

Water in the envelopes and disks around young high-mass stars

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Received 28 July 2005 / Accepted 21 October 2005

ABSTRACT

Single-dish spectra and interferometric maps of (sub-)millimeter lines of H_2^{18}O and HDO are used to study the chemistry of water in eight regions of high-mass star formation. The spectra indicate HDO excitation temperatures of ~ 110 K and column densities in an $11''$ beam of $\sim 2 \times 10^{14} \text{ cm}^{-2}$ for HDO and $\sim 2 \times 10^{17} \text{ cm}^{-2}$ for H_2O , with the $N(\text{HDO})/N(\text{H}_2\text{O})$ ratio increasing with decreasing temperature. Simultaneous observations of CH_3OH and SO_2 indicate that 20–50% of the single-dish line flux arises in the molecular outflows of these objects. The outflow contribution to the H_2^{18}O and HDO emission is estimated to be 10–20%. Radiative transfer models indicate that the water abundance is low ($\sim 10^{-6}$) outside a critical radius corresponding to a temperature in the protostellar envelope of ~ 100 K, and “jumps” to $\text{H}_2\text{O}/\text{H}_2 \sim 10^{-4}$ inside this radius. This value corresponds to the observed abundance of solid water and together with the derived HDO/ H_2O abundance ratios of $\sim 10^{-3}$ suggests that the origin of the observed water is evaporation of grain mantles. This idea is confirmed in the case of AFGL 2591 by interferometer observations of the HDO $1_{10-1_{11}}$, H_2^{18}O $3_{13-2_{20}}$ and SO_2 $12_{0,12-11_{1,11}}$ lines, which reveal compact ($\varnothing \sim 800$ AU) emission with a systematic velocity gradient. This size is similar to that of the 1.3 mm continuum towards AFGL 2591, from which we estimate a mass of $\sim 0.8 M_{\odot}$, or $\sim 5\%$ of the mass of the central star. We speculate that we may be observing a circumstellar disk in an almost face-on orientation.

Key words. ISM: molecules – molecular processes – stars: formation – astrochemistry

1. Introduction

Water is a cornerstone molecule in the oxygen chemistry of dense interstellar clouds and a major coolant of warm molecular gas¹. In the surroundings of embedded protostars, water can be formed by three very different mechanisms. In cold (~ 10 K) protostellar envelopes, water may be formed in the gas phase by ion-molecule chemistry, through dissociative recombination of H_3O^+ . Simultaneously, on the surfaces of cold dust grains, O and H atoms may combine to form water-rich ice mantles. These mantles will evaporate when the grains are heated to ~ 100 K, either by protostellar radiation or by grain sputtering in outflow shocks. Third, in gas with temperatures ≥ 250 K, reactions of O and OH with H_2 drive all gas-phase oxygen into water. Such high temperatures may occur very close to the star due to radiation, or further out in outflow shocks. The water molecule thus offers the opportunity to study the relative importance of each of these types of chemistry in the protostellar environment.

There has been considerable controversy about the water abundance around high-mass protostars. Observations of the

H_2O $6 \mu\text{m}$ bending mode with ISO-SWS have revealed abundant water ($\text{H}_2\text{O}/\text{H}_2 \sim 10^{-5}$ – 10^{-4}) in absorption toward several high-mass protostars (Boonman & van Dishoeck 2003). The absorption data do not tell us the location of the H_2O along the line of sight, except that the high excitation temperatures (~ 300 – 500 K) imply an origin in warm gas. In contrast, observations of the o- H_2O ground state line at 557 GHz with SWAS of the same sources indicate much lower abundances ($\text{H}_2\text{O}/\text{H}_2 \sim 10^{-7}$ – 10^{-6} ; Snell et al. 2000). The narrow line width indicates an origin in the envelopes rather than the outflows of the sources, but the data have too low angular resolution (several arcminutes) for more detailed statements. Boonman et al. (2003) performed a simultaneous analysis of ISO-SWS, ISO-LWS and SWAS data and inferred a water abundance jump in the inner envelope by four orders of magnitude for several high-mass YSOs.

Locating the water around protostars requires observations at high spatial and spectral resolution, which presently can only be done from the ground. Most ground-based observations of H_2O have targeted the maser lines at 22 and 183 GHz (e.g., Cernicharo et al. 1990). However, the anomalous excitation of these lines makes it hard to derive column densities from such data, which may in any case not be representative of the surrounding region. The only thermal water lines that

¹ This paper uses the word “water” to denote the chemical species, and the notations H_2O , H_2^{18}O and HDO to denote specific isotopologues.

Table 1. Source sample.

Source	RA (1950) (h m s)	Dec (1950) (° ' ")	L ($10^4 L_{\odot}$)	d (kpc)	$N(\text{H}_2)^a$ (10^{23} cm^{-2})
W3 IRS5	02 21 53.1	+61 52 20	17	2.2	2.3
AFGL 490	03 23 38.9	+58 36 33	0.2	1	2.0
W33A	18 11 43.7	-17 53 02	10	4	6.2
AFGL 2136	18 19 36.6	-13 31 40	7	2	1.2
AFGL 2591	20 27 35.8	+40 01 14	2	1	2.3
S140 IRS1	22 17 41.1	+63 03 42	2	0.9	1.4
NGC 7538 IRS1	23 11 36.7	+61 11 51	13	2.8	6.5
NGC 7538 IRS9	23 11 52.8	+61 10 59	4	2.8	3.3

^a Column density in a $15''$ beam.

Table 2. Observed transitions.

