

Early spectral evolution of Nova Sagittarii 2004 (V5114 Sagittarii)^{★,★★,★★★}

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ABSTRACT

Aims. We present optical and near-infrared spectral evolution of the Galactic nova V5114 Sgr (2004) during few months after the outburst.

Methods. We use multi-band photometry and line intensities derived from spectroscopy to put constraints on the distance and the physical conditions of the ejecta of V5114 Sgr.

Results. The nova showed a fast decline ($t_2 \approx 11$ days) and spectral features of Fe II spectroscopic class. It reached $M_V = -8.7 \pm 0.2$ mag at maximum light, from which we derive a distance of 7700 ± 700 kpc and a distance from the galactic plane of about 800 pc. Hydrogen and oxygen mass of the ejecta are measured from emission lines, leading to $\sim 10^{-6}$ and $10^{-7} M_\odot$, respectively. We compute the filling factor of the ejecta to be in the range $0.1-10^{-3}$. We found the value of the filling factor to decrease with time. The same is also observed in other novae, then giving support to the idea that nova shells are not homogeneously filled in, rather being the material clumped in relatively higher density blobs less affected by the general expanding motion of the ejecta.

Key words. stars: novae, cataclysmic variables – stars: individual: V5114 Sgr

1. Introduction

Nova Sgr 2004 (V5114 Sgr) was independently discovered by Nishimura (2004) and Liller (2004) on Mar. 15.8 UT and Mar. 17.3 UT, respectively. West (2004) provided precise coordinates RA = $18^{\text{h}}19^{\text{m}}32^{\text{s}}.29$, Dec = $-28^{\circ}36'35''.7$ (gal. coord. $l = 3^{\circ}.9$ $b = -6^{\circ}.3$). Early spectroscopy on Mar. 18.3 UT by Della Valle et al. (2004) confirmed this object to be a classical nova caught near maximum light. Here, we present photometric and spectroscopic observations taken at McDonald Observatory, Cerro Tololo Inter-American Observatory (CTIO), ESO-La Silla, and Lick Observatory. The

spectra from ESO-La Silla have been obtained as part of the Target of Opportunity campaign for the observation of Classical Novae in the Galaxy and in the Magellanic Clouds. The paper is organized as follows: in Sect. 2 we present the analysis of the photometric and spectroscopic data, in Sects. 3 and 4 we analyze the evolution of the light curve and spectra of V5114 Sgr, in Sect. 5 we discuss the effects of interstellar absorption. In Sect. 6 we derive the physical parameters of the nova ejecta. Summary and conclusions are given in Sect. 7.

2. Observations and data reduction

Optical photometry has been carried out during 8 nights in March and April 2004 with the 0.8-m telescope at McDonald Observatory. Another observing run has been carried out during 14 nights in June and July with the Small and Moderate Aperture Research Telescope System (SMARTS) 1.0-m telescope at CTIO. Optical and infrared photometry has been also carried out with the ANDICAM dual-channel imager on the SMARTS 1.3-m at CTIO during 25 nights from March to

* Based on observations collected at the European Southern Observatory, La Silla, Chile, Cerro Tololo Inter-American Observatory (CTIO), McDonald Observatory and Lick Observatory.

** Tables 2–4 and Appendix are only available in electronic form at <http://www.aanda.org>

*** Spectroscopic material used in this paper is available in electronic form at the CDS via anonymous ftp to cdsarc.u-strasbg.fr (130.79.128.5) or via <http://cdsweb.u-strasbg.fr/cgi-bin/qcat?J/A+A/459/875>

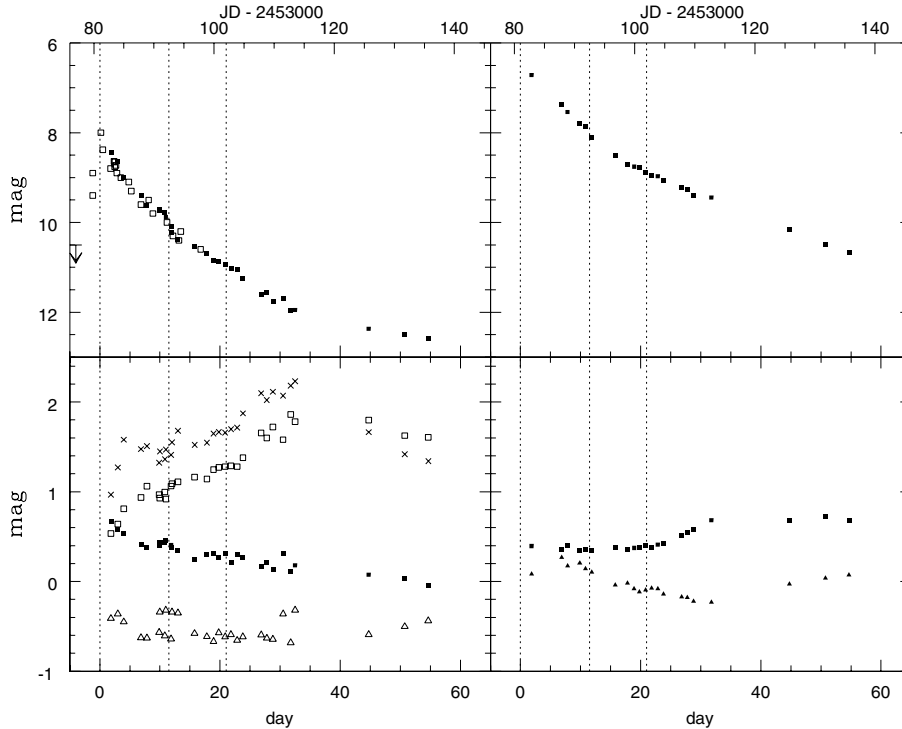


Fig. 1. *Upper left panel:* V band light curve. Filled symbols represent our data points, empty symbols represent IAUC data. Vertical dotted lines represent maximum light, t_2 , and t_3 . Data points are affected by errors $\lesssim 0.04$ mag in all bands but in U , where the error is 0.1 mag. *Lower left panel:* evolution of different colors: filled squares represent $B - V$, empty triangles $U - B$, crosses $V - I$ and empty squares $V - R$. *Upper right panel:* J band light curve. *Lower right panel:* evolution of near-infrared colors: triangles represent $J - H$ and squares $H - K$.

August. One night of observations has been carried out with a 0.3-m telescope equipped with an SBIG ST7E CCD camera, in Exmouth, Australia. A log of photometric observations is given in Table A.1.

Spectra at maximum and during the early decline have been obtained with FEROS (Kaufer et al. 1999) with a resolution $R \sim 48\,000$ and spectral range 4000–9000 Å. Spectrophotometric standard stars have not been observed each night, and in this case the spectra have been corrected with an “average response curve”. This procedure can introduce an uncertainty on the flux measurement up to 50%. Flux-determination is affected by underdetermined uncertainty because FEROS is a fiber-fed spectrograph that was not equipped with an atmospheric distortion corrector at the time of these observations. Indetermination is due to the fact that the observations are carried out guiding on the V -band image of the star that is differently displaced (due to atmospheric refraction) in the other bands. During our analysis, fluxes were corrected in order to match the observed magnitudes.

An independent spectroscopic follow up has been carried out with the RC spectrograph on the SMARTS 1.5-m telescope at CTIO. A spectrophotometric standard star (either LTT 4364 or Feige 110) has been observed each night to remove the instrumental signature. Standard reduction has been carried out with an author’s written IDL routine.

An IR spectrum was taken June 22, 2004 UT at Lick Observatory using the Aerospace Corporation’s Near-Infrared and Visible Imaging Spectrograph (NIRIS). The standard star used was HR 6836.

All spectra have been analyzed with the onedspec package in IRAF. Line fluxes have been measured by the integration of the line profile and not by Gaussian fitting. Full width at half maximum (FWHM) of lines have been measured also via Gaussian fitting but show no significant difference from direct

Table 1. Observed photometric properties for V5114 Sgr.

t_0	JD = 2 453 081.556
t_2	11 days
t_3	21 days
$m_{V,\max}$	8.0 mag
$(B - V)_{\max}$	0.66 mag
$(B - V)_{t_2}$	0.38 mag
$(B - V)_{15d}$	0.25 mag

measure. A complete log of our spectroscopic observation is reported in Table A.2.

3. Light curve

The optical and near-infrared light curves of V5114 Sgr are shown in Fig. 1. The light curves have been derived using both our photometric data and photometry available in the literature (IAUC 8306, 8307, 8310). The V light curve shows that V5114 Sgr reached $V = 8.0$ mag on Mar. 17.17 UT (MJD = 53 081.556). West (2004) noted that nothing was visible at the same position in the red Digitized Sky Survey. After considering that the DSS limiting magnitude is ~ 21 mag, we can infer that the outburst amplitude was $\gtrsim 13$ mag which is consistent with values observed for other novae with about the same rate of decline (see Warner 1995).

