

Elemental abundances in the atmosphere of clump giants^{*,**}

T. V. Mishenina¹, O. Bienaymé², T. I. Gorbaneva¹, C. Charbonnel^{3,4}, C. Soubiran⁵, S. A. Korotin¹, and V. V. Kovtyukh¹

¹ Astronomical Observatory of Odessa National University and Isaac Newton Institute of Chile, Odessa Branch, Shevchenko Park, 65014 Odessa, Ukraine

e-mail: tamar@deneb1.odessa.ua

² Observatoire Astronomique de l'Université Louis Pasteur, 11 rue de l'Université, 67000 Strasbourg, France

³ Geneva Observatory, 1290 Sauverny, Switzerland

⁴ LATT CNRS UMR 5572, 14 Av. E. Belin, 31400 Toulouse, France

⁵ Observatoire Aquitain des Sciences de l'Univers, CNRS UMR 5804, BP 89, 33270 Floirac, France

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ABSTRACT

Aims. The aim of this paper is to provide the fundamental parameters and abundances for a large sample of local clump giants with a high accuracy. This study is a part of a big project, in which the vertical distribution of the stars in the Galactic disc and the chemical and dynamical evolution of the Galaxy are being investigated.

Methods. The selection of clump stars for the sample group was made applying a colour-absolute magnitude window to nearby Hipparcos stars. The effective temperatures were estimated by the line depth ratio method. The surface gravities ($\log g$) were determined by two methods (the first one was the method based on the ionization balance of iron and the second one was the method based on fitting of the wings of the Ca I 6162.17 Å line). The abundances of carbon and nitrogen were obtained from the molecular synthetic spectrum, and the Mg and Na abundances were derived using the non-LTE approximation. The “classical” models of stellar evolution without atomic diffusion and rotation-induced mixing were employed.

Results. The atmospheric parameters (T_{eff} , $\log g$, $[\text{Fe}/\text{H}]$, V_r) and Li, C, N, O, Na, Mg, Si, Ca, and Ni abundances in 177 clump giants of the Galactic disc were determined. The underabundance of carbon, overabundance of nitrogen, and “normal” abundance of oxygen were detected. A small sodium overabundance was found. A possibility of a selection of the clump giants based on their chemical composition and the evolutionary tracks was explored.

Conclusions. The theoretical predictions based on the classical stellar evolution models are in good agreement with the observed surface variations of the carbon and nitrogen just after the first dredge-up episode. The giants show the same behaviour of the dependencies of O, Mg, Ca, and Si (α -elements) and Ni (iron-peak element) abundances vs. $[\text{Fe}/\text{H}]$ as dwarfs do. This allows us to use such abundance ratios to study the chemical and dynamical evolution of the Galaxy.

Key words. stars: abundances

1. Introduction

Low-mass stars (i.e., below $\sim 2.3 M_{\odot}$) climb the red giant branch (hereafter RGB) with a degenerate He core whose mass increases until it reaches a critical value of about $0.45 M_{\odot}$ at the top of the RGB. At this point, helium ignites in a series of flashes removing this degeneracy. The star then becomes a “clump” giant, which undergoes the central He burning. All the low-mass stars have similar core masses at the beginning of He burning, and hence similar luminosities. Due to this fact, the red giants at this evolutionary stage exhibit a specific feature in the colour-magnitude diagram (CMD), called the “clump”. According to evolutionary tracks and in agreement with the observation of open clusters, all giants older than about 1 Gyr fall in the clump (Girardi 1999).

Clump giants are especially interesting in two aspects. Their intrinsic brightness combined with the numerous and sharp features of their spectra make them good tracers for Galactic

kinematics and chemistry. They are also very useful in clarifying the advanced evolutionary stages of the low mass stars. Until now, the comprehensive investigation of a large sample of these stars has not yet been carried out.

This work is a part of our study of the Galactic disc surface mass density (Siebert et al. 2003; Bienaymé et al. 2006) of the properties of both thin and thick discs, and the abundance trends in the solar neighbourhood (Soubiran et al. 2003; Mishenina et al. 2004; Soubiran & Girard 2005). In this part of the project, the local clump giants observed with high spectral resolution and high signal-to-noise ratios (S/N) serve as reference stars. Here we take the opportunity to provide strong observational constraints for the theory of stellar evolution. As a matter of fact, the chemical composition of the clump giant atmospheres reflects both the chemical composition of the prestellar matter and nucleosynthesis and mixing processes inside the stars. Therefore, for successful Galactic studies we need to know the abundances of which chemical elements are not affected by the mixing processes. On the other hand, we can use the elements with abundances affected by the mixing processes to distinguish different stages of the stellar evolution, in particular the clump phase.

* Based on spectra collected with the ELODIE spectrograph at the 1.93-m telescope of the Observatoire de Haute Provence (France).

** Tables A.1–A.4 are only available in electronic form at <http://www.edpsciences.org>

While investigating the clump giants, we faced the problem of their selection, since this region of the CMD is also occupied by the stars of the ascending giant branch. The differentiation between the first-ascending RGB stars and the “clump” stars is rather complicated. Even for open cluster stars, it is very difficult to establish, with a good level of certainty, which stars from the group under investigation are the real clump ones (Pasquini et al. 2004). We can solve the problem of the correct selection using the extended observational data on the clump giants, selected with photometric criteria. However, is it possible to identify the clump giants based on additional criteria, including the data about their chemical composition?

When the star moves towards the RGB, the superficial convection zone deepens and the nuclearly processed material penetrates into the atmosphere, changing its chemical composition. During this so-called first dredge-up phase, the surface abundances of Li, C, N, and Na, together with the $^{12}\text{C}/^{13}\text{C}$ ratio, are being altered. The effect depends both on the stellar mass and metallicity (see Charbonnel 1994). Typically, the surface abundance of carbon decreases by $\sim 0.1\text{--}0.2$ dex, and that of nitrogen increases by 0.3 dex or more (Iben 1991).

Despite a large dispersion, the abundances of CNO elements and their isotopes observed previously in the giants of the solar metallicity (Lambert & Ries 1981; Kjørgaard et al. 1982) were found to be in a good agreement with theoretical predictions (Iben 1991). We reinvestigate this problem with a larger sample of stars that have been observed with higher quality.

For the giants and supergiants of solar metallicity, the sodium overabundance has been found in many studies (Cayrel de Strobel et al. 1970; Korotin & Mishenina 1999; Boyarchuk et al. 2001; Andrievsky et al. 2002). As has been shown in some papers (Mashonkina et al. 1993; Korotin & Mishenina 1999), the Na overabundance cannot only be explained by deviations from LTE. In this paper, we look into this problem using recent theoretical considerations and track calculations. The aim of this paper is to provide the fundamental parameters and abundances for a large sample of local clump giants, determined with a high accuracy, and to use these data to 1) probe the vertical distribution of the stars in the Galactic disc; 2) investigate the Galactic chemical evolution; and 3) explore the possibility of selecting the clump giants based on their elemental abundances.

In Sect. 2, we described the photometric selection criteria for the clump giants and the detail of our spectroscopic observations. In Sect. 3, the determination of the atmospheric parameters is presented. In Sect. 4, the chemical abundances are determined. In Sect. 5, we discussed the behaviour of each element in our sample with metallicity. In Sect. 6, we provided the analysis of the signs of the first dredge-up in stars under current study.

2. Selection of the stars and observations

The clump stars analysed in this paper were selected using a colour-absolute magnitude window applied to nearby Hipparcos stars. Stars were selected from the Hipparcos catalogue according to the following criteria:

$$\begin{aligned} \pi &\geq 10 \text{ mas} \\ \delta_{\text{ICRS}} &\geq -20 \text{ deg} \\ 0.8 &\leq B - V \leq 1.2 \\ 0 &\leq M_V \leq 1.6 \end{aligned}$$

where π is the Hipparcos parallax and δ_{ICRS}^1 the declination. The Johnson $B - V$ colour was transformed from the *Tycho2* $B_T - V_T$ colour, applying Eq. (1.3.20) from ESA (1997):

$$B - V = 0.850 (B_T - V_T).$$

The absolute magnitude M_V was computed using the V magnitude transformed from the Hipparcos H_p magnitude to the Johnson system with the equation calibrated by Harmanec (1998).

Among the nearly 400 giants selected, about a half of them were observed: the first priority group contained stars, for which radial velocities were not determined, and their $[\text{Fe}/\text{H}]$ were not accurate or were old. The known spectroscopic binaries were excluded.

The list was completed with a few clump stars having distances in the range of 100–200 pc, for which we were expecting the low metallicity. Most of the clump stars with $B - V$ between 0.75 and 0.8, even lacking previously published metallicities, were also included in the list and observed. Finally, we have included six *Tycho-2* stars that are distant clump stars located at about 500 pc from the Galactic plane toward the North Galactic Pole. They were previously identified from the low S/N spectra (Bienaymé et al. 2006).

The spectra of the studied stars were obtained using the facilities of the 1.93 m telescope of the Haute-Provence Observatoire (France) equipped with the échelle spectrograph ELODIE. The resolving power was 42 000, the region of the wavelengths was 4400–6800 Å, and the signal-to-noise ratio was about 130–230 (at 5500 Å). The initial processing of the spectra (image extraction, cosmic particles removal, flatfielding, etc.) was carried out following Katz et al. (1998). Further processing of the spectra (continuum level location, measurement of the equivalent widths, etc.) was performed using the software package DECH20 (Galazutdinov 1992). The equivalent widths were measured using the Gaussian fitting.

In Fig. 1 we have shown the comparison of our EW s measured in the spectrum of HD 180711 with those reported by Boyarchuk et al. (1996). In Boyarchuk et al. (1996), the spectra of program stars were obtained using the 2.6 m telescope at the Crimean Astrophysical Observatory (Ukraine) with a coude échelle spectrograph. The reciprocal dispersion of those spectra was 3 Å/mm. An agreement between two independent EW systems appears to be good, as one can see in Fig. 1. $\Delta EW = EW(\text{our}) - EW(\text{Boyarchuk}) = 0.23 \pm 4.5 \text{ m}\text{\AA}$.

The basic characteristics of studied stars are given in Table A.1. The spectral classes Sp , magnitudes $V(\text{Simbad})$ were taken from the SIMBAD database, the magnitude $V(\text{Hipparcos})$ was transformed from the Hipparcos H_p to the Johnson system, the parallaxes π were taken from the Hipparcos catalogue (ESA 1997), and the M_V were calculated. The 6 fainter stars of the sample are not part of the Hipparcos catalogue. Their spectral type come from SIMBAD. The V magnitudes were extracted from the General Catalogue of Photometric Data by Mermilliod et al. (1997). Absolute magnitudes M_V have been estimated from the TGMET software.

¹ The Hipparcos star positions are expressed in the ICRS (see <http://www.iers.org/iers/earth/icrs/icrs.html>).

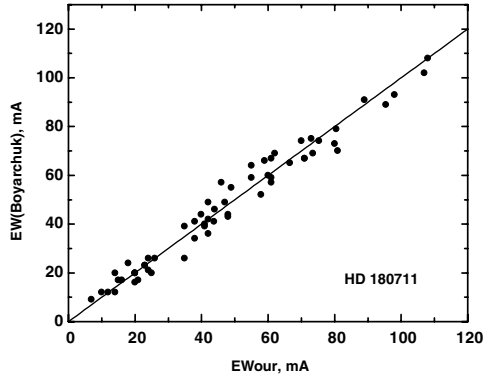


Fig. 1. Comparison between EWs measured from the spectrum of HD 180711 and those from Boyarchuk et al. (1996). *Upper panel:* EWs from Boyarchuk et al. (1996); *lower panel:* our determinations.

3. Atmospheric parameters

3.1. Effective temperature T_{eff}

The temperatures were determined with the very high level of accuracy ($\sigma = 10\text{--}15$ K) using the line depth ratios. The spectral lines of high and low excitation potentials respond differently to the change in effective temperature (T_{eff}). Therefore, the ratio of their depths (or equivalent widths) is a very sensitive temperature indicator. This technique allows one to determine T_{eff} with an exceptional level of accuracy. The method used is based on the ratio of the central depths of two lines having very different functional T_{eff} dependences. This method is independent of the interstellar reddening and only marginally dependent on individual characteristics of the stars, such as rotation, microturbulence, and metallicity. NLTE effects will most likely affect ratios of high- and low-excitation line strengths, and ratios between different chemical elements. Perhaps these effects and also those of varying individual stellar chemical abundance can be reduced by the statistics of a large number of different line ratios. The zeropoint is well established and is based on a large number of independent measurements from the literature; it would be unlikely that the error on the zeropoint is larger than 20–50 K.

We used a set of 100 line ratio – T_{eff} relations obtained in Kovtyukh et al. (2006), with the mean random error of a single calibration being 65–95 K (45–50 K in most cases and 90–95 K in the least accurate cases). The use of $\sim 70\text{--}100$ calibrations per spectrum reduces the uncertainty to 5–25 (1σ) K. This precision indicates that these 100 calibrations are weakly sensitive to non-LTE effects, metallicity, surface gravity, micro- and macroturbulence, rotation, and other individual stellar parameters. These relations have been calibrated with the reference stars in common with Gray & Brown (2001) – 21 stars, $\sigma = 27$ K; Blackwell & Lynas-Gray (1998) – 18 stars, $\sigma = 81$ K; Alonso et al. (1999) – 14 stars, $\sigma = 47$ K; Strassmeier & Schordan (2000) – 20 stars, $\sigma = 71$ K; and Soubiran et al. (1998) – 103 stars, $\sigma = 106$ K (see Fig. 2). For the majority of stars, we obtained an internal error smaller than 20 K.

After the accurate effective temperatures have been determined, the other atmospheric parameters were found iteratively.

3.2. Surface gravity $\log g$, microturbulent velocity V_t , and metallicity $[\text{Fe}/\text{H}]$

Because clump giants have similar luminosity, but different initial masses, their gravities cannot be correctly determined using their parallaxes and mass data. These $\log g$ values are affected

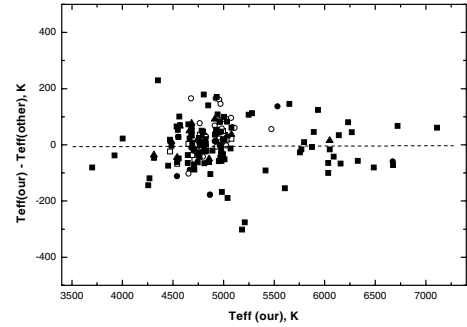


Fig. 2. Comparison between the temperatures derived in the present work and those derived by Gray & Brown (2001) – *open squares*, Blackwell & Lynas-Gray (1998) – *solid circles*, Alonso et al. (1999) – *solid triangles*, Strassmeier & Schordan (2000) – *open circles*, and Soubiran et al. (1998) – *solid squares*.

by ± 0.3 dex, if the stellar mass is based on an assumption of mass of $2.2 M_{\odot}$. We used the two following methods of spectroscopic determination of the gravity $\log g$: 1) using the iron ionization equilibrium assumption, where the average iron abundance determined from Fe I lines and Fe II lines must be identical, and 2) from the wing fitting of the Ca I 6162 Å line. For the method of the ionization equilibrium, we used iron lines with $EW < 120$ mÅ; the wing profiles for such lines practically do not depend on damping constants, but they are sensitive to microturbulent velocity V_t and metallicity $[\text{Fe}/\text{H}]$. Therefore, we take adopted T_{eff} , and then we obtain the parameters ($\log g$, V_t , and $[\text{Fe}/\text{H}]$) iteratively. Two or three steps were enough to get a good convergence.