Species	Transition $J_{K_a K_c}$	Frequency MHz	E_{up} K	A_{ul} s^{-1}	Telescope	Beam "
HDO	$1_{10}-1_{11}$	80578.3	47	1.3×10^{-6}	IRAM 30 m	30
HDO	$3_{12}-2_{21}$	225896.7	168	1.3×10^{-5}	IRAM 30 m	11
HDO	$2_{11}-2_{12}$	241561.6	95	1.2×10^{-5}	JCMT 15 m	21
HDO	$1_{01}-0_{00}$	464924.5	22	1.7×10^{-4}	JCMT 15 m	12
H_2^{18}O	$3_{13}-2_{20}$	203407.5	204	4.9×10^{-6}	IRAM 30 m	12
SO_2	$12_{0,12}-11_{1,11}$	203391.6	70	8.1×10^{-5}	IRAM 30 m	12
CH_3OH	$5_{-1}-4_0 \text{ E}$	84521.2	40	2.0×10^{-6}	IRAM 30 m	30

can be studied from the ground are the $3_{13}-2_{20}$ line of H_2^{18}O at 203 GHz (Phillips et al. 1978; Jacq et al. 1988), and several HDO lines. These lines were used by Gensheimer et al. (1996) and Helmich et al. (1996) to estimate envelope-averaged abundances of H_2O and HDO around several young high-mass stars. Advances in sensitivity and resolution allow us to consider lower-luminosity objects closer than the Sun than before, and also enable us to study abundance variations with position in the envelope. This paper presents new observations of these lines toward sources that have been studied previously with ISO and SWAS, including the first published interferometric observations of the H_2^{18}O line (and in fact of any non-masing water line). The sources are eight deeply embedded high-mass stars, with luminosities of 2×10^3 – $2 \times 10^5 L_{\odot}$, distances of 1–4 kpc, and H_2 column densities of 1 – $7 \times 10^{23} \text{ cm}^{-2}$, as listed in Table 1. Single-dish mapping of dust continuum and molecular line emission at (sub-)millimeter wavelengths indicates envelope masses of 30–1100 M_{\odot} within radii of 0.09–0.36 pc (Van der Tak et al. 2000b). The sources drive powerful outflows as revealed by mid-infrared and millimeter-wave observations of CO and HCO^+ (Mitchell et al. 1991; Hasegawa & Mitchell 1995a). The unique aspect of this source sample is its high mid-infrared brightness, which allows us to compare its (sub-)millimeter emission with solid state data for the chemistry, and with rovibrational absorption lines for the geometry.

This paper is organized as follows. Section 2 describes the observations, and Sect. 3 their direct results. Section 4 describes modeling of the data with a radiative transfer program. Section 5 discusses the results of the observations and the models in the context of a disk/outflow geometry for these sources. Section 6 concludes the paper with an outlook toward future opportunities in this field.

2. Observations

Table 2 summarizes spectroscopic parameters of the observed lines, and gives the relevant telescope and its $FWHM$ beam size at that frequency. With $E_{\text{up}} \approx 200$ K, the $3_{13}-2_{20}$ line is the lowest-lying transition of H_2^{18}O that can be observed from the ground. We use this line to measure the abundance of H_2O in the warm inner envelopes of the sources. The HDO lines cover the range of excitation energies from 20 to 200 K, and are used to constrain the excitation and chemical state of the gas, in particular its deuterium fractionation. The SO_2 and CH_3OH lines have comparable excitation requirements, and are used to measure the effects of shock chemistry (SO_2) and ice evaporation (CH_3OH). The difference in Einstein A -coefficients of the lines is mostly due to the ν^3 dependence: all the lines have transition dipole moments of a few Debye.

2.1. Single-dish observations

Observations of lines of H_2^{18}O , HDO, SO_2 and CH_3OH in the 80–225 GHz range were made with the 30-m telescope of the Institut de Radio Astronomie Millimétrique (IRAM)² on Pico Veleta, Spain, in May 2003. The front ends were the facility receivers A100, B100, A230 and B230, and the backend was the Versatile Spectral Assembly (VESPA) autocorrelator. The five lines were observed simultaneously with a spectral resolution of 0.1–0.3 km s^{-1} . Integration times are 60–180 min (on+off) using double beam switching with a throw of $180''$.

² IRAM is an international institute for research in millimeter astronomy, co-funded by the Centre National de la Recherche Scientifique (France), the Max Planck Gesellschaft (Germany) and the Instituto Geografico Nacional (Spain).

System temperatures were 100–150 K at 3 mm and 300–600 K at 1.3 mm wavelength. Data were calibrated onto T_{MB} scale by multiplying by η_f/η_b , where the forward efficiency η_f is 0.95 at 3 mm and 0.91 at 1.3 mm, and the main beam efficiency η_b is 0.78 at 3 mm and 0.57 at 1.3 mm wavelength. The spectra have noise levels per 0.25 km s⁻¹ channel of $T_{\text{MB}} = 10$ –15 mK at 80 GHz and 20–30 mK at 215 GHz.

Additional observations of HDO lines at 225, 241 and 464 GHz toward selected sources were carried out with the James Clerk Maxwell Telescope (JCMT)³ on Mauna Kea, Hawaii, in 1995–1997. These data were taken as part of a spectral line survey program, and have a lower spectral resolution and signal to noise ratio than the IRAM spectra. The facility receivers A2 and C2 were used as front ends and the Dutch Autocorrelation Spectrometer (DAS) as back end. The JCMT has a main beam efficiency of 0.65 at 1.3 mm and 0.53 at 0.6 mm wavelength. Integration times are 30 min at 241 GHz and 40 min at 464 GHz, resulting in rms noise levels in T_{MB} per 625 kHz channel of ≈ 40 mK at 241 GHz and ≈ 200 mK at 464 GHz.

