

Optical *BVI* imaging and H I synthesis observations of the dwarf irregular Galaxy ESO 364-G029

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ABSTRACT

As part of an effort to enlarge the number of well-studied Magellanic-type galaxies, we obtained broadband optical imaging and neutral hydrogen radio synthesis observations of the dwarf irregular galaxy ESO 364-G029. The optical morphology characteristically shows a bar-like main body with a one-sided spiral arm, an approximately exponential light distribution, and offset photometric and kinematic centers. The H I distribution is mildly asymmetric and, although slightly offset from the photometric center, roughly follows the optical brightness distribution, extending to over 1.2 Holmberg radii (where $\mu_B = 26.5$ mag arcsec⁻²). In particular, the highest H I column densities closely follow the bar, one-arm spiral, and a third optical extension. The rotation is solid-body in the inner parts but flattens outside of the optical extent. The total H I flux $F_{\text{HI}} = 23.1 \pm 1.2$ Jy km s⁻¹, yielding a total H I mass $M_{\text{HI}} = (6.4 \pm 1.7) \times 10^8 M_\odot$ (for a distance $D = 10.8 \pm 1.4$ Mpc) and a total H I mass-to-blue-luminosity ratio $M_{\text{HI}}/L_B = (0.96 \pm 0.14) M_\odot/L_{B,\odot}$ (distance independent). The H I data suggest a very complex small-scale H I structure, with evidence of large shells and/or holes, but deeper observations are required for a detailed study. Follow-up observations are also desirable for a proper comparison with the Large Magellanic Cloud, where despite an optical morphology very similar to ESO 364-G029 the H I bears little resemblance to the optical.

Key words. galaxies: individual: ESO 364-G029 – galaxies: irregular – galaxies: photometry – galaxies: kinematics and dynamics – galaxies: structure – galaxies: ISM

1. Introduction

In the currently popular bottom-up galaxy and structure formation scenarios, studies of gas-rich low-metallicity dwarf galaxies are invaluable, as these objects must have dominated the Universe in the past (e.g. White & Frenk 1991; Kauffmann et al. 1993; Cole et al. 1994; but see also Chiu et al. 2001). They are also expected to be more uniformly distributed than their larger counterparts (e.g. Dekel & Silk 1986) and may in fact still be the most common type of object in the Universe (e.g. Mateo 1998; Im et al. 1999; Driver 1999).

The Magellanic-type spirals (Sm; with rotational symmetry and some spiral structure) and Magellanic irregulars (Im; asymmetric with no spiral structure) are at the transition between fully fledged spirals and true dwarf irregulars (de Vaucouleurs 1956, 1959). With respect to the latter two groups of galaxies, Sm and Im galaxies have intermediate physical properties such as rotation velocity, nuclear concentration, colour, and neutral hydrogen content. Although the prototype of the class, the Large Magellanic Cloud (LMC; classified as SB(s)m in NED¹), is amongst the best studied galaxies (e.g. Westerlund 1997; van der Marel et al. 2002; Kim et al. 2003; Brüns et al. 2005;

van der Marel 2006, and references therein), the general properties of Magellanic systems are surprisingly poorly known. The classical reference on the subject remains that of de Vaucouleurs & Freeman (1972). Odewahn (1991, 1994) provides more modern studies and argues that it is a misconception to consider the LMC and other galaxies like it as irregular. Magellanic-type galaxies are characterized by an asymmetric spiral arm connected at one end to a high surface brightness bar. The bar center is often offset from the center of the galaxy as defined by the outer optical isophotes, and the arm normally has a clumpy appearance, presumably due to triggered star formation. The arm can sometimes be followed completely around the galaxy and, in some cases, small “embryonic” arms are present at the ends of the bar (de Vaucouleurs 1955).

Various simulations have shown that one-armed morphology can develop from a strong tidal encounter with a companion (e.g. Byrd et al. 1986; Howard & Byrd 1990). The models generally predict only a short-lived one-armed spiral structure, implying that a companion galaxy should generally be present close to any Magellanic galaxy. Using a sample from the Third Reference Catalog of Bright Galaxies (RC3 1991, hereafter RC3), Odewahn (1994) demonstrated that close companions are indeed observed in almost every case. As for the LMC itself, its

¹ NASA/IPAC Extragalactic Database.

interaction with both the Milky Way and the Small Magellanic Cloud (SMC) is most likely at the origin of its one-arm structure and large-scale disturbed morphology (e.g. Putman et al. 1998, 2003, and references therein).

The kinematics of the warp in the HI distribution of the Magellanic-type galaxy NGC 3109 and the nearby Antlia dwarf galaxy also suggests that the galaxies had a mild interaction about 1 Gyr ago (Barnes & de Blok 2001). Based on the morphology and kinematics, Bush & Wilcots (2004) draw a similar conclusion for the neighbours NGC 4618 and NGC 4625. However, not all Magellanic-type galaxies have an obvious neighbour for a recent interaction. Bekki & Chiba (2006) suggest that the morphologies of these apparently isolated galaxies may be explained by a recent interaction with an (optically) invisible companion. The HI survey performed by Doyle et al. (2005), however, shows that evidence for the existence of these “dark galaxies” is marginal.

In order to address the asymmetry issue, Wilcots & Prescott (2004) performed an HI survey of a sample of 13 Magellanic spiral galaxies with apparent optical companions. In their study, they find that only four of these have confirmed HI-detected neighbours. Their study also indicates that the presence of companions near NGC 2537 and UGC 5391 has no effect on the morphology of these galaxies, and that Magellanic spirals are no more asymmetric than a random sample in the field. The latter conclusion is supported by HI observations of the interacting Magellanic spirals NGC 4618 and NGC 4625 (Bush & Wilcots 2004). On the other hand, in their HI study of (possibly) interacting Magellanic galaxies in the M 81 group, Bureau et al. (2004) find no large-scale tidal feature and no intergalactic HI cloud near the Magellanic dwarf M81dwA, and remark that this galaxy could be a tidal dwarf galaxy (rather than a perturbed “regular” dwarf), resulting from the debris of a past encounter with Holmberg II.

Magellanic dwarf galaxies are also known to contain a large number of shells and holes in their interstellar medium. These structures are created by winds and supernova explosions of the most massive stars in star forming regions. Unlike in spiral systems, these structures are long-lived in dwarf galaxies, due to solid-body rotation and a lack density of waves. Well-known dwarfs exhibiting shells and holes include the LMC (SB(s)m; Kim et al. 1998), IC 2574 (SAB(s)m; Walter & Brinks 1999) and Holmberg II (Im; Puche et al. 1992; but see also Rhode et al. 1999; Bureau & Carignan 2002).

A significant amount of neutral hydrogen is known to be present beyond the optical extent of Magellanic-type galaxies, often up to 2–3 Holmberg radii (e.g. NGC 925; Pisano et al. 1998; Pisano & Wilcots 2000; DDO 43; Simpson et al. 2005), occasionally reaching up to 6 Holmberg radii (NGC 4449; Hunter et al. 1998). The HI distribution of several Magellanic dwarfs is known to exhibit further irregularities, such as HI loops surrounding the galaxy (NGC 4618; Bush & Wilcots 2004) and external spurs or blobs (e.g. NGC 5169; Mühle et al. 2005; UGCA 98; Stil et al. 2005). An S-shaped distortion in the HI velocity field, possibly indicating counterrotation, has been observed in NGC 4449 (Hunter et al. 1998) and DDO 43 (Simpson et al. 2005).

The asymmetries and irregularities in the HI distribution of Magellanic-type galaxies are often associated with a high star formation rate, such as for IC 10 (e.g. Wilcots & Miller 1998; see also Thurow & Wilcots 2005). Using H α observations, Wilcots & Thurow (2001) find that, in several regions of the Magellanic spiral NGC 4214, the velocity of the ionized gas is 50–100 km s⁻¹ higher than that of the HI, indicating massive

star formation. The study of NGC 4449 by Hunter et al. (2000) compares H α , H $_2$, HI and near-infrared emission, and suggests that the different regions of this galaxies are in different stages of star formation.

As part of an effort to increase the number of Magellanic dwarfs with detailed observations, we focus in this paper on the relatively nearby Magellanic dwarf irregular galaxy ESO 364-G029 (PGC 018396). ESO 364-G029 was first reported by Holmberg et al. (1978), who classified it as a “dwarf irregular or emission nebulae” using the ESO (B) Atlas. By combining optical data from the UK Schmidt plates with Parkes HI observations, Longmore et al. (1982) concluded that ESO 364-G029 is a dwarf irregular galaxy, although Karachentseva & Karachentsev (2000) designate the object as “dwarf irregular or reflection nebula” in their optical search for nearby dwarf galaxy candidates in the Southern hemisphere. ESO 364-G029 was identified (and rejected) as a possible companion of the nearby dwarf galaxy NGC 2188 by Domgörgen et al. (1996). Zimmermann et al. (2001) reported a possible X-ray counterpart to ESO 364-G029 using ROSAT data. ESO 364-G029 appears in a handful of other surveys, such as the Nearby Optical Galaxies Catalogue (Giuricin et al. 2000), the Catalog of Nearby Galaxies (Karachentsev et al. 2004), the HIPASS² catalogue (Huchtmeier et al. 2001; Meyer et al. 2004; Koribalski et al. 2004), and the surveys for optical/HI counterparts of Paturel et al. (2005) and Doyle et al. (2005). Although the general properties (such as luminosity and HI content) of ESO 364-G029 are reasonably well constrained by the above-mentioned surveys, the detailed morphology and kinematics of the galaxy were not studied in detail previously. General properties of ESO 364-G029 are listed in Table 1 and an optical image from the Digitized Sky Survey (DSS) is shown in Fig. 1.