The second method is motivated by the fact that the Ca I line is strong in giants, and therefore its wing profile is sensitive to the gas pressure in a stellar atmosphere, and to the surface gravity. The use of the Ca I triplet lines (6102, 6122, 6162 Å) as indicators of surface gravity was proposed for dwarfs and subgiants by Edvardsson (1988), and analysed by Cayrel et al. (1996). Cayrel et al. (1996) explored the possible influence of errors on the profiles of these lines. They found that a change of 10% of the damping constants has a negligible influence, a change of 15% becomes more or less detectable. The effective temperatures were also varied by 100 K, and no alteration of profiles was detected. NLTE effects are important in the core of the line, but negligible in the wings. Recently, Affer et al. (2005) used this method for K dwarfs and subgiants. We applied this method to giants to assess our gravity determination in case of iron ionization equilibrium. The Ca I 6162 Å line, which was recommended by Katz et al. (2003) as a best luminosity indicator among the triplet lines, was used. We estimated the influence of atmospheric parameter uncertainties on accuracy of the gravity determination. The Ca I wings are not sensitive to the microturbulent velocity V_t . A change by 30 K in T_{eff} brings the errors in $\log g$ about 0.05–0.10 (for $T_{\text{eff}} = 5000$ K and 4500 K, respectively). To determine the calcium abundance value, we used weaker Ca I lines, which are presumed to be less affected by the damping and the microturbulent velocity.

The departures from LTE in the computation of the wing profiles for these lines are negligible for dwarfs and subgiants (Cayrel et al. 1996), but in the case of giants it may elevate the level of uncertainty in $\log g$ up to 0.2 dex. The total error of the $\log g$ determination for giants is about 0.2–0.3 dex. An example of the line wing fitting for HD 180711 is given in Fig. 3. The values of $\log g$, obtained by two methods, are given in Table A.2. The mean difference $\Delta(\log g(\text{Ca}) - \log g(\text{IE}))$ is

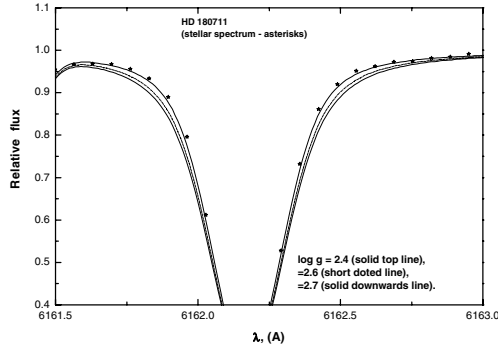


Fig. 3. Derivation of $\log g$ from the Ca I 6162 Å profile for HD 180711.

$-0.01 \sigma = \pm 0.09$ The results of these two methods applications are in good agreement. The value of surface gravity for each star was obtained by keeping the condition of the ionization equilibrium between the Fe I and Fe II species, and these values were used for abundance determination.

The value of the microturbulent velocity V_t is determined by the standard method from a condition of independence of the iron abundances determined from the given line of Fe I upon its equivalent width EW . The accuracy of V_t determination is $\Delta V_t = \pm 0.2 \text{ km s}^{-1}$.

The $[\text{Fe}/\text{H}]$ metallicity is obtained from the abundance determined from Fe I lines. (In this paper we use the customary spectroscopic notation $[X/Y] = \log_{10}(N_X/N_Y)_{\text{star}} - \log_{10}(N_X/N_Y)_{\odot}$).

In Table A.2 we give the mean T_{eff} , the number of calibrations used (N), the errors of the mean σ , $\log g$, V_t , and metallicities $[\text{Fe}/\text{H}]_I$ and $[\text{Fe}/\text{H}]_{II}$, which were determined from Fe I and Fe II lines, respectively.

4. Determination of chemical abundances

We employ the grid of stellar atmospheres from Kurucz (1993) to compute abundances of Li, C, N, O, Na, Mg, Si, Ca, and Ni. The choice of the model was made using the standard interpolation on T_{eff} and $\log g$. The abundance analysis of Si, Ca, Ni, and Fe has been done in the LTE approximation (Kurucz's WIDTH9 code) using the measured equivalent widths of these elements' lines and the solar oscillator strengths (Kovtyukh & Andrievsky 1999). Abundances of Fe, Si, Ca, Ni were derived from a differential analysis relatively to the Sun's data (see discussion in Mishenina et al. 2004). In Table A.3, the relative-to-solar Fe, Si, Ca, and Ni abundances and individual errors are given.

4.1. The Li abundance

The Li abundances in program stars were obtained by fitting synthetic spectra to the observational profiles. We used the STARSPEX LTE spectral synthesis code developed by Tsymbal (1996). Considering a wide range of temperatures and metallicities of our sample stars, a special effort was put into a compilation of a full list of atomic and molecular lines close to the ^7Li 6707 Å line (Mishenina & Tsymbal 1997). In Fig. 4, the comparison was made of the observed and the calculated spectra of HD 90633 for different lithium abundances $\log A(\text{Li}) = 1.0, 1.85, \text{ and } 2.1$, where $\log A(X) = 12 + \log(N_X/N_H)$. The derived values of $\log A(\text{Li}) > 0.5$ dex are given for 24 stars in Table A.4. We consider the lithium abundance of about 0.5 dex as a lower limit of the reliable determination. The comparison of our results

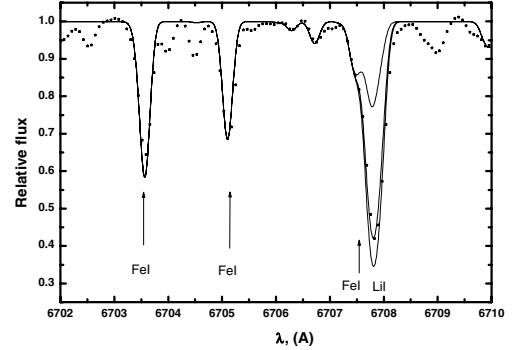


Fig. 4. Comparison between the observed spectrum for HD 90633 and the synthetic one with the Li abundances $\log A(\text{Li}) = 1.0, 1.85, \text{ and } 2.1$.

with values found in the literature (Brown et al. 1989) shows a good agreement $\Delta \log A(\text{Li})_{\text{Brown}} - \log A(\text{Li})_{\text{our}} = -0.01 \pm 0.13$ (for 8 common stars).

4.2. CNO abundances

The abundances of carbon, nitrogen, and oxygen are determined by the method of synthetic spectrum using the STARSPEX code (Tsymbal 1996). The spectrum of a molecule C_2 at 5630 Å (head of a band $\text{C}_2(0,1)$ of the Swan system $d^3\Pi_g - a^3\Pi_u$) was used to derive the carbon abundance. The nitrogen abundance was determined from the spectrum of a molecule CN at 6330 Å and 6470 Å (red system $\text{A}^2\Pi - \text{X}^2\Sigma$, heads of bands CN (5,1) and (6,2)). The wavelengths and parameters of molecular lines (including $\log gf$) were taken from Kurucz (1993), and they were corrected using the technique proposed by Kuznetsova & Shavrina (1996). For these spectral regions, the contribution from blended lines of other systems of a molecule C_2 , molecule CN, and (NH, OH, CH, MgH, and SiH) was estimated. The lines from our list were compared to the lines from the solar spectrum, using the atmosphere model from the Kurucz grid. The contribution from the blended lines in the solar spectrum appears to be insignificant (except for lines of a molecule CN in the regions of a molecule C_2 and line [OI] 6300.3 Å). For the regions of a molecule C_2 and [OI] 6300.3 Å line, the CN lines were included in the final line list. The calculation was carried out with the dissociation potentials $D_0(\text{C}_2) = 6.15 \text{ eV}$ and $D_0(\text{CN}) = 7.76 \text{ eV}$. In Fig. 5, a comparison between the synthetic spectrum and the observed one near 5630 Å is shown. In Table A.4, the abundances of carbon, nitrogen, and oxygen are given with the scale $\log A(\text{H}) = 12$. Below, we use the values relative to solar and iron abundances ($[\text{C}, \text{N}, \text{O}/\text{Fe}]$). The solar C, N, O abundances are determined by fitting the synthetic and solar spectra. As solar spectra, we used those of the Moon and asteroids that were obtained with the spectrograph ELODIE. The adopted solar values of $\log \log A(\text{C})$, $\log A(\text{N})$, and $\log A(\text{O})$ are the following: 8.55, 7.97, and 8.70, respectively.

4.3. NLTE abundances of magnesium and sodium

In the spectra of cool giants, the lines of sodium and magnesium are strong enough ($EW > 200 \text{ mÅ}$), therefore one can expect a significant deviation from LTE. For determination of the abundances of Na and Mg, we used an NLTE approximation. Four lines of Na I and 9 lines of Mg I were considered.

NLTE abundances of Mg and Na were determined with the help of a modified version of the MULTI code (Carlsson 1986)

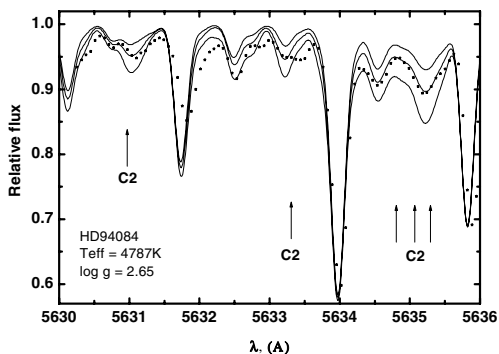


Fig. 5. Comparison between the observed spectrum for HD 94084 and the synthetic one with the C abundances $\log A(\text{C}) = 8.36, 8.46,$ and 8.56 .

described in Korotin et al. (1999a,b). In such a modified version, in particular, additional opacity sources from the ATLAS9 code (Kurucz 1993) were included. This was done to calculate the continuum opacity more precisely, and to take into account the absorption by a great number of spectral lines (especially within the region of the near UV). This allowed us to more accurately calculate the intensity distribution in the region 900–1500 Å. In turn, this significantly affects the determination of the radiative rates of bound – free transitions. A simultaneous solution of the radiative transfer and statistical equilibrium equations has been performed in the approximation of a complete frequency redistribution for all the lines. All the NLTE calculations were also based on the Kurucz grid of atmospheric models.

4.4. Parameters of sodium and magnesium atoms

The model of the sodium atom described by Sakhbullin (1987) has been modified (see Korotin & Mishenina 1999). It consists of 27 levels of Na I and the ground level of Na II. We considered the radiative transitions between the first 20 levels of Na I and the ground level of Na II. Transitions between the remaining levels were only used in the equations of particle number conservation. Finally, 46 *bound – bound* and 20 *bound – free* transitions were included in the linearization procedure. For 34 transitions the radiative rates were fixed.

We employed the model of the magnesium atom consisting of 97 levels: 84 levels of Mg I, 12 levels of Mg II, and a ground state of Mg III. Within the described system of the magnesium atom levels, we considered the radiative transitions between the first 59 levels of Mg I and the ground level of Mg II. Transitions between the rest levels were not taken into account, and they were used only in the equations of particle number conservation. For detail see Mishenina et al. (2004).

The difference between synthetic and observed spectra becomes visible if the sodium and magnesium abundances are changed by about 0.05 dex. The difference between sodium and magnesium abundances derived under the LTE assumption and for the NLTE case is within an interval of 0.10–0.15 dex. As an example, for better comparison we have shown the LTE line profile in Figs. 6–9.

In Table A.4, NLTE abundances of Na and Mg are given in the scale where $\log A(\text{H}) = 12$.

4.5. Abundance determination errors

The metallicities $[\text{Fe}/\text{H}]$ for the giants have been determined. These determinations were based on the iron abundance

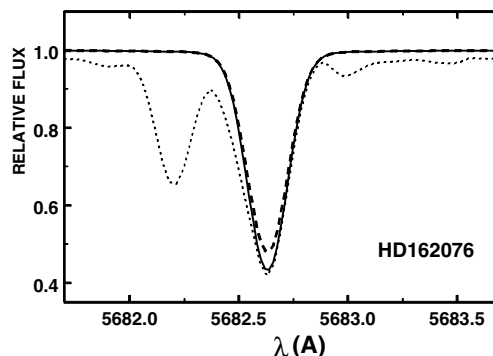


Fig. 6. NLTE profile fitting for HD 162076 (for the line Na I 5683 Å) and the LTE profile (dashed line).

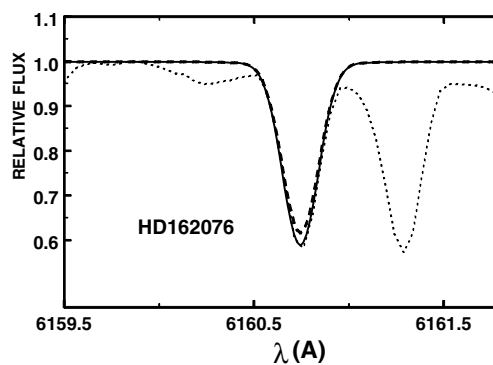


Fig. 7. Same as Fig. 6, but for the lines Na I 6164, 6160.

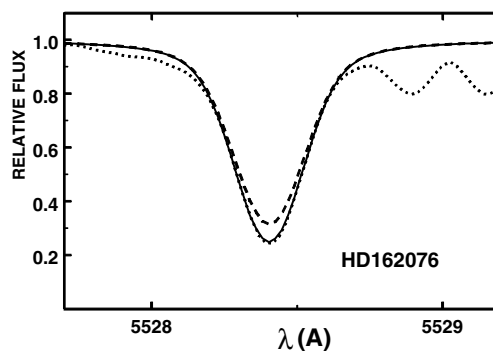


Fig. 8. Same as Fig. 6, but for the lines Mg I 4703.

values derived from Fe I lines. For this purpose, we used from 100 to 170 lines, depending on the temperature of the star: for the cooler stars, the number of iron lines was lower. The typical line-to-line scatter for Fe I is 0.11 dex s.d. The abundances of silicon, calcium, and nickel have been determined from 12 to 22 lines of Si I, 8 to 10 lines of Ca I, and 15 to 20 lines of Ni I. Typical standard deviations of the abundances derived from a single line of these elements are 0.12, 0.14, and 0.10, respectively.

Several factors may influence the abundance determination. Among them are: 1) the accuracy of the model parameters; 2) the equivalent width measurements; 3) the quality of the synthetic spectrum adjustment; and 4) internal errors of the method used. Concerning the last factor, one can notice that somewhat different abundance results can be obtained if one uses the LTE or NLTE approximations, 1D-, 2D-, or 3D atmosphere models. There are also uncertainties in atomic constants. The use of the differential method minimizes these determination errors. Uncertainties that are attributed to observed spectrum are the

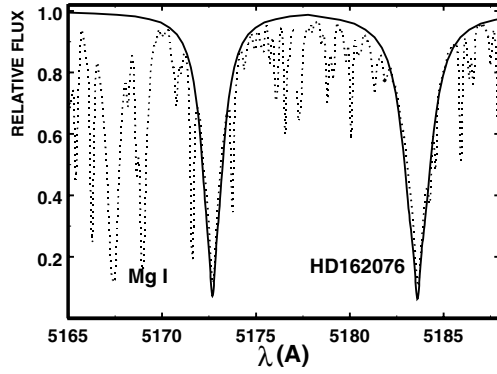


Fig. 9. Same as Fig. 6, but for the lines Mg I 5173, 5184.

following. A change in equivalent width of 2 mÅ corresponds to a change in abundance of about 0.03–0.06 dex for Fe I, Si I, Ca I, and Ni I. The fitting procedure between synthetic and observed spectra in the case of C_2 lines produces uncertainty of about 0.02 dex. In other cases (N, Na, Mg), it is about 0.05 dex. The value of total uncertainty due to the choice of the stellar parameters is shown in Table 1. The atmospheric parameters were changed by $\Delta T_{\text{eff}} = +100$ K, $\Delta \log g = +0.2$, $\Delta [\text{Fe}/\text{H}] = -0.25$ for $[\text{Fe}/\text{H}] < 0$ and $\Delta [\text{Fe}/\text{H}] = +0.1$ for $[\text{Fe}/\text{H}] > 0$, $\Delta V_t = +0.2$ km s $^{-1}$.

As one can see from Table 1, the total uncertainty reaches 0.18–0.24 dex for iron abundance determined from Fe II species, and 0.10–0.12 dex in the case of the Fe I species. For C, N, and O abundances, such uncertainties are: 0.13–0.23, 0.09–0.13, and 0.09–0.21, respectively. In the case of carbon, the maximal error takes place for the cooler stars. For oxygen, the uncertainty is caused by the choice of the model metallicity for metal-deficient stars. For Na, Mg, and Si abundances, the total error is about 0.08–0.11, and for Ca and Ni it is about 0.11–0.16. The microturbulence uncertainty supplies the largest uncertainty to the Fe I iron abundance.