The nova decreased by two magnitudes from maximum in $t_2 = 11$ days and by three magnitudes in $t_3 = 21$ days. Adopting the maximum magnitude versus rate of decline (MMRD) relation by Della Valle & Livio (1995), V5114 Sgr achieved an absolute V magnitude at maximum light of $M_V = -8.7 \pm 0.2$ mag. Photometric properties are summarized in Table 1. The typical photometric errors are smaller than 0.04 mag in all bands but in

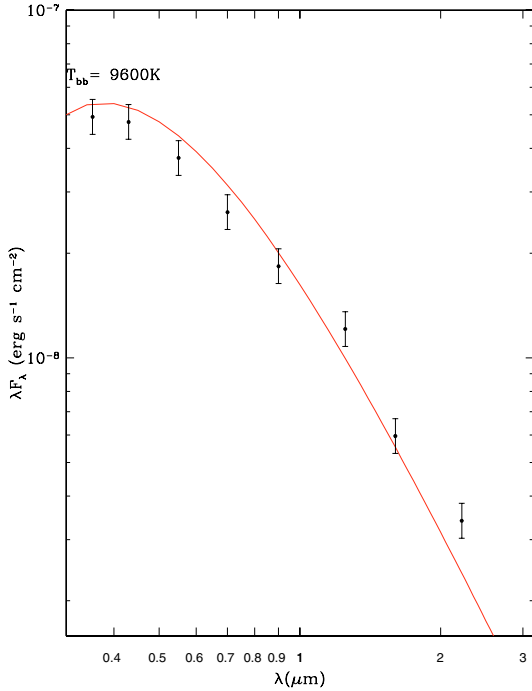


Fig. 2. Spectral energy distribution of V5114 Sgr one day after maximum. The magnitudes in the various bands have been corrected for interstellar extinction.

$U \sim 0.1$ mag ($J - H$) color changes its slope after \sim day 40. This could be in principle due to formation of dust in the ejecta but an inspection of the B and V light curves rules out “DQ Herculis” behavior, therefore the reddened color is likely due to variations in intensity of emission lines in the NIR. The spectral energy distribution (SED) one day after maximum appears to be well fitted by a blackbody at $T = 9600$ K as shown in Fig. 2.

4. Spectral evolution

Spectroscopic observations started immediately after discovery. Line identification for FEROS spectra (see Figs. 3 and 4) is given in Table 2 while in Table 3 we show line identification for spectra taken at CTIO (Williams et al., in preparation).

The first spectrum (phase +1¹) was dominated by Balmer, Fe II and O I emission lines. This behavior characterized the nova as a typical “Fe II” type object, according to the Cerro Tololo classification (Williams et al. 1991, 1994). P-Cyg profiles were clearly visible in Balmer lines as well as in Fe II, O I, and Na I lines. P-Cyg profiles were double, thus suggesting the presence of two expanding systems with velocities (obtained by averaging of measurements of Balmer lines) of 1400 ± 50 km s⁻¹ and 850 ± 30 km s⁻¹.

Eight days after maximum the spectrum was still dominated by low ionization species. The double P-Cyg profiles were still clearly visible and the velocities (as derived from both the P-Cyg profiles and FWHM) were increasing. The emission lines started developing a flat topped profile.

On April 9 (phase +23), we observe the 4640 Å emission band together with N II and N III, although Fe II emission lines were still present. The O I $\lambda 8446$ emission line was more intense than the H β one and showed a flatter profile.

By April 18 (phase +32) P-Cyg profiles had disappeared and Fe II emission lines were fading and forbidden and high

excitation lines strengthened. The intensity of [O III] $\lambda 4363$ indicated that V5114 Sgr entered the auroral phase, described in Williams et al. (1991, 1994). Fluxes of Balmer lines, that had decreased very slowly until this moment, started to decrease faster (see Fig. 7). The FWHM of Balmer lines reached a plateau (2000 ± 100 km s⁻¹).

The O I $\lambda 8446$ and He I $\lambda 5876$ lines show flat topped profiles while the hydrogen lines have a clearly asymmetric profile (the red side being more prominent than the blue one). At this stage the O I $\lambda 8446$ emission line reached its maximum intensity.

By May 13 (phase +57) the hydrogen lines turn to flat topped profiles (like oxygen) while nitrogen lines were still rounded. It has been noted in the past (see Payne-Gaposchkin 1957) that different line profiles observed at the same stage indicate that the emission lines originate in different layers of the ejecta. Flat topped profiles originate from optically thin spherical shells while rounded profiles are related to optically thick winds. The NIR part of the spectrum (observed only on June 22, see Fig. 6, line identification given in Table 5) showed prominent Paschen and Brackett lines as well as oxygen and nitrogen lines. Common but unknown lines (1.10, 1.19, 1.55 and 2.10 μ m) were present in this spectrum (see Venturini et al. 2004). Tentative identifications for these lines with van Hoof’s line list² are given in Table 6. Few suggested identifications have already been suggested by Rudy et al. (2002) during analysis of lines of V723 Cas.

In September the spectrum was dominated by O I $\lambda\lambda$ 4959–5007 lines. All lines showed saddle-shaped profiles. The O I $\lambda 8446$ line had almost disappeared. H α was clearly strongly blended with N II.

The overall evolution of expansion velocities measured from P-Cyg absorption is shown in Fig. 8. Cassatella et al. (2004) have shown that the evolution of P-Cyg profiles as measured in the UV of Nova Cyg 1992 can be modelled with an exponential law $v(t) = v_\infty - (v_0 - v_\infty)e^{-t/\tau}$, where $v(t)$ is the velocity t days after maximum, v_0 is the expansion velocity at maximum light, v_∞ is the asymptotic velocity and τ is a time scale similar to t_3 . The best fit to the data gives $v_\infty = 2100 \pm 50$ km s⁻¹ and 1500 ± 50 km s⁻¹, $v_0 = 1300 \pm 50$ km s⁻¹ and 800 ± 50 km s⁻¹ and $\tau = 18 \pm 1$ day. Both v_∞ and velocity measurements based on the FWHM of emission lines during the nebular stage are important quantities that should be considered in view of future observations aimed at determining the distance to V5114 Sgr, via nebular expansion parallax.

5. Reddening and distance

We have used several methods to estimate the reddening of V5114 Sgr. This piece of information is crucial to measure the distance to the object. The results of our various estimates of the interstellar reddening are summarised in Table 7.

We could not use the commonly used NaD lines because they are saturated. Nevertheless it has been shown by Munari & Zwitter (1997) that alternatively to these lines, the KI $\lambda 7699$ line can be used. The advantage to use this line is that it has a less steep curve of growth, thus giving the opportunity to measure higher reddenings. After applying this method to V5114 Sgr, we found $E_{B-V} = 0.45$ mag. Colors at maximum are a widely used empirical indicator of interstellar reddening. Van Den Bergh & Younger (1987) showed that novae at maximum have an intrinsic $B - V$ of 0.23 ± 0.06 mag. The same authors showed also that novae, two magnitudes below maximum, have an intrinsic

¹ Phase +1 = 1 day after maximum light.