All single-dish spectra have been reduced with the CLASS package, developed at IRAM⁴. Linear baselines were subtracted and the spectra were smoothed once and calibrated onto T_{MB} scale. We estimate a calibration uncertainty of 30% for the 80 GHz IRAM and 240 GHz JCMT data, and of 50% for the 230 GHz IRAM and 460 GHz JCMT data, due to higher atmospheric opacity. The estimated pointing uncertainty for all single-dish data is 3'' rms.

2.2. Interferometric observations

The IRAM interferometer on Plateau de Bure (France) consists of six 15-m antennas on North-South and East-West baselines. We used this instrument to map the HDO 1₁₀–1₁₁, H₂¹⁸O 3₁₃–2₂₀ and SO₂ 12_{0,12}–11_{1,11} lines and continuum at 80.6 and 204.9 GHz toward AFGL 2591. Due to tuning problems, only five antennas took 1.3 mm data in *C*-configuration on December 6, 2003; these problems were solved before the *D*-array observations on May 15–16, 2004. The correlator was configured to produce “narrow” 80 MHz and “broad” 160 MHz windows, with one of each centered on the HDO and H₂¹⁸O lines. The number of channels is 128 per window. The SO₂ line falls in the 160 MHz window of the H₂¹⁸O line. The continuum bandwidth is 640 MHz at 81 GHz and twice as much at 205 GHz, where tuning is double side band. Antenna gains and phases were monitored by observing 2013+370 and 2005+403 for 2 min every 20 min. The total observing time was 13.2 h of good weather in *C*-array and 10 h of excellent weather in *D*-array. The combined dataset has baselines ranging from the antenna shadowing limit out to 309 m, corresponding to an angular scale of 2.2'' at 81 GHz and 0.85'' at 205 GHz. Data reduction was performed at the IRAM headquarters in Grenoble,

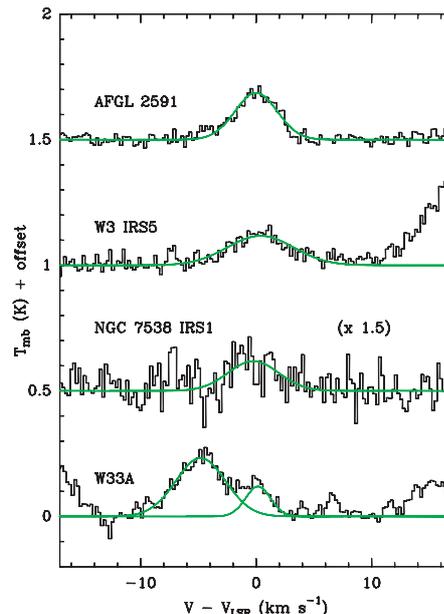


Fig. 1. Spectra of the H₂¹⁸O 3₁₃–2₂₀ transition, obtained with the IRAM 30 m telescope. The line at $V = +16$ km s⁻¹ is the SO₂ 12_{0,12}–11_{1,11} transition; the lines in the W33A spectrum at $V = -5$ and -17 km s⁻¹ are due to CH₃OCH₃.

using the GILDAS software. Bandpass was checked on 3C273 and 2145+067. Flux calibration was performed on MWC 349, assuming $S_\nu = 1.0$ Jy at 81 GHz and 2.0 Jy at 205 GHz.

3. Results

3.1. Single-dish spectra

With the IRAM 30 m, we have detected H₂¹⁸O in four objects (Fig. 1). The HDO 80 GHz line is detected in the same four objects, and the 225 GHz line in three of them. The CH₃OH line is seen in all eight sources, and the SO₂ line in all but one. Tables 3 and 4 list the integrated intensities and widths of the lines, obtained through Gaussian fits to the profiles. Note that for strong lines, calibration dominates the uncertainty on the line flux, while spectral noise dominates for weak lines. The spectra also show a few unexpected lines. In W33A, the 4₀₄–3₀₃ line of formamide (NH₂CHO) at 84542.4 MHz was detected with $T_{\text{MB}} = 42$ mK and $\Delta V = 5.6$ km s⁻¹. Next to the H₂¹⁸O line, the 3₃₀–2₂₁, 3₃₁–2₂₁ and 3₃₀–2₂₀ lines of dimethyl ether (CH₃OCH₃) at 203 420, 203 410 and 203 384 MHz are detected with $T_{\text{MB}} = 0.21, 0.23$ and 0.25 K and $\Delta V = 4.6, 4.4$ and 4.9 km s⁻¹. The 3₃₀–2₂₀ line of CH₃OCH₃ is also detected toward NGC 7538 IRS1, with $T_{\text{MB}} = 0.12$ K and $\Delta V = 3.7$ km s⁻¹.

With the JCMT, we have detected three HDO lines in AFGL 2591, and one in NGC 7538 IRS1 (Fig. 2; Table 5). Upper limits (1σ) of $T_{\text{MB}} = 0.24, 0.35$ and 0.17 K on 0.625 MHz channels were obtained for the 464 GHz line toward W33A, AFGL 2136 and S140 IRS1. For the 225 GHz line, upper limits of $T_{\text{MB}} = 38$ mK were found for S140 IRS1 and NGC 7538 IRS9. Helmich et al. (1996) report a tentative detection of the 464 GHz line toward W3 IRS5 and upper limits

³ The JCMT is operated by the Joint Astronomy Centre, on behalf of the Particle Physics and Astronomy Research Council of the UK, The Netherlands Organization for Scientific Research and the National Research Council of Canada.