5. Abundance trends with metallicity

5.1. The C, N, O abundance

The abundance ratios $[\text{C}/\text{Fe}]$, $[\text{N}/\text{Fe}]$, and $[\text{O}/\text{Fe}]$ for each star in our set are plotted against $[\text{Fe}/\text{H}]$ in Figs. 10, 11, and 13. For the whole sample of giants, the average values of the abundances of these elements are the following: $\langle [\text{C}/\text{Fe}] \rangle = -0.23 \pm 0.08$, $\langle [\text{O}/\text{Fe}] \rangle = 0.08 \pm 0.16$, and $\langle [\text{N}/\text{Fe}] \rangle = 0.25 \pm 0.09$. For the stars of metallicity $[\text{Fe}/\text{H}] > -0.3$ they are: $\langle [\text{C}/\text{Fe}] \rangle = -0.24 \pm 0.07$, $\langle [\text{O}/\text{Fe}] \rangle = 0.04 \pm 0.12$, and $\langle [\text{N}/\text{Fe}] \rangle = 0.24 \pm 0.08$. For stars near the solar metallicity $-0.01 < [\text{Fe}/\text{H}] < 0.01$, they are: $\langle [\text{C}/\text{Fe}] \rangle = -0.28 \pm 0.05$, $\langle [\text{O}/\text{Fe}] \rangle = 0.02 \pm 0.08$, and $\langle [\text{N}/\text{Fe}] \rangle = 0.21 \pm 0.07$. These averaged abundance ratios agree well with evolutionary model predictions of Iben (1991), who showed that stars should have decreased carbon and increased nitrogen abundances in their atmospheres.

Thus, our data (see Fig. 10) exhibit a clear anticorrelation between $[\text{C}/\text{Fe}]$ and $[\text{Fe}/\text{H}]$. In Fig. 12, we compare our determinations (open circles) of the carbon abundance with those of Lambert & Ries (1981), Kjrgaard & Gustafsson (1982) for all stars. The mean values obtained in these works are: $\langle [\text{C}/\text{Fe}] \rangle = -0.22 \pm 0.21$ (Lambert & Ries 1981) and $\langle [\text{C}/\text{Fe}] \rangle = -0.31 \pm 0.30$ (Kjrgaard & Gustafsson 1982). They are within the error limits of determinations with our value $\langle [\text{C}/\text{Fe}] \rangle = -0.23 \pm 0.08$, but in our case we have smaller scatter. The dependence of $[\text{C}/\text{Fe}]$

Table 1. Abundance determination errors. Parameter variation and corresponding uncertainty in abundance determination ($\Delta T_{\text{eff}} = +100$ K, $\Delta \log g = +0.2$, $\Delta [\text{Fe}/\text{H}] = -0.25$ for $[\text{Fe}/\text{H}] < 0$ and $\Delta [\text{Fe}/\text{H}] = +0.1$ for $[\text{Fe}/\text{H}] > 0$, $\Delta V_t = +0.2$ km s $^{-1}$).

Element	ΔT_{eff}	$\Delta \log g$	$\Delta [\text{Fe}/\text{H}]$	ΔV_t	Δtot
HD 161178					
4789/2.2/-0.24/1.3					
FeI	0.06	0.01	-0.03	-0.10	0.12
FeII	-0.10	0.08	-0.19	-0.08	0.24
LiI	0.14	0.00	0.02	-0.01	0.14
CI	-0.12	0.05	0.03	0.00	0.13
NI	0.08	0.03	-0.10	0.00	0.13
OI	0.02	0.09	-0.19	0.00	0.21
NaI	0.08	-0.03	0.02	-0.07	0.11
MgI	0.07	-0.02	-0.04	-0.06	0.10
SiI	-0.02	0.03	-0.08	-0.04	0.10
CaI	0.11	0.00	0.05	-0.07	0.14
NiI	0.05	0.03	-0.06	-0.08	0.12
HD 17361					
4646/2.5/0.12/1.5					
FeI	0.03	0.03	0.03	-0.09	0.10
FeII	-0.11	0.17	0.08	-0.05	0.22
LiI	0.16	-0.01	0.00	0.00	0.16
CI	-0.13	0.15	-0.01	0.00	0.20
NI	0.03	0.07	0.05	0.00	0.09
OI	-0.04	0.13	0.05	0.00	0.14
NaI	0.09	0.00	-0.01	-0.07	0.11
MgI	0.05	0.01	0.03	-0.06	0.08
SiI	-0.05	0.09	0.04	-0.05	0.12
CaI	0.09	-0.01	-0.01	-0.07	0.11
NiI	0.03	0.07	0.05	-0.11	0.14
HD 27697					
4975/2.65/0.11/1.4					
FeI	0.07	0.00	0.00	-0.08	0.11
FeII	-0.07	0.09	-0.13	-0.06	0.18
LiI	0.12	-0.01	0.07	-0.01	0.14
CI	-0.10	0.07	-0.19	0.00	0.23
NI	0.08	0.03	0.06	0.00	0.10
OI	0.02	0.09	0.00	0.00	0.09
NaI	0.07	-0.02	0.00	-0.05	0.09
MgI	0.06	-0.02	0.03	-0.03	0.08
SiI	-0.01	0.03	-0.05	-0.04	0.07
CaI	0.09	-0.01	0.02	-0.06	0.11
NiI	0.05	0.02	-0.05	-0.14	0.16

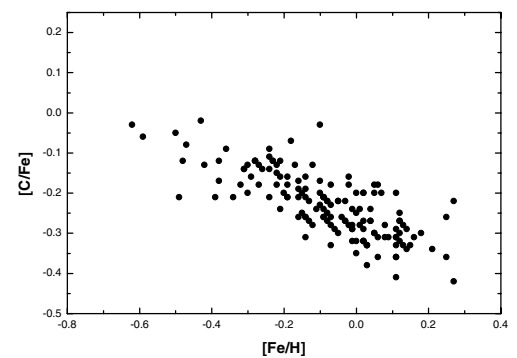


Fig. 10. Carbon abundance $[\text{C}/\text{Fe}]$ vs. $[\text{Fe}/\text{H}]$.

on $[\text{Fe}/\text{H}]$ (see Fig. 12) is clearly observed only for our clump giant sample; however, it is probably not a feature of the clump stars. The same behaviour of $[\text{C}/\text{Fe}]$ versus $[\text{Fe}/\text{H}]$ was discovered in the disc dwarfs (Bensby & Feltz 2006; Reddy et al. 2006), but the average values of $[\text{C}/\text{Fe}]$ are different for dwarfs and for giants. Therefore, we can conclude that the observed trend is not a peculiar feature of the clump giants, since it reflects the general tendency of the C abundance decreasing with $[\text{Fe}/\text{H}]$

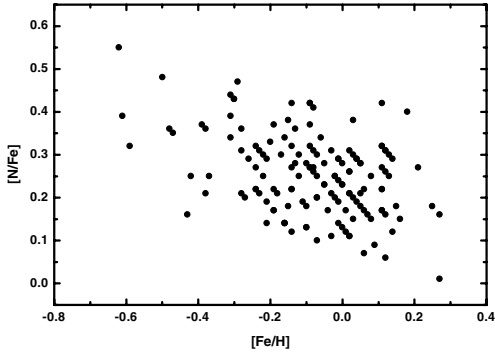


Fig. 11. Nitrogen abundance $[N/Fe]$ vs. $[Fe/H]$.

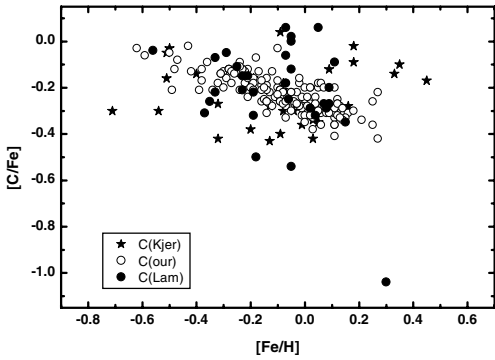


Fig. 12. Comparison of our carbon abundances (open circles) and those of Kjrgaard et al. (1982) (asterisks) and Lambert & Ries (1981) (filled circles).

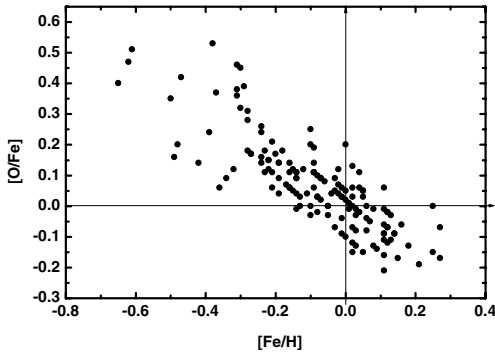


Fig. 13. Oxygen abundance $[O/Fe]$ vs. $[Fe/H]$.

increasing. We have detected it because our $[C/Fe]$ are obtained with a smaller scatter, compared to other similar works.

We have found some (not distinctive) dependence between $[N/Fe]$ and $[Fe/H]$ (see Fig. 11), but quite large scatter for our nitrogen data prevents us from making a definitive conclusion. An analysis of the C and N abundances within the frameworks of the evolutionary models is presented in Sect. 6.

The abundance of oxygen increases with the metallicity decrease (see Fig. 13). This behaviour is similar to α -element behaviour in the dwarf stars. This confirms the well-known fact that the relative-to-iron abundance of the α -elements decreases when the metallicity increases. This is connected with the growing contribution from the SNe I stars to the iron enrichment. Obviously, the determined oxygen abundance in giants can be used in an investigation of Galactic evolution.

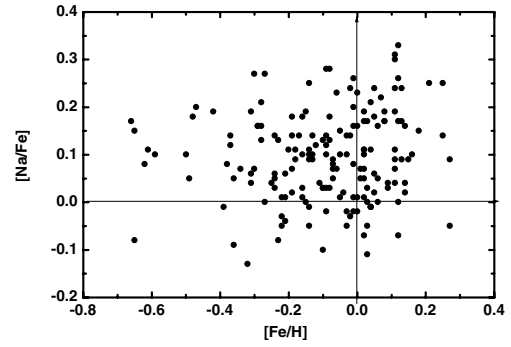


Fig. 14. Sodium abundance $[Na/Fe]$ vs. $[Fe/H]$.

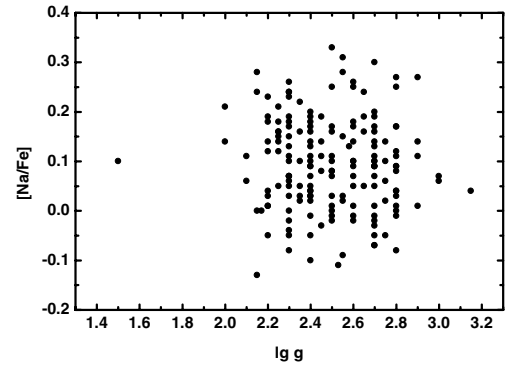


Fig. 15. Sodium abundance $[Na/Fe]$ vs. $\log g$.

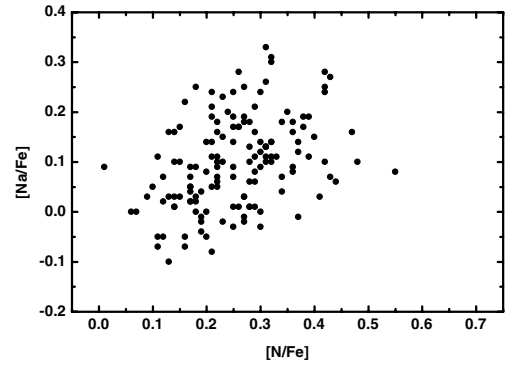


Fig. 16. Sodium abundance $[Na/Fe]$ vs. nitrogen abundance $[N/Fe]$.

5.2. The Na abundance

We found a small Na overabundance of about 0.1 dex (see Fig. 14), and we also established that there is no visible dependence of $[Na/Fe]$ upon $\log g$ (see Fig. 15). Some overabundance of sodium can be the sign of the NeNa cycle operation. Nevertheless, an absence of any dependence between $[Na/Fe]$ and $\log g$ does not support this supposition. On the other hand, this can be the result of a restricted region of the $\log g$ we considered. Additionally, there is a correlation between $[Na/Fe]$ and $[N/Fe]$ (see Fig. 16). We notice that the behaviour of $[Na/Fe]$ vs. $[Fe/H]$ is not similar for giants and for dwarfs (Soubiran & Girard 2005). We will consider the sodium abundance below in Sect. 6.

5.3. The α -element and Ni abundances

The behaviour of Mg, Ca, and Si (α -elements) and Ni (iron-peak element) (see Figs. 17–20) abundances vs. $[Fe/H]$ in giants is

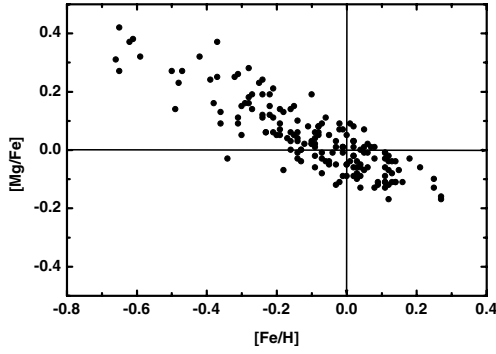


Fig. 17. Magnesium abundance $[Mg/Fe]$ vs. $[Fe/H]$.

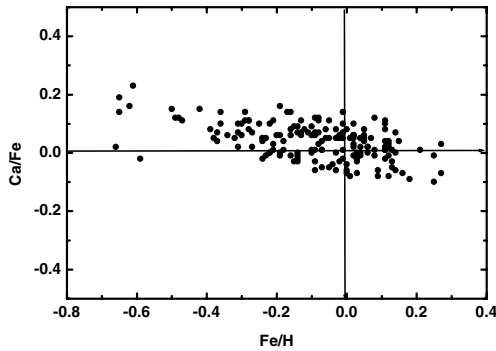


Fig. 18. Calcium abundance $[Ca/Fe]$ vs. $[Fe/H]$.

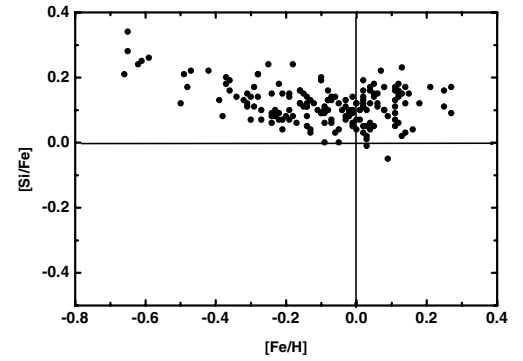


Fig. 19. Silicon abundance $[Si/Fe]$ vs. $[Fe/H]$.

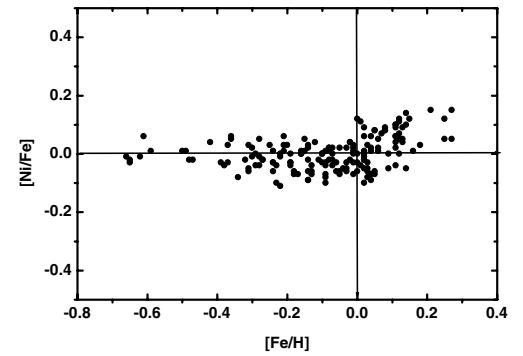


Fig. 20. Nickel abundance $[Ni/Fe]$ vs. $[Fe/H]$.

the same as in dwarfs (Soubiran & Girard 2005). This allows us to use abundances of these elements to study the chemical and dynamical evolution of the Galaxy.

6. Abundance variations due to the first dredge-up

Due to surface abundance modifications during the first dredge-up episode, clump giants do not exhibit the chemical pattern that they inherited at their birth. In this section, we compare our data with first dredge-up theoretical predictions.