This paper is organized as follows. Optical observations and surface photometry of ESO 364-G029 are presented in Sect. 2 while HI spectral line-imaging is discussed in Sect. 3. In Sect. 4 we constrain the distance to ESO 364-G029. The discussion in Sect. 5 focuses on a comparison with the LMC, the importance of interactions, and the possible presence of HI shells and holes in ESO 364-G029. Finally, our main conclusions are briefly summarized in Sect. 6.

2. Optical data

2.1. Observations and data reduction

Optical imaging observations were obtained on the nights of 9–11 April 1997 using the 1-m telescope of Siding Spring Observatory (SSO) in Australia. A total of 30 broadband images combining short and long exposures were obtained in the *B*, *V* and *I* Cousins filter bands. A 2048 × 2048 thinned TEK CCD with 15 μ m pixels was used in direct imaging, yielding a spatial sampling of 0''.61 pix⁻¹. A 1644 × 1000 CCD subsection with 100 pixels virtual overscan was read, for a total field-of-view of 15''.5 × 10''.0. The seeing was moderate (1''.7–2''.4) but sufficient for basic photometry and comparison with our HI observations.

The data reduction was carried out using standard procedures in IRAF³ (Tody 1986, 1993). All the images were overscan-subtracted, bias-subtracted using bias scans obtained

² HI Parkes All-Sky Survey:

<http://www.atnf.csiro.au/research/multibeam/release/>

³ IRAF is distributed by the National Optical Astronomy Observatories (NOAO), which are operated by the Association of Universities for Research in Astronomy (AURA), Inc., under cooperative agreement with the National Science Foundation (NSF).

Table 1. Basic properties of the dwarf irregular galaxy ESO 364-G029. The left and middle columns list the different quantities and their values; the right column lists corresponding references – (1) NED; (2) HyperLEDA; (3) This paper; (4) Karachentsev et al. (2004).

Quantity	Value	R.
Morphological type	IB(s)m	1
Right ascension (J2000)	$\alpha = 06^{\text{h}}05^{\text{m}}45.^{\text{s}}2$	1
Declination (J2000)	$\delta = -33^{\circ}04'51''$	1
Galactic longitude	$l = 239.^{\circ}47$	1
Galactic latitude	$b = -23.^{\circ}36$	1
Heliocentric radial velocity	$V_{\odot} = 786 \pm 11 \text{ km s}^{-1}$	2
(from HI measurements)	$V_{\odot} = 784 \pm 2 \text{ km s}^{-1}$	3
Scale length (B)	$h_B = 50'' \pm 5''$	3
Radius (25 B mag arcsec $^{-2}$)	$R_{25} = 1:29 \pm 0:09$	2
	$R_{25} = 1:27 \pm 0:03$	3
Radius (26.5 B mag arcsec $^{-2}$)	$R_{\text{Ho}} = 2:30 \pm 0:08$	3
Inclination	$i = 70.^{\circ}5$	2
Axial ratio (25 B mag arcsec $^{-2}$)	$q_{25} = 0.51 \pm 0.06$	2
	$q_{25} = 0.54 \pm 0.04$	3
Asymmetry	$Q_{\text{asym}} = 1.7\%$	3
Position angle (25 B mag arcsec $^{-2}$)	PA = $55.^{\circ}4$	2
	PA = $62.^{\circ} \pm 4.^{\circ}$	3
Central surface brightness	$\mu_B(0) = 23.3 \text{ mag arcsec}^{-2}$	3
Apparent total B mag	$B_T = 13.81 \pm 0.22 \text{ mag}$	2
	$B_T = 13.8 \pm 0.1 \text{ mag}$	3
Corrected apparent B mag	$B_c = 13.8 \pm 0.22 \text{ mag}$	2
	$B_c = 13.6 \pm 0.1 \text{ mag}$	3
Corrected absolute B mag	$M_B = -16.44 \text{ mag}$	2
	$M_B = -16.6 \pm 0.3 \text{ mag}$	3
Distance	$D = 8.02 \text{ Mpc}$	2
	$D = 10.8 \pm 1.4 \text{ Mpc}$	3
	$D = 7.7 \text{ Mpc}$	4
Scale	$1' = 2.3 \text{ kpc}$	2
	$1' = 3.1 \pm 0.4 \text{ kpc}$	3
	$1' = 2.2 \text{ kpc}$	4
Total HI mass	$M_{\text{HI}} = (6.4 \pm 1.7) \times 10^8 M_{\odot}$	3
HI mass-to-blue-luminosity ratio	$M_{\text{HI}}/L_B = 0.96 \pm 0.14 M_{\odot}/L_{B,\odot}$	3

every night, and flatfielded using twilight flatfields. Dark current corrections were not necessary due to the low dark current levels and the relatively short exposure times. The cosmetic of the CCD is good, with only one bad column and few bad pixels. All target images in a given filter were then registered and combined using a pixel rejection algorithm, convolving all images to a common (largest) seeing in each filter. Because ESO 364-G029 is faint but lies in a region with many bright foreground stars, all pixels with counts above 20% of the saturation level were rejected in the long exposure images, to prevent excessive bleeding. Short exposures fill in the rejected regions, easing the masking and interpolation of bright foreground stars (see below).

Three sets of carefully selected and corrected Landolt $UBVRI$ standards (Landolt 1992; Bessell 1995) were used for photometric calibration. Since we do not require precise photometry, we adopted the averaged extinction coefficients for the SSO 1-m telescope. The resulting combined and calibrated images are shown in Fig. 2.

2.2. Surface photometry

Standard ellipse fitting using the *Ellipse* and *Isophote* packages in IRAF (see Jedrzejewski 1987) was used to derive the azimuthally-averaged surface brightness profile of ESO 364-G029 in each filter. For each semi-major axis a , the ellipses are characterized by the central coordinates (x_0, y_0) , the

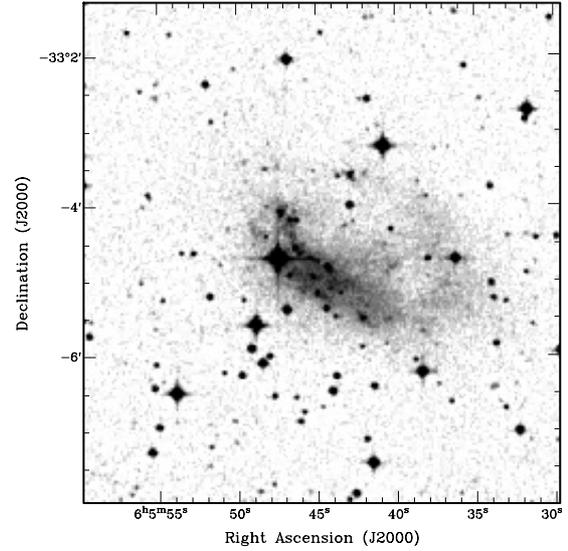


Fig. 1. Optical image of ESO 364-G029 from the Digitized Sky Survey (DSS); cf. our BVI images in Fig. 2.

ellipticity ϵ (defined as $\epsilon \equiv 1 - b/a$, where a and b are the semi-major and semi-minor axis, respectively) and the position angle PA, measured from North to East.

The derived surface brightness profiles should be used with caution, for the following reasons. First, ESO 364-G029's irregular and lopsided morphology renders any azimuthally-averaged profile at best a modest representation of its two-dimensional light distribution. Second, many bright foreground stars are present in the field of ESO 364-G029. Particularly troublesome is the bright star TYC 7075-383-1 located near the North-East edge of the central bar at $(\alpha = 06^{\text{h}}05^{\text{m}}47.^{\text{s}}6, \delta = -33^{\circ}04'41''.5; \text{J2000})$, which has $B = 12.1 \pm 0.1$ and $V = 11.4 \pm 0.1$ mag (Kharchenko 2001). The most elegant way to remove the foreground stars is to subtract the stellar profiles from the image. However, the brightest stars in the field are saturated and show bleeding effects even in our shortest exposures, so we masked out the brightest foreground stars, and interpolated the surface brightness in the masked regions. As the images in each of the three filters have a comparable resolution, the same regions were used for all filters.

Mainly because of ESO 364-G029's irregular morphology, it was difficult to obtain reliable and stable fits for the whole range of possible semi-major axes ($3''$ to $155''$). Our results are shown in Figs. 3–5. The I -band images are less deep because of the relatively bright background. In order to get meaningful colors, the ellipses were fitted on the B image only and imposed on the V and I band images. One can see that the ellipticity and position angle of the ellipses were successfully fitted over a limited range of semi-major axes only, $15'' \lesssim a \lesssim 83''$. For lack of a better prescription, they are assumed constant outside of those limits.