² Version 2.04 <http://www.pa.uky.edu/~peter/atomic/>

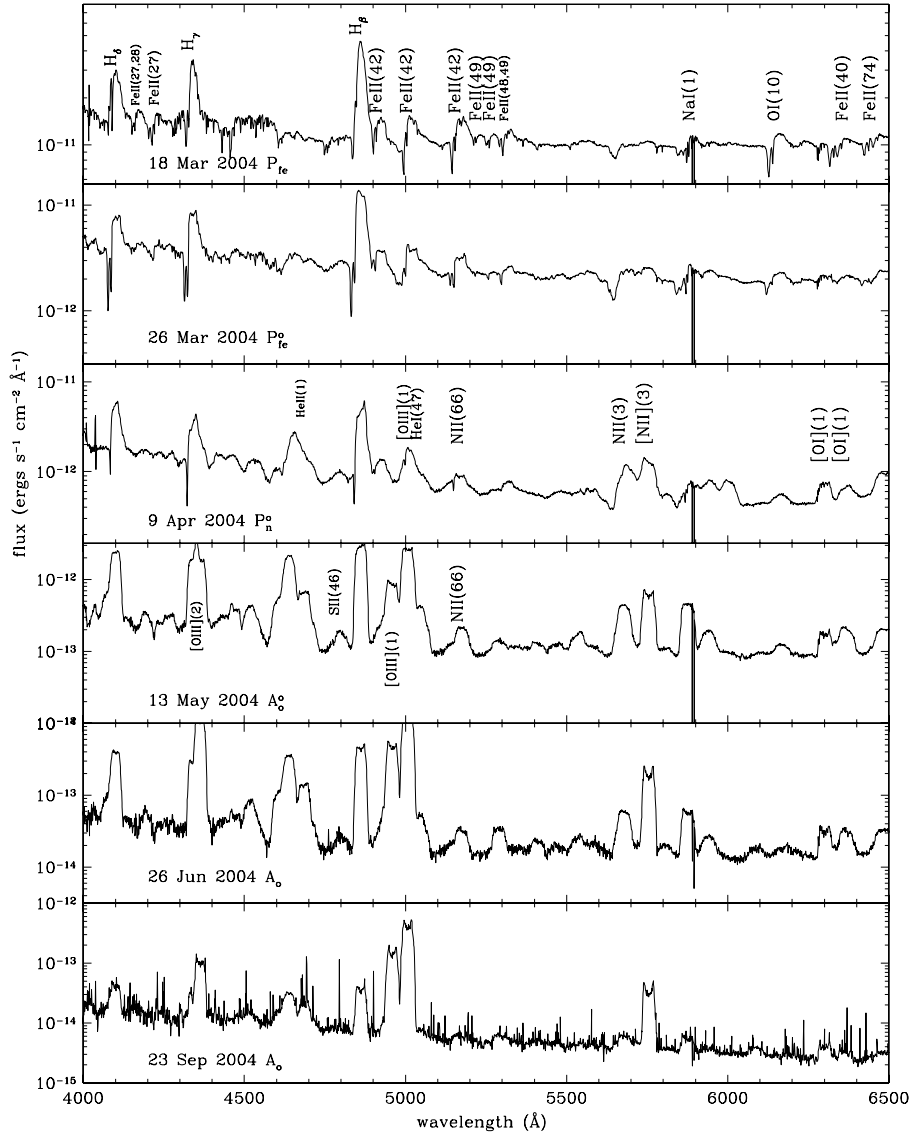


Fig. 3. Dereddened spectra of V5114 Sgr (blue part). Fluxes are in logarithmic scale to show the less intense lines.

$B - V = -0.02 \pm 0.04$ mag. As from Table 1, V5114 Sgr had $(B - V)_{\max} = 0.66$ mag and $(B - V)_{12} = 0.38$ mag. This leads to $E_{B-V} = 0.43 \pm 0.06$ mag and $E_{B-V} = 0.40 \pm 0.04$ mag, respectively. From Schlegel et al. (1998) dust maps of the Milky Way, we have derived $E_{B-V} = 0.58$ mag. Relations based on emission-line ratios are considered the best ones (see Williams 1994), unfortunately they can be used only at very late stages when the nova is in the optically thin phase. As an example, if we compute the reddening through OI $\lambda 8446$ and OI $\lambda 113164$ (following Rudy et al. 1991) we find $E_{B-V} = 1.2$ mag, which is very different from the reddening estimates derived with other methods (see Table 7). Williams (1994) suggested that if the HeII $\lambda \lambda 4686, 10124$ and the HI $\lambda \lambda 4861, 10049$ are optically thin, then they can be used to measure the absorption to the nova via:

$$E_{(B-V)_{\text{HeII}}} = 1.01 \log (4.1 F_{10124} / F_{4686})$$

and

$$E_{(B-V)_{\text{HI}}} = 1.08 \log (17 F_{10049} / F_{4861}).$$

The observed flat-topped profiles suggest that these lines likely originate in the ejected shell and therefore they may be optically

thin. Therefore we can derive $E_{B-V} = 0.57$ mag and $E_{B-V} = 0.65$ mag, respectively. A last attempt has been carried out making use of the fact that the HeI triplet ratios $\lambda \lambda 5876, 4471$ and $\lambda \lambda 10830, 4471$ are rather insensitive to density. Robbins (1968) showed that $\frac{F_{5876}}{F_{4471}} \sim 2.9$ for different values of density and temperature, while, from Osterbrock (1989), $\frac{F_{10830}}{F_{4471}} = 4.4$ for typical values of temperature and density. From the observed HeI $\lambda \lambda 5876, 4471$ lines we find zero extinction. This value is unlikely to be correct and Ferland (1977) advised against the use of these two lines because of the small baseline and the faintness of the HeI $\lambda 4471$ line. On the other hand, the HeI $\lambda \lambda 4471, 10830$ lines point towards $E_{B-V} = 0.9$ mag thus supporting rather high extinction. In the following we adopt $E_{B-V} = 0.6 \pm 0.3$ mag.

After assuming the maximum magnitude obtained in Sect. 3 with the MMRD ($M_V = -8.7 \pm 0.2$ mag) and the average absorption derived above, we derive a distance to V5114 Sgr of 9000 ± 900 pc. A complementary estimate of the distance has been obtained using the Buscombe-DeVaucouleurs relation (all novae show the same magnitude 15 days after maximum) as from Capaccioli et al. (1989). This second estimate leads

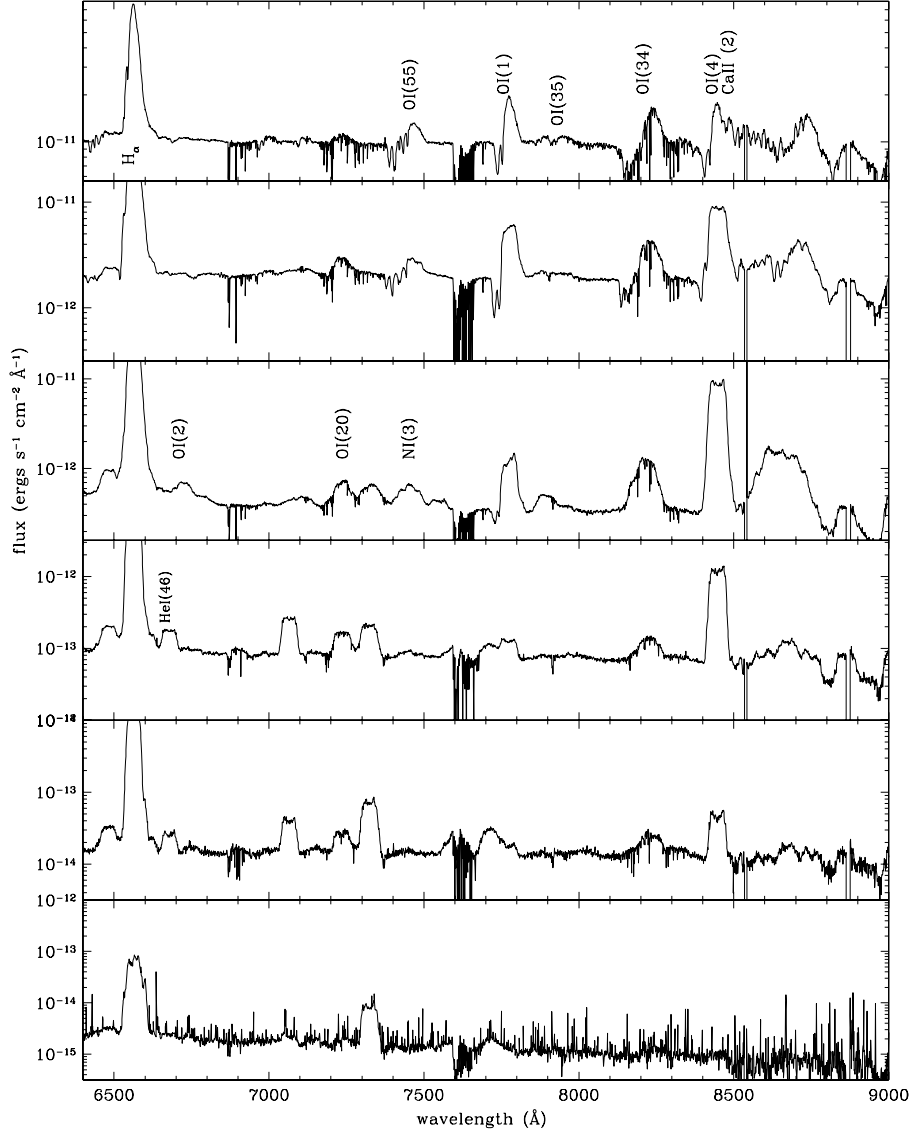


Fig. 4. Dereddened spectra of V5114 Sgr (red part). Fluxes are in logarithmic scale to show the less intense lines.

to 6000 ± 1200 pc. After taking the average weight, we get $d = 7700 \pm 700$ pc. For $l = 3^\circ.9$ and $b = -6^\circ.3$, we find the nova to be located 770 ± 70 pc above the galactic plane

6. Physical parameters

Analysis of dereddened line fluxes is the only way to derive physical parameters of nova ejecta (e.g. masses and temperatures). Since ejecta are still evolving toward the nebular stage, the line ratios are not the ones expected from atomic transition probabilities. For example, looking at Fig. 7 the $H\alpha/H\beta$ ratio converges toward the theoretical value only after phase ~ 100 .

