0.1	0.67	0.60	322 ± 15	337 ± 15	277 ± 15	262 ± 15	262
PAH2	11.26	7.7 ± 0.1	1.60	0.90	337 ± 15	360 ± 15	300 ± 15	285 ± 15	285
NeII	12.27	6.7 ± 0.2	0.67	0.60	390 ± 15	412 ± 15	300 ± 15	307 ± 15	307

4.2. Radiative transfer

We use the 3-dimensional Monte Carlo radiative transfer code RADMC (Dullemond & Dominik 2004a), for which a module to treat the emission from quantum-heated PAH molecules has been included. This module will be described in detail in Dullemond et al. (in prep.), but a rough description has been given by Pontoppidan et al. (2006). The code RADMC (Dullemond & Dominik 2004a; Pontoppidan & Dullemond 2005) solves the temperature structure of the disk in a Monte-Carlo way using a variant of the algorithm of Bjorkman & Wood (1997). This Monte-Carlo code also produces the source terms for scattering, in the isotropic-scattering approximation.

4.3. Dust model and PAH properties

The dust is a mixture of grains in thermal equilibrium and transiently heated PAHs.

- **Thermal grains:** the grains are composed of graphite and silicate with optical constants from Draine (1985). They have a MRN (Mathis et al. 1977) size distribution ($n(a) \propto a^{-3.5}$) with a size between $0.01 \mu\text{m}$ and $0.3 \mu\text{m}$. Since no silicate emission features have been detected in HD 97048 (see Sect. 2), we have considered the hypothesis of thermally decoupling the carbon and silicate grains (see Sect. 5.3). The optical to mid-IR opacity ratio of silicate is much smaller than that of carbon. When the grains are thermally decoupled, the low temperature of the silicates produces very weak features and the emergent flux will be dominated by emission from the carbon grains. But the absence of the silicate feature could also be due to a geometrical effect, to a lower silicate abundance, and/or larger silicate grains in the inner regions.
- **Transiently heated grains (PAH):** we can explain the observed PAH spectra, at least of the isolated HAe stars, with PAH abundances and qualitative properties similar to those of PAHs in the ISM (Habart et al. 2004a). In the ISM, PAHs are made up of a few tens up to a few hundreds of carbon atoms; for simplicity, we take only one PAH size in our model, $N_C = 100$. The hydrogen to carbon ratio is $H/C = f_H \times (6/N_C)^{0.5}$ (case of compact symmetric PAHs, see Omont 1986) with f_H – the hydrogenation fraction of the molecule. Here, we consider f_H equal to 1 (i.e., essentially fully hydrogenated PAHs) or 0.5 (partially hydrogenated PAHs). We take the absorption cross section from Li & Draine (2001) based on both laboratory data and astrophysical spectra. They consider the bands at 3.3, 8.6, 11.3, 11.9 and $12.7 \mu\text{m}$ from vibrational modes of the aromatic C-H bond; the strong bands at 6.2 and $7.7 \mu\text{m}$ due to vibrations of the aromatic C-C bonds; and a few weak features probably caused by C-C bending modes at 16.4, 18.3,

21.2 and $23.1 \mu\text{m}$. With respect to the charge, in order to keep the model simple, we assume that all PAHs are neutral. Simple determination of the ratio between the photoionisation rate of the grain to the electron-grain recombination rate suggests that this is probably the case in the outer regions ($R \geq 150 \text{ AU}$) of a disk heated by a typical HAe star (see Habart et al. (2004b) for more details). Finally, we do not take into account photo-destruction of PAHs in a strong FUV radiation field.

In summary, we adopt a dust model with large thermal grains of graphite and silicate and small transiently heated aromatic particles. The silicate abundance in dust is $[\text{Si}/\text{H}] = 3 \times 10^{-5}$, and the total carbon abundance in dust is $[\text{C}/\text{H}] = 2.2 \times 10^{-4}$. Of this 10% are in PAH and 90% in large grains.

5. Comparison between model and observations

In this section, we compare the model's predictions to the observations. In the following, we first describe the adopted stellar and disk parameters for HD 97048 and then discuss the results of the calculation and the comparison with the observations.