The light profiles are approximately exponential over most of the optical extent ($a \lesssim 135''$), falling off with roughly $1.38 \text{ mag arcmin}^{-1}$. A break to a steeper fall-off is seen at $a \approx 135''$. A fit to the B profile within that region yields a radial scale length $h_B = 50'' \pm 5''$ and an (uncorrected) extrapolated central surface brightness $\mu_B(0) = 23.3 \text{ mag arcsec}^{-2}$. The standard criterion for defining low surface brightness galaxies is for the extrapolated central surface brightness (based on the outer exponential disk) to be equal or fainter than $23.0 \text{ mag arcsec}^{-2}$ in B (Impey & Bothun 1997). Following this definition, ESO 364-G029 would qualify as a low surface

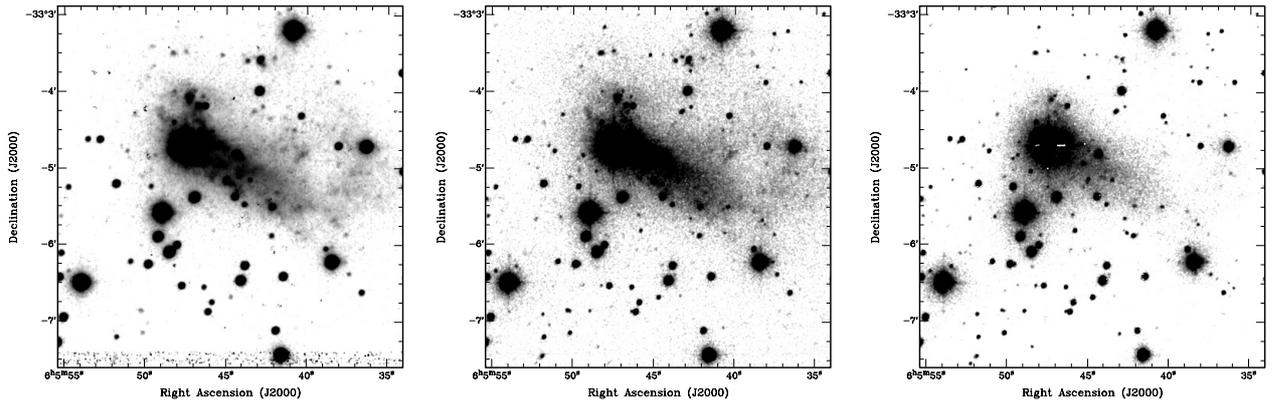


Fig. 2. Logarithmic grayscale images of ESO 364-G029 in the Cousins B (left), V (middle) and I (right) filters. The central bar and one spiral arm characterize ESO 364-G029 as a Magellanic-type dwarf galaxy.

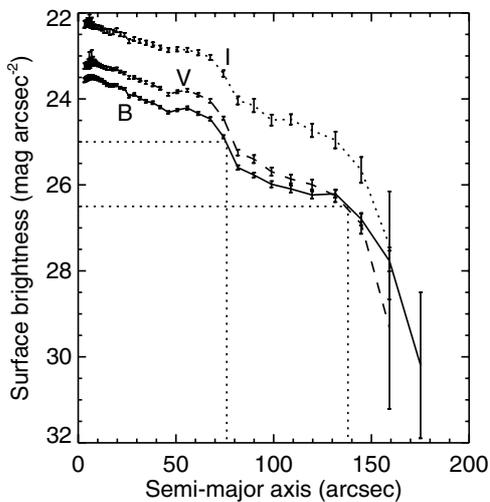


Fig. 3. The azimuthally-averaged surface brightness profiles of ESO 364-G029 in the B band (solid curve), the V band (dashed curve) and the I band (dotted curve). The error bars indicate the 1σ formal errors. Systematic errors due to the irregular morphology of ESO 364-G029 are probably much larger. The dotted lines indicate the semi-major axis of ESO 364-G029 at the $B = 25$ and 26.5 mag arcsec $^{-2}$ isophotes, respectively: $R_{25} = 1:27 \pm 0:03$ and $R_{H0} = 2:30 \pm 0:08$.

brightness galaxy. The B surface brightness profile also suggests a radius at the 25 mag arcsec $^{-2}$ isophote of $R_{25} = 1:27 \pm 0.03$, identical to the values listed in HyperLEDA⁴ and the RC3. The Holmberg radius (defined at the 26.5 mag arcsec $^{-2}$ isophote in B ; see Holmberg 1958) is $R_{H0} = 2:30 \pm 0:08$.

Over most of the disk $B - V = 0.4 \pm 0.08$ mag, slightly bluer in the very inner and outer parts. This is marginally bluer than the LMC ($B - V = 0.51$ mag; HyperLEDA) but similar to NGC 4618 (SBm; Odewahn 1991). The $V - I$ color increases from 0.85 ± 0.08 mag in the inner parts to 1.3 ± 0.1 mag in the outer parts, again similar to NGC 4618.

Constructing growth curves from the surface brightness profiles shown in Fig. 3, the extrapolated total apparent magnitudes are $B_T = 13.8 \pm 0.1$, $V_T = 13.5 \pm 0.1$ and $I_T = 12.4 \pm 0.15$ mag, in good agreement with the values listed in the RC3 ($B_T = 13.58 \pm 0.21$ mag), HyperLEDA ($B_T = 13.81 \pm 0.22$ mag) and The Surface Photometry Catalogue of the ESO-Uppsala Galaxies ($B_T = 13.6 \pm 0.1$ mag; Lauberts & Valentijn 1989). Correcting for Galactic extinction following

Schlegel et al. (1998), but neglecting the inclination-dependent internal absorption, we obtain the following corrected total apparent magnitudes: $B_c \approx 13.6 \pm 0.1$, $V_c \approx 13.4$ and $I_c \approx 12.3$ mag.

3. HI radio synthesis data

3.1. Observations and data reduction

Radio synthesis data of ESO 364-G029 were obtained with four different configurations of the Australia Telescope Compact Array (ATCA). Full 12 h syntheses were obtained with the 375, 750C and 750E arrays, while shorter observations (3–4 h on-source) were also obtained twice in the 1.5D array. Antenna CA06, providing the longest possible baselines, was not used in the 375, 750C and 750E arrays. Both HI line and continuum data (centered at 1380 MHz) were obtained simultaneously in two polarizations, but only the line data will be discussed here. In the 375, 750C and 750E configurations, 512 channels in each polarization were used, each 1.65 km s $^{-1}$ wide, for a total bandwidth of 844 km s $^{-1}$. In the 1.5D configuration, 256 channels of 6.60 km s $^{-1}$ width were used in each polarization for a total coverage of 1689 km s $^{-1}$. These and other relevant observing parameters are listed in Table 2.

The data were reduced using standard procedures in MIRIAD⁵ (Sault et al. 1995) and visualized with KARMA (Gooch 1995). The strong source PKS 1934-638 (Reynolds 1994) was observed once briefly in each configuration for use as a flux calibrator. PKS 0614-349, located $1^{\circ}26'$ to the South-East of ESO 364-G029, was observed roughly every 50 min for use as a complex gain and bandpass calibrator.

The (ATCA-specific) self-interference channels were first removed, after which the line data were transformed to heliocentric rest frame and Hanning smoothed. Bad data, mostly due to short duration interference, was then masked out manually and the flux density, complex gains, and bandpass were calibrated. The continuum emission was subtracted in the uv plane using a linear fit to line-free channels on each side of ESO 364-G029's emission (excluding bandpass-affected channels), and the uv data were then imaged to create dirty cubes.

Both uniform- and natural-weighted cubes were created and a $33' \times 33'$ box (equal to the primary beam diameter) was imaged. In each case the pixel size was chosen to sample the

⁴ HyperLEDA: <http://leda.univ-lyon1.fr>

⁵ MIRIAD was developed by the Australia Telescope National Facility (ATNF) which is part of the Commonwealth Scientific and Industrial Research Organization (CSIRO) of Australia.

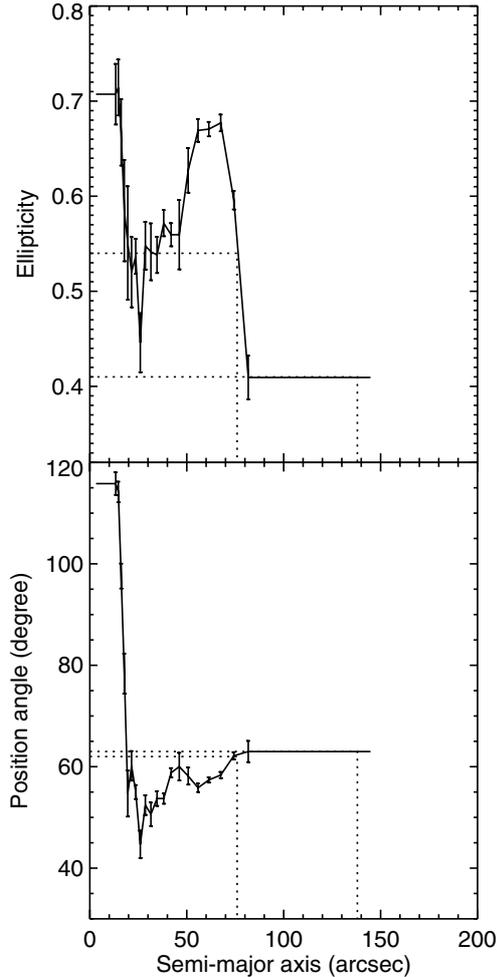


Fig. 4. Ellipticity ϵ (*top*) and position angle PA (*bottom*) profiles of ESO 364-G029, derived with the azimuthally-averaged B surface brightness profile shown in Fig. 3. Due to the irregular morphology of ESO 364-G029 it was not possible to successfully fit ellipses at radii larger than $83''$. ϵ and PA are assumed to be constant for radii larger than this value. The error bars indicate 1σ formal errors. Systematic errors due to the irregular morphology of ESO 364-G029 are probably much larger. The dotted lines in both panels indicate the ellipticity and position angle at radii R_{25} and R_{H_0} .

synthesized beam with 2 to 3 pixels. The four datacubes used in our analysis are listed in Table 3, and are hereafter abbreviated as NH (naturally-weighted, high-sensitivity datacube), UH (uniformly weighted, high-sensitivity datacube), NL (naturally weighted, low-sensitivity datacube) and UL (uniformly weighted, low-sensitivity datacube). For datacubes UH and NH we merge the 375, 750C and 750E arrays only, while for datacubes UL and NL we combine the data for all arrays. Note that the low-sensitivity datasets (NL and UL) have a significantly higher spatial resolution. In the analysis below we use the the high-sensitivity datasets for global properties and the high-resolution datasets for the small-scale morphology. Table 3 lists the root-mean-square (rms) noise per channel and synthesized beam for each dataset. When considering the noise levels reached, one should remember that the channel width for the 1.5D observations was four times as large as that for the other arrays, although these observations were regridded onto the narrower 1.65 km s^{-1} channels when combined with the shorter arrays.