A different explanation (also at later stages) is invoked for the case of [OI] $\lambda\lambda 6300, 6364$. This is a well known example of lines that does not respect the theoretical ratio $\sim 3:1$. Williams (1994) interpreted this as due to large optical depth in the 6300 Å line, and showed that the optical depth of that line can be derived from

$$\frac{j_{6300}}{j_{6364}} = \frac{1 - e^{-\tau}}{1 - e^{-\tau/3}}. \quad (1)$$

Following this argument, we find τ_{6300} to be in the range 1.6–6.1, in good agreement with the values exhibited by other novae. The optical depth, together with the ratio of [OI] $\lambda\lambda 6300, 5577$ can be used to determine electron temperature of the zones of the ejecta where [OI] lines are formed through the formula:

$$T_e = \frac{11\,200}{\log [43\tau / (1 - e^{-\tau}) \times F_{\lambda 6300} / F_{\lambda 5577}]}. \quad (2)$$

We find T_e to be between 3700 K and 6000 K.

The knowledge of the optical depth and electron temperature allows us to estimate the mass of oxygen in the ejecta using the 6300 Å line:

$$M_{OI} = 152 d_{\text{kpc}}^2 \exp \frac{22850}{T_e} \times 10^{1.05E(B-V)} \frac{\tau}{1 - e^{-\tau}} F_{\lambda 6300} M_{\odot}. \quad (3)$$

We find $M_{OI} = 1.9 \times 10^{-5} - 2.4 \times 10^{-7} M_{\odot}$.

Electron densities can be determined adopting the temperatures computed above and [OIII] line ratios as in Osterbrock (1989)

$$\frac{j_{4959} + j_{5007}}{j_{4363}} = 7.73 \frac{e^{3.29 \times 10^{-4} / T_e}}{1 + 4.5 \times 10^{-4} \frac{N_e}{T_e^{1/2}}}. \quad (4)$$

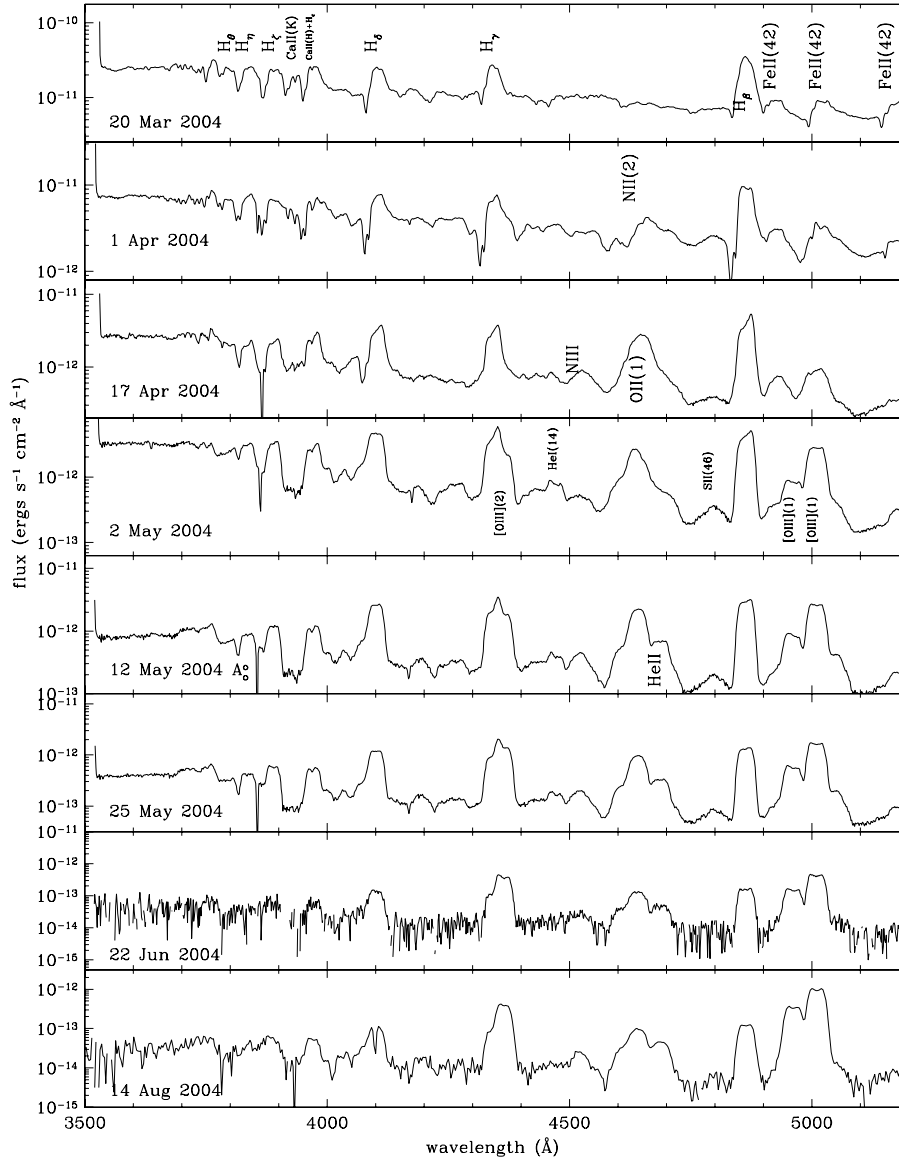


Fig. 5. Dereddened spectra of V5114 Sgr taken at CTIO and covering the region 3500–5200 Å.

The values we obtain are in the range 10^7 – 10^9 cm^{-3} , close to the upper limit of the critical densities to give rise to nebular and auroral lines. This is an indication that these lines likely arise from regions characterized by a relatively high density.

Hydrogen mass can be derived following Mustel & Boyarchuk (1970). The $H\alpha$ line is given by

$$j_{H\alpha} = g_{\alpha} \times n_e^2 \times \epsilon \times V \quad (5)$$

where g_{α} is the emission coefficient, n_e is the electron density, ϵ is the so-called “filling factor” (a measure of the clumpiness of the ejecta) and V is the volume that can be expressed as

$$V = 4\pi \times R^2 \times \delta \quad (6)$$

where the radius R can be obtained by simple expansion $v_{\text{exp}} \times \Delta t$ and $\delta = R v_{\text{ther}}/v_{\text{exp}}$ is assumed to be representative of the thickness of the shell. Deriving v_{exp} from the FWHM of Balmer lines (~ 2000 km s^{-1}) and $v_{\text{ther}} = \sqrt{3kT/m_p}$ (for $T_e \sim 5 \times 10^3$ K, $v_{\text{ther}} \sim 10$ km s^{-1}) it follows that $\delta \sim 0.01 \times R$ (cf. $v_{\text{ther}}/v_{\text{exp}} \sim 0.05$ for GK Per, Mustel & Boyarchuk 1970). On the other hand this is true only for the case of a very thin shell instantaneously

ejected. In fact, the shell ejection process occurs on a much longer time scale, of the order of t_3 . Therefore the thickness of the shell can be computed, as order of magnitude, by the ratio between $t_3 \times v_{\text{exp}}$ and $t_{\text{neb}} \times v_{\text{exp}}$, being t_{neb} the time needed by the nova to reach the nebular stage, i.e. $\delta \sim 20$ days/100 days = $0.2 \times R$. This is consistent with spectroscopic measurements obtained by Della Valle et al. (1997) for FH Ser ($\delta = 0.5 \times R$) and Humason (1940) for DQ Her ($\delta = 0.5 \times R$).

Solving Eq. (5) for ϵ ,

$$\epsilon = \frac{j_{H\alpha} d^2}{g_{\alpha} n_e^2 V} \quad (7)$$

we find that the filling factor spans a range of 7.1×10^{-2} – 7.9×10^{-4} . Finally, we can determine from

$$M_H = n_e m_H 4\pi R^3 \epsilon 0.2 \quad (8)$$

values of M_H to be $\sim 3.0 \times 10^{-5}$ – $1.1 \times 10^{-6} M_{\odot}$.

The derived physical parameters are summarized in Tables 8 and 9.