5.1. Stellar and disk parameters

A Kurucz model spectrum is taken for the central star with $T_{\text{eff}} = 10\,000 \text{ K}$. The stellar parameters ($L_\star = 32 L_\odot$, $M_\star = 2.5 M_\odot$) were chosen from stellar evolutionary tracks by Siess et al. (2000) for an age of 3 Myr (Lagage et al. 2006). Once the stellar flux has been reddened, the resulting SED is in agreement with the photometry extracted from Hillenbrand et al. (1992) in the UV and near-IR. To correct for extinction, we used the method of Cardelli et al. (1989) where we adopt $A_V = 1.24$ (van den Ancker et al. 1998) and $E(B - V) = 0.36$ (Davies et al. 1991; The et al. 1986).

The surface density is taken equal to $\Sigma = \Sigma_0 (r/R_{\text{in}})^{-q}$ with q the power law index equal to $3/2$ inferred for the solar nebula (Weidenschilling 1977) and $\Sigma_0 = 444 \text{ g cm}^{-2}$ a minimum value deduced from VISIR observations (Lagage et al. 2006). The inner radius R_{in} is at the dust (silicate) evaporation radius at 0.4 AU from the central star and the outer radius is at 370 AU as deduced from observations (Lagage et al. 2006). Finally, as the disk is vertically optically thick at the wavelength of observation, a minimum mass of $0.01 M_\odot$ (gas + dust) can be derived (Lagage et al. 2006). We take this minimum disk mass as a first guess.

5.2. Disk structure

Figure 4 shows the change with radius of the pressure (H_p) and photospheric (H_s) scale height. The disk flares with H_s

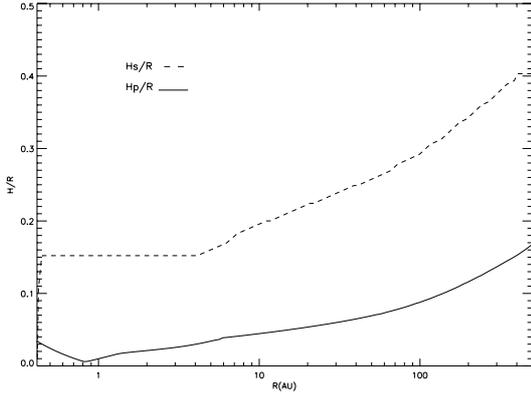


Fig. 4. Structure of the disk for the template model. The surface (dashed line) and pressure (full line) scale height of the flared disk versus the distance from the central star.

increasing with radius as $R^{9/7}$ (Chiang & Goldreich 1997). At a radius of 135 AU, $H_s = 51$ AU.

In a previous work, we used the VISIR PAH band image to constrain the parameters of the disk structure around HD 97048 (Lagage et al. 2006). Using a very simple model, Lagage et al. (2006) measured a flaring index of $1.26^{+0.05}_{-0.05}$ and $H_s = 51.3^{+0.7}_{-3.3}$ AU at 135 AU. Both values are very close to those expected from our hydrostatic, radiative equilibrium models of passive flared disks (Chiang & Goldreich 1997; Dullemond et al. 2001).

Concerning the structure of the inner region, VISIR does not have enough spatial resolution to constrain it. The puffed inner rim and the shadow region lie within the central pixel of VISIR (1 pixel is equivalent to 13.5 AU for a distance of 180 pc) and the two effects² compensate each other in terms of resulting mid-IR emission. As soon as the disk is vertically optically thick in the mid-IR, the change of the slope of the surface density in the model no longer has any influence on the structure of the disk or its mid-IR emission and cannot be constrained here.

5.3. Spectral energy distribution

Figure 5 shows the calculated spectrum in the 0.1–1300 μm range for a star/disk system inclined by 43 degrees and situated at 180 pc. At short wavelengths between 0.1 to ~ 1 μm , the emission comes from the stellar photosphere. The near-IR emission (around 3 μm) is induced by the puffed inner rim. The region that emits between 5 and 8 μm corresponds to the shadow region. In this part, the disk is not in sight of the star itself, but it receives flux from the inner rim, which is sufficiently strong to keep the disk up. At larger radii, the flaring disk reappears from the shadow, and produces the rise observed in the spectrum around 20–30 μm . Concerning the dust spectroscopic features, the PAH emission features are clearly visible, and some of them at 3.3, 6.2, 7.7 and 11.3 μm are very strong. Strong ratios between 30 and 10 μm and PAH features show that the disk intercepts a large fraction of the energy in the outer part and are evidence of a flared disk geometry. For geometrically flat disks, PAH features are predicted to be very weak, even when PAH with standard properties are present on the disk surface (Habart et al. 2004). The infrared emission under the narrow PAH features is mostly due to the large thermal grains which are very hot in the inner regions. For the model in which carbon and