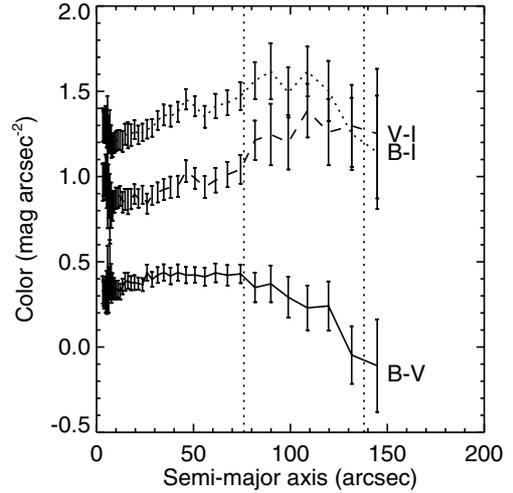


Fig. 5. Color profiles of ESO 364-G029, derived from the azimuthally-averaged surface brightness profiles shown in Fig. 3: $B-V$ (*solid curve*), $V-I$ (*dashed curve*) and $B-I$ (*dotted curve*). The error bars indicate the 1σ formal errors. Systematic errors due to the irregular morphology of ESO 364-G029 are probably much larger. Vertical dotted lines are plotted at radii R_{25} and R_{H_0} .

Table 2. Parameters of the ATCA HI observations. Observations were obtained in four different configurations, and combined (see Sect. 3) for analysis.

ATCA array	Date of obs.	Maximum baseline
375	1997 Oct. 3–4	459 m
750C	1997 Oct. 25–26	750 m
750E	1998 May 30	643 m
1.5D	1997 Mar. 8, 10	3000 m

Property	1.5D array	375, 750C, 750E arrays
Primary beam	$HPBW \approx 33'$	$HPBW \approx 33'$
Correlator	FULL_8_512-128	FULL_4_1024-128
Channel width	31.3 kHz (6.60 km s^{-1})	7.8 kHz (1.65 km s^{-1})
Bandwidth	8.00 MHz (1689 km s^{-1})	4.00 MHz (844 km s^{-1})
Flux calibrator	PKS 1934-638	PKS 1934-638
Phase calibrator	PKS 0614-349	PKS 0614-349

Note: Antenna CA06 was not used with the 375, 750C and 750E arrays.

The Clark Clean algorithm (Clark 1980) was used for the deconvolution, cleaning an area tightly encompassing the emission in each channel. The maps were cleaned until the total cleaned flux stopped increasing, typically at a depth of 2.5 times the rms noise per channel, which was sufficient to remove all side-lobes. The clean components were then restored and added to the residuals using a Gaussian clean beam as listed in Table 3 (and resulting from a fit to the inner parts of the dirty beam). Channel maps of datacube UL (with the best spatial resolution) are shown in Fig. 6, where each map displayed is the average of 5 channels to increase sensitivity.

3.2. Global HI profile

To maximize the flux recovered, the (masked) NH datacube was used to derive the global profile of ESO 364-G029, which is shown in Fig. 7. The linewidth (uncorrected for instrumental broadening) at 20% of the peak is $\Delta V_{20}^{\text{obs}} = 89 \pm 2 \text{ km s}^{-1}$, while

Table 3. Properties of the four datacubes used in our analysis. For the analysis of the detailed morphology of ESO 364-G029 we use the long-baseline datacubes NL and UL, which have a high spatial resolution. For the derivation of the global structure of ESO 364-G029 we use the datacubes NH and UH, which contain only the short baseline observations. The latter datasets are therefore deeper than the former, but have lower spatial resolutions.

Datacube	Properties	Arrays used	Pixel size (arcsec)	Weighting	rms noise per channel (mJy beam ⁻¹)	Beam properties <i>FWHM</i> (arcsec)	φ (deg)
Datacube NH	high sensitivity, low resolution	375, 750C and 750E	16 × 16	natural	2.0	132 × 70	-1
Datacube UH	high sensitivity, low resolution	375, 750C and 750E	16 × 16	uniform	3.3	72 × 45	0
Datacube NL	low sensitivity, high resolution	All arrays	4 × 4	natural	1.5	58 × 38	+8
Datacube UL	low sensitivity, high resolution	All arrays	4 × 4	uniform	2.6	21 × 10	+30

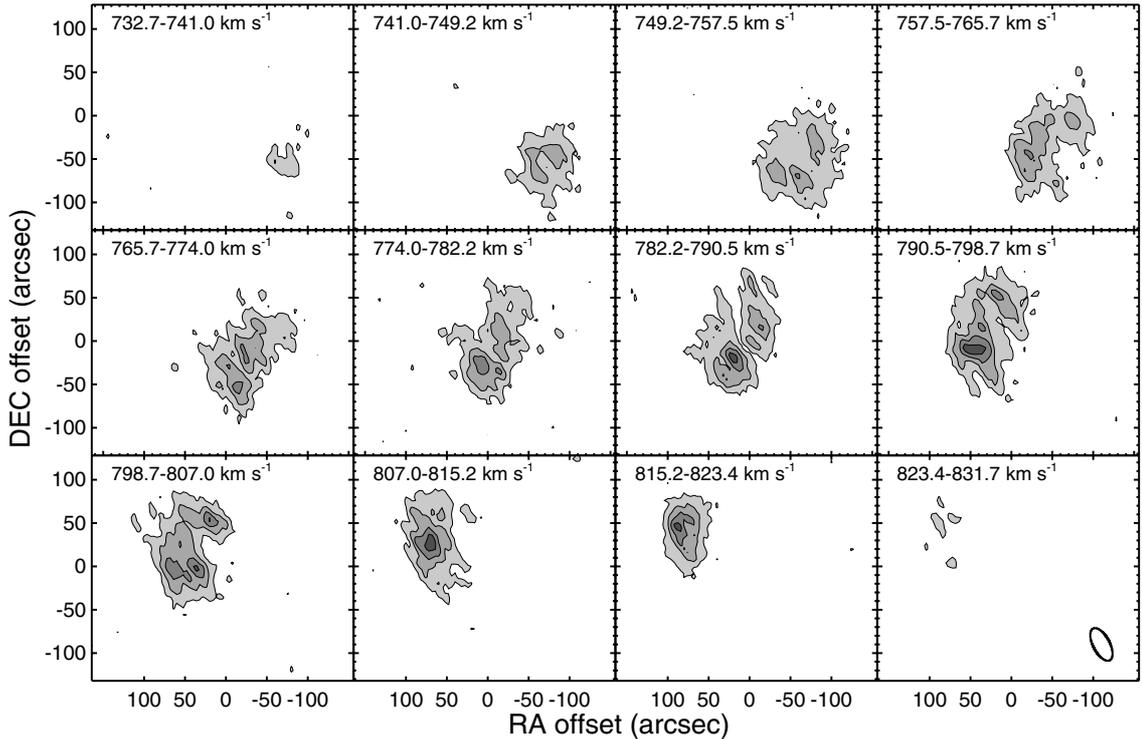


Fig. 6. HI channel maps of ESO 364-G029 for datacube UL (see Table 3). Each map is the average of 5 channels to increase the sensitivity (8.25 km s⁻¹ per map). Contour levels are at 3, 6, 9 and 12 times the rms noise in each map (1.2 mJy beam). The heliocentric velocity range of each map is indicated in the top-left corner. The synthesized beam is 21'' × 10'' with a position angle of 30°, and is indicated in the bottom-right panel.

that at 50% $\Delta V_{50}^{\text{obs}} = 73 \pm 2$ km s⁻¹. We adopt a systemic velocity, taken to be the midpoint of the velocities at 20% of the peak, of $V_{\text{sys}} = 784 \pm 2$ km s⁻¹. An intensity-weighted mean of the entire profile gives $\langle V \rangle = 785 \pm 2$ km s⁻¹, consistent with the former value. Longmore et al. (1982) obtained $\langle V \rangle = 790 \pm 5$ km s⁻¹ and $\Delta V_{50} = 85 \pm 10$ km s⁻¹ from Parkes Telescope data ($HPBW = 15'$). Both values are in good agreement with our observations, especially considering that the slightly lower spectral resolution of Longmore et al. (1982) may have led to a mild overestimation of ΔV_{50} . Huchtmeier et al. (2001) list $\langle V \rangle = 787 \pm 3$ km s⁻¹ and $\Delta V_{50} = 65$ km s⁻¹ from HIPASS, also using Parkes.

After integration over the full profile in Fig. 7, we find a total HI flux $F_{\text{HI}} = 23.1 \pm 1.2$ Jy km s⁻¹. Longmore et al. (1982) obtained 36 ± 4 Jy km s⁻¹ but Huchtmeier et al. (2001) list 13.7 Jy km s⁻¹, while Koribalski et al. (2004) and Meyer et al. (2004) respectively quote 17.6 ± 2.5 and 17.2 Jy km s⁻¹ from HIPASS data (the latter HIPASS measurements were obtained under the assumption that ESO 364-G029 is a point source). It is thus unclear from those if our synthesis observations are missing

any short spacing flux or not. We reanalyzed the HIPASS archive observations ourselves, summing the data in a 40' × 40' box, and we find a total HI flux of 21.8 ± 1.0 Jy km s⁻¹, still lower than the result of Longmore et al. (1982) but consistent with our own ATCA measurement. This HIPASS global profile is also shown in Fig. 7.