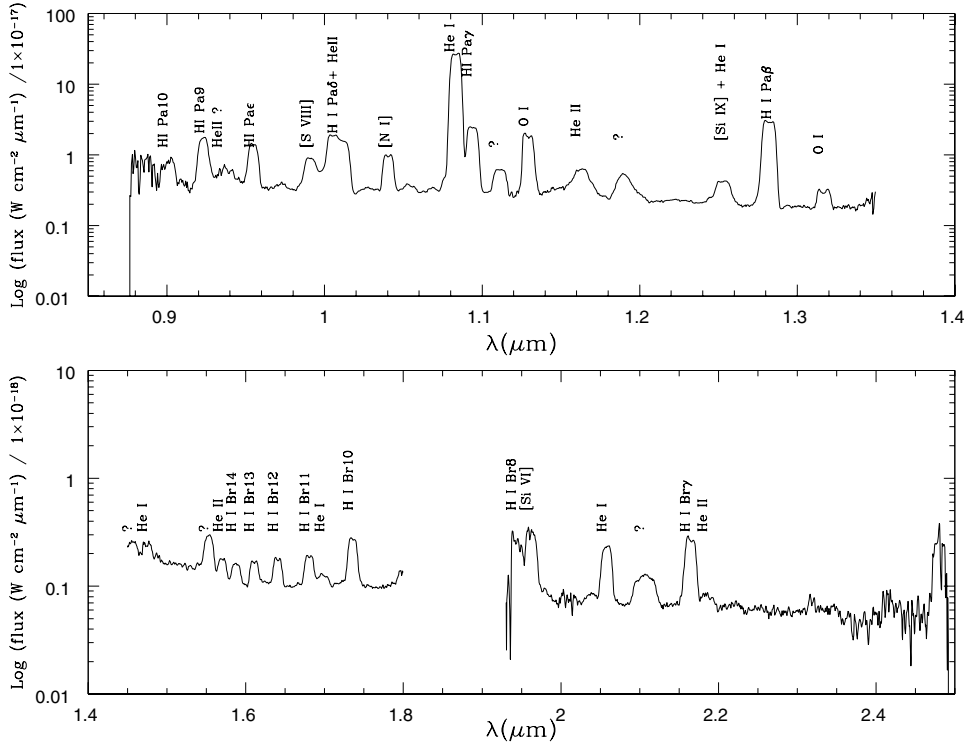


Fig. 6. Dereddened IR spectrum of V5114 Sgr observed with NIRIS at Lick Observatory on June 22, 2004 UT. Upper panel shows the “blue” part of the spectrum and the lower panel shows the “red” part.

Table 5. Reddening-corrected near infra red line fluxes (in $\text{erg s}^{-1} \text{cm}^{-2}$) observed with NIRIS at Lick Observatory on June 22, 2004, UT.

wavelength	ID	flux
0.9015	H I Pa10	1.615E-13
0.9229	H I Pa9	6.132E-12
0.9381		1.026E-13
0.9545	H I Paε	5.604E-13
0.9913	[S VIII]	4.641E-13
1.0049 + 1.0124	H I Paδ+He II	1.841E-12
1.0400	[N I]	4.165E-13
1.0534		4.482E-14
1.0830	He I	1.665E-11
1.0938	H I Paγ	1.421E-12
1.1114	?	2.438E-13
1.1287	O I	1.076E-12
1.1626	He II ?	3.708E-13
1.1911	?	2.826E-13
1.2528	[Si IX]+He I	2.144E-13
1.2818	H I Paβ	2.230E-12
1.3164	O I	1.044E-13
1.4567		5.686E-14
1.4760	He II	5.382E-14
1.5528	?	1.671E-13
1.5719	He II	3.884E-14
1.5881	H I Br14	4.318E-14
1.6109	H I Br13	6.792E-14
1.6407	H I Br12	7.936E-14
1.6806	H I Br11	1.056E-13
1.7002	He I	3.356E-14
1.7362	H I Br10	2.157E-13
1.9440		2.207E-13
1.9621		3.045E-13
2.0383		1.653E-14
2.0581	He I	2.128E-13
2.1068		1.501E-13
2.1655	H I Bry	3.019E-13

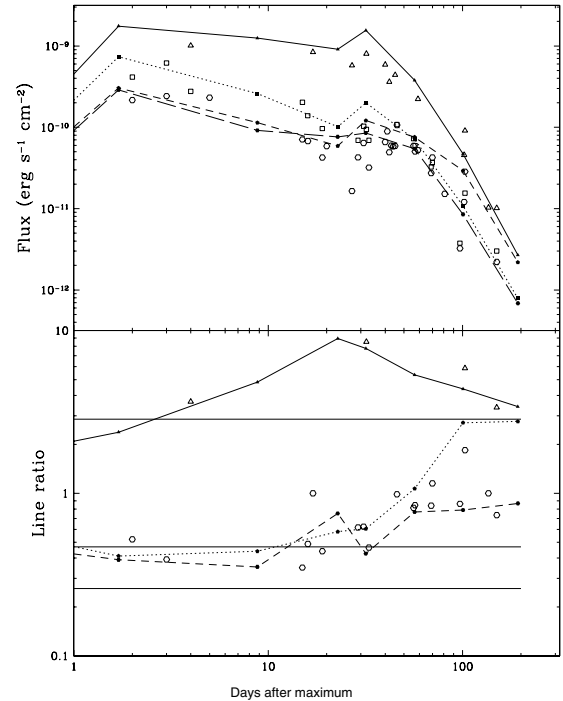


Fig. 7. Upper panel: evolution of Balmer lines fluxes. Triangles represent $H\alpha$, squares $H\beta$, pentagons $H\gamma$ and hexagons $H\delta$. Filled symbols represent data from FEROS, empty symbols from the RC spectrograph at the 1.5-m telescope at CTIO. Lines connecting the FEROS symbols are meant to show the trend of the data. Lower panel: evolution of ratios of Balmer lines divided by the $H\beta$ flux (symbols refer to the same lines as in the upper panel). Horizontal lines represent the Balmer lines ratios for case B as from Osterbrock (1989) (from top to bottom $H\alpha/H\beta$, $H\gamma/H\beta$ and $H\delta/H\beta$).

Table 6. Tentative identification of previously unidentified NIR lines.

Observed wavelength (μm)	Rest frame wavelength	Element	Transition
1.1114	11112.4	C I]	E1
	11112.43	C I	E1
	11119.98	C I	E1
	11112.	N I	E1
	11114.	N I]	E1
	11114.	N I	E1
	11116.	N I	E1
	11117.	N I]	E1
	11117.	N I]	E1
	11115.3	N II	E1
	11116.00	N II]	E1
	11120.10	N II	E1
	11108.6	N III]	E1
	11116.48	O II	E1
	11116.27	O III]	E1
1.1911	11905.516	He II	E1
	11907.4	C I]	E1
	11915.35	C I	E1
	11905.6	C II	E1
	11906.1	C IV	E1
	11908.3	C IV	E1
	11907.	N I]	E1
	11910.2	N I]	E1
	11910.2	N I	E1
	11911.005	O I]	E1
	11911.090	O I]	E1
	11914.54	O III]	E1
	1.4567	14566.14	C I]
14564.9		C II	E1
14568.7		C II	E1
14560.7		C III	E1
14560.7		C III	E1
14563.964		O I]	E1
14564.136		O I]	E1
1.5528		15514.13	He I
	15524.5	C I	E1
	15527.71	C I]	E1
	15531.3	C I	E1
	15528.085	N I]	E1
	15524.371	O I	E1
	15524.388	O I]	E1
2.0383	20373.28	He II	E1
	20379.9	N I	E1
	20389.	N I	E1
	20387.63	O II	E1
	20389.2	O II	E1
	2.1068	21067.	C III
21067.		C III	E1
21068.		C III	E1
21068.		C III	E1
21061.		C IV	E1
21067.		N I]	E1
21067.2		N I	E1
21067.6		N I	E1
21068.0		N I	E1

7. Summary and conclusions

V5114 Sgr is an Fe II nova that occurred in the bulge of the Milky Way. The rate of decline characterizes V5114 Sgr as a borderline object between the fast ($t_2 < 12$ days) and slow ($t_2 > 12$ days) classes (Della Valle & Livio 1998). These authors have shown that He/N and Fe IIb (i.e. Fe II *broad*) novae belong

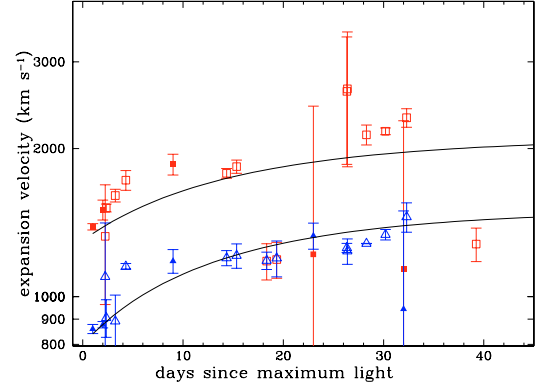


Fig. 8. Velocities of the ejecta. Measures are derived as average of P-Cyg profiles of Balmer lines. Squares refer to the faster absorption system and triangles to the slower one. Filled symbols refer to FEROS observations and empty symbols to CTIO observations. Continuous lines are exponential description of the evolution (see text).