6.1. Stellar models

We computed evolution models from the pre-main sequence up to the AGB phase for stars with initial masses of 1.0, 1.5, 2.0, 2.5, and 3.0 M_{\odot} , and for three values of $[Fe/H]$, i.e., -0.293 , 0, and $+0.252$, using the code STAREVOL (i.e., Siess et al. 2000; Palacios et al. 2006). These are “classical” models, i.e., they do not take into account atomic diffusion and rotation-induced mixing. The nuclear reaction rates are those of NACRE (Angulo et al. 1999). For the radiative opacities, we use the OPAL tables above 8000 K (Iglesias & Rogers 1996), and at lower temperatures, the atomic and molecular opacities of Alexander & Ferguson (1994). The conductive opacities are computed from a modified version of the Iben (1975) fit to the Hubbard & Lampe (1969) tables for non-relativistic electrons, and from Itoh et al. (1983) and Mitake et al. (1984) for relativistic electrons. The equation of state is described in detail in Siess et al. (2000) and accounts for the non-ideal effects due to Coulomb interactions and pressure ionization. The standard mixing length theory is used to model the convection $\alpha_{MLT} = 1.75$ calibrated for the solar case. Neither overshooting, nor undershooting is considered for convection. The atmosphere is treated in the gray approximation and integrated up to an optical depth $\tau \approx 5 \times 10^{-3}$. Mass loss

is considered during the whole evolution and follows the Reimer (1975) empirical relation.

6.2. Comparison with theoretical predictions

In Figs. 21–23 we show the corresponding evolutionary tracks together with the positions of the sample stars in the HR diagram. As expected, the objects appear to be slightly more massive, on average, when one moves to higher metallicity.

Also shown in these figures are the predictions for the surface abundance variations of C, N, and Na as a function of effective temperatures along the RGB, together with the corresponding observational data. Stars with $[Fe/H]$ below -0.15 are compared with the $[Fe/H] = -0.293$ tracks, those with $[Fe/H]$ between -0.15 and $+0.12$ are compared with the $[Fe/H] = 0$ tracks, and the more metallic ones with the $[Fe/H] = +0.252$ tracks. As can be seen, the region of the clump in the evolutionary tracks overlaps the region where the first dredge-up ceases.

Our finding about the nitrogen abundance is in good agreement with the prediction of the canonical theory of evolution for first dredge-up phase.

In the case of carbon, we show two sets of tracks for the two more metallic subsamples: one assuming an initial $[C/Fe]$ equal to solar, while the other is obtained by simply shifting the previous ones by -0.15 and -0.20 dex, respectively. These quantities correspond to the values of the upper envelope of $[C/Fe]$ at the corresponding $[Fe/H]$ (see Fig. 10). Again, the models explain the observational pattern well.

Regarding sodium, the observed dispersion is higher than the theoretical one. Numerous overabundances are observed, especially for the more metal-rich subsample. This cannot be attributed to an extra-mixing process because any additional processing of the envelope of the giant would also lead to further

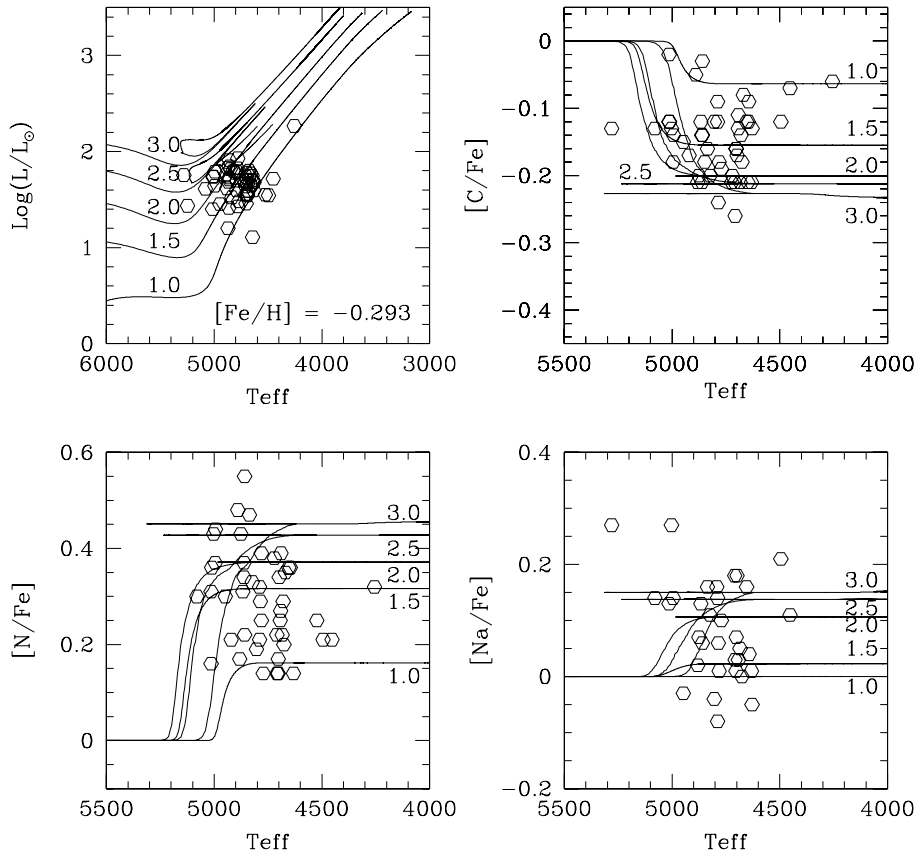


Fig. 21. Comparison of the theoretical tracks and predictions for surface abundance variations with observations. Position of our target stars with $[\text{Fe}/\text{H}] < -0.15$ in the H-R diagram.

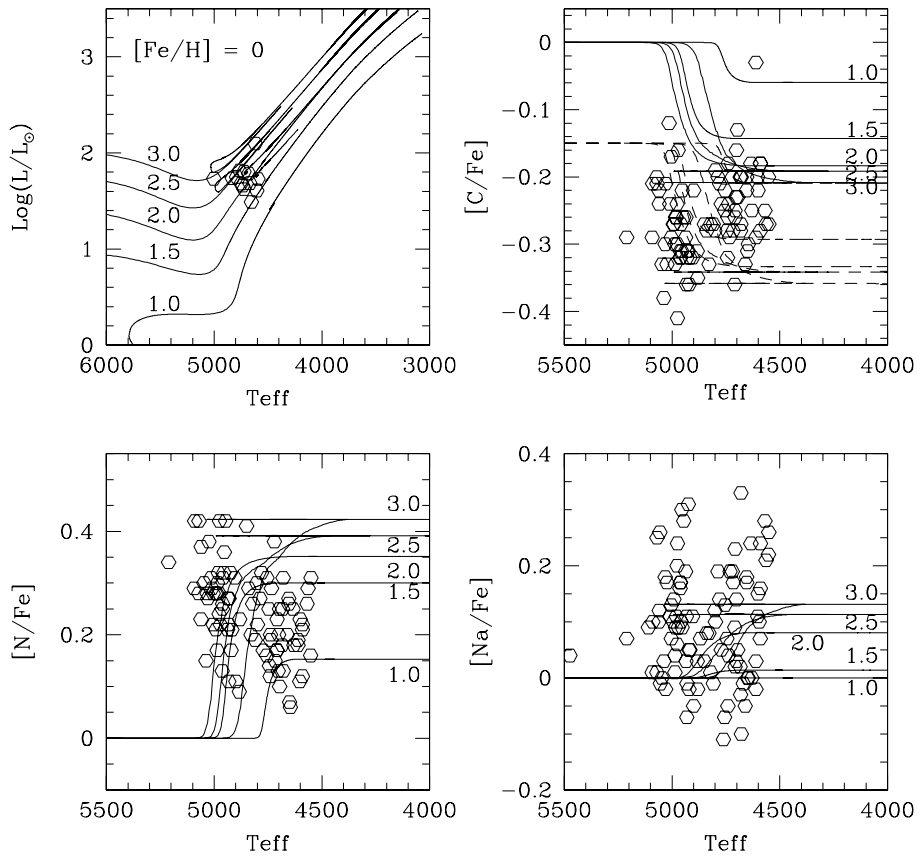


Fig. 22. Same as Fig. 21 for $[\text{Fe}/\text{H}] = 0$. *Top right panel:* the solid lines are those assuming an initial $[\text{C}/\text{Fe}]$ equal to solar, while the dotted lines are obtained by simply shifting the previous ones by -0.15 .

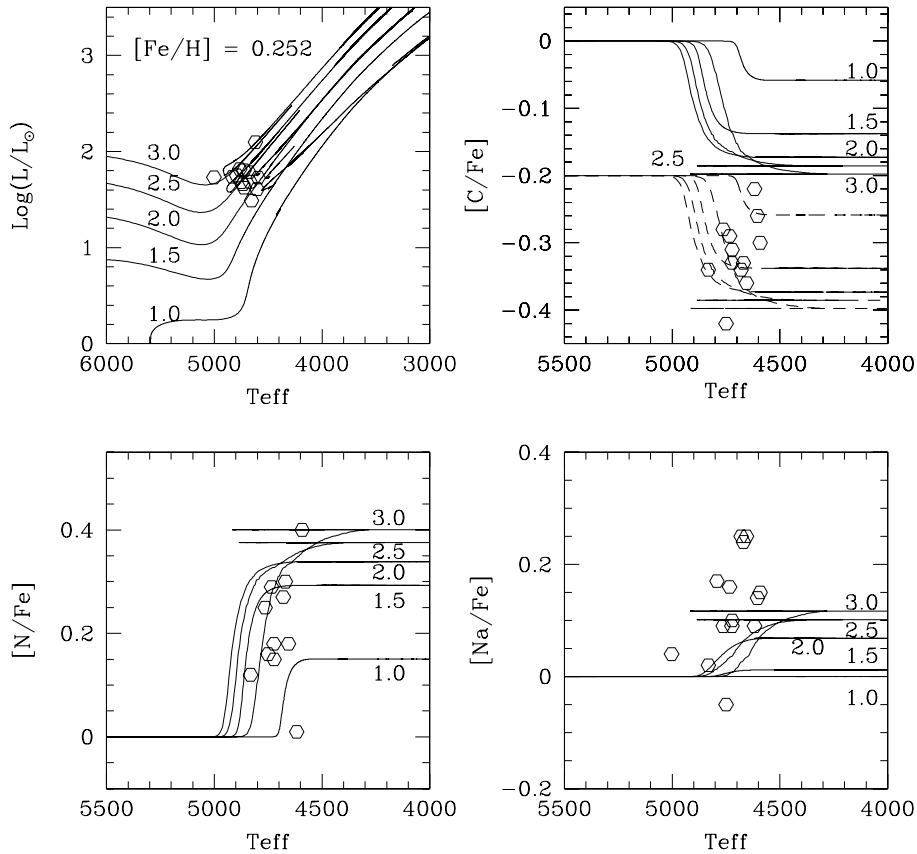


Fig. 23. Same as Fig. 22 for $[\text{Fe}/\text{H}] = 0.252$ and $[\text{C}/\text{Fe}] = -0.20$.

changes in the C and N abundances, which are not observed in our sample. One possibility to remove part of the discrepancy could lie in the rates that intervene in the NeNa cycle. For the reaction that forms sodium, $^{22}\text{Ne}(p,\gamma)^{23}\text{Na}$, the new rate calculated by Hale et al. (2001), is slightly smaller (for the central temperature of the models on the main sequence, i.e., below ~ 20 million degrees) than the NACRE prescription used in the present computations. Thus it would, not favour a stronger dredge-up of sodium. On the other hand, the present uncertainty of the $^{23}\text{Na}(p,\gamma)^{24}\text{Mg}$ and $^{23}\text{Na}(p,\alpha)^{20}\text{Ne}$ reactions is still large (Hale et al. 2004). The possibility that the initial sodium abundance was higher for some stars also exists. The disc dwarfs show some dispersion of the Na abundance, and this is confirmed, for example, by Mishenina et al. (2003) and Edvardsson et al. (1993) (especially for $[\text{Fe}/\text{H}] > 0$).

In the central regions of the main sequence stars, ^{16}O is partially converted into ^{14}N . However, the convective envelope hardly reaches the O-depleted region during the first dredge-up for the mass range considered here (see, for example, Fig. 1 of Charbonnel 1994). As a consequence, surface O variations are not expected in our sample stars (Fig. 24). In our sample of giants, we observe an O/Fe versus Fe/H variation (Fig. 13) similar to the variation observed with dwarfs (see Figs. 4 and 10 in Soubiran & Girard 2005).

6.3. The Li abundance

According to the theory, the surface Li abundance decreases with respect to its value at the end of the main sequence (MS) by a factor of 30 to 60, depending on the stellar mass and metallicity (Iben 1991). Starting from the present interstellar medium

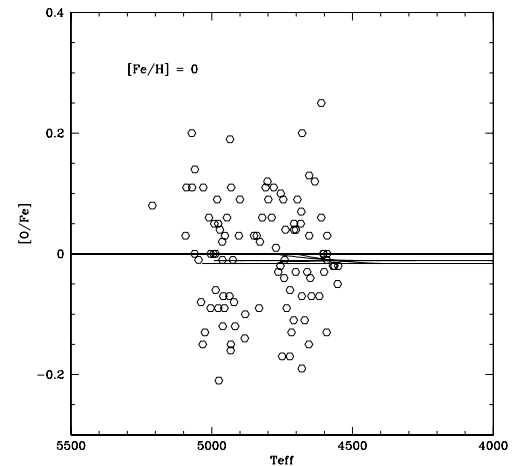


Fig. 24. Predictions for surface abundance variations of O during the first dredge-up for the $[\text{Fe}/\text{H}] = 0$ models, and comparison with the observations.

abundance of $\log N(\text{Li}) = 3.3$, we thus expect the Li values that lie around 1.5 after the first dredge-up, as shown in Fig. 25. Let us insist on the fact that these are “classical” predictions that do not take the effect of non-standard transport processes such as those induced by rotation into account, and that are thus not able to explain the Li patterns observed in low-mass main sequence stars (Charbonnel & Talon 2005; see the review by Deliyannis et al. 2000). As such, we have the right to expect lower values of lithium for our giants, as proves to be true by observation (Brown et al. 1989; Mallik 1999).

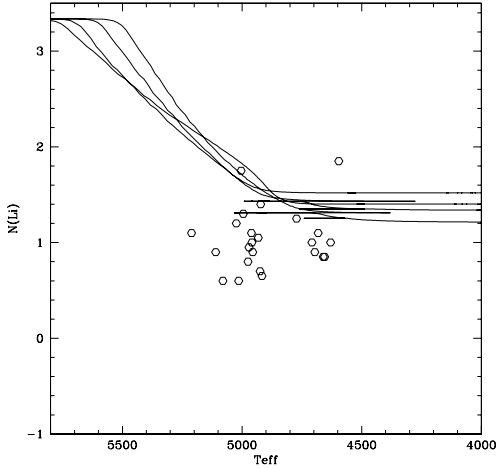


Fig. 25. Li abundance $\log A(\text{Li})$ vs. T_{eff} for our sample stars compared with the theoretical predictions for the tracks at $[\text{Fe}/\text{H}] = 0$. The line with the strongest Li depletion corresponds to the $1.0 M_{\odot}$ model.

The observed Li abundances versus the effective temperatures for the giants studied here are depicted in Fig. 25.

In the case of low-mass stars ($M_{\text{star}} < 2.2\text{--}2.5 M_{\odot}$, HD 8733, 15453, 42341, 46374, 90633, 117304, 136138, 139254, 148604, 171994, and 192836) that undergo some extramixing at the RGB bump (Charbonnel et al. 1998; Charbonnel & Balachandran 2000), the fact that we see some Li certainly indicates that these objects are RGB stars that have not yet reached the bump. Otherwise, their Li would have been destroyed.

In the case of more massive stars that do not undergo this extramixing because they do not go through the bump, Li should be consistent with the standard post dredge-up values predicted by the models.

It is most likely that the Li abundance cannot be used as the criterion to segregate the clump giants from RGB giants.