Additional uncertainties in the HI flux include the possibility of HI self-absorption and the existence of dense gas clumps associated with the apparently diffuse HI gas (Grenier et al. 2005). Given that $N_{\text{HI}}/\tau \approx T_s \Delta v$, where N_{HI} is the column density, τ the opacity, T_s the spin temperature, and τ and T_s are correlated (Dickey & Lockman 1990), it is not easy to calculate a definite upper limit for N_{HI} even from the line analysis. Simplifications can however be made using the isophotal axis ratio as a measure of the disk inclination (Giovanelli et al. 1994). The correction factor, $(a/b)^{0.12}$, for the disk of ESO 364-G029 is about 1.1, so this effect is likely small; we have not corrected for HI self-absorption above.

The ATCA global profile shown in Fig. 7 appears asymmetric, with more HI at higher velocities. This behaviour is common

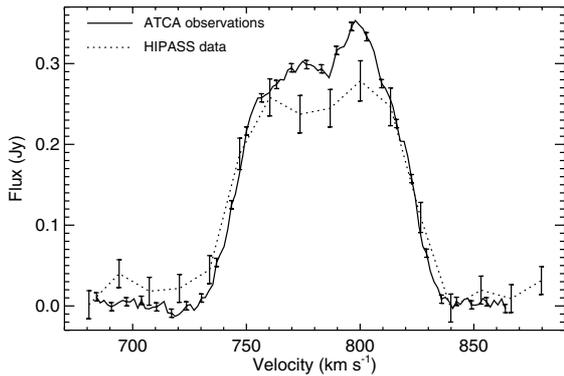


Fig. 7. The global HI profile of ESO 364-G029 (*solid line*), obtained from datacube NH by integrating the individual channel maps over the area containing emission. The total HI flux is $F_{\text{HI}} = 23.1 \pm 1.2 \text{ Jy km s}^{-1}$. The HIPASS global profile summed over a $40' \times 40'$ area is overlotted (*dotted line*). For clarity, the error bars for our ATCA data are only plotted for each fourth channel.

among late-type dwarfs (e.g. Richter & Sancisi 1994; Swaters et al. 2002). Schoenmakers et al. (1997) have developed a complex algorithm to quantify asymmetry using Fourier techniques. In our analysis, however, we restrict ourselves to the standard moment maps, and quantify the importance of the asymmetry with the parameter

$$Q_{\text{asym}} \equiv \frac{\left| \int_{-\infty}^{V_{\text{sys}}} I(v) dv - \int_{V_{\text{sys}}}^{\infty} I(v) dv \right|}{\int_{-\infty}^{\infty} I(v) dv} = 1.7\%, \quad (1)$$

where $I(v)$ is the flux as a function of frequency v . In the expression above, we adopted integration limits of 720 and 860 km s^{-1} , respectively (as for the mean velocity and total fluxes above). The global profile asymmetry of ESO 364-G029 is thus mild and the HIPASS data show an even milder asymmetry.

3.3. HI moments and rotation curve

Moment maps were derived in the standard manner from the cleaned cubes of the datasets listed in Table 3. Figures 8–10 show the total HI map, the velocity field and the velocity dispersion field for the high-resolution dataset, both uniform and natural-weighted. All significant emission is contiguous so a mask was derived from the zeroth moments and applied to the others.

The total HI maps (Fig. 8) show that the HI roughly follows the optical light, although slightly offset to the North-East, and has the same asymmetries. The map from datacube UL shows very well that the HI is concentrated along the main bar and the one-arm spiral, although there is a third component intermediate between those two in terms of position, length, and mass. While this third component has a stellar counterpart where it joins with the bar, this does not appear to be the case over its entire length. The velocity fields (Fig. 9) reveal solid-body rotation over the entire optical extent, a common feature among dwarf irregular galaxies (e.g. de Blok et al. 1996; Swaters 1999). Interestingly, no kinematic lopsidedness is present and the one-arm spiral has no obvious influence on the kinematics, although this may simply be related to the poor spatial resolution of most maps. The rotation curve appears to flatten at the largest radii and a kink in the kinematic major-axis suggests the presence of a warp in the outer parts. The velocity dispersion is 8–16 km s^{-1} over most

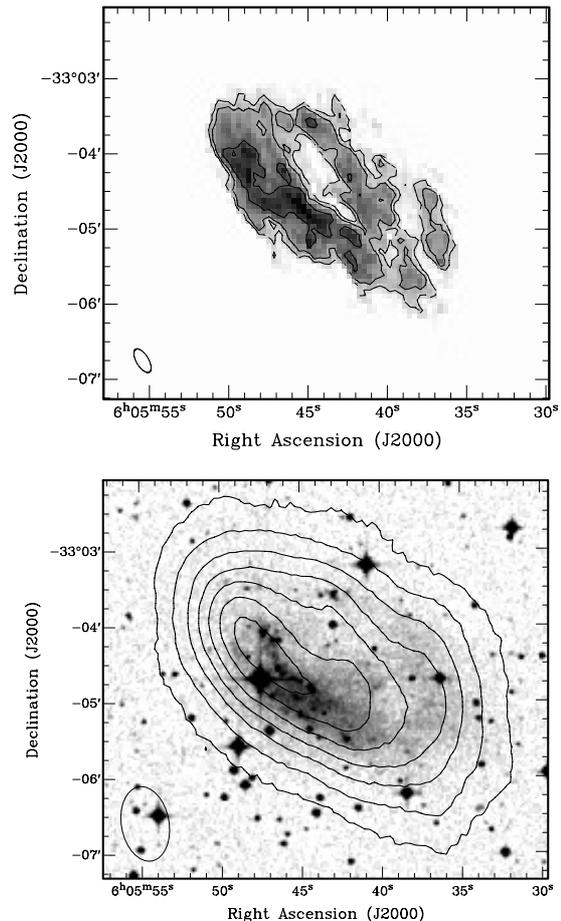


Fig. 8. Total HI maps of ESO 364-G029 for the low sensitivity (high resolution) datacubes. *Top:* UL isodensity contours overlaid on a grayscale representation. The contours represent 1, 2 and 4 times the faintest contour of $50 \text{ mJy beam}^{-1} \text{ km s}^{-1}$ or $2.65 \times 10^{20} \text{ cm}^{-2}$. *Bottom:* NL contours overlaid on a DSS image. The contours represent 1, 5, 10, 15, 20, 25 and 30 times the faintest contour of $80 \text{ mJy beam}^{-1} \text{ km s}^{-1}$ or $0.4 \times 10^{20} \text{ cm}^{-2}$. The contours generally increase toward the center and the synthesized beam is indicated in the bottom-left corner of each map.

of the galaxy (Fig. 10), the higher values being reached over the optical bar and the third component.

The GIPSY⁶ (van der Hulst et al. 1992; Vogelaar & Terlouw 2001) implementation of ROTCUR (Begeman 1987) was used for rotation curve fitting. The target was divided into concentric tilted rings, each approximately half a beam width wide, and at first all six parameters describing each ring were left free: the center coordinates (x_0, y_0) , systemic velocity V_{sys} , inclination i , position angle of the kinematic major-axis ϕ and rotation velocity V_{rot} . After a few trials with the velocity field of the NL datacube, the systemic velocity was fixed to $V_{\text{sys}} = 784 \text{ km s}^{-1}$, and the center coordinates were fixed at $(\alpha = 6^{\text{h}}05^{\text{m}}43^{\text{s}}.4, \delta = -33^{\circ}04'31''; \text{J2000})$.

As the inclination is only loosely constrained by the observations, we fixed it at the HyperLEDA value of $i = 70.5$. This value for the inclination depends on the assumed thickness of ESO 364-G029, but is consistent with the values returned by ROTCUR in our first trials, and is likely a good approximation of the true inclination of ESO 364-G029. For an oblate galaxy with semi-major axis a and semi-minor axis c , the relation between

⁶ Groningen Image Processing System.

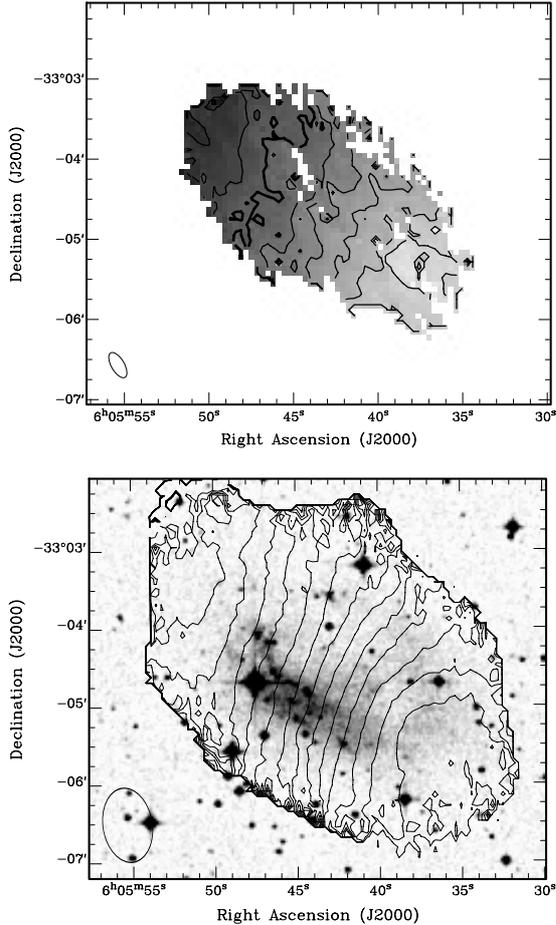


Fig. 9. Same as Fig. 8 but for the velocity fields of ESO 364-G029. The isovelocity contours increase toward the North-East and cover the velocity range $750\text{--}820\text{ km s}^{-1}$ in both panels. The contours are spaced by 10 km s^{-1} (*top*) and 5 km s^{-1} (*bottom*).

the intrinsic thickness $q_0 \equiv c/a$ and the projected axis ratio q is given by

$$\cos^2 i = \frac{q^2 - q_0^2}{1 - q_0^2} \quad (2)$$

(e.g. Bottinelli et al. 1983). For an intrinsic thickness $q_0 = 0.2$ (Holmberg 1946), an inclination $i = 70:5$ implies a projected axis ratio $q = 0.38$, consistent with our observations at large radii (e.g. Figs. 4 and 8). In any case, the rotation velocities can easily be scaled for other inclination values.