⁴ <http://www.iram.fr/IRAMFR/GILDAS/>

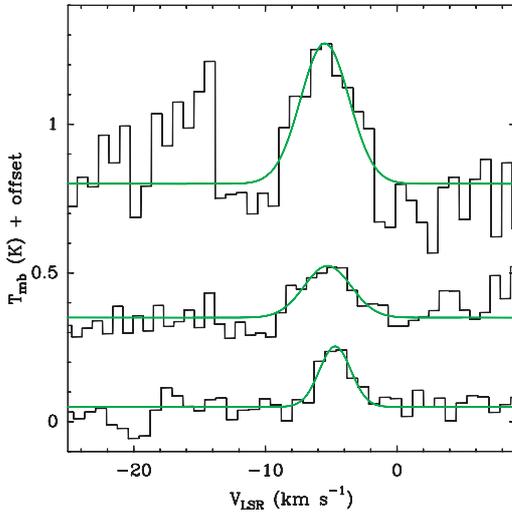
Table 3. Line fluxes (K km s^{-1}) or 1σ upper limits (mK) observed with the IRAM 30 m. Numbers in brackets are uncertainties in units of the last decimal place. The errors do not take calibration into account but only spectral noise.

Source	HDO $1_{10}-1_{11}$	CH ₃ OH $5_{-1}-4_0$ E	H ₂ ¹⁸ O $3_{13}-2_{20}$	SO ₂ $12_{0,12}-11_{1,11}$	HDO $3_{12}-2_{21}$
W3 IRS5	0.14(2)	0.17(2)	0.84(7)	19.44(8)	<36
AFGL 490	<13	0.39(2)	<24	<24	<31
W33A	0.66(3)	7.69(3)	0.46(2)	4.73(2)	4.06(6)
AFGL 2136	<9	0.50(1)	<21	0.85(3)	<27
AFGL 2591	0.15(1)	1.51(1)	0.86(3)	4.01(3)	0.59(3)
S140 IRS1	<10	1.41(1)	<22	2.26(3)	<23
NGC 7538 IRS1	0.26(3)	2.64(2)	0.43(9)	2.09(8)	1.57(8)
NGC 7538 IRS9	<12	2.07(2)	<22	0.62(4)	<29

Table 4. Widths ($FWHM$ in km s^{-1}) of the lines observed with the IRAM 30 m. Numbers in brackets are uncertainties in units of the last decimal place.

Source	HDO $1_{10}-1_{11}$	CH ₃ OH $5_{-1}-4_0$ E	H ₂ ¹⁸ O $3_{13}-2_{20}$	SO ₂ $12_{0,12}-11_{1,11}$	HDO $3_{12}-2_{21}$
W3 IRS5	3.2(5)	2.7(4)	6.8(10)	6.67(4) ^a	...
AFGL 490	...	2.5(2) ^a
W33A	4.6(2)	4.67(2) ^a	3.8(1)	6.1(1) ^a	5.06(9) ^a
AFGL 2136	...	2.79(7) ^a	...	4.2(2)	...
AFGL 2591	3.3(4)	2.87(3) ^a	4.3(2)	5.33(5)	3.2(2)
S140 IRS1	...	2.68(3) ^a	...	2.67(4) ^a	...
NGC 7538 IRS1	3.6(5)	3.21(3) ^a	5.2(16)	6.1(3)	3.8(2)
NGC 7538 IRS9	...	2.94(3) ^a	...	5.3(4)	...

^a Line core only; also wings visible.

**Fig. 2.** Observations of HDO lines toward AFGL 2591 with the JCMT. Top to bottom: 465 GHz, 241 GHz, 225 GHz lines. The bottom two spectra have been multiplied by 2 and all spectra are vertically offset for clarity. The feature at $V_{\text{LSR}} = -15 \text{ km s}^{-1}$ in the upper spectrum is due to the CH₃OH $18_{8,10}-18_{9,10}$ E line.

on the 225 and 241 GHz lines; Schreyer et al. (2002) set an upper limit to the 464 GHz emission from AFGL 490. The JCMT spectra of HDO have a lower signal to noise ratio and spectral resolution than the IRAM 30 m spectra, and the line positions and widths in Table 5 are too uncertain to extract kinematic information.

Table 5. JCMT observations of HDO.

Line	$\int T_{\text{MB}} dV$ K km s^{-1}	V_{LSR} km s^{-1}	ΔV km s^{-1}	T_{MB} K
AFGL 2591				
$3_{12}-2_{21}$	0.308(48)	-4.70(22)	2.83(52)	0.10
$2_{11}-2_{12}$	0.394(80)	-5.27(45)	4.24(88)	0.09
$1_{01}-0_{00}$	2.19(41)	-5.48(43)	4.35(81)	0.47
NGC 7538 IRS1				
$3_{12}-2_{21}$	0.74(12)	-58.43(28)	3.57(62)	0.20

The central velocities and the widths of the HDO lines in the 30 m spectra (Table 4) are consistent with the values for the molecular envelopes of these objects, derived from C¹⁷O and C³⁴S spectra (Van der Tak et al. 2000b). In contrast, the H₂¹⁸O lines are 30–90% broader than the HDO lines in the same sources (Table 4), which may be an indication that part of the H₂¹⁸O emission arises in the molecular outflows of these sources. Only in W33A, the fitted width of the H₂¹⁸O line is smaller than that of the HDO lines, but for this source, the H₂¹⁸O line is blended with other lines (see Fig. 1), making its width hard to measure.