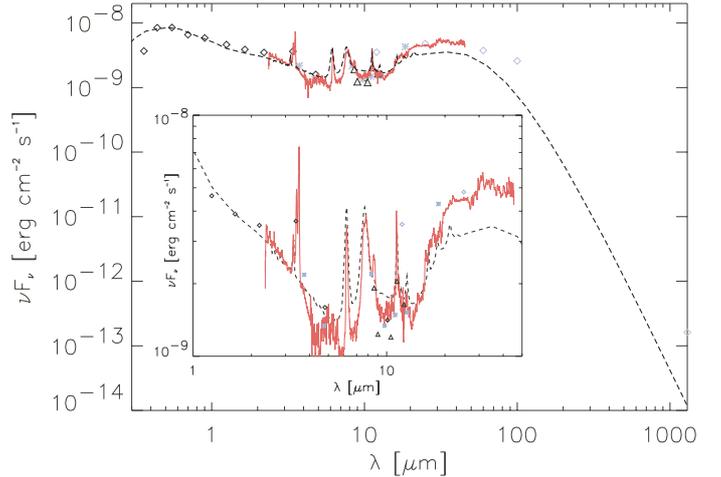


Fig. 5. SED calculated with the template model (dashed line). The disk is inclined by 43 degrees and the system is situated at a distance of 180 pc. A Kurucz model spectrum is taken for the central star with $T_{\text{eff}} = 10000$ K. The luminosity is $L_{\star} = 32 L_{\odot}$ and the mass $M_{\star} = 2.5 M_{\odot}$. The disk is flared with a total mass of $0.01 M_{\odot}$, $R_{\text{in}} = 0.41$ AU and $R_{\text{out}} = 370$ AU. Full red line shows ISO SWS spectrum of HD 97048. Points of photometry are taken from Hillenbrand et al. (1992) (open diamond black), IRAS (open diamond blue), Prusti et al. (1994) (blue crosses) and VISIR measurements (black triangle).

silicate grains are thermally coupled, it is possible to see, for example, the strong broad feature due to silicate emission peaking at about 10 μm . On the other hand, for the model in which the grains are decoupled, one can see (Fig. 5) that this feature almost disappears.

In Fig. 5, we also compare the predicted SED to that of HD 97048 constructed with different sets of data. Photometric points from the visible to the mid-IR were taken from Hillenbrand et al. (1992) and Prusti et al. (1994). To that we have added IRAS and VISIR photometric measurements. Spectroscopic observations were obtained with ISO-SWS (Short Wavelength Spectrometer) (Acke & van den Ancker 2004). The model reproduces correctly the global shape of the observed SED of HD 97048. The agreement between predicted slopes and absolute fluxes from the near- to mid-IR waves with observed ones is acceptable (differences ≤ 20 –40%). Moreover, our model also reproduces well the observed intensity of the most commonly observed PAH features, i.e., 3.3, 6.2, 7.7 and 11.3 μm , especially if we correct for contamination by the associated reflection nebula (about 40%, see Sect. 3). This is in agreement with previous comparisons made by Habart et al. (2004) between flared disk model results and observed PAH emission features with ISO and ground-based telescopes of some thirty HAeBe stars, including HD 97048.

Nevertheless, there is some mismatch between the predicted and observed PAH spectra; the strength of the 8.6 μm feature is, for example, significantly underestimated by the model. However, this is not surprising considering the uncertainties on the PAH absorption cross section (Li & Draine 2001) and our simple hypothesis that PAH are characterized by a single charge state, size or hydrogenation parameter along the disk. This is unlikely to be the case and is briefly discussed in this section. Also, HD 97048 presents peculiar strong features that peak at 3.43 and 3.53 μm not predicted by our model. Several studies have proposed attributing these features to surface C-H stretching modes on nanodiamond particles (e.g., Guillois et al. 1999; Van Kerckhoven et al. 2002; Sheu et al. 2002; Jones et al. 2004).

² Increasing the height of the inner rim augments the shadow region.