With the values of (x_0, y_0) , V_{sys} and i are fixed, we fit for the kinematic position angle ϕ and rotation velocity V_{rot} as a function of semi-major axis position a using datacube NL. The fitted position angle is consistent with $\phi = 60^\circ \pm 5^\circ$ over the whole semi-major axis range. The resulting rotation curve is shown in Fig. 11, overlaid on a position-velocity diagram taken along the major-axis. As the HI column density decreases rapidly at large radii ($a \gtrsim 2:3$), we obtain our data points at these radii from the high-sensitivity datacube NH. The rotation curve is typical of late-type systems, rising in an almost solid-body manner in the central parts. The maximum rotational velocity of $V_{\text{rot}} \sin i \approx 40\text{ km s}^{-1}$ is reached at $a \approx 2'$ (the optical extent), corresponding to a projected distance of $6.3 \pm 0.8\text{ kpc}$, and the rotation curve flattens out at larger radii. The data hint at a decline in the rotational velocity at radii $a \gtrsim 2:7$, though deeper HI observations are necessary to confirm this. One should however remember

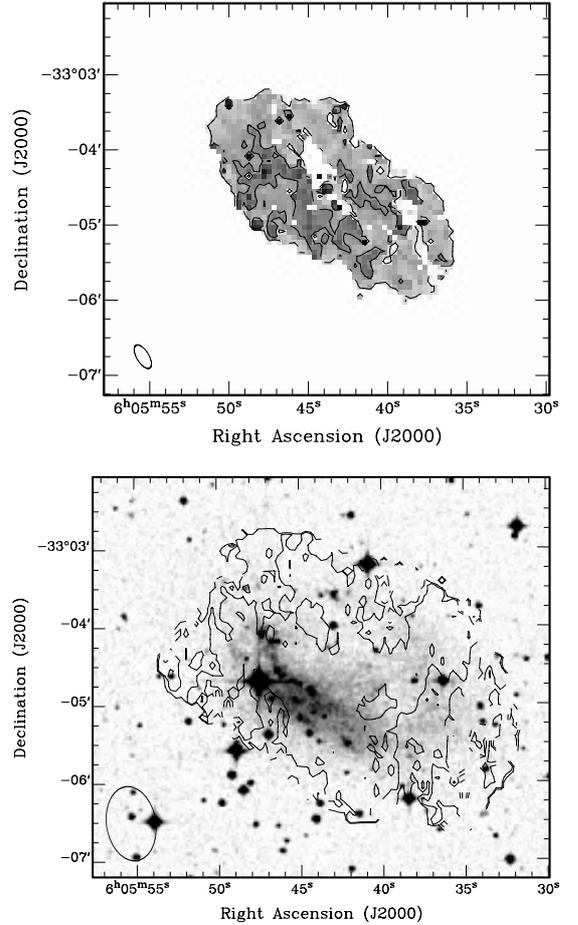


Fig. 10. Same as Fig. 8 but for the velocity dispersion fields of ESO 364-G029. The contours generally increase toward the center and represent 2 and 6 km s^{-1} (*top*) (unreliable) and 8 and 12 km s^{-1} (*bottom*).

that ROTCUR assumes axisymmetry (like `Ellipse` for the surface photometry; see Sect. 2.2), which may not be the case in ESO 364-G029. Our results may thus be systematically biased.

4. Distance determination and total fluxes

No precise determination of the distance to ESO 364-G029 is available in the literature (e.g. through Cepheids or the tip of the red-giant branch), but it can be estimated in several ways. Adopting $H_0 = 70\text{ km s}^{-1}$ and a Local Group infall toward the Virgo cluster of 208 km s^{-1} , HyperLEDA lists a distance of 8.02 Mpc . For our analysis we adopt the slightly more refined linear Virgocentric flow model of Schechter (1980). We derive a distance to ESO 364-G029 of $D = 10.8 \pm 1.4\text{ Mpc}$ (random errors only), adopting a Virgo distance of $D_{\text{Virgo}} = 16.1 \pm 1.2\text{ Mpc}$ (random errors only; Kelson et al. 2000), an observed Virgo redshift of $V_{\text{Virgo}} = 1079\text{ km s}^{-1}$ (Ebeling et al. 1998), a Local Group Virgocentric infall velocity of $257 \pm 32\text{ km s}^{-1}$ (Aaronson et al. 1982) and a Virgocentric density contrast $\gamma = 2$ (representative of Abell clusters; Peebles 1976). An angular separation of $1'$ then corresponds to a projected distance $3.1 \pm 0.4\text{ kpc}$.

At this adopted distance, the total HI flux $F_{\text{HI}} = 23.1 \pm 1.2\text{ Jy km s}^{-1}$ (see Sect. 3.1) corresponds to a total HI mass $M_{\text{HI}} = (6.4 \pm 1.7) \times 10^8 M_\odot$. From the corrected total apparent magnitudes listed in Sect. 2.2 (corrected for Galactic extinction but not internal absorption), we derive the following corrected absolute magnitudes: $M_B = -16.6 \pm 0.3$, $M_V = -16.8 \pm 0.3$ and

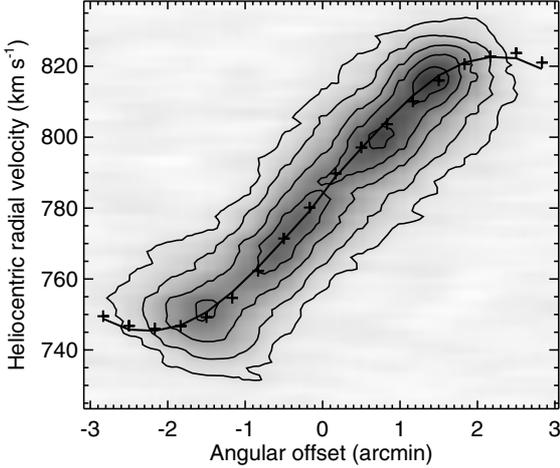


Fig. 11. The rotation curve of ESO 364-G029 as derived from the tilted-ring model described in the text, overlaid on a grayscale of the major-axis position-velocity diagram (datacube NL; beam size $58'' \times 38''$). The data points at radii larger than $2/3$ are obtained from datacube NH (beam size $132'' \times 70''$), which has a higher sensitivity. The solid curve shows the fit to the whole velocity field, while the crosses show the fit to the approaching side (*left*) and receding side (*right*) of the galaxy only. Contour levels are overlaid at 7.5, 20, 40, 60 and 80 mJy beam^{-1} . Rotational velocities are *not* corrected for the uncertain inclination. For an inclination $i = 70:5$ (HyperLEDA), the rotation curve must be scaled up by 6%.

$M_I = -17.9 \pm 0.3$ mag. This is more than a magnitude fainter than the absolute magnitude of the LMC listed in HyperLEDA ($M_B = -17.9$ mag). Combining the optical and HI measurements, we obtain a (distance independent) total HI mass-to-blue-luminosity ratio $M_{\text{HI}}/L_B = (0.96 \pm 0.14) M_{\odot}/L_{B,\odot}$, roughly five times larger than that of the LMC based on HyperLEDA's entries ($M_{\text{HI}}/L_B = 0.18 M_{\odot}/L_{B,\odot}$). This is consistent with ESO 364-G029 being of later type than the LMC, although the total HI flux of the LMC is highly dependent on the area considered, given the presence of the Magellanic bridge (e.g. Putman et al. 1998; Brüns et al. 2005).

5. Discussion

5.1. Structure of ESO 364-G029 and a comparison to the LMC

Figures 1–2 reveal an irregular but characteristic optical morphology for Magellanic-type dwarf galaxies. Indeed, the main body of ESO 364-G029 exhibits a bar-like morphology, from which protrudes a one-arm spiral circling the galaxy by at least 90° , possibly more. This structure is sharpest in the bluer bands, suggesting that it is associated with relatively young stars and possibly ongoing star formation. Narrow-band $H\alpha$ imaging would be particularly useful to confirm this latter point. Despite the irregular morphology, however, the azimuthally-averaged optical profiles are approximately exponential up to $a \approx 135''$, after which they fall off more rapidly (see Fig. 3).