Evidence for a contribution to the observed emission by outflows is even more pronounced in the SO₂ and CH₃OH spectra, which have higher signal-to-noise than those of HDO and H₂¹⁸O. The profiles of SO₂ in three sources and of CH₃OH in seven show low-level emission at high velocities (Table 4). We have fitted these profiles with the sum of two independent Gaussians: a narrow one corresponding to the “envelope”

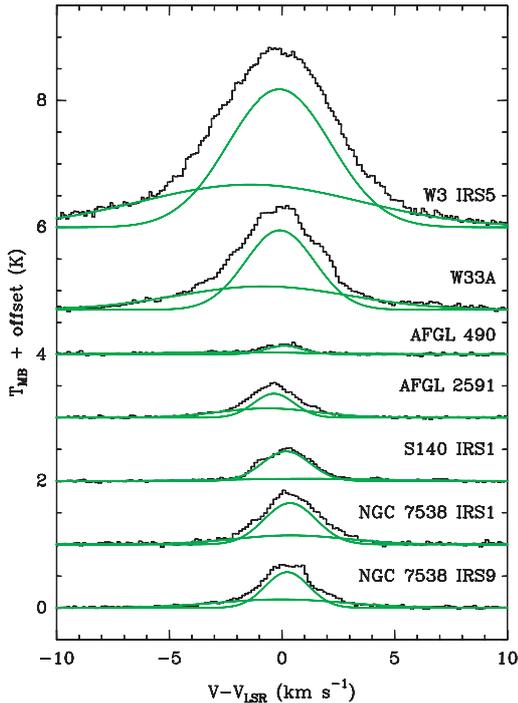


Fig. 3. Line profiles of CH₃OH, observed with the IRAM 30 m, with two-component fits overlotted. For W3 IRS5, the SO₂ line is plotted instead of the CH₃OH line.

component seen in C¹⁷O and C³⁴S, and a broader one (Fig. 3) which we attribute to the outflow. We find widths for the broad components between 4.7 km s⁻¹ in AFGL 2591 and 12.8 km s⁻¹ in W3 IRS5; the source-averaged width is 7.1 km s⁻¹. The broad component is blueshifted from the narrow one in all sources except S140 IRS1, which ties in with the tendency for blueshifted mid-infrared absorption of CO (Mitchell et al. 1991), which also arises in outflows. The fraction of the line flux carried by the broad CH₃OH component ranges from 16% in S140 IRS1 to 58% in AFGL 2591, and is 39% on average. These fractions are comparable with the values of 10–50% found for CS, SO and SO₂ in these sources (Van der Tak et al. 2003) which may depend on excitation energy or beam size. Previous CH₃OH spectra of these sources which had lower signal to noise ratios and lower spectral resolution did not show line wings, although they did show wings for two other, similar objects (Van der Tak et al. 2000a).

The line profiles of H₂¹⁸O and HDO do not have high enough signal to noise ratios to perform two-component fits. Therefore we assume that the outflow contribution estimated for CS, SO, SO₂ and CH₃OH also holds for H₂¹⁸O and HDO, which have similar excitation requirements and optical depths. However, instead of a double Gaussian fit, the outflow contribution may be estimated as the line flux at high velocities only. The result is somewhat less than half of that from a double Gaussian fit, or ~10–20%. This value is probably consistent with the fraction of ~50% estimated for the o-H₂¹⁶O ground state line observed with SWAS (Snell et al. 2000; Boonman et al. 2003), given the different excitation requirements.

The presence of significant amounts of CH₃OH in the outflow raises the question of its desorption from grain mantles.

Toward young low-mass protostars, the methanol emission of transitions with low energy levels is dominated by outflow gas (Bachiller et al. 2001), indicating that grain mantle desorption by shocks is at work. Higher excitation methanol lines are broad toward some low-mass protostars and narrow towards others, indicating that shocks or thermal desorption may dominate in the warmer regions of the envelope (Maret et al. 2005; Jørgensen et al. 2005). The present data show that for high-mass protostars, radiation and shocks both have a relevant role in releasing ice mantles from dust grains too.

3.2. Excitation of HDO

For the sources where several HDO lines are detected, we have estimated the excitation temperature using rotation diagrams. This method, described in detail by Blake et al. (1987) and Helmich et al. (1994), assumes that the lines are optically thin and describes the molecular excitation by a single temperature, the “rotation temperature”. However, this temperature is only meaningful if all the data refer to the same physical volume. The beam sizes of our observations range from 12 to 30″, and it is important to consider the effect of beam dilution. Indeed, if the HDO emission from W33A, AFGL 2591 and NGC 7538 IRS1 were extended on scales as large as 30″, the upper-state column densities of the higher-excitation lines would be larger than those of the low-excitation lines, implying an infinite or negative excitation temperature. Since such non-thermal excitation is unlikely, the data must be corrected for the effect of a finite source size.

The size of the HDO emission can be estimated for AFGL 2591 and NGC 7538 IRS1, where the 225 GHz line has been measured both with the IRAM 30 m and the JCMT. The emission is about twice as bright in the IRAM 30 m beam, suggesting a compact source size. Therefore we assume a source size of 12″ for the HDO in W33A, AFGL 2591 and NGC 7538 IRS1, and correct the 80 GHz line fluxes upward by (30/12)². This factor is much larger than any plausible optical depth effect on the 80 GHz line, especially since the 464 GHz line, which lies lower in excitation, is expected to have a larger optical depth. Statistical equilibrium calculations indeed indicate optical depths of ~10⁻² for the excitation temperatures and brightness levels of HDO found here.