Table 7. Measured values for reddening of V5114 Sgr.

method	reference	E_{B-V}
$(B - V)_{\text{max}}$	van den Bergh & Younger (1987)	0.43
$(B - V)_{t_2}$	van den Bergh & Younger (1987)	0.40
K I $\lambda 7699$	Munari & Zwitter (1997)	0.45
O I emission line ratios	Rudy et al. (1991)	1.23
Dust maps	Schlegel et al. (1998)	0.58
He II $\lambda 10124, 4686$	Williams (1994)	0.57
H I $\lambda 10049, 4861$	Williams (1994)	0.65
He I $\lambda 10830, 4471$	Osterbrock 1989	0.9
He I $\lambda 5876, 4471$	Robbins (1968)	0

to the fast class and are preferentially concentrated towards the galactic disc, i.e. at small z above the galactic plane ($z < 200$ pc), while Fe II novae belong to both slow and fast classes and are observed both in the disk and in the bulge, extending up to ~ 1 kpc. V5114 Sgr located at about 0.8 kpc above the galactic plane in the direction of the galactic bulge is an Fe II nova which does not represent an exception to this scenario (Della Valle et al. 1992).

Spectroscopic observations showed a dramatic change in the overall appearance of the spectrum of the nova at about 30 days after maximum light (in coincidence with the entrance in the “auroral phase”): permitted lines start fading at a different rate, P-Cyg profiles disappear, and velocities reach a “plateau” phase and $U - V$ and $J - K$ colors increase. The lack of detection of [Ne III] line hints for a “standard” evolution of the nebular spectrum (see Williams et al. 1994).

The very high values of optical depth in the O I $\lambda 6300$ suggest very high densities for the zones of the ejecta where these lines are formed. We have derived the filling factor in the range $\sim 7.1 \times 10^{-2}$ to 7.9×10^{-4} . Comparing these values with other values reported in the literature (see Table 10), two facts emerge: a) the filling factors in nova ejecta are definitely smaller than 1, likely close to 0.1, during the early stages; b) these values decrease by 1–3 orders of magnitude with time. This fact indicates that the volume of the expanding shell (computed with Eq. (6)) increases with time more rapidly than the volume actually occupied by most of the ejected material. In other words, the decreasing trend exhibited by the filling factor suggests that the ejected matter tends to remain clumped in sub-structures having higher density than the average density characterizing the expanding shell. This is consistent with the [O I] $\lambda 6300, 6364$ ratio < 2 (Table 8), which suggests that [O I] lines should be formed in

Table 8. Physical parameters for V5114 Sgr: O I $\lambda\lambda 6300/6363$ line ratios, optical depth in $\lambda 6300$ (τ_{6300}), electron temperature, O I mass, [O III] line ratios and electron densities.

Date	j_{6300}/j_{6363}	τ	T_e (K)	$M_{O I}(M_{\odot})$	$\frac{j_{4959}+j_{5007}}{j_{4363}}$	$N_e(\text{cm}^{-3})$
Jun. 26	1.74	2.10	4390	3.02E-6	1.60	1.28E9
Jun. 28	1.15	6.09	3659	1.87E-5	2.37	3.53E9
Aug. 14	1.93	1.60	5743	2.36E-7	5.69	7.01E8
Sep. 26	1.17	5.69	6022	5.73E-7	23.79	1.30E7

Table 9. Physical parameters for V5114 Sgr: ϵ and the hydrogen mass.

Date	ϵ	$M_H (M_{\odot})$
Jun. 26	7.08E-2	3.02E-5
Jun. 28	7.04E-2	2.91E-5
Aug. 14	3.97E-3	6.49E-6
Sep. 26	7.92E-4	1.10E-6

Table 10. Comparison between filling factors in literature.

Nova	Filling factor value	Reference
T Pyx	10^{-2} – 10^{-5}	Shara et al. (1997)
Nova Vel 1999	0.1	Della Valle et al. (2002)
Nova SMC 2001	0.1 – 10^{-3}	Mason et al. (2005)
Nova LMC 2002	10^{-2} – 10^{-4}	Mason et al. (2005)
V5114 Sgr	0.1 – 10^{-3}	this paper

very dense small blobs of neutral material embedded within the ionized shell (see paragraph 6). However, recently Williams & Mason (2006) proposed an alternative interpretation of this behavior, assuming that the [O I] lines arise in regions of high magnetic field, and their intensity and profile are modified by Quadratic Zeeman Effect.

Computed oxygen and hydrogen masses are in the ranges 1.9×10^{-5} – $2.4 \times 10^{-7} M_{\odot}$ and 3.0×10^{-5} – $1.1 \times 10^{-6} M_{\odot}$. This high mass ratio is close to the upper limit for classical novae shown in Warner (1995).

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Online Material

Table 2. V5114 Sgr reddening-corrected emission line fluxes (in $\text{erg s}^{-1} \text{cm}^{-2}$) from FEROS spectra.

Ident. wavelength	Mar. 18	Mar. 19	Mar. 26	Apr. 9	Apr. 18	May 13	Jun. 26	Sep. 23
He I 4026	–	–	3.061–11	–	–	2.158E–12	4.457E–13	3.158E–13
[S II] 4076	–	–	–	–	–	1.193E–11	2.827E–12	–
H δ 4101	3.440E–10	4.820E–10	1.925E–10	2.202E–11	6.402E–10	4.634E–11	1.266E–11	2.982E–13
Fe II(27,28) 4173/78	7.030E–11	7.445E–11	3.684E–11	1.603E–12	–	–	–	–
Fe II(27) 4233	6.490E–11	4.450E–11	2.704E–11	1.258E–12	–	–	–	–
Fe II(27) 4273	–	–	2.450E–11	1.247E–12	3.213E–12	2.785E–12	–	–
H γ 4340	4.055E–10	4.516E–10	2.401E–10	1.791E–11	8.476E–11	4.197E–11	1.059E–11	2.033E–13
[O III](2) 4363	–	–	–	–	–	4.012E–11	4.249E–11	6.652E–13
Fe II(27) 4417	–	–	7.224E–12	3.159E–12	–	–	–	–
He I 4438	–	–	5.393E–12	–	–	–	–	–
He I(14) 4472	–	–	2.768E–11	3.384E–12	5.158E–12	3.990E–12	6.806E–13	–
N III 4517	–	–	–	4.260E–12	1.085E–11	5.969E–12	2.921E–12	–
Fe II(37) 4629	2.678E–11	2.365E–11	–	–	–	–	–	–
N II(5)+N III(2)+O II(1)4601/34/49	–	–	–	2.074E–11	1.243E–10	6.591E–11	1.843E–11	2.831E–13
Fe II(37) 4665	2.742E–11	1.917E–11	4.826E–11	–	–	–	–	–
He II(1) 4686	–	–	–	–	–	1.490E–11	5.556E–12	1.517E–13
S II(46) 4792	–	–	–	1.819E–12	2.729E–12	2.431E–12	5.134E–13	–
H β 4861	8.489E–10	9.439E–10	5.244E–10	2.819E–11	1.386E–10	7.492E–11	1.889E–11	2.655E–13
Fe II(42) 4924	1.495E–10	1.741E–10	6.124E–11	4.031E–12	1.429E–11	–	–	–
[O III](1) 4959	–	–	–	–	–	2.165E–11	1.692E–11	3.913E–12
[O III](1) 5007	–	–	–	5.972E–11	3.967E–11	7.599E–11	5.095E–11	1.191E–11
Fe II(42) 5018	2.633E–10	2.612E–10	1.252E–10	–	–	–	–	–
Fe II(52) 5169	2.121E–10	2.060E–10	5.872E–11	2.034E–11	9.650E–12	4.177E–12	9.209E–13	–
Fe II(49) 5234	8.313E–11	3.980E–11	7.826E–12	–	–	–	–	–
Fe II(49) 5276	5.449E–11	4.539E–11	1.102E–11	3.140E–12	–	–	–	–
??? 5290	–	–	–	–	–	1.432E–12	7.422E–13	–
Fe II(49) 5317	–	–	–	8.981E–12	7.669E–12	–	–	–
Fe II(48) 5337	8.597E–11	1.102E–10	2.566E–11	–	–	–	–	–
He II(2) 5412	–	–	–	–	–	6.440E–13	2.642E–13	–
[O I](1) 5577	–	–	1.878E–11	5.771E–12	2.789E–12	6.216E–13	2.222E–13	<8.498E–14
N II(3) 5667/76/77/86	–	–	–	3.282E–11	1.801E–11	1.203E–11	1.983E–12	1.057E–13
[N II](3) 5755	–	2.952E–11	3.544E–11	4.246E–11	4.583E–11	1.703E–11	7.536E–12	9.420E–13
He I(11) 5876	–	–	–	–	1.412E–11	1.090E–11	1.697E–12	8.809E–14
Na I(1) 5890	2.423E–11	1.510E–11	1.978E–11	1.501E–11	–	–	–	–
N II(28) 5940	–	1.924E–11	2.267E–11	1.761E–11	2.880E–12	3.253E–12	5.892E–13	–
??? 6004	–	–	–	1.961E–11	8.580E–12	–	–	–
O I(10) 6159	2.423E–11	3.686E–11	1.887E–11	6.049E–12	3.257E–12	1.310E–12	–	–
Fe II(74) 6247	1.022E–11	1.440E–11	4.418E–12	–	–	–	–	–
[O I] 6300	–	–	8.803E–12	1.152E–11	7.271E–12	2.719E–12	7.685E–13	2.507E–14
[O I](1) 6364	–	–	–	4.192E–12	3.995E–12	2.585E–12	4.408E–13	2.138E–14
Fe II(40) 6370	9.104E–12	1.834E–11	1.324E–11	4.727E–12	–	–	–	–
N II(8) 6482	–	–	2.507E–11	1.983E–11	1.035E–11	4.192E–12	8.092E–13	1.470E–14
H α 6563	4.634E–10	1.124E–9	1.598E–9	1.152E–9	9.801E–10	3.601E–10	5.965E–11	8.597E–13
He I(46) 6678	–	–	9.063E–12	–	6.050E–12	2.689E–12	5.309E–13	–
O I(2) 6726	9.166E–12	7.611E–12	1.125E–11	1.971E–11	7.864E–12	–	–	–
O I(21) 7002	1.212E–11	9.505E–12	9.552E–12	1.224E–12	–	–	–	–
He I(10) 7065	–	–	–	–	1.023E–11	6.891E–12	1.257E–12	1.106E–14
??? 7113	–	9.514E–12	–	–	3.531E–12	–	–	–
C II(3) 7231/36	1.195E–12	4.984E–12	3.402E–11	1.951E–11	5.706E–12	3.637E–12	6.099E–13	–
[O II](2) 7319/20/30/31	–	–	–	1.881E–11	1.913E–11	5.491E–12	3.060E–12	9.895E–14
N I(3)+O I(55) 7468 + 7476	3.876E–11	7.413E–11	5.497E–11	1.793E–11	6.388E–12	–	–	–
??? 7546	–	–	–	–	2.686E–12	1.298E–12	–	–
O I(1) 7772/4/5	7.000E–11	1.579E–10	2.013E–10	4.869E–11	1.936E–11	2.504E–12	–	–
Mg II(8) 7896	–	–	–	9.620E–12	5.092E–12	–	–	–
O I(35) 7947	2.086E–11	1.997E–11	–	–	2.286E–12	8.680E–13	–	–
O I(34) 8223	7.846E–11	1.374E–10	1.194E–10	6.634E–11	2.467E–11	3.605E–12	1.050E–12	–
O I(4) 8446/7	9.500E–11	1.307E–10	4.585E–10	5.154E–10	3.448E–10	4.844E–11	2.062E–12	–
Ca II 8498	2.659E–11	4.547E–11	–	–	–	–	–	–
Ca II 8662?	8.537E–12	1.371E–11	–	–	–	–	–	–