In our total HI map of datacube UL, the HI distribution has a radial extent of at least $2/8$ (corresponding to a column density of $0.4 \times 10^{20} \text{ cm}^{-2}$; see Fig. 8), thus extending to at least $2.2 R_{25}$ or $1.2 R_{H_0}$. This is fairly typical of dwarf irregular galaxies (e.g. Hunter 1997). The HI is also asymmetric, the column density falling off more steeply to the South-East than the North-West, and the absolute peak being slightly offset to the

North-East. Both effects are common in late-type dwarf galaxies (see, e.g., Swaters et al. 2002).

The HI distribution is also slightly offset compared to the optical in all our maps, being shifted toward the North-East, and the HI kinematic center does not correspond to the optical center as defined by the surface photometry. The latter displacement is about $23''$, corresponding to a projected distance of 1.2 ± 0.2 kpc. This effect is present in a number of Magellanic-type galaxies (e.g. de Vaucouleurs & Freeman 1972). In particular, there are strong indications in the LMC that the kinematic center derived from the HI (e.g. Kim et al. 1998) is very different from that of the main disk, as traced by the distribution of red giant branch (RGB) and asymptotic giant branch (AGB) stars (van der Marel 2001) or the kinematics of intermediate-age carbon stars (van der Marel et al. 2002; Olsen & Massey 2007), while the morphological centers derived from those latter tracers are in rough agreement with each other (Weinberg & Nikolaev 2001). Evidence is also gathering that the LMC is intrinsically elongated (van der Marel & Cioni 2001; van der Marel 2006), presumably due to the tidal force from the Milky Way (van der Marel 2001; Mastropietro et al. 2005). For a distant galaxy, this would translate in different kinematic and photometric major axes. Given the uncertainties in the position angle, however, no significant difference is seen between the photometric and kinematic position angle profiles of ESO 364-G029 (Figs. 4 and 9).

The rotation curve of ESO 364-G029 rises smoothly and reaches a value of about $V_{\text{rot}} \sin i = 40 \text{ km s}^{-1}$ in the outer parts (Fig. 11). In comparison, the rotation of the LMC shows a sudden decline after its peak velocity $V_{\text{rot}} \approx 63 \text{ km s}^{-1}$ at $a = 2.4$ kpc (Kim et al. 1998). An inclination of about 40° is expected if ESO 364-G029 is to have the same rotation as the LMC. This however appears unlikely given the relatively high apparent ellipticity observed (Figs. 4 and 8). Even the lowest optical ellipticity measured ($\epsilon_{\text{min}} = 0.4$; see Fig. 4) is inconsistent with such a low inclination, for any intrinsic disk thickness (e.g. Bottinelli et al. 1983). As a disk with an intrinsically high ellipticity is unlikely, it is almost certain that ESO 364-G029 is less massive than the LMC over the range of radii probed. Interestingly however, the total HI mass of ESO 364-G029 is about a third higher than that of the LMC ($\approx (4.8 \pm 0.2) \times 10^8 M_{\odot}$; Staveley-Smith et al. 2003), resulting in a significantly larger HI mass-to-blue-luminosity ratio $M_{\text{HI}}/L_B = (0.96 \pm 0.14) M_{\odot}/L_{B,\odot}$ for ESO 364-G029 (see Sect. 4).

Despite a number of similarities between the optical and HI content of ESO 364-G029 and the LMC, one large and surprising difference exists. The medium- and large-scale HI structure of the LMC is very regular, bearing little resemblance to the optical structure and lacking any obvious trace of the main bar (e.g. Putman et al. 1998; Kim et al. 2003). As stated above, however, the HI distribution of ESO 364-G029 is highly asymmetric and roughly follows its optical morphology. In particular, and as best shown in the HI map (Fig. 8), the HI is concentrated along the main bar and appears to follow the one-arm spiral to the North-West, in addition to a third intermediate component (also with an optical counterpart, at least partially). We have no explanation for these differences in the HI morphology; further observations in ATCA 1.5 km configurations are needed to characterize the HI morphology of ESO 364-G029 better. In this respect, it would be interesting to see if the optical morphology of ESO 364-G029 remains the same at longer wavelengths or if, like the LMC, the preponderance of the bar decreases and the smoothness of the disk increases in the near-infrared (e.g. van der Marel 2001; Meixner et al. 2006). Also of interest is

whether the HI differences arise from internal processes (e.g. supernovae feedback) or from external factors (e.g. tidal interactions). The former appears to dominate the HI structure of the LMC on small scales (Kim et al. 1998; Kim et al. 2003) and the latter on large scales (Putman et al. 1998; Staveley-Smith et al. 2003). However, this could be different for ESO 364-G029.

Interestingly, the velocity dispersion maps of ESO 364-G029 and the LMC are also somewhat different. In ESO 364-G029, the highest HI velocity dispersions are observed along the bar and the third component. In the LMC, the HI velocity dispersion is higher along the Eastern edge of the 30 Doradus region and near the dense molecular clouds extending to the South (Kim et al. 2003). Both regions are turbulent due to the combined effects of shocks from ram pressure on the Galactic halo and the very active star formation near 30 Doradus. Star formation is probably also at the origin of the high HI velocity dispersions observed near the North-West end of the LMC bar and Constellation III (e.g. Efremov & Elmegreen 1998).

5.2. Comparison to other dwarf irregulars

The optical and HI distributions of ESO 364-G029 are typical of those of a Magellanic dwarf galaxy: a prominent bar and a pronounced single spiral arm in the optical, an HI distribution extending to about 1.2 Holmberg radii and a kinematic center offset from the photometric center. Furthermore, our HI data suggest the presence of numerous shells and holes (see Sect. 5.4).

The HI kinematic center of ESO 364-G029 is offset by a projected distance of 1.2 ± 0.2 kpc from the optical center (as derived from surface photometry). This offset is known to be present in other Magellanic-type systems (e.g. de Vaucouleurs & Freeman 1972) such as NGC 925, which has an offset of ≈ 1 kpc (Pisano et al. 1998), and the LMC (see Sect. 5.1). Additional features in the HI distribution have been detected in Magellanic dwarfs, such as loops surrounding the galaxy (NGC 4618; Bush & Wilcots 2004), external spurs or blobs (e.g. NGC 5169; Mühle et al. 2005; UGC 98; Stil et al. 2005) and S-shaped distortions in the HI velocity field (e.g. NGC 4449; Hunter et al. 1998; DDO 43; Simpson et al. 2005). Such features, often associated with a recent or ongoing interaction, are not apparent in the HI distribution of ESO 364-G029. The origin of the third component identified in HI between the optical bar and one-arm spiral also remains unclear. Shell- and hole-like structures in the HI distribution, as observed in ESO 364-G029 (see Sect. 5.4), are observed in practically all dwarfs. These structures are commonly thought to be the result of stellar winds and supernova explosions (e.g. Weaver et al. 1977; McCray & Kafatos 1987; Puche et al. 1992; Kim et al. 1999; but see also Rhode et al. 1999; Bureau & Carignan 2002; Dib & Burkert 2005).

We find no significant difference between the photometric and kinematic position angle of ESO 364-G029. Hunter et al. (2000) find a similar result in their HI study of UGC 199. Many other Magellanic dwarfs, however, show a significant disagreement between the morphological and kinematic axes, such as NGC 1156 (Hunter et al. 2002) and DDO 26 (Hunter & Wilcots 2002). DDO 43 even has a kinematic axis nearly perpendicular to its morphological axis (Simpson et al. 2005). Hunter et al. (2002) suggest that such apparent inconsistencies may be explained by the presence of an inclined bar.

The asymmetry we observe in the global HI profile of ESO 364-G029 is a common characteristic of Magellanic-type galaxies, often attributed to a recent or ongoing interaction. Bush & Wilcots (2004), for example, analyse the HI morphology of the interacting Magellanic spiral galaxies NGC 4618 and

NGC 4625, particularly in relation to the interaction between them. Their HI observations show a loop-like structure around NGC 4618 (the most massive of the two), indicating that the outer gas of this galaxy is strongly perturbed by the recent interaction. NGC 4625 and the inner part of NGC 4618, however, appear unaffected by the interaction. They find an asymmetry ratio $A = 1.0$ for the HI profile of NGC 4618 and $A = 1.29$ for NGC 4625, where A is defined as the ratio between the areas under the HI profile at velocities smaller and greater than the systemic velocity (Haynes et al. 1998). Using the same algorithm, we find an asymmetry ratio $A = 1.1$ for ESO 364-G029, which is bracketed by the values for NGC 4618 and NGC 4625. However, a similar amount of asymmetry is detected in non-interacting Magellanic spirals (Wilcots & Prescott 2004); an asymmetry measurement alone cannot confirm or reject the hypothesis of a recent interaction.

ESO 364-G029 exhibits solid-body rotation at small radii, reaches its maximum rotational velocity of $V_{\text{rot}} \sin i \approx 40 \text{ km s}^{-1}$ at a projected distance of ≈ 6.4 kpc from its center, and flattens out at large radii. This behaviour is similar to that of other Magellanic dwarfs, e.g. IC 10 ($\approx 35 \text{ km s}^{-1}$ at 1 kpc; Wilcots & Miller 1998), DDO 43 ($\approx 30 \text{ km s}^{-1}$ at 2 kpc; Simpson et al. 2005) and NGC 4618 ($\approx 50 \text{ km s}^{-1}$ at 5 kpc; Bush & Wilcots 2004), though variations in these parameters clearly exist. The HI velocity dispersion of ESO 364-G029 is $8\text{--}16 \text{ km s}^{-1}$ over most of the galaxy (Fig. 10), which is very similar to that of other dwarfs. The Magellanic dwarfs mentioned above exhibit a somewhat higher velocity dispersion at several locations (including the bar), indicating recent or ongoing star formation. For ESO 364-G029, the largest velocity dispersion is reached over the optical bar and the third component, which also likely indicates recent star formation in these regions.