Table 6 lists the assumed sizes for all sources where HDO has been detected, and the resulting HDO excitation temperatures. The data do not rule out source sizes <12″, and the assumed size may be regarded as an upper limit. Smaller source sizes would influence all lines equally, though, and not change the excitation temperature estimates. In the case of W3 IRS5, the only firm detection of HDO is the 80 GHz line. The observational limits on the 225, 241 and 464 GHz lines indicate $T_{\text{ex}} = 85\text{--}115$ K, but do not constrain the source size.

The excitation temperatures found for HDO may be used as a first clue to its chemical origin by comparison with SO₂ (a product of shock chemistry), CH₃OH (a product of ice evaporation), and C₂H₂ (a product of hot gas-phase reactions). The excitation temperatures of HDO are similar to the values of 50–200 K derived for CH₃OH

Table 6. Sizes and excitation temperatures of the HDO emission, derived from the single-dish observations.

Source	Size "	T_{ex} K
W3 IRS5	...	85–115
W33A	12	110 ± 58
AFGL 2591	12	117 ± 57
NGC 7538 IRS1	12	108 ± 56

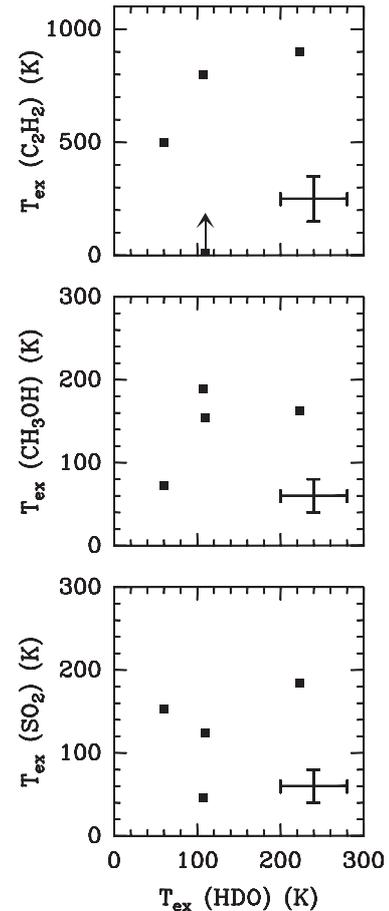
(Van der Tak et al. 2000a) and SO_2 (Van der Tak et al. 2003), measured in (sub-)millimeter emission in 14–18'' beams. These excitation temperatures are lower limits to the kinetic temperature of the emitting gas, but this probably has little effect on the comparison with HDO since the molecules have similar dipole moments. The excitation temperatures are much lower than the values of 500–1000 K measured in mid-infrared absorption of C_2H_2 (Lahuis & van Dishoeck 2000) as expected for two reasons. First, the C_2H_2 molecule does not have a permanent dipole moment, and its excitation temperature directly reflects the kinetic temperature of its surroundings. Second, C_2H_2 is seen in absorption, which in these centrally condensed envelopes preferentially probes smaller radii, where the temperature is higher. The excitation temperatures of HCN 14 μm and H_2O 6.2 μm absorption of 300–500 K for these sources (Boonman & van Dishoeck 2003) are between the values for HDO emission and C_2H_2 absorption, suggesting that the two effects contribute about equally.

The HDO, CH_3OH , SO_2 and C_2H_2 in these objects are thus located in warm (several 100 K) gas, such as found in the inner envelopes and outflow shocks of protostars. The exact temperature is hard to estimate because of the above caveats, so that we have searched for trends instead. Even here, the data appear inconclusive (Fig. 4). Further discussion of the origin of the HDO is deferred until after the radiative transfer model in Sect. 5.

3.3. Column densities of HDO and H_2O

Knowing the excitation conditions, we derive molecular column densities from the observed line strengths. For HDO, these follow directly from the rotation diagrams. For H_2O , these come from the H_2^{18}O data, assuming optically thin H_2^{18}O emission with the same excitation temperature as the HDO in that source. We use an oxygen isotope ratio of $^{16}\text{O}/^{18}\text{O} = 500$ and assume an ortho/para ratio of 3 for H_2O , as expected for warm gas. The resulting column densities (Table 7) are uncertain by a factor of ≈ 2 , mainly through the uncertain excitation temperature. The sensitivity of $N(\text{H}_2\text{O})$ to T_{ex} is such that increasing T_{ex} from 110 to 220 K increases the derived column density just slightly, while decreasing T_{ex} from 110 to 60 K almost doubles it and further lowering T_{ex} leads to implausibly high H_2O column densities (Fig. 5). In addition, if the source size is smaller than the 12'' assumed above, the column density estimates would increase.

For the four sources where HDO and H_2^{18}O were not detected, the noise levels of the spectra imply limits on the

**Fig. 4.** Excitation temperatures of HDO plotted versus values of SO_2 (bottom), CH_3OH (middle) and C_2H_2 (top). The low T_{ex} of C_2H_2 in W33A may be affected by a high continuum optical depth at 14 μm .

column densities. Assuming $T_{\text{ex}} = 110$ K and $\Delta V = 3.0$ km s^{-1} , the 3σ limits in an 11'' beam are $N(\text{HDO}) < 6.5 \times 10^{13}$ cm^{-2} and $N(\text{H}_2\text{O}) < 6.4 \times 10^{16}$ cm^{-2} , which for both species is a factor of ≈ 2 below the weakest detection. In the cases of AFGL 490, S140 IRS1 and NGC 7538 IRS9, the non-detection of HDO and H_2^{18}O emission ties in with non-detections of H_2O 6.2 μm absorption (Boonman & van Dishoeck 2003; Schreyer et al. 2002) and low column densities of warm H_2 as traced by ^{13}CO 4.7 μm absorption (Mitchell et al. 1990, 1995). In contrast, AFGL 2136 does have high H_2O and H_2 column densities measured in mid-infrared absorption, but this gas must be compact, since the H_2 column density from (sub-)millimeter data is low (Van der Tak et al. 2003).