Table 3. V5114 Sgr reddening-corrected emission line fluxes (in $\text{erg s}^{-1} \text{cm}^{-2}$) from CTIO spectra with wavelength range larger than 2000 Å and including the H α region.

Ident. wavelength	Mar. 21	Apr. 18	Jun. 28	Aug. 14
He + Ca II 3968/70	–	–	5.709E–12	8.794E–13
H δ 4101	–	–	1.823E–11	1.910E–12
H γ 4340	–	–	5.539E–11	1.857E–12
[O III](2) 4363	–	–	–	8.853E–12
N III 4517	–	–	1.232E–12	2.403E–13
4640	–	–	1.277E–11	3.021E–12
He II(1) 4686	–	–	3.334E–12	1.060E–12
H β 4861	3.295E–10	1.046E–10	1.077E–11	3.130E–12
Fe II(42) 4924	4.610E–11	1.226E–11	–	–
[O III](1) 4958	–	–	1.55E–11	1.295E–11
[O III](1) 5007	–	–	5.69E–11	3.748E–11
Fe II(42) 5018	1.751E–10	5.317E–11	–	–
Fe II(52) 5169	9.638E–11	9.411E–12	7.794E–13	2.891E–13
Fe II(49) 5276	4.055E–12	–	–	–
??? 5296	–	–	7.250E–13	3.543E–13
Fe II(48) 5337	2.020E–11	5.505E–12	–	–
He II(2) 5412	–	–	2.154E–13	1.123E–13
??? 5474	–	–	9.443E–14	–
5540	1.432E–11	–	2.917E–13	–
[O I](1)5577	–	–	1.505E–13	<2.360E–13
N II(3)5667/76/77/78	1.870E–11	2.218E–11	1.975E–12	4.354E–13
[N II](3) 5755	1.869E–11	4.821E–11	8.219E–12	4.477E–12
He I(11) 5876	–	1.764E–11	2.054E–12	4.346E–13
Na I(1) 5890	2.015E–11	–	–	–
N II(28) 5940	1.520E–11	1.197E–12	6.378E–13	1.047E–13
??? 6004	–	8.634E–12	–	–
O II(10) 6159	3.056E–11	3.412E–12	–	–
Fe II(74) 6247	9.383E–12	–	–	–
[O I]6300	2.858E–12	7.739E–12	6.599E–13	2.445E–13
[O I]6363 + FeII(40)6370	1.239E–11	3.287E–12	5.745E–13	1.265E–13
N II(8) 6482	7.461E–12	5.879E–12	3.771E–13	1.263E–13
H α 6563	1.240E–9	9.902E–10	5.747E–11	1.283E–11
He I(46) 6678	1.571E–12	3.450E–12	4.168E–13	1.020E–13
O I(2) 6726	6.578E–12	4.497E–12	–	–
He I(10) 7065	–	1.346E–11	1.167E–12	–
C II(3) 7231/36	–	6.288E–12	5.851E–13	–
[O II](2) 7319/20/30/31	–	2.147E–11	2.717E–12	–
N I(3)+O I(55) 7468 + 7476	5.187E–11	6.109E–12	–	–
??? 7586	–	–	1.717E–13	–
??? 7630	–	–	7.579E–14	–
??? 7731	–	–	9.544E–13	–
O I(1) 7772/4/5	2.315E–10	2.231E–11	–	–
Mg II(8) 7896	–	5.863E–12	–	–
O I(34) 8223	1.302E–10	2.734E–11	3.385E–13	–
O I(4) 8446	2.825E–10	4.794E–10	9.237E–13	–
??? 8704	–	–	6.435E–13	–
??? 9030	–	–	2.261E–13	–
??? 9245	–	–	9.574E–13	–
??? 9556	–	–	2.905E–13	–

Table 4. V5114 Sgr reddening-corrected emission line fluxes (in $\text{erg s}^{-1} \text{cm}^{-2}$) from CTIO spectra shown in Fig. 5.

Ident. wavelength	Mar. 20	Apr. 1	Apr. 17	May 2	May 12	May 25	Jun. 22	Aug. 14
??? 3765	1.553E-10	4.039E-11	1.1117E-11	3.114E-12	1.402E-11	1.089E-11	6.005E-12	–
H10 3798	9.596E-11	3.257E-11	–	6.295E-12	–	–	–	–
H9 3835	2.030E-10	8.256E-11	2.219E-11	2.242E-11	7.722E-12	6.032E-12	–	–
H8 3889	2.412E-10	9.473E-11	3.930E-11	4.155E-11	1.984E-11	1.480E-11	7.604E-12	1.146E-12
Ca II(K)3933	1.137E-10	–	–	–	–	–	–	–
H ϵ + Ca II 3968/70	3.719E-10	1.144E-10	4.848E-11	5.301E-11	2.259E-11	2.156E-11	5.386E-12	8.293E-13
H δ 4101	3.644E-10	1.142E-10	7.861E-11	1.079E-10	8.113E-11	5.493E-11	1.760E-12	3.798E-12
Fe II(27,28) 4173/78	5.184E-11	9.433E-12	–	–	–	–	–	–
Fe II(27) 4233	3.419E-11	–	–	–	–	–	–	–
H γ 4343	4.241E-10	1.465E-10	8.000E-11	2.172E-10	5.656E-11	4.077E-11	1.341E-11	1.857E-12
[O III]4363	–	–	–	8.706E-11	5.334E-11	6.423E-11	4.605E-11	8.853E-12
4422	–	1.413E-11	–	–	–	–	–	–
4468	–	2.312E-11	–	1.029E-11	–	2.687E-12	–	–
N III4517	–	–	–	3.199E-12	–	5.006E-12	2.617E-12	2.403E-13
4640	–	1.275E-10	1.202E-10	8.310E-11	9.007E-11	5.906E-11	2.167E-11	3.021E-12
He II(1)4686	–	–	–	–	2.039E-11	1.518E-11	6.264E-12	1.060E-12
H β 4861	8.560E-10	2.517E-10	1.402E-10	1.138E-10	1.028E-10	6.833E-11	2.161E-11	3.130E-12
Fe II(42) 4924	7.125E-11	3.407E-11	1.484E-11	–	–	–	–	–
[O III]4959	–	–	–	1.961E-11	2.661E-11	3.037E-11	2.255E-11	1.295E-11
[O III]5007	–	–	–	8.073E-11	1.008E-10	9.709E-11	6.511E-11	3.748E-11
Fe II(42) 5018	1.581E-10	1.007E-10	2.767E-11	–	–	–	–	–
Fe II(52) 5169	9.383E-11	2.288E-11	6.805E-12	3.897E-12	5.086E-12	2.296E-12	–	–

Appendix A: Logs of the observations**Table A.1.** Measured magnitudes during photometric campaigns (described in text).