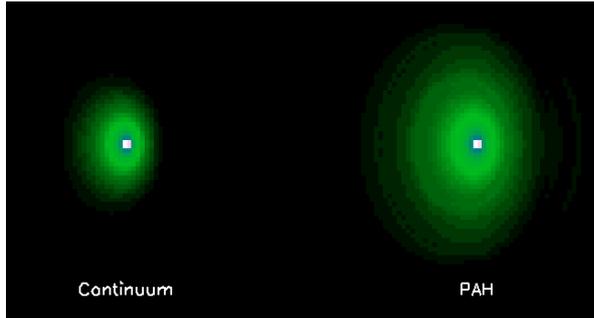


Fig. 6. The non PSF-convolved predicted image for our template model in the adjacent continuum (at $10.5 \mu\text{m}$ on the left) and in the PAH band at $11.3 \mu\text{m}$ (on the right).

Because of the good match between laboratory and observed spectra, this identification appears convincing.

The absence of the small grains of silicate features in the HD 97048 spectra is intriguing since they are the most abundant dust species in interstellar space. However, in ISO spectra, Acke & van den Ancker (2004) reported non detection of silicate features for 16 objects out of 46, showing that the absence of the silicate feature is a common phenomenon among HAeBe stars. The silicate emission at $10 \mu\text{m}$ arises from grains with a size of $0.1 \mu\text{m}$ thermally heated in the inner region ($1 < r < 10 \text{ AU}$) from direct and/or indirect irradiation by the central star. The absence of this emission could probably result from various effects concerning either the dust properties either the disk geometry in the inner region. Here, we have considered the thermal decoupling hypothesis but it could also be due to less silicate abundance in the inner disk. Globally decreasing the silicate abundance is not a solution since the 20 micron feature shows the presence of silicates in the $50 \text{ AU} < R < 100 \text{ AU}$ parts of the disk. We have also tested different modeled structures of a disk and we have shown that it is possible to explain the absence of the silicate emission at $10 \mu\text{m}$ with a geometrical effect. But at the moment, the structure is calculated self-consistently and the aim of the paper is not to construct a structure by hand that could fit all the data of HD 97048 but that is not physical. This issue could not be investigated with the resolution of VISIR observations.

Concerning the longer wavelength, we find that the submm flux is too low by a factor of 70 compared to observations (Henning & Launhardt 1998). This shows that a large reservoir of large grains (around a millimeter size) must exist in the outer regions of the disk. These grains would naturally reside close to the midplane and therefore do not affect the shape of the disk. They are therefore not within the focus of this paper, and for this reason we do not include them in our model. If we would have included them, they would only affect these long-wavelength fluxes, because the disk has a flaring geometry (see e.g. Dullemond & Dominik 2004a, Figs. 6 and 7).

5.4. Imaging

Figure 6 shows the modelled disk image in the PAH band at $11.3 \mu\text{m}$ and the adjacent continuum. These emissions both originate from the optically thin surface of the disk but the PAH emission is much more extended than the adjacent continuum. The continuum reaches 50% of its integrated intensity at a very small radius (about 2 AU), while the PAH feature does so at large radii (about 100 AU). This behaviour reflects the different excitation mechanisms of grains in thermal equilibrium and transiently heated PAHs. Only grains very near the star are warm enough to

emit at $11 \mu\text{m}$ whereas PAHs farther away can be excited and emit in the $11.3 \mu\text{m}$ feature. The emission at $11.3 \mu\text{m}$ (band + continuum) in the central part of the disk ($< 2 \text{ AU}$) is dominated by the thermal emission from silicate and carbon grains, whereas in the outer regions of the disk, the emission is dominated by the PAH feature. Moreover, among the PAH features, the $11.3 \mu\text{m}$ is one of the most extended spatially. Indeed, as discussed in Habart et al. (2004b), the features at shorter wavelengths are stronger in the inner part of the disk (where PAHs are hotter because multiphoton events occur where the radiation field is most intense), decreasing rapidly in the outer cold part. On the other hand, the features at longer wavelengths are more extended.

Figure 7 shows the brightness emission profiles of the PAH feature and the adjacent continuum obtained by convolving the model with the corresponding PSF observed by VISIR. The model predicts a spatial distribution of the PAH emission very similar to that observed. The predicted *FWHM* and wing extension are very close to the observed ones (differences $\leq 20\%$). Moreover, our disk model predicts, as in the observations, an asymmetry in the east/west direction, which results from the inclination of a flared disk that is optically thick at the observed wavelength. This acceptable agreement gives strong support for the physics underlying our flared disk model. The disk parameter that most affects the PAH emission is the disk flaring angle, which determines at each radius the fraction of FUV intercepted by the disk surface. Lower values of the flaring could be caused by a variety of reasons, for example if the dust settles toward the disk midplane (Dullemond & Dominik 2004b). If the disk is less flared, the PAH emission that directly tracks the illumination of the disk surface will be strongly reduced in the outer disk region. Less flared disks will have less extended and weaker PAH emission features. This will be particularly true for the features at long wavelengths, such as the $11.3 \mu\text{m}$ one, which have a large contribution from the outer disk. In the extreme case of a fully self-shadowed disk, the PAH feature strengths should decrease by orders of magnitude and the spatial distribution should be similar to that of the adjacent continuum.