The derived molecular column densities may be used as a clue to the source geometry by comparing them with values measured in mid-infrared absorption. The H_2O column densities in Table 7 are consistent with the values from ISO 6.2 μm absorption (Boonman & van Dishoeck 2003) to within factors of a few. This situation is similar to that for CO and dust, where (sub-)millimeter data indicate column densities ~ 3 – $5\times$ higher than mid-infrared data (Van der Tak et al. 2000b). Therefore, if these species are spherically distributed around the central star, this region must be extended on the scales of the single-dish beams (diameter $\gtrsim 1''$, corresponding to $\gtrsim 2000$ AU at

Table 7. Column densities of HDO and H₂O in an 11'' beam, derived from the single-dish observations.

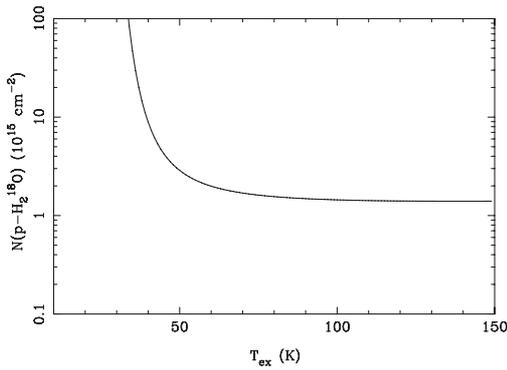
Source	HDO 10 ¹⁴ cm ⁻²	H ₂ O 10 ¹⁷ cm ⁻²	H ₂ O ^a 10 ¹⁷ cm ⁻²	\bar{T} ^d K	HDO/H ₂ O 10 ⁻⁴
W3 IRS5	0.3–0.6 ^b	2.6–5.0 ^c	3	33	1
W33A	6.9	1.4	<8	20	70
AFGL 2591	2.3	3.0	4	28	8
NGC 7538 IRS1	2.7	1.3	<5	25	30

^a From ISO 6.2 μ m absorption (Boonman & van Dishoeck 2003) in a pencil beam.

^b Values for source sizes of 30'' and 12''.

^c Values for $T_{\text{ex}} = 110$ and 60 K.

^d Mass-weighted envelope temperature from Van der Tak et al. (2000b).

**Fig. 5.** Sensitivity of the derived $p\text{-H}_2^{18}\text{O}$ column density to the adopted excitation temperature, for a line flux of 1.0 K km s⁻¹.

$d = 2$ kpc), Alternatively, these molecules do not have spherically symmetric distributions. In contrast, for SO₂ and HCN, mid-infrared absorption lines indicate $\sim 100\times$ higher column densities than (sub-)millimeter emission lines (Keane et al. 2001; Boonman et al. 2001). These molecules must have distributions as compact as $\lesssim 0.1$ ($\lesssim 200$ AU) which may or may not be spherical. Additional constraints on the source geometry come from the interferometer data (Sect. 3.5).

The $N(\text{HDO})/N(\text{H}_2\text{O})$ ratios (Table 7, Col. 6) are consistent with the limits on solid-state HDO/H₂O obtained for these same sources (Dartois et al. 2003), and similar to the values measured for “hot core” type regions (Jacq et al. 1990; Gensheimer et al. 1996). The ratios correspond to enhancements of the HDO/H₂O ratio over the elemental abundance ratio ($\text{D}/\text{H} = 1.5 \times 10^{-5}$; Linsky 1998) of 7 for W3 IRS5, 50 for AFGL 2591, 200 for NGC 7538 IRS1 and 470 for W33A. The enhancement level shows a correlation with the mass-weighted average envelope temperature \bar{T} of these sources (Van der Tak et al. 2000b), listed in the fifth column of Table 7, in the sense that colder sources have higher HDO/H₂O ratios. The decrease of the HDO enhancement from 470 in W33A ($\bar{T} = 20$ K) to 7 in W3 IRS5 ($\bar{T} = 33$ K) is the combined result of a decrease in $N(\text{HDO})$ by a factor of 16 and an increase in $N(\text{H}_2\text{O})$ by a factor of 4. The HDO/H₂O ratios in Table 7 are 10–100 \times higher than the equilibrium value at the estimated gas temperatures (\sim few 100 K), which suggests that the HDO (and H₂O) molecules are a remnant from an earlier, colder phase of the protostars (like in the case of low-mass protostars; see Sect. 5.3). A natural explanation is that HDO (and H₂O)

molecules are sublimated from the grain mantles in the warm region where the dust temperature exceeds 100 K, the ice sublimation temperature. In this scenario, the coldest sources would be also the youngest, where gas-phase reactions occurring in the warm region containing the sublimated ices have had less time to bring the HDO/H₂O ratio back down to the equilibrium ratio at ≥ 100 K.