JD	<i>U</i>	<i>B</i>	<i>V</i>	<i>R</i>	<i>I</i>	<i>J</i>	<i>H</i>	<i>K</i>
2453 083.25	8.70	9.11	8.45	7.91	7.48	6.72	6.63	6.24
2453 084.00	8.87	9.23	8.65	8.01	7.38	–	–	–
2453 084.25	8.74	9.24	8.73	8.07	7.55	–	–	–
2453 084.50	8.87	9.23	8.65	8.01	7.38	–	–	–
2453 085.00	9.08	9.53	9.00	8.19	7.42	–	–	–
2453 085.50	9.08	9.53	9.00	8.19	7.42	–	–	–
2453 088.25	9.19	9.81	9.40	8.46	7.93	7.37	7.10	6.74
2453 089.25	9.38	10.01	9.63	8.57	8.12	7.54	7.37	6.97
2453 091.25	9.54	10.10	9.70	8.74	8.38	7.79	7.59	7.25
2453 091.50	9.83	10.17	9.73	8.80	8.28	–	–	–
2453 092.25	9.61	10.22	9.78	8.79	8.42	7.86	7.72	7.35
2453 092.50	10.02	10.34	9.88	8.96	8.41	–	–	–
2453 093.25	9.85	10.49	10.08	9.02	8.67	8.11	8.01	7.66
2453 093.50	10.26	10.60	10.22	9.13	8.67	–	–	–
2453 094.50	10.38	10.73	10.39	9.28	8.71	–	–	–
2453 097.25	10.20	10.78	10.53	9.37	9.01	8.52	8.56	8.18
2453 099.25	10.39	11.00	10.70	9.55	9.15	8.70	8.72	8.36
2453 100.50	10.49	11.16	10.84	9.59	9.19	8.76	8.84	8.46
2453 101.25	10.57	11.14	10.87	9.60	9.20	8.78	8.89	8.52
2453 102.25	10.63	11.24	10.93	9.65	9.27	8.88	8.97	8.57
2453 103.25	10.64	11.23	11.02	9.73	9.32	8.95	9.02	8.64
2453 104.25	10.69	11.34	11.05	9.77	9.33	8.97	9.05	8.64
2453 105.25	10.90	11.51	11.24	9.86	9.37	9.06	9.20	8.77
2453 108.25	11.17	11.77	11.60	9.94	9.50	9.23	9.40	8.89
2453 109.25	11.14	11.77	11.56	9.96	9.54	9.26	9.44	8.89
2453 110.25	11.25	11.90	11.76	10.04	9.64	9.39	9.61	9.03
2453 112.00	11.64	12.00	11.69	10.11	9.62	–	–	–
2453 113.25	11.39	12.07	11.95	10.10	9.77	9.44	9.68	8.99
2453 114.00	11.81	12.13	11.95	10.17	9.72	–	–	–
2453 126.25	11.85	12.44	12.37	10.57	10.70	10.16	10.19	9.50
2453 132.25	12.03	12.53	12.50	10.87	11.08	10.48	10.45	9.72
2453 136.25	12.10	12.53	12.58	10.97	11.24	10.67	10.59	9.91
2453 156.27				<i>unfiltered</i>	12.36			
2453 186.25	14.35	14.11	13.67	13.20	13.60	–	–	–
2453 187.00	14.39	14.16	13.68	13.25	13.66	–	–	–
2453 188.25	14.37	14.17	13.68	13.26	13.58	–	–	–
2453 189.25	14.44	14.19	13.77	13.40	–	–	–	–
2453 190.25	14.46	14.21	13.74	13.34	13.76	–	–	–
2453 191.25	14.51	14.25	13.76	13.40	13.81	–	–	–
2453 192.00	14.53	14.26	13.77	13.42	13.82	–	–	–
2453 193.00	14.56	14.31	13.82	13.46	13.82	–	–	–
2453 194.25	14.59	14.33	13.83	13.50	13.90	–	–	–
2453 195.25	14.62	14.33	13.83	13.53	13.95	–	–	–
2453 196.25	14.66	14.40	13.86	13.58	13.99	–	–	–
2453 197.25	14.69	14.39	13.87	13.61	14.00	–	–	–
2453 198.25	14.68	14.43	13.88	13.62	14.04	–	–	–
2453 200.00	14.78	14.50	13.92	13.70	14.08	–	–	–
2453 537.83	>18.6	18.81	17.49	17.69	17.92	–	–	–
2453 541.87	18.70	18.77	17.50	17.63	18.02	–	–	–

Table A.2. Log of the spectroscopic observations for V5114 Sgr.

Date (UT)	Instrument	Exp.time (s)	Wavelength range (Å)	Resolution (or scale)
Mar. 18.3	FEROS	240	4000–9000	48 000
Mar. 19.3	FEROS	400	4000–9000	48 000
Mar. 19.3	SMARTS	360	4000–5000	0.77
Mar. 19.3	SMARTS	360	4000–5000	0.77
Mar. 20.3	SMARTS	270	3500–5300	1.5
Mar. 21.4	SMARTS	120	4800–9500	5.6
Mar. 22.4	SMARTS	600	3900–4500	0.6
Mar. 26.4	FEROS	632	4000–9000	48 000
Apr. 1.4	SMARTS	360	3500–5300	1.5
Apr. 2.3	SMARTS	360	4000–4900	0.77
Apr. 3.3	SMARTS	300	5900–7700	1.5
Apr. 5.3	SMARTS	360	3500–5300	1.5
Apr. 6.4	SMARTS	360	3900–4550	0.56
Apr. 9.4	FEROS	900	4000–9000	48 000
Apr. 13.3	SMARTS	360	5600–7000	1.1
Apr. 13.4	SMARTS	360	3800–4550	0.56
Apr. 13.4	SMARTS	360	4050–4750	0.56
Apr. 15.4	SMARTS	540	4000–5000	0.77
Apr. 16.3	SMARTS	720	3500–5300	1.48
Apr. 17.3	SMARTS	360	3500–5300	1.48
Apr. 18.4	FEROS	900	4000–9000	48 000
Apr. 18.4	SMARTS	300	4800–9500	5.6
Apr. 18.4	SMARTS	180	4800–9500	5.6
Apr. 19.4	SMARTS	480	3800–5600	1.48
Apr. 26.2	SMARTS	450	5600–6950	1.1
Apr. 26.4	SMARTS	540	3850–4550	0.56
Apr. 27.3	SMARTS	450	3870–4550	0.56
Apr. 28.2	SMARTS	600	3870–4550	0.56
Apr. 28.4	SMARTS	600	5650–7000	1.1
Apr. 29.2	SMARTS	600	3870–4550	0.56
Apr. 30.2	SMARTS	450	3870–4550	0.56
May 1.2	SMARTS	450	3870–4550	0.56
May 1.4	SMARTS	450	5600–7000	1.1
May 2.2	SMARTS	360	3500–5300	1.48
May 12.4	SMARTS	360	3500–5300	1.48
May 13.3	FEROS	1500	4000–9000	48 000
May 13.4	SMARTS	360	3500–5300	1.48
May 14.3	SMARTS	360	3500–5300	1.48
May 15.1	SMARTS	450	5650–7000	1.1
May 15.4	SMARTS	450	3900–4550	0.56
May 25.3	SMARTS	720	3500–5300	1.48
May 26.3	SMARTS	720	4000–5000	0.77
Jun. 6.2	SMARTS	720	4050–4750	0.56
Jun. 22	NIRIS	?	4500–25000	?
Jun. 22.2	SMARTS	360	3500–5300	1.48
Jun. 26.2	FEROS	3600	4000–9000	48 000
Jun. 27.1	SMARTS	450	5650–7000	1.1
Jun. 27.4	SMARTS	720	3850–4550	0.56
Jun. 28.1	SMARTS	360	4800–9600	5.67
Jul. 15.1	SMARTS	600	3550–5300	1.48
Jul. 31.2	SMARTS	540	5650–7000	1.10
Aug. 14.2	SMARTS	600	3450–6900	2.88
Sep. 26.2	FEROS	7200	4000–9000	48 000