6.4. Determination of the evolutionary status

The stars considered in this study have been selected as clump stars according to photometric criteria. Nevertheless, our sample could be contaminated by ascending giant branch stars that cohabit with clump stars in the considered region of the CMD (see Figs. 21–23). It would thus be interesting to perform a more subtle separation of the clump giants from the whole sample of the stars.

We aimed to do this selection by comparing the abundances of individual stars with theoretical predictions of stellar evolution models. This is, however, a very difficult task since the clump overlaps the region where the first dredge-up ceases in the evolutionary tracks. Despite this difficulty, we checked the status of each sample star individually following the procedure described below.

We first attributed a mass and evolutionary status to each object by comparing its position in the HRD with the theoretical tracks. The values of obtained masses for our target stars are given in Table A1. Again, the stars with $[\text{Fe}/\text{H}]$ below -0.15 are compared with the $[\text{Fe}/\text{H}] = -0.293$ models, those with $[\text{Fe}/\text{H}]$ between -0.15 and $+0.12$ are compared with the $[\text{Fe}/\text{H}] = 0$ models, and the more metallic ones with the $[\text{Fe}/\text{H}] = +0.252$ ones. From this first iteration, 125 of our sample stars were identified as possible RHB or clump stars, 4 as supergiants, 38 as probable RGBs, and 2 as probable AGB stars.

Then, for each star, we checked whether its nitrogen abundance was compatible with the model predictions for the corresponding stellar mass previously attributed. The carbon abundance was used only as a cross-check because of the variation it presents as a function of metallicity (see Sect. 5.1). Among the 125 possible RGB/clump stars, 15 objects have no N determination and 32 objects appear to still be undergoing the first dredge-up dilution, as indicated by their $[\text{N}/\text{Fe}]$ and effective temperature. For the others, we made the following distinctions: (i) 38 stars are found to have completed their first dredge-up, but present slightly lower (by 0.05 to 0.2 dex) N abundances than predicted values for the corresponding stellar mass. 3 stars have N overabundances by ~ 0.2 dex. We consider, however, that these slight discrepancies are not significant because of the observational error bars on the effective temperature and on the abundance determination. Moreover, part of the small discrepancy can be accounted for by the fact that not all the stars have the exact metallicity of the theoretical tracks they are compared with. (ii) 16 giants have a N abundance in good agreement with the post-dredge predictions for the given stellar mass. Both the (i) and (ii) stars would be preferentially identified as RGB stars according to their effective temperature, although they are still good clump candidates. (iii) 21 stars could be selected as clump giants according to both their N abundance in good agreement with the post-dredge predictions and their effective temperature.

We have thus reliably selected 21 clump giants, plus 54 clump candidates and about 100 usual giants that show all the signs of first dredge-up. Unfortunately, we have to state that some uncertainty exists in the separation of the clump giants, if we rely only on the evolutionary tracks and elemental abundances that are sensitive to the stellar evolution.

An important conclusion of the present study is that the theoretical predictions of the classical models do account well for the observed surface variations of both carbon and nitrogen during the first dredge-up episode.

However, it should be noted that the considered analysis and the results depend on the accuracy of the determination of the chemical composition and the theoretical preconditions used.

7. Conclusions

We have performed the detailed analysis of the atmospheric parameters and the abundances of some elements in 177 giant stars.

The stars analysed in this study have been selected as clump stars according to the photometric criteria. We have estimated the possibility of defining the evolutionary status of these giants on the basis of evolutionary tracks and their measured element abundances that are modified during stellar evolution. We reliably selected 21 clump giants, about 54 clump candidates, and about 100 usual giants that show all the signs of a first dredge-up.

The determined C, N, and Na abundances in our program stars reflect the CNO- and NaNe cycle operation in the giant stars.

The O, Mg, Ca, and Si (α -elements) and Ni (iron-peak element) abundances in giants show trends similar to those observed in dwarfs which allows us to use these abundances to study the chemical and dynamical evolution of the Galaxy.

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Online Material

Appendix A:**Table A.1.** The basic characteristics of studied stars.

Star	Sp	V (Simbad)	V (Hipparcos)	π (0.001)	M_V	Masses
HD 2910	K0III	5.347	5.38	12.69	0.432	2.2
HD 4188	K0III	4.775	4.77	15.54	0.331	2.5
HD 4482	G8II	5.515	5.51	12.45	0.646	2.5
HD 5395	G8III-IV	4.632	4.62	15.48	0.202	2.0
HD 6319	K2III:	6.193	6.20	10.01	0.696	1.7
HD 6482	K0III	6.101	6.09	10.58	0.780	2.0
HD 7106	K0,5III	4.523	4.51	20.11	0.563	2.0
HD 7578	K1III	6.050	6.04	10.36	0.647	2.0
HD 8207	K0III	4.875	4.87	16.68	0.550	2.5
HD 8599	G8III	6.167	6.17	12.17	1.176	1.2
HD 8733	K0	6.450	6.44	10.71	1.262	2.0
HD 9408	G9III	4.692	4.68	15.96	0.303	2.0
HD 10975	K0III	5.947	5.94	10.57	0.705	1.5
HD 11559	K0III	4.621	4.61	17.11	0.471	2.7
HD 11749	K0III	5.698	5.69	10.18	0.257	2.0
HD 11949	K0IV	5.712	5.70	13.25	0.861	1.0
HD 15453	K2III	6.098	6.09	11.4	0.913	1.7
HD 15755	K0III	5.846	5.84	12.54	0.813	1.5
HD 15779	G3III	5.364	5.36	12.27	0.414	2.5
HD 16247	K0III:	5.819	5.81	11.18	0.549	1.0
HD 16400	G5III:	5.656	5.65	10.29	0.334	2.5
HD 17361	K1,5III	4.510	4.52	18.06	0.292	2.0
HD 18885	G6III:	5.834	5.84	11.55	0.693	2.2
HD 19270	K3III	5.648	5.64	10.21	0.240	2.5
HD 19787	K2III	4.350	4.35	19.44	0.405	2.7
HD 19845	G9III	5.921	5.93	10.47	0.684	2.5
HD 20791	G8.5III	5.690	5.70	11.23	0.630	2.5
HD 25602	K0III-IV	6.320	6.31	10.21	0.894	1.2
HD 25604	K0III	4.354	4.36	18.04	0.207	2.7
HD 26546	K0III	6.091	6.09	11.42	0.939	2.0
HD 26659	G8III	5.477	5.47	10.74	0.397	3.0
HD 26755	K1III	5.727	5.72	12.38	0.679	1.5
HD 27348	G8III	4.933	4.93	14.42	0.423	3.0
HD 27371	G8III	3.654	3.65	21.17	-0.043	3.0
HD 27697	G8III	3.746	3.77	21.29	0.070	3.0
HD 28292	K2III	4.971	4.96	16.78	0.458	1.0
HD 28305	K0III	3.540	3.53	21.04	-0.185	3.0
HD 28307	G7III	3.847	3.84	20.66	0.100	3.0
HD 30557	G9III	5.642	5.64	10.16	0.287	2.5
HD 31444	G6/G8III	5.726	5.71	11.3	0.719	2.2
HD 33419	K0III	6.118	6.11	10.37	0.735	2.0
HD 33618	K2III-IV	6.147	6.15	10.61	0.738	1.5
HD 34200	G5	6.384	6.39	8.5	0.749	2.5
HD 34559	G8III	4.956	4.96	15.83	0.652	2.5
HD 35369	G8III	4.144	4.13	18.71	0.167	3.0
HD 37638	G5III:	6.168	6.17	10.26	0.957	2.5
HD 39070	G8III	5.490	5.49	10.51	0.364	3.0
HD 39910	K2III:	4.874	5.87	10.71	-0.497	3.0
HD 40020	K2III	5.891	5.89	10.62	0.536	2.2
HD 40801	K0III	6.092	6.08	12.47	1.106	1.0
HD 42341	K2III	5.563	5.56	15.53	1.023	2.0
HD 43023	G8III	5.835	5.83	10.36	0.603	2.5
HD 45415	G9III	5.553	5.55	11.15	0.360	2.2
HD 46374	K2III:	5.562	5.57	12.75	0.598	2.0
HD 46758	G5	7.156	7.15	5.15	0.411	2.2
HD 47138	G8/K0III	5.704	5.71	11.79	0.838	2.5
HD 47366	K1III:	6.117	6.11	11.75	1.044	1.2
HD 48432	K0III	5.349	5.34	15.69	0.941	1.5
HD 50904	G5	6.539	6.53	7.02	0.444	2.5
HD 53329	G8IV	5.557	5.55	10.68	0.399	2.0
HD 54810	K0III	4.920	4.91	15.45	0.378	1.5
HD 55280	K2III	5.200	5.20	16.88	0.840	1.5
HD 56891	K0	6.640	6.66	7.86	0.656	2.0
HD 58207	G9III+..	3.793	3.78	25.9	0.453	2.5
HD 59686	K2III	5.450	5.45	10.81	0.123	2.0

Table A.1. continued.

Star	Sp	V (Simbad)	V(Hipparcos)	π (0.001)	M_V	Masses
HD 60294	K2III	5.941	5.93	11.63	0.717	1.5
HD 60318	K0III	5.348	5.34	10.78	0.189	3.0
HD 60341	K0III	5.637	5.64	11.15	0.364	2.0
HD 60986	K0III	5.888	5.58	10.65	0.744	2.5
HD 61363	K0III	5.596	5.58	10.01	0.184	1.5
HD 61935	K0III	3.930	3.94	22.61	0.284	2.5
HD 62141	K0III	6.244	6.25	10.63	1.058	2.5
HD 62345	G8III	3.570	3.57	22.73	0.109	3.0
HD 63798	G5	6.500	6.49	8.43	0.825	2.5
HD 64152	K0III	5.619	5.62	11.9	0.681	2.5
HD 64967	G8IV	6.567	6.57	10.1	1.218	1.2
HD 65066	K0III	6.033	6.03	11.08	0.887	2.0
HD 65345	K0III	5.297	5.30	12.33	0.430	2.5
HD 67539	G5	6.502	6.48	5.87	-0.072	2.0
HD 68312	G8III	5.351	5.36	10.32	0.151	2.5
HD 68375	G8III	5.560	5.55	11.18	0.526	3.0
HD 70523	K0III	5.720	5.71	11.4	0.501	1.2
HD 71088	G8III	5.889	5.89	10.12	0.562	2.5
HD 73017	G8IV	5.665	5.66	13.55	0.853	0.9
HD 74794	K0III:	5.698	5.70	11.71	0.574	2.0
HD 75506	K0III	5.160	5.15	11.91	0.175	2.0
HD 75958	G6III	5.571	5.57	10.59	0.450	2.7
HD 76291	K1IV	5.733	5.72	14.21	0.891	0.8
HD 76813	G9III	5.233	5.23	10.19	-0.006	3.0
HD 78235	G8III	5.427	5.42	12.56	0.646	2.7
HD 79181	G8III	5.727	5.72	10.85	0.535	1.5
HD 80546	K3III	6.180	6.16	10.3	0.714	1.5
HD 81688	K0III-IV	5.413	5.40	11.33	0.272	1.5
HD 82969	G5	6.412	6.41	10.17	1.119	1.5
HD 83240	K1IIIvar	5.012	5.00	14.45	0.333	2.0
HD 83371	K0	6.601	6.59	5.61	-0.027	2.2
HD 86513	G9III:	5.753	5.75	10.05	0.331	2.2
HD 90633	K2III	6.320	6.32	10.41	0.873	1.5
HD 93291	G4III:	5.488	5.49	11.34	0.481	3.0
HD 93875	K2III	5.564	5.57	12.39	0.491	1.7
HD 94084	K2III	6.445	6.44	10.35	1.106	2.0
HD 94402	G8III	5.454	5.45	10.45	0.246	3.0
HD 94497	G7III:	5.734	5.73	10.69	0.413	1.5
HD 95808	G7III...	5.513	5.51	10.24	0.234	3.0
HD 98366	K0III:	5.908	5.90	10.48	0.544	2.0
HD 100696	K0III	5.203	5.19	13.49	0.481	1.5
HD 101484	K1III	5.266	5.26	14.04	0.693	2.5
HD 102928	K0IV	5.635	5.62	12.59	0.639	1.0
HD 103605	K1III	5.838	5.83	10.34	0.386	2.0
HD 103912	G7IVw...	8.390	8.35	5.55	1.744	0.9
HD 104783	G5III	9.160	9.16	2.77	1.162	2.2
HD 106714	K0III	4.938	4.93	13.12	0.192	3.0
HD 108381	K2IIICN+	4.350	4.35	19.18	0.285	2.6
HD 109053	G8III	9.220	9.21	-	0.595	-
HD 110024	G9III	5.487	5.49	11.39	0.427	2.5
HD 113321	K0	9.388	9.35	-	0.798	-
HD 113997	K0	9.275	9.32	-	0.489	-
HD 114357	K3III	6.020	6.01	10.89	0.640	1.5
HD 116292	K0III	5.367	5.36	10.2	0.068	3.0
HD 117304	K0III	5.656	5.65	11.99	0.538	2.0
HD 117566	G2.5IIIB	5.750	5.74	11.15	0.842	2.5
HD 119126	G9III	5.633	5.63	10.13	0.256	2.5
HD 120084	G7III:	5.910	5.91	10.24	0.600	2.5
HD 120164	K0III+...	5.507	5.51	10.73	0.222	2.5
HD 120420	K0III	5.620	5.61	10.49	0.242	1.5
HD 136138	G5IV	5.694	5.68	11.23	0.638	2.0
HD 138852	K0III-IV	5.750	5.74	10.24	0.428	2.0
HD 139254	K0III	5.796	5.79	11.97	0.724	2.0
HD 139329	K0III	5.829	5.82	11.17	0.596	1.2
HD 143553	K0III:	6.812	5.82	13.62	1.980	1.0
HD 146388	K3III	5.714	5.72	10.31	0.332	2.0
HD 148604	G5III/IV	5.675	5.66	12.18	0.843	2.5

Table A.1. continued.

Star	Sp	V (Simbad)	V (Hipparcos)	π (0.001)	M_V	Masses
HD 152224	K0III	6.168	6.16	10.36	0.768	1.0
HD 153956	K1III:	6.050	6.04	10.99	0.726	1.5
HD 155970	K1III	6.006	5.98	10.63	0.682	1.5
HD 156874	K0III	5.688	5.68	10.24	0.377	2.5
HD 159353	K0III:	5.691	5.68	10.2	0.356	2.5
HD 161178	G9III	5.881	5.87	10.17	0.505	1.5
HD 162076	G5IV	5.700	5.69	13.04	0.952	2.5
HD 166578	K0	6.683	6.67	5.83	0.138	1.5
HD 168653	K1III:	5.969	5.96	11.48	0.758	1.0
HD 170693	K1,5III	4.833	4.82	10.28	-0.916	1.0
HD 171994	G8IV	6.316	6.31	11.14	1.251	1.5
HD 175743	K1III	5.706	5.69	11.84	0.586	2.0
HD 176408	K1III	5.676	5.67	11.4	0.405	1.7
HD 176598	G8III	5.632	5.62	10.39	0.420	3.0
HD 180711	G9III	3.082	3.07	32.54	0.252	2.0
HD 185644	K1III	5.302	5.30	13.31	0.386	1.5
HD 187739	K0III	5.905	5.88	10.46	0.503	1.0
HD 188119	G8III	3.830	3.84	22.4	0.272	2.0
HD 192787	K0III	5.711	5.70	10.86	0.578	2.5
HD 192836	K1III	6.127	6.11	10.95	0.901	2.0
HD 195330	K1/K2III	6.121	6.10	10.25	0.764	1.5
HD 196134	K0III-IV	6.513	6.50	10.28	1.131	1.7
HD 198431	K1III	5.880	5.87	13.06	0.876	1.0
HD 199870	G8III	5.562	5.55	12.32	0.680	2.5
HD 204771	K0III	5.231	5.22	14.58	0.699	2.5
HD 206005	K0	6.070	6.07	10.22	0.656	1.2
HD 207130	K0III	5.182	5.18	13.19	0.373	2.5
HD 208111	K2III	5.716	5.71	11.22	0.429	2.0
HD 211006	K2III	5.878	5.87	13.08	0.897	1.5
HD 212496	G8,5IIb	4.430	4.42	19.51	0.380	1.0
HD 214567	G8II	5.849	5.84	8.55	0.195	3.0
HD 215030	G9III	5.941	5.93	10.08	0.505	1.2
HD 215721	G8II	5.260	5.24	12.26	0.345	1.5
HD 216131	M2III	3.513	3.51	27.95	0.432	2.7
HD 216228	K0III	3.510	3.50	28.27	0.298	1.5
HD 218031	K0IIIb	4.650	4.64	18.20	0.478	1.2
HD 219418	G5III	6.830	6.81	5.58	0.363	2.7
HD 219916	K0III	4.868	4.75	15.48	0.577	2.7
HD 221345	K0III	5.220	5.22	13.09	0.315	1.5
HD 221833	K0	6.476	6.47	10.51	1.055	1.5
HD 225197	K0III	5.781	5.78	11.29	0.598	2.5
HD 225216	K1III	5.691	5.68	10.30	0.300	2.0
BD+22 2606	K5	9.450	9.42	-	1.148	-
BD+25 2555	G7III	9.162	9.14	-	0.594	-
BD+28 2250	G5	9.370	9.34	-	1.394	-

Table A.2. Parameters of atmospheres of studied stars.