Doyle et al. (2005) present a search for optical counterparts of HIPASS sources. Out of the 3618 optical counterparts found, 151 are galaxies with morphological type IB(s)m, including ESO 364-G029. From their (total) HI fluxes and B -band magnitudes we calculated the HI mass-to-blue-luminosity ratio M_{HI}/L_B for each galaxy. Magellanic dwarf irregular galaxies have a median value $\langle M_{\text{HI}}/L_B \rangle = 0.97 M_{\odot}/L_{B,\odot}$, with a standard deviation of $0.74 M_{\odot}/L_{B,\odot}$. Thus, ESO 364-G029 (with $M_{\text{HI}}/L_B = (0.96 \pm 0.14) M_{\odot}/L_{B,\odot}$) is a fairly typical Magellanic dwarf galaxy in this respect as well.

5.3. Group membership and origin of the asymmetries

Tidal interactions in disk galaxies are known to lead to strongly asymmetric spiral features (e.g. Byrd et al. 1986; Olson & Kwan 1990). The structure of M51-like systems, for example, is often explained by tidal interactions with a close neighbour (e.g. Howard & Byrd 1990; Howard et al. 1993; Salo & Laurikainen 2000). While this mechanism need not be unique, it is interesting to question whether the large-scale morphology of ESO 364-G029 is consistent with a (trans)formation through tidal interactions, or whether another mechanism needs to be invoked. Certainly, the break in ESO 364-G029's light profile seen at $a \approx 135''$ is consistent with an interaction, since tidal features typically have sharp boundaries.

Tully (1988) lists ESO 364-G029 as being part of a small group of 3 galaxies in the Dorado Cloud, which also includes NGC 2090 and NGC 2188. Table 4 lists all galaxies with a known redshift within a radius of 5° of ESO 364-G029, and with a relative radial velocity less than 500 km s^{-1} , as found in NED. Only one additional galaxy is found, AM 0605-341, although a number of galaxies without redshift are also present in

Table 4. Basic properties of galaxies near ESO 364-G029. This table lists all galaxies with a projected distance differing by less than 5° and a redshift differing by less than 500 km s^{-1} from ESO 364-G029. Only galaxies with a known redshift are included. From left to right, the columns list the galaxy name, spectral type, total apparent B magnitude and the angular separation, projected separation (assuming a distance of 10.8 Mpc), and relative radial velocity from ESO 364-G029.

Name	Type	B_T (mag)	$\Delta\theta$	ΔR (kpc)	ΔV_{sys} (km s^{-1})
ESO 364-G029	IB(s)m	13.8	–	–	–
AM 0605-341	SBdm	14.1	70.2	221	-19
NGC 2188	SB(s)m	12.1	82.5	259	-37
NGC 2090	SA(rs)b	12.0	244.0	768	+137

the 5° region. There are thus a number of candidates for interaction with ESO 364-G029.

NGC 2090 is the least certain and most distant member of the group containing ESO 364-G029 (among the galaxies with a measured redshift). Its distance was measured using the Hubble Space Telescope (HST) and the Cepheid period-luminosity relation, yielding a distance of $12.3 \pm 0.9 \pm 0.9$ Mpc (random and systematic errors, respectively; Phelps et al. 1998). Although slightly higher, this is consistent within the errors with our adopted distance for ESO 364-G029, a comforting fact given that our estimate for ESO 364-G029 was itself higher than other published values (see Sect. 4).

The edge-on galaxy NGC 2188 is rather peculiar, both the HI and H α emission bending away from the disk, leading to a crescent shape not unlike that expected from ram pressure stripping (Domgörgen et al. 1996). The HI distribution is also highly asymmetric with respect to the equatorial plane. The velocity field shows many peculiarities, with apparent rotation about two axes. Although not bullet-proof evidence for interaction, these properties hint at least to a disturbed past. Given the large separation between ESO 364-G029 and NGC 2188, however, any interaction with ESO 364-G029 would have occurred a very long time ago (>1.25 Gyr for a relative velocity of 200 km s^{-1}). The evidence for a recent interaction involving ESO 364-G029 thus remains marginal.

5.4. HI shells and holes

Although the current observations are best suited to study the medium and large-scale structure of the HI, studies of the small-scale structure in the interstellar medium (ISM) are also of interest. In particular, the radiation and mechanical energy produced by stellar winds and supernova explosions are generally thought to give rise to a very dynamic multi-phase ISM, through the interacting and continually evolving cavities created. For recent work on the latter topic, see, e.g., Korpi et al. (1999) and de Avillez & Berry (2001). Such shells and holes have indeed been observed in nearby large spiral galaxies such as M 31 (e.g. Brinks & Bajaja 1986) and M 33 (Deul & den Hartog 1990). Dwarf galaxies are however better targets for studying these phenomena. The low gravitational potential well of dwarf galaxies, the lack of shear due to solid-body rotation and the absence of density waves all facilitate the expansion and survival of the shells created. Some of the best examples include the LMC (SB(s)m; Kim et al. 1998), IC 2574 (SAB(s)m; Walter & Brinks 1999) and Holmberg II (Im; Puche et al. 1992; but see also Rhode et al. 1999; Bureau & Carignan 2002).

The channel maps of datacube UL (Fig. 6) and the corresponding moment maps (Figs. 8–10) of ESO 364-G029 reveal

much small-scale structure, suggesting that expanding shells or holes may be present in ESO 364-G029. Most obvious are the two large HI gaps, the first one between the main bar and the so-called third component, centered at ($\alpha = 06^{\text{h}}05^{\text{m}}44^{\text{s}}.5$, $\delta = -33^\circ04'10''$; J2000), and the second one between the third component and the one-arm spiral, itself perhaps composed of two smaller holes centered respectively at ($\alpha = 06^{\text{h}}05^{\text{m}}38^{\text{s}}.7$, $\delta = -33^\circ04'50''$) and ($\alpha = 06^{\text{h}}05^{\text{m}}39^{\text{s}}.9$, $\delta = -33^\circ05'20''$). These holes raise the interesting possibility that the third component and the one-arm spiral may not be caused by large-scale dynamical processes, but may instead simply represent gas concentrations at the edges of large cavities, emptied by the supernova explosions and stellar winds associated with past bursts of star formation. Their optical counterparts would then trace the location of secondary (i.e. triggered and more recent) star formation.

Despite its relatively large distance, ESO 364-G029 is thus a prime target for studies of star formation feedback and self-propagating star formation. However, due to the short integrations with the 1.5D array and the shallowness of the high-resolution maps, it is hard to convincingly argue that any single structure is surely due to a (centrally-located) localized energetic phenomenon. Better quality maps and longer observations with 1.5 km or longer arrays are necessary for a proper study of the small-scale structure of the ISM. As mentioned before, narrow-band H α imaging would be useful to see how current and recent sites of star formation relate to the HI morphology and kinematics.

6. Summary and conclusions

We presented the analysis of optical imaging observations and HI radio synthesis observations of the dwarf irregular galaxy ESO 364-G029. The optical BVI imaging data reveal a morphology characteristic of Magellanic-type spirals and irregulars, with a large dominant bar and a one-sided spiral or tidal arm, although the absolute magnitude of ESO 364-G029 (for a distance $D = 10.8 \pm 1.4$ Mpc) is more than a magnitude fainter than that of the LMC. While poorly-defined, the azimuthally-averaged surface brightness profiles show an exponential disk with a possible break at large radii and colors typical of late-type disk galaxies.

The radio synthesis observations reveal an HI disk extending well outside of the optical extent. The total HI mass is $M_{\text{HI}} = (6.4 \pm 1.7) \times 10^8 M_\odot$, yielding a (distance independent) HI mass-to-blue-luminosity ratio $M_{\text{HI}}/L_B = (0.96 \pm 0.14) M_\odot/L_{B,\odot}$, significantly more HI-rich than the LMC. The latter value is at the high end of the distribution for late-type spirals (Sc–Sd), but is typical of late-type dwarfs (Roberts & Haynes 1994; Swaters et al. 2002; Doyle et al. 2005). The large-scale HI distribution is also asymmetric and roughly follows the optical light distribution, although slightly offset from it. The HI distribution thus appears consistent with that expected from tidal interactions, but most evidence for a past encounter is circumstantial. Our highest spatial resolution maps show that the highest column densities and velocity dispersions are reached over the central bar, the one-arm spiral, and over a third component which also has an optical counterpart (at least partially). Despite a similar optical morphology, this is opposite to what is observed in the LMC, where the HI distribution bears little resemblance to that of the stars. The two galaxies may thus have formed differently, or may simply be in a different evolutionary stage. The rotation curve of ESO 364-G029 is solid-body over the optical extent but flattens out at large radii, reaching $V_{\text{rot}} \sin i \approx 40 \text{ km s}^{-1}$. The inclination remains poorly constrained by observations, but is consistent

with the HyperLEDA value $i \approx 70:5$. Our high-resolution observations also hint at a complex HI structure, reminiscent of that expected from stellar winds and supernova explosions. Deeper and higher-resolution optical and HI observations, as well as H α observations, are however necessary to properly characterize the small-scale HI morphology and kinematics, and to establish any relationship to star formation.

The observations presented in this paper provide important information on the large-scale optical and HI structure of ESO 364-G029, and provide a solid foundation for further investigation of its small-scale structure, particularly in HI with longer baseline observations. The ultimate goal is that studies of individual nearby Magellanic galaxies will yield information on their formation and evolution as a class, but will also strengthen our understanding of the LMC itself.

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