As predicted by the model, the observed spatial extension of the $11.3 \mu\text{m}$ feature is much larger than that observed for the PAH feature at $3.3 \mu\text{m}$, extended on a scale of (several) 10 AU (Habart et al. 2004a). This has the interesting implication that PAHs appears to be present over a large range of radii; in other words, PAHs can survive over a wide range of physical conditions. Concerning the spatial extent of the $10 \mu\text{m}$ adjacent continuum emission, the model predicts that it is slightly broader than the PSF but still agrees with the observations.

There are several complications that we have neglected. The most obvious is that we have assumed that PAHs can be characterized by a single size, hydrogenation and charge state. This is unlikely to be the case, and one can expect variations as a function of radius and depth in the disk. For example, PAHs are likely to be more positively ionized in the inner disk regions. Moreover, processes such as photo-evaporation or coagulation could affect the abundance and size of PAHs. In order to obtain some insight into the specific PAH properties, one needs spatial information of several band strength ratios. We are developing this study for a forthcoming paper.

6. Summary and conclusions

In Lagage et al. (2006), we were able to constrain for the first time the flaring geometry of a disk around an intermediate-mass young star HD 97048. These results were based on a very simple

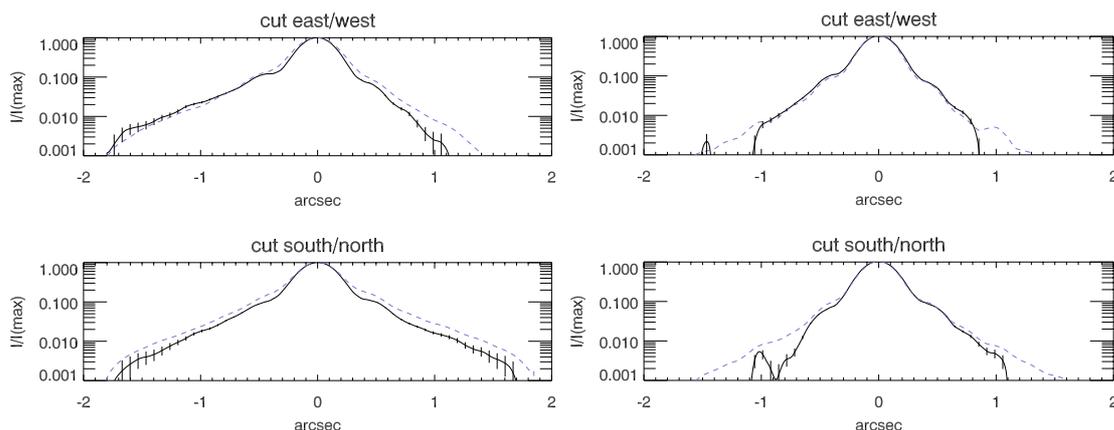


Fig. 7. *Left*, cut in north/south and east/west (*up panel*) for the template model of the PAH band at $11.3 \mu\text{m}$ convolved with VISIR PSF (dashed line) compared to the observation (full line). *Right*, the same as on the left with the adjacent continuum at $10.5 \mu\text{m}$. The errors drawn are $1-\sigma$ errors rms.

model making several assumptions such as surface PAH emission, an optically thick disk, and a power-law function for the surface height and the intensity. In the present paper, using a full radiative transfer model based on predicted disk geometries assuming hydrostatic equilibrium (Dullemond et al. 2001; Dominik et al. 2003; Dullemond & Dominik 2004a), we could:

- justify the hypothesis made in the previous paper (Lagage et al. 2006), and therefore confirm the results, e.g., the images calculated with a different code (Pinte et al. 2006);
- show that both SED and spatial distribution of the PAH emission and the adjacent continuum are in very good agreement with the model predictions for common disk parameters;
- take the general agreement between observations and predictions as a strong support for the physical considerations underlying our flared disk model.

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