Star	T_{eff} , K	$\log g(\text{E})$	$\log g(\text{Ca})$	V_t , km s $^{-1}$	[Fe/H] $_{\text{I}}$	[Fe/H] $_{\text{II}}$
HD 2910	4756	2.70	2.40	1.5	0.12	0.11
HD 4188	4809	2.70	2.60	1.5	0.04	0.06
HD 4482	4917	2.65	2.65	1.4	0.02	0.03
HD 5395	4849	2.15	2.15	1.3	-0.32	-0.31
HD 6319	4650	2.30	2.30	1.3	0.06	0.07
HD 6482	4738	2.40	2.50	1.4	-0.11	-0.10
HD 7106	4684	2.55	2.45	1.5	0.05	0.06
HD 7578	4680	2.50	2.50	1.4	0.12	0.13
HD 8207	4750	2.75	2.60	1.5	0.27	0.30
HD 8599	4781	2.50	2.40	1.1	-0.22	-0.25
HD 8733	4932	2.70	2.80	1.2	0.02	0.03
HD 9408	4804	2.30	2.40	1.5	-0.21	-0.20
HD 10975	4881	2.20	2.40	1.5	-0.19	-0.16
HD 11559	4977	3.00	2.90	1.5	0.05	0.03
HD 11749	4679	2.40	2.40	1.5	-0.10	-0.06
HD 11949	4708	2.30	2.30	1.2	-0.16	-0.15
HD 15453	4696	2.40	2.40	1.3	-0.07	-0.07
HD 15755	4611	2.30	2.30	1.2	-0.01	-0.02
HD 15779	4821	2.70	2.60	1.5	0.02	0.05
HD 16247	4629	2.20	2.30	1.4	-0.22	-0.20
HD 16400	4840	2.50	2.50	1.35	-0.01	0.01
HD 17361	4646	2.50	2.45	1.5	0.12	0.11
HD 18885	4722	2.50	2.60	1.4	0.16	0.16
HD 19270	4723	2.40	2.30	1.45	0.15	0.16
HD 19787	4832	2.75	2.65	1.5	0.14	0.15
HD 19845	4933	2.80	2.70	1.3	0.11	0.11
HD 20791	4986	2.80	2.60	1.2	0.11	0.13
HD 25602	4693	2.40	2.40	1.15	-0.42	-0.45
HD 25604	4764	2.70	2.60	1.5	0.13	0.14
HD 26546	4743	2.25	2.15	1.3	-0.01	-0.01
HD 26659	5178	2.90	3.00	1.2	-0.13	-0.11
HD 26755	4630	2.20	2.20	1.3	-0.06	-0.03
HD 27348	5003	2.80	2.70	1.2	0.14	0.14
HD 27371	4955	2.70	2.60	1.4	0.11	0.10
HD 27697	4975	2.65	2.55	1.4	0.11	0.07
HD 28292	4453	2.10	2.10	1.5	-0.18	-0.17
HD 28305	4925	2.55	2.45	1.4	0.11	0.11
HD 28307	4961	2.70	2.75	1.3	0.12	0.08
HD 30557	4829	2.45	2.45	1.35	-0.07	-0.05
HD 31444	5080	2.75	2.75	1.2	-0.17	-0.13
HD 33419	4708	2.30	2.30	1.4	0.00	0.05
HD 33618	4590	2.30	2.25	1.4	0.05	0.05
HD 34200	5055	2.80	2.80	1.3	0.04	0.06
HD 34559	5010	2.90	2.90	1.2	0.04	0.06
HD 35369	4931	2.40	2.40	1.4	-0.14	-0.14
HD 37638	5093	2.80	2.80	1.3	-0.01	0.01
HD 39070	5047	2.80	2.80	1.15	0.03	0.05
HD 39910	4618	2.60	2.60	1.35	0.27	0.26
HD 40020	4670	2.30	2.40	1.5	0.13	0.14
HD 40801	4703	2.20	2.20	1.05	-0.21	-0.23
HD 42341	4655	2.60	2.80	1.4	0.25	0.22
HD 43023	4994	2.40	2.40	1.3	-0.13	-0.12
HD 45415	4762	2.30	2.30	1.3	0.03	0.02
HD 46374	4661	2.30	2.30	1.5	0.03	0.03
HD 46758	5003	2.90	2.80	1.3	-0.30	-0.32
HD 47138	5211	3.00	3.00	1.2	-0.06	-0.04
HD 47366	4772	2.60	2.60	1.2	-0.16	-0.13
HD 48432	4836	2.65	2.65	1.3	-0.29	-0.31
HD 50904	4953	2.70	2.70	1.3	-0.13	-0.08
HD 53329	5012	2.80	2.80	1.2	-0.38	-0.40
HD 54810	4669	2.40	2.40	1.4	-0.47	-0.49
HD 55280	4654	2.25	2.25	1.3	-0.08	-0.07
HD 56891	4709	2.40	2.40	1.4	0.11	0.07
HD 58207	4799	2.35	2.35	1.4	-0.14	-0.09
HD 59686	4654	2.40	2.40	1.4	0.02	0.03
HD 60294	4569	2.15	2.15	1.3	-0.08	-0.07
HD 60318	4962	2.80	2.80	1.2	0.08	0.11

Table A.2. continued.

Star	T_{eff} , K	$\log g(\text{E})$	$\log g(\text{Ca})$	V_t , km s $^{-1}$	[Fe/H] $_{\text{I}}$	[Fe/H] $_{\text{II}}$
HD 60341	4634	2.15	2.15	1.4	-0.02	0.01
HD 60986	5057	2.60	2.60	1.3	-0.01	0.00
HD 61363	4785	2.10	2.10	1.1	-0.21	-0.17
HD 61935	4780	2.40	2.30	1.3	-0.09	-0.09
HD 62141	4971	2.80	2.70	1.3	-0.14	-0.13
HD 62345	5032	2.60	2.50	1.2	0.05	0.09
HD 63798	5004	2.50	2.50	1.3	-0.10	-0.08
HD 64152	4977	2.70	2.70	1.3	-0.01	-0.01
HD 64967	4864	2.55	2.55	1.5	-0.65	-0.64
HD 65066	4868	2.60	2.60	1.5	0.02	0.04
HD 65345	4963	2.70	2.60	1.3	0.00	0.03
HD 67539	4781	2.45	2.45	1.2	-0.61	-0.63
HD 68312	5090	2.70	2.70	1.3	-0.09	-0.10
HD 68375	5071	2.90	2.90	1.3	0.00	0.02
HD 70523	4642	2.20	2.10	1.4	-0.25	-0.26
HD 71088	4900	2.70	2.70	1.3	-0.03	-0.01
HD 73017	4693	2.30	2.40	1.2	-0.66	-0.64
HD 74794	4701	2.25	2.25	1.4	-0.02	-0.01
HD 75506	4876	2.50	2.30	1.3	-0.30	-0.31
HD 75958	5030	2.70	2.80	1.3	-0.09	-0.07
HD 76291	4495	2.00	2.20	1.3	-0.28	-0.30
HD 76813	5060	2.80	2.80	1.4	-0.09	-0.07
HD 78235	5070	2.80	2.80	1.3	-0.14	-0.14
HD 79181	4867	2.40	2.40	1.2	-0.28	-0.25
HD 80546	4601	2.25	2.35	1.3	-0.05	-0.03
HD 81688	4789	2.30	2.30	1.3	-0.23	-0.21
HD 82969	4948	2.70	3.00	1.2	-0.22	-0.20
HD 83240	4682	2.45	2.45	1.3	-0.02	0.01
HD 83371	4861	2.60	2.50	1.3	-0.39	-0.39
HD 86513	4755	2.30	2.30	1.4	-0.08	0.07
HD 90633	4596	2.30	2.30	1.3	0.02	0.03
HD 93291	5061	2.75	2.75	1.3	-0.05	0.05
HD 93875	4590	2.25	2.25	1.4	0.06	0.01
HD 94084	4787	2.65	2.45	1.4	0.11	0.11
HD 94402	5004	2.70	2.70	1.4	0.11	0.13
HD 94497	4702	2.30	2.30	1.3	-0.19	-0.18
HD 95808	4946	2.55	2.55	1.4	-0.09	-0.11
HD 98366	4702	2.40	2.40	1.2	-0.10	-0.11
HD 100696	4862	2.40	2.30	1.4	-0.31	-0.31
HD 101484	4991	2.70	2.70	1.3	-0.03	-0.03
HD 102928	4654	2.35	2.25	1.4	-0.28	-0.28
HD 103605	4611	2.35	2.35	1.4	-0.10	-0.08
HD 103912	4870	2.80	2.80	1.1	-0.65	-0.68
HD 104783	5247	2.55	2.55	1.5	-0.36	-0.34
HD 106714	4935	2.50	2.50	1.2	-0.09	-0.05
HD 108381	4680	2.50	2.50	1.3	0.21	0.17
HD 109053	4921	2.50	-	1.4	-0.38	-
HD 110024	4921	2.70	2.70	1.4	0.06	0.08
HD 113321	4739	2.10	-	1.4	-0.07	-
HD 113997	4697	1.70	-	1.5	-0.12	-
HD 114357	4551	2.30	2.50	1.5	0.12	0.11
HD 116292	4922	2.60	2.60	1.5	-0.03	-0.01
HD 117304	4630	2.17	2.35	1.3	-0.15	-0.12
HD 117566	5475	3.15	3.15	1.35	0.09	0.08
HD 119126	4802	2.25	2.25	1.35	-0.12	-0.11
HD 120084	4883	2.55	2.65	1.5	0.09	0.08
HD 120164	4746	2.30	2.30	1.5	-0.07	-0.08
HD 120420	4676	2.15	2.30	1.25	-0.27	-0.24
HD 136138	4995	2.60	2.80	1.5	-0.19	-0.19
HD 138852	4859	2.30	2.30	1.4	-0.24	-0.21
HD 139254	4708	2.35	2.35	1.4	-0.04	-0.02
HD 139329	4690	2.30	2.20	1.4	-0.31	-0.33
HD 143553	4644	2.30	2.15	1.0	-0.36	-0.36
HD 146388	4731	2.45	2.45	1.4	0.08	0.08
HD 148604	5110	2.80	2.80	1.1	-0.18	-0.16
HD 152224	4685	2.25	2.15	1.4	-0.24	-0.25
HD 153956	4604	2.45	2.45	1.5	0.25	0.24

Table A.2. continued.

Star	T_{eff} , K	$\log g(\text{E})$	$\log g(\text{Ca})$	V_r , km s $^{-1}$	[Fe/H] _I	[Fe/H] _{II}
HD 155970	4717	2.50	2.50	1.3	0.08	0.09
HD 156874	4881	2.50	2.50	1.3	0.00	0.02
HD 159353	4850	2.40	2.40	1.35	-0.08	-0.08
HD 161178	4789	2.20	2.40	1.3	-0.24	-0.24
HD 162076	4959	2.70	2.70	1.3	-0.03	0.01
HD 166578	4859	2.50	2.55	1.4	-0.62	-0.64
HD 168653	4632	2.20	2.40	1.4	-0.16	-0.16
HD 170693	4256	1.50	1.50	1.25	-0.59	-0.61
HD 171994	5014	2.70	2.70	1.2	-0.23	-0.22
HD 175743	4669	2.50	2.40	1.4	0.04	0.04
HD 176408	4564	2.25	2.25	1.5	0.04	0.03
HD 176598	5024	2.80	2.80	1.3	0.03	0.04
HD 180711	4824	2.40	2.40	1.3	-0.20	-0.17
HD 185644	4591	2.40	2.40	1.20	0.01	-0.01
HD 187739	4649	2.30	2.10	1.2	-0.34	-0.37
HD 188119	4993	2.75	2.60	1.1	-0.31	-0.28
HD 192787	4987	2.60	2.60	1.35	-0.05	-0.97
HD 192836	4772	2.60	2.55	1.35	0.01	0.00
HD 195330	4792	2.40	2.50	0.6	-0.29	-0.30
HD 196134	4741	2.40	2.50	1.3	-0.14	-0.12
HD 198431	4524	2.00	2.10	1.2	-0.37	-0.39
HD 199870	4937	2.70	2.70	1.3	0.02	-0.01
HD 204771	4904	2.70	2.70	1.2	0.02	-0.01
HD 206005	4709	2.20	2.20	1.2	-0.19	-0.19
HD 207130	4792	2.60	2.50	1.2	0.13	0.12
HD 208111	4592	2.30	2.10	1.5	0.18	0.21
HD 211006	4553	2.35	2.35	1.4	0.07	0.05
HD 212496	4646	2.30	2.30	1.2	-0.48	-0.49
HD 214567	4981	2.50	2.50	1.3	-0.14	-0.13
HD 215030	4723	2.35	2.35	1.25	-0.49	-0.49
HD 215721	4890	2.40	2.40	1.3	-0.50	-0.52
HD 216131	4984	2.70	2.70	1.2	-0.07	-0.07
HD 216228	4698	2.40	2.40	1.4	-0.16	-0.13
HD 218031	4692	2.20	2.30	1.4	-0.24	-0.22
HD 219418	5281	2.80	2.90	1.3	-0.27	-0.29
HD 219916	5038	2.80	2.80	1.2	0.03	0.01
HD 221345	4664	2.20	2.10	1.4	-0.37	-0.41
HD 221833	4603	2.30	2.30	1.4	0.02	0.03
HD 225197	4734	2.50	2.40	1.4	0.14	0.13
HD 225216	4720	2.20	2.20	1.4	-0.15	-0.13
BD+22 2606	4680	2.10	-	1.3	-0.26	-
BD+25 2555	5014	2.80	-	1.5	-0.43	-
BD+28 2250	4630	2.00	-	1.2	-0.69	-

Table A.3. Si, Ca, and Ni abundances.

Star	[Fe/H]	σ	[Si/Fe]	σ	[Ca/Fe]	σ	[Ni/Fe]	σ
HD 2910	0.12	0.11	0.16	0.10	0.02	0.16	0.05	0.13
HD 4188	0.04	0.13	0.16	0.13	0.02	0.07	0.01	0.16
HD 4482	0.02	0.10	0.12	0.08	0.05	0.06	-0.10	0.09
HD 5395	-0.32	0.06	0.13	0.10	0.05	0.08	-0.02	0.06
HD 6319	0.06	0.09	0.11	0.09	0.02	0.09	0.01	0.08
HD 6482	-0.11	0.07	0.10	0.07	0.07	0.10	-0.02	0.09
HD 7106	0.05	0.11	0.18	0.10	0.05	0.12	0.08	0.12
HD 7578	0.12	0.13	0.14	0.05	0.03	0.14	0.12	0.12
HD 8207	0.27	0.12	0.17	0.11	0.03	0.15	0.05	0.14
HD 8599	-0.22	0.05	0.09	0.05	0.10	0.13	-0.01	0.07
HD 8733	0.02	0.06	0.03	0.08	0.06	0.09	-0.02	0.09
HD 9408	-0.21	0.10	0.15	0.08	0.03	0.07	0.06	0.08
HD 10975	-0.19	0.12	0.10	0.10	-0.01	0.06	-0.03	0.12
HD 11559	0.05	0.14	0.15	0.16	0.06	0.11	-0.07	0.14
HD 11749	-0.10	0.15	0.20	0.15	0.00	0.11	0.00	0.17
HD 11949	-0.16	0.07	0.12	0.06	0.14	0.08	0.00	0.09
HD 15453	-0.07	0.13	0.14	0.09	0.01	0.02	0.00	0.03
HD 15755	-0.01	0.09	0.14	0.09	-0.02	0.11	0.03	0.08
HD 15779	0.02	0.13	0.13	0.12	0.02	0.07	-0.05	0.16
HD 16247	-0.22	0.16	0.18	0.08	0.00	0.14	0.00	0.13
HD 16400	-0.01	0.08	0.05	0.07	0.05	0.10	0.01	0.12
HD 17361	0.12	0.13	0.18	0.10	-0.01	0.18	0.07	0.15
HD 18885	0.16	0.09	0.04	0.04	-0.07	0.09	0.01	0.11
HD 19270	0.15	0.16	0.15	0.10	0.04	0.12	0.12	0.09
HD 19787	0.14	0.13	0.17	0.13	0.07	0.11	-0.05	0.16
HD 19845	0.11	0.09	0.17	0.10	0.02	0.07	0.10	0.12
HD 20791	0.11	0.08	0.13	0.12	-0.02	0.10	0.02	0.12
HD 25602	-0.42	0.12	0.22	0.03	0.15	0.11	0.04	0.08
HD 25604	0.13	0.14	0.23	0.09	-0.03	0.21	0.04	0.19
HD 26546	-0.01	0.12	0.10	0.12	0.14	0.13	0.02	0.15
HD 26659	-0.13	0.09	0.03	0.09	0.06	0.12	-0.04	0.06
HD 26755	-0.06	0.11	0.14	0.10	0.08	0.13	-0.03	0.11
HD 27348	0.14	0.11	0.12	0.11	-0.06	0.10	0.14	0.08
HD 27371	0.11	0.10	0.07	0.12	0.10	0.12	-0.04	0.12
HD 27697	0.11	0.09	0.07	0.11	0.08	0.12	0.06	0.09
HD 28292	-0.18	0.16	0.24	0.11	0.01	0.10	-0.06	0.13
HD 28305	0.11	0.09	0.09	0.11	0.11	0.12	0.09	0.11
HD 28307	0.12	0.13	0.06	0.10	0.04	0.12	0.00	0.12
HD 30557	-0.07	0.11	0.07	0.08	0.01	0.09	0.02	0.13
HD 31444	-0.17	0.10	0.06	0.13	0.14	0.13	-0.07	0.08
HD 33419	0.00	0.10	0.12	0.09	0.08	0.06	0.12	0.10
HD 33618	0.05	0.11	0.14	0.09	0.04	0.12	0.08	0.13
HD 34200	0.04	0.08	0.04	0.07	0.00	0.06	-0.07	0.04
HD 34559	0.04	0.09	0.06	0.10	0.05	0.11	-0.09	0.08
HD 35369	-0.14	0.12	0.05	0.10	-0.02	0.08	-0.02	0.11
HD 37638	-0.01	0.12	0.06	0.09	0.01	0.02	-0.07	0.06
HD 39070	0.03	0.07	0.01	0.07	-0.02	0.01	-0.03	0.06
HD 39910	0.27	0.10	0.09	0.13	-0.07	0.05	0.15	0.10
HD 40020	0.13	0.11	0.15	0.14	-0.05	0.02	0.05	0.11
HD 40801	-0.21	0.07	0.07	0.06	0.11	0.09	0.03	0.07
HD 42341	0.25	0.12	0.16	0.14	-0.01	0.01	0.12	0.10
HD 43023	-0.13	0.12	0.04	0.06	0.11	0.10	-0.07	0.08
HD 45415	0.03	0.12	-0.01	-	-0.07	0.10	0.03	0.12
HD 46374	0.03	0.13	0.10	0.09	-0.02	0.12	-0.05	0.14
HD 46758	-0.30	0.09	0.14	0.10	0.06	0.10	-0.01	0.09
HD 47138	-0.06	0.06	0.03	0.09	0.05	0.07	-0.06	0.07
HD 47366	-0.16	0.11	0.12	0.07	0.06	0.09	0.00	0.09
HD 48432	-0.29	0.10	0.17	0.08	0.14	0.08	-0.04	0.10
HD 50904	-0.13	0.09	0.13	0.10	0.07	0.10	-0.06	0.10
HD 53329	-0.38	0.10	0.08	0.10	0.05	0.06	-0.04	0.07
HD 54810	-0.47	0.13	0.22	0.11	0.11	0.07	-0.02	0.07
HD 55280	-0.08	0.12	0.13	0.09	0.07	0.06	0.01	0.12
HD 56891	0.11	0.11	0.11	0.14	0.01	0.07	0.10	0.12
HD 58207	-0.14	0.11	0.12	0.09	0.09	0.08	0.03	0.10
HD 59686	0.02	0.11	0.19	0.09	0.03	0.11	0.06	0.16
HD 60294	-0.08	0.13	0.16	0.11	0.03	0.05	0.02	0.11
HD 60318	0.08	0.11	0.10	0.11	-0.01	0.07	0.08	0.12

Table A.3. continued.

Star	[Fe/H]	σ	[Si/Fe]	σ	[Ca/Fe]	σ	[Ni/Fe]	σ
HD 60341	-0.02	0.11	0.10	0.08	0.05	0.14	0.04	0.13
HD 60986	-0.01	0.10	0.10	0.10	0.10	0.09	-0.03	0.08
HD 61363	-0.21	0.30	0.04	0.07	0.01	0.08	0.01	0.11
HD 61935	-0.09	0.13	0.10	0.07	0.12	0.13	0.01	0.07
HD 62141	-0.14	0.11	0.14	0.12	-0.03	0.12	-0.09	0.10
HD 62345	0.05	0.11	0.05	0.11	0.08	0.07	-0.06	0.11
HD 63798	-0.10	0.11	0.09	0.10	0.09	0.07	-0.03	0.11
HD 64152	-0.01	0.10	0.09	0.09	0.07	0.07	0.00	0.07
HD 64967	-0.65	0.13	0.34	0.07	0.14	0.12	-0.02	0.12
HD 65066	0.02	0.11	0.11	0.09	0.06	0.10	0.00	0.12
HD 65345	0.00	0.11	0.12	0.10	-0.04	0.08	0.00	0.11
HD 67539	-0.61	0.08	0.25	0.11	0.23	0.10	0.06	0.06
HD 68312	-0.09	0.12	0.06	0.10	0.06	0.10	-0.07	0.10
HD 68375	0.00	0.11	0.07	0.10	-0.06	0.05	-0.03	0.10
HD 70523	-0.25	0.11	0.24	0.08	0.10	0.12	0.03	0.11
HD 71088	-0.03	0.12	0.07	0.06	-0.06	0.06	-0.03	0.11
HD 73017	-0.66	0.10	0.21	0.09	0.02	0.04	-0.01	0.08
HD 74794	-0.02	0.12	0.09	0.12	0.06	0.09	0.04	0.15
HD 75506	-0.30	0.10	0.07	0.10	0.10	0.10	0.02	0.11
HD 75958	-0.09	0.12	0.09	0.10	-0.06	0.02	-0.10	0.08
HD 76291	-0.28	0.10	0.21	0.11	0.08	0.10	-0.01	0.08
HD 76813	-0.09	0.08	0.10	0.07	-0.03	0.05	-0.08	0.10
HD 78235	-0.14	0.11	0.11	0.10	0.07	0.11	-0.09	0.09
HD 79181	-0.28	0.09	0.14	0.08	0.11	0.13	0.05	0.10
HD 80546	-0.05	0.13	0.12	0.09	0.11	0.14	0.02	0.10
HD 81688	-0.23	0.10	0.10	0.09	-0.01	0.07	-0.04	0.09
HD 82969	-0.22	0.09	0.07	0.10	0.00	0.09	-0.11	0.09
HD 83240	-0.02	0.13	0.08	0.10	-0.03	0.09	-0.02	0.12
HD 83371	-0.39	0.07	0.13	0.11	0.08	0.12	-0.03	0.11
HD 86513	-0.08	0.13	0.10	0.12	0.12	0.08	-0.02	0.14
HD 90633	0.02	0.12	0.14	0.09	-0.01	0.10	-0.02	0.09
HD 93291	-0.05	0.11	0.00	0.07	0.05	0.12	-0.06	0.06
HD 93875	0.06	0.10	0.15	0.14	0.02	0.10	0.05	0.10
HD 94084	0.11	0.10	0.16	0.11	0.01	0.11	0.04	0.11
HD 94402	0.11	0.10	0.05	0.09	0.04	0.14	0.06	0.16
HD 94497	-0.19	0.13	0.15	0.10	0.00	0.11	-0.03	0.09
HD 95808	-0.09	0.12	0.11	0.07	0.07	0.04	-0.03	0.12
HD 98366	-0.1	0.12	0.09	0.07	0.06	0.11	0.00	0.11
HD 100696	-0.31	0.13	0.11	0.09	0.07	0.04	-0.06	0.09
HD 101484	-0.03	0.12	0.13	0.06	0.05	0.09	0.02	0.11
HD 102928	-0.28	0.12	0.21	0.13	0.11	0.11	-0.04	0.14
HD 103605	-0.10	0.12	0.19	0.06	0.01	0.11	0.00	0.13
HD 103912	-0.65	0.12	0.28	0.10	0.19	0.08	-0.03	0.11
HD 104783	-0.36	0.10	0.19	0.08	0.10	0.12	0.06	0.09
HD 106714	-0.09	0.10	0.00	0.06	0.01	0.08	-0.04	0.12
HD 108381	0.21	0.09	0.17	0.13	0.01	0.11	0.15	0.14
HD 109053	-0.38	0.14	0.10	0.08	0.08	0.12	-0.03	0.09
HD 110024	0.06	0.10	0.12	0.08	0.00	0.08	0.02	0.12
HD 113321	-0.07	0.12	0.09	0.12	0.04	0.10	0.04	0.13
HD 113997	-0.12	0.15	0.08	0.14	0.18	0.19	-0.02	0.12
HD 114357	0.12	0.13	0.15	0.11	-0.08	0.10	0.11	0.14
HD 116292	-0.03	0.11	0.08	0.07	0.00	0.07	-0.06	0.12
HD 117304	-0.15	0.12	0.15	0.13	-0.03	0.09	0.01	0.10
HD 117566	0.09	0.10	-0.05	0.10	-0.06	0.14	-0.05	0.08
HD 119126	-0.12	0.11	0.12	0.10	0.08	0.12	0.04	0.08
HD 120084	0.09	0.10	0.08	0.10	-0.08	0.10	0.00	0.09
HD 120164	-0.07	0.13	0.13	0.08	-0.05	0.04	-0.04	0.13
HD 120420	-0.27	0.12	0.10	0.07	0.07	0.07	-0.02	0.10
HD 136138	-0.19	0.13	0.14	0.09	0.06	0.07	-0.04	0.09
HD 138852	-0.24	0.12	0.09	0.09	0.04	0.05	-0.03	0.08
HD 139254	-0.04	0.13	0.10	0.09	-0.04	0.12	-0.05	0.13
HD 139329	-0.31	0.12	0.15	0.08	0.10	0.07	-0.05	0.10
HD 143553	-0.36	0.10	0.16	0.10	0.14	0.08	0.05	0.10
HD 146388	0.08	0.12	0.10	0.09	0.12	0.07	0.09	0.08
HD 148604	-0.18	0.11	0.08	0.10	0.04	0.10	-0.07	0.08
HD 152224	-0.24	0.15	0.08	0.07	0.05	0.11	-0.06	0.09
HD 153956	0.25	0.13	0.11	0.11	-0.10	0.10	0.05	0.14

Table A.3. continued.

Star	[Fe/H]	σ	[Si/Fe]	σ	[Ca/Fe]	σ	[Ni/Fe]	σ
HD 155970	0.08	0.10	0.17	0.08	0.01	0.10	0.08	0.11
HD 156874	0.00	0.12	0.04	0.13	-0.07	0.03	-0.06	0.12
HD 159353	-0.08	0.11	0.13	0.08	0.11	0.09	0.02	0.09
HD 161178	-0.24	0.09	0.06	0.10	0.05	0.08	0.01	0.08
HD 162076	-0.03	0.11	0.11	0.09	0.08	0.08	-0.05	0.09
HD 166578	-0.62	0.09	0.24	0.09	0.16	0.05	-0.01	0.08
HD 168653	-0.16	0.13	0.16	0.11	-0.01	0.09	0.00	0.14
HD 170693	-0.59	0.08	0.26	0.12	-0.02	0.06	0.01	0.11
HD 171994	-0.23	0.10	0.08	0.08	0.05	0.04	-0.10	0.09
HD 175743	0.04	0.13	0.12	0.08	-0.01	0.12	0.06	0.15
HD 176408	0.04	0.14	0.17	0.11	0.05	0.14	0.02	0.15
HD 176598	0.03	0.10	0.05	0.06	0.00	0.12	-0.08	0.10
HD 180711	-0.20	0.11	0.08	0.08	0.06	0.10	0.03	0.08
HD 185644	0.01	0.11	0.07	0.05	-0.08	0.13	-0.04	0.09
HD 187739	-0.34	0.13	0.14	0.08	0.06	0.12	-0.08	0.12
HD 188119	-0.31	0.10	0.12	0.10	0.02	0.10	0.03	0.13
HD 192787	-0.05	0.12	0.04	0.06	-0.05	0.07	-0.07	0.12
HD 192836	0.01	0.12	0.10	0.08	0.05	0.11	0.11	0.06
HD 195330	-0.29	0.11	0.11	0.12	0.11	0.15	0.00	0.13
HD 196134	-0.14	0.09	0.11	0.07	0.00	0.12	0.02	0.10
HD 198431	-0.37	0.09	0.18	0.10	0.04	0.11	-0.03	0.10
HD 199870	0.02	0.13	0.16	0.09	0.10	0.14	0.09	0.14
HD 204771	0.02	0.10	0.05	0.07	0.06	0.12	-0.01	0.07
HD 206005	-0.19	0.11	0.07	0.06	0.16	0.15	0.00	0.11
HD 207130	0.13	0.11	0.02	0.08	0.01	0.12	0.09	0.14
HD 208111	0.18	0.11	0.12	0.06	-0.09	0.13	0.03	0.16
HD 211006	0.07	0.10	0.22	0.09	0.05	0.12	0.07	0.14
HD 212496	-0.48	0.09	0.17	0.08	0.12	0.08	-0.02	0.05
HD 214567	-0.14	0.12	0.08	0.10	-0.01	0.04	-0.06	0.11
HD 215030	-0.49	0.12	0.21	0.10	0.12	0.10	0.01	0.07
HD 215721	-0.50	0.09	0.12	0.09	0.15	0.10	0.01	0.08
HD 216131	-0.07	0.13	0.06	0.10	0.04	0.10	-0.02	0.12
HD 216228	-0.16	0.14	0.16	0.10	0.08	0.11	0.01	0.13
HD 218031	-0.24	0.09	0.15	0.11	-0.02	0.08	-0.03	0.12
HD 219418	-0.27	0.09	0.07	0.10	0.02	0.09	-0.02	0.05
HD 219916	0.03	0.11	0.02	0.07	-0.03	0.06	-0.06	0.12
HD 221345	-0.37	0.13	0.20	0.10	0.07	0.04	0.03	0.13
HD 221833	0.02	0.12	0.11	0.08	0.06	0.13	0.01	0.16
HD 225197	0.14	0.13	0.03	0.12	0.00	0.10	0.10	0.16
HD 225216	-0.15	0.10	0.11	0.06	0.09	0.13	0.05	0.10
BD+22 2606	-0.26	0.12	0.10	0.07	0.00	0.05	0.02	0.12
BD+25 2555	-0.43	0.10	0.22	0.05	0.17	0.09	-0.02	0.11
BD+28 2250	-0.69	0.13	0.29	0.11	0.38	0.09	0.00	0.14

Table A.4. The abundances of lithium, carbon, nitrogen, oxygen, sodium, and magnesium in atmospheres of studied stars.

Star	T_{eff}	$\log g$	(Li/H)	(C/H)	(N/H)	(O/H)	(Na/H)	(Mg/H)
HD 2910	4756	2.70	–	8.42	8.25	8.80	6.30	7.45
HD 4188	4809	2.70	–	8.32	8.20	8.85	6.28	7.41
HD 4482	4917	2.65	0.65	8.25	8.20	8.60	6.32	7.55
HD 5395	4849	2.15	–	8.05	–	8.50	5.80	7.43
HD 6319	4650	2.30	–	8.30	8.10	8.72	6.31	7.55
HD 6482	4738	2.40	–	8.20	8.05	8.63	6.18	7.47
HD 7106	4684	2.55	–	8.40	8.20	8.80	6.32	7.56
HD 7578	4680	2.50	–	8.40	8.40	8.75	6.70	7.58
HD 8207	4750	2.75	–	8.35	8.40	8.80	6.47	7.60
HD 8599	4781	2.50	–	8.15	8.00	8.63	6.04	7.47
HD 8733	4932	2.70	1.05	8.25	8.10	8.57	6.20	7.50
HD 9408	4804	2.30	–	8.22	7.95	8.70	6.00	7.35
HD 10975	4881	2.20	–	8.15	7.95	8.65	6.08	7.40
HD 11559	4977	3.00	–	8.30	8.30	8.80	6.36	7.54
HD 11749	4679	2.40	–	8.25	7.80	8.80	6.05	7.42
HD 11949	4708	2.30	–	8.13	7.95	8.60	6.12	7.48
HD 15453	4696	2.40	0.90	8.25	8.00	8.72	6.23	7.52
HD 15755	4611	2.30	–	8.25	8.15	8.75	6.22	7.56
HD 15779	4821	2.70	–	8.30	8.25	8.78	6.28	7.43
HD 16247	4629	2.20	–	8.20	–	8.63	5.98	7.43
HD 16400	4840	2.50	–	8.26	8.25	8.72	6.32	7.40
HD 17361	4646	2.50	–	8.37	8.15	8.75	6.37	7.55
HD 18885	4722	2.50	–	8.40	8.28	8.80	6.51	7.55
HD 19270	4723	2.40	–	8.30	8.30	8.68	6.49	7.58
HD 19787	4832	2.75	–	8.35	8.23	8.75	6.41	7.53
HD 19845	4933	2.80	–	8.30	8.30	8.65	6.47	7.50
HD 20791	4986	2.80	–	8.33	8.25	8.75	6.40	7.48
HD 25602	4693	2.40	–	8.08	7.80	–	6.02	7.40
HD 25604	4764	2.70	–	8.40	8.35	8.80	6.47	7.52
HD 26546	4743	2.25	–	8.22	8.10	8.65	6.40	7.50
HD 26659	5178	2.90	–	–	–	–	6.26	7.33
HD 26755	4630	2.20	–	–	–	–	6.42	7.55
HD 27348	5003	2.80	–	–	–	–	6.43	7.53
HD 27371	4955	2.70	0.90	8.35	8.40	8.72	6.66	7.60
HD 27697	4975	2.65	0.80	8.25	8.50	8.60	6.60	7.58
HD 28292	4453	2.10	–	8.30	8.00	8.70	6.18	7.45
HD 28305	4925	2.55	0.70	8.35	8.40	8.80	6.67	7.60
HD 28307	4961	2.70	1.10	8.35	8.35	8.70	6.54	7.50
HD 30557	4829	2.45	–	8.15	8.10	8.65	6.26	7.35
HD 31444	5080	2.75	0.60	8.25	8.10	8.60	6.22	7.37
HD 33419	4708	2.30	–	8.30	8.20	8.75	6.48	7.55
HD 33618	4590	2.30	–	8.42	8.23	8.78	6.54	7.62
HD 34200	5055	2.80	–	–	–	–	6.28	7.49
HD 34559	5010	2.90	–	8.35	8.30	8.80	6.40	7.45
HD 35369	4931	2.40	–	8.10	8.10	8.67	6.10	7.35
HD 37638	5093	2.80	–	8.25	8.25	8.72	6.25	7.48
HD 39070	5047	2.80	–	8.25	8.30	8.72	6.28	7.49
HD 39910	4618	2.60	–	8.60	8.25	8.90	6.61	7.61
HD 40020	4670	2.30	–	8.30	8.40	8.72	6.62	7.59
HD 40801	4703	2.20	–	8.18	7.90	8.55	6.05	7.50
HD 42341	4655	2.60	0.85	8.44	8.40	8.80	6.75	7.65
HD 43023	4994	2.40	–	8.15	8.12	8.57	6.22	7.37
HD 45415	4762	2.30	–	–	–	–	6.17	7.43
HD 46374	4661	2.30	0.85	8.25	8.20	8.70	6.23	7.43
HD 46758	5003	2.90	–	8.22	8.10	8.72	6.22	7.25
HD 47138	5211	3.00	1.10	8.20	8.25	8.72	6.26	7.40
HD 47366	4772	2.60	–	8.20	7.95	8.65	6.19	7.37
HD 48432	4836	2.65	–	8.25	8.15	8.80	6.12	7.37
HD 50904	4953	2.70	–	8.20	8.20	8.60	6.21	7.37
HD 53329	5012	2.80	–	8.25	7.95	8.85	5.95	7.28
HD 54810	4669	2.40	–	8.20	7.85	8.65	5.98	7.30
HD 55280	4654	2.25	–	8.25	8.16	8.65	6.35	7.48
HD 56891	4709	2.40	–	8.30	8.25	8.70	6.45	7.55
HD 58207	4799	2.35	–	8.22	8.05	8.65	6.21	7.46
HD 59686	4654	2.40	–	8.37	8.25	8.85	6.44	7.55
HD 60294	4569	2.15	–	8.22	8.15	8.60	6.45	7.47
HD 60318	4962	2.80	–	8.32	8.30	8.77	6.50	7.45

Table A.4. continued.

Star	T_{eff}	$\log g$	(Li/H)	(C/H)	(N/H)	(O/H)	(Na/H)	(Mg/H)
HD 60341	4634	2.15	–	8.35	8.20	8.80	6.47	7.55
HD 60986	5057	2.60	–	–	–	–	6.50	7.50
HD 61363	4785	2.10	–	8.10	8.05	8.60	6.10	7.40
HD 61935	4780	2.40	–	8.25	8.20	8.72	6.30	7.43
HD 62141	4971	2.80	0.95	8.25	8.05	8.60	6.20	7.42
HD 62345	5032	2.60	–	8.30	8.30	8.60	6.48	7.48
HD 63798	5004	2.50	1.75	8.28	8.15	8.60	6.28	7.42
HD 64152	4977	2.70	–	8.30	8.20	8.60	6.44	7.50
HD 64967	4864	2.55	–	–	–	8.45	5.75	7.12
HD 65066	4868	2.60	–	–	–	–	6.34	7.46
HD 65345	4963	2.70	–	8.23	8.10	8.72	6.41	7.41
HD 67539	4781	2.45	–	–	7.75	8.60	5.75	7.27
HD 68312	5090	2.70	–	8.25	8.30	8.72	6.26	7.42
HD 68375	5071	2.90	–	8.35	8.25	8.90	6.26	7.45
HD 70523	4642	2.20	–	–	–	–	6.04	7.48
HD 71088	4900	2.70	–	8.30	8.05	8.76	6.17	7.35
HD 73017	4693	2.30	–	–	–	–	5.76	7.15
HD 74794	4701	2.25	–	8.37	8.15	8.72	6.37	7.57
HD 75506	4876	2.50	–	8.20	8.10	8.85	6.02	7.35
HD 75958	5030	2.70	–	8.25	8.15	8.72	6.14	7.35
HD 76291	4495	2.00	–	8.22	7.90	8.60	6.18	7.50
HD 76813	5060	2.80	–	8.22	8.25	8.75	6.28	7.45
HD 78235	5070	2.80	–	8.20	8.25	8.67	6.36	7.40
HD 79181	4867	2.40	–	8.20	8.00	8.70	6.10	7.38
HD 80546	4601	2.25	–	8.28	8.15	8.62	6.35	7.50
HD 81688	4789	2.30	–	8.20	7.85	8.65	5.94	7.33
HD 82969	4948	2.70	–	8.18	8.05	8.60	6.00	7.40
HD 83240	4682	2.45	1.10	8.25	8.20	8.75	6.20	7.37
HD 83371	4861	2.60	–	8.15	7.95	8.55	5.85	7.35
HD 86513	4755	2.30	–	8.25	8.20	8.72	6.30	7.50
HD 90633	4596	2.30	1.85	–	–	–	6.37	7.60
HD 93291	5061	2.75	–	8.28	8.15	8.65	6.30	7.41
HD 93875	4590	2.25	–	8.43	8.25	8.76	6.47	7.58
HD 94084	4787	2.65	–	8.42	8.35	8.87	6.55	7.50
HD 94402	5004	2.70	–	8.30	8.30	8.72	6.43	7.52
HD 94497	4702	2.30	–	8.20	7.95	8.65	6.13	7.38
HD 95808	4946	2.55	–	8.20	8.30	8.67	6.44	7.47
HD 98366	4702	2.40	–	8.22	8.00	8.57	6.18	7.43
HD 100696	4862	2.40	–	8.22	8.00	8.77	5.98	7.30
HD 101484	4991	2.70	–	8.25	8.15	8.72	6.36	7.50
HD 102928	4654	2.35	–	8.25	8.05	8.73	6.13	7.40
HD 103605	4611	2.35	–	8.42	8.05	8.85	6.18	7.59
HD 103912	4870	2.80	–	–	–	–	5.52	7.27
HD 104783	5247	2.55	–	–	–	–	5.80	7.23
HD 106714	4935	2.50	–	8.20	8.15	8.80	6.19	7.39
HD 108381	4680	2.50	–	8.42	8.60	8.72	6.71	7.65
HD 109053	4921	2.50	–	8.00	7.80	8.50	5.93	7.38
HD 110024	4921	2.70	–	8.25	8.20	8.68	6.36	7.50
HD 113321	4739	2.10	–	8.30	8.10	8.72	6.23	7.62
HD 113997	4697	1.70	–	8.30	8.10	8.60	6.41	7.35
HD 114357	4551	2.30	–	8.40	8.40	8.80	6.63	7.60
HD 116292	4922	2.60	1.40	–	–	–	6.20	7.48
HD 117304	4630	2.17	1.00	8.15	8.00	8.60	6.10	7.40
HD 117566	5475	3.15	–	–	–	–	6.38	7.48
HD 119126	4802	2.25	–	8.15	8.15	8.70	6.25	7.40
HD 120084	4883	2.55	–	8.33	8.15	8.65	6.37	7.47
HD 120164	4746	2.30	–	8.20	8.15	8.72	6.25	7.40
HD 120420	4676	2.15	–	8.10	7.90	8.60	5.98	7.37
HD 136138	4995	2.60	1.30	8.18	8.15	8.55	6.20	7.36
HD 138852	4859	2.30	–	8.17	7.95	8.72	6.07	7.38
HD 139254	4708	2.35	1.00	8.25	8.10	8.70	6.23	7.45
HD 139329	4690	2.30	–	–	8.05	8.75	6.13	7.45
HD 143553	4644	2.30	–	8.10	–	8.40	5.94	7.27
HD 146388	4731	2.45	–	–	–	–	6.52	7.59
HD 148604	5110	2.80	0.90	–	–	–	6.16	7.25
HD 152224	4685	2.25	–	8.10	7.95	8.62	6.06	7.37
HD 153956	4604	2.45	–	8.54	–	8.95	6.64	7.62

Table A.4. continued.

Star	T_{eff}	$\log g$	(Li/H)	(C/H)	(N/H)	(O/H)	(Na/H)	(Mg/H)
HD 155970	4717	2.50	–	8.35	8.20	8.65	6.50	7.55
HD 156874	4881	2.50	–	8.20	8.20	8.60	6.23	7.45
HD 159353	4850	2.40	–	8.20	8.30	8.65	6.20	7.50
HD 161178	4789	2.20	–	8.22	8.05	8.60	6.15	7.50
HD 162076	4959	2.70	1.00	8.25	8.25	8.60	6.32	7.42
HD 166578	4859	2.50	–	8.10	7.95	8.55	5.71	7.25
HD 168653	4632	2.20	–	8.18	7.95	8.60	6.10	7.40
HD 170693	4256	1.50	–	8.00	7.70	–	5.76	7.23
HD 171994	5014	2.70	0.60	8.20	8.05	8.58	6.15	7.33
HD 175743	4669	2.50	–	–	–	–	6.28	7.48
HD 176408	4564	2.25	–	8.32	8.30	8.72	6.50	7.54
HD 176598	5024	2.80	1.20	8.25	8.38	8.60	6.45	7.47
HD 180711	4824	2.40	–	8.15	8.10	8.67	6.16	7.35
HD 185644	4591	2.40	–	8.28	8.10	8.70	6.33	7.60
HD 187739	4649	2.30	–	8.00	–	8.45	5.98	7.13
HD 188119	4993	2.75	–	8.20	8.30	8.85	6.00	7.28
HD 192787	4987	2.60	–	8.20	8.20	8.65	6.21	7.40
HD 192836	4772	2.60	1.25	8.32	8.15	8.72	6.31	7.47
HD 195330	4792	2.40	–	–	–	–	6.12	7.10
HD 196134	4741	2.40	–	8.15	7.95	8.55	6.06	7.33
HD 198431	4524	2.00	–	–	7.85	8.70	6.02	7.38
HD 199870	4937	2.70	–	8.28	8.20	8.65	6.46	7.53
HD 204771	4904	2.70	–	8.25	8.30	8.75	6.38	7.55
HD 206005	4709	2.20	–	8.15	8.00	8.60	6.24	7.45
HD 207130	4792	2.60	–	–	–	–	6.55	7.55
HD 208111	4592	2.30	–	8.43	8.75	8.75	6.58	7.65
HD 211006	4553	2.35	–	8.42	8.20	8.72	6.54	7.58
HD 212496	4646	2.30	–	8.06	7.85	8.42	5.95	7.25
HD 214567	4981	2.50	–	8.15	8.15	8.65	6.22	7.39
HD 215030	4723	2.35	–	8.00	–	8.37	5.81	7.15
HD 215721	4890	2.40	–	8.10	8.00	8.55	5.85	7.27
HD 216131	4984	2.70	–	8.22	8.20	–	6.27	7.42
HD 216228	4698	2.40	–	8.22	8.15	8.68	6.27	7.34
HD 218031	4692	2.20	–	8.20	8.00	8.70	6.04	7.45
HD 219418	5281	2.80	–	8.15	8.35	8.60	6.25	7.42
HD 219916	5038	2.80	–	8.00	8.15	8.65	6.31	7.45
HD 221345	4664	2.20	–	–	–	–	6.00	7.50
HD 221833	4603	2.30	–	8.37	8.00	8.72	6.38	7.50
HD 225197	4734	2.50	–	8.37	8.40	8.75	6.55	7.60
HD 225216	4720	2.20	–	8.20	8.20	8.67	6.29	7.50
BD+22 2606	4680	2.10	–	8.15	8.00	8.60	6.28	7.49
BD+25 2555	5014	2.80	–	8.10	7.70	9.00	5.86	7.44
BD+28 2250	4630	2.00	–	–	–	8.75	5.76	7.24