3.4. Interferometric continuum images

Figure 6 presents continuum maps of AFGL 2591, made by gridding and Fourier transforming the IRAM interferometer data and deconvolving with the Clean algorithm. Using uniform weight, the size ($FWHM$) of the synthesized beam is $(1.33 \times 1.07)''$ at PA = 47° at 205 GHz and $(4.60 \times 3.38)''$ at PA = 74° at 81 GHz. The rms noise level of the maps is 0.34 mJy/beam at 81 GHz and 1.3 mJy/beam at 205 GHz.

The 205 GHz map shows three sources, and Table 8 lists their properties, derived by fitting two-dimensional Gaussians to the u, v data. The strongest, Eastern source (VLA3) coincides with the “dust peak” and the infrared source AFGL 2591. The South-Western source VLA1 is a compact H II region which dominates the field at frequencies $\lesssim 100$ GHz. The weakest, North-Western source coincides with feature “VLA2” in low-frequency (< 10 GHz) VLA maps. Our 81 GHz map (Fig. 6, bottom) shows the same three features, but their emission is blended due to the lower angular resolution. These results are consistent with previous mapping at similar frequencies with the OVRO interferometer (Van der Tak et al. 1999). However, the Bure data have higher sensitivity, making these the first detections of the NW and SW H II regions at frequencies > 200 GHz. On the other hand, the OVRO data have higher resolution, so that the infrared source and the H II regions are well separated at 87 GHz. Source VLA2 is visible in the OVRO 87 GHz data, but not firmly detected.

To study the physical nature of the continuum sources, we calculate their millimeter-wave spectral indices γ , defined as $S_\nu \propto \nu^\gamma$. To do so we combine the 87 GHz data from OVRO with the 205 GHz data from IRAM. The uv coverage of these telescopes is similar, and the effect of “missing flux” on their comparison is probably small. The result is $\gamma = 0.0 \pm 0.05$ for VLA1, $\gamma = +2.1 \pm 0.2$ for VLA2 and $\gamma = +2.7 \pm 0.1$ for VLA3. These values indicate optically thin free-free emission for VLA1, optically thick dust or free-free emission in VLA2,

Table 8. Positions, deconvolved sizes, and strengths of continuum sources in the AFGL 2591 region, detected with the interferometer. Numbers in brackets denote uncertainties in units of the last decimal.

Component ^a	RA (J2000) hh mm ss	Dec (J2000) dd mm ss	Major axis arcsec	Minor axis arcsec	Pos. angle deg.	Flux density Jy
<i>81 GHz</i>						
VLA1	20:29:24.5507(10)	40:11:15.250(10)	2.64(4)	2.36(6)	-22(8)	0.061(1)
VLA2	20:29:24.4795(50)	40:11:22.460(46)	0.009(1)
VLA3	20:29:24.8917(30)	40:11:19.687(28)	0.016(2)
<i>205 GHz</i>						
VLA1	20:29:24.5721(28)	40:11:14.669(29)	2.13(9)	1.29(9)	90(4)	0.061(2)
VLA2	20:29:24.4394(49)	40:11:23.505(51)	3.24(17)	2.21(16)	90(5)	0.065(4)
VLA3	20:29:24.8694(3)	40:11:19.498(5)	1.08(2)	0.85(1)	00(3)	0.194(1)

^a Nomenclature from Trinidad et al. (2003).

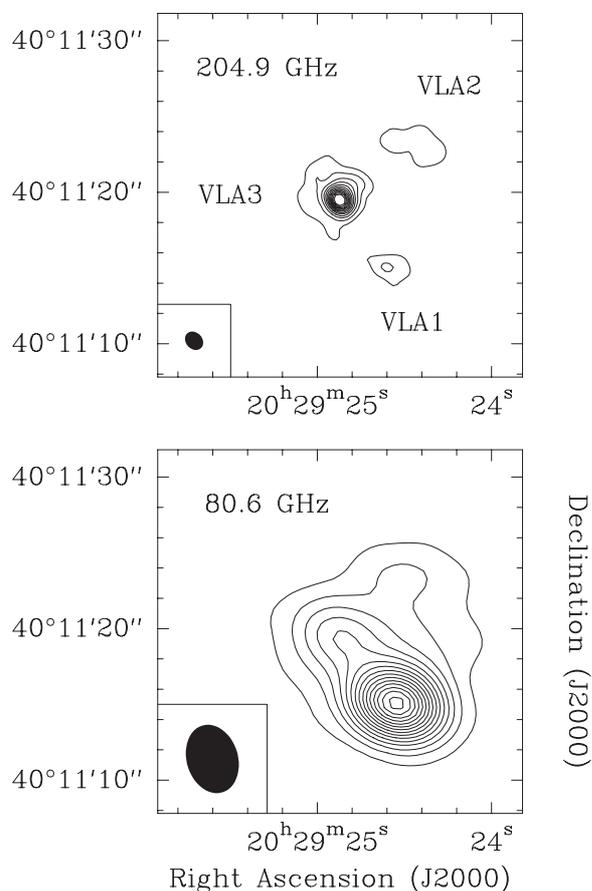


Fig. 6. Maps of the continuum emission of AFGL 2591 made with the IRAM interferometer. Contours are drawn every 8 mJy/beam at 205 GHz (*top*) and every 3 mJy/beam at 81 GHz (*bottom*). Source nomenclature and synthesized beam size are indicated.

and optically thin dust emission in VLA3. The brightness temperatures of the 205 GHz sources are 0.7 K for VLA1, 0.3 K for VLA2 and 6.1 K for VLA3. These values are much lower than the expected physical temperatures of either ionized gas or dust, and indicate either a low optical depth or a small filling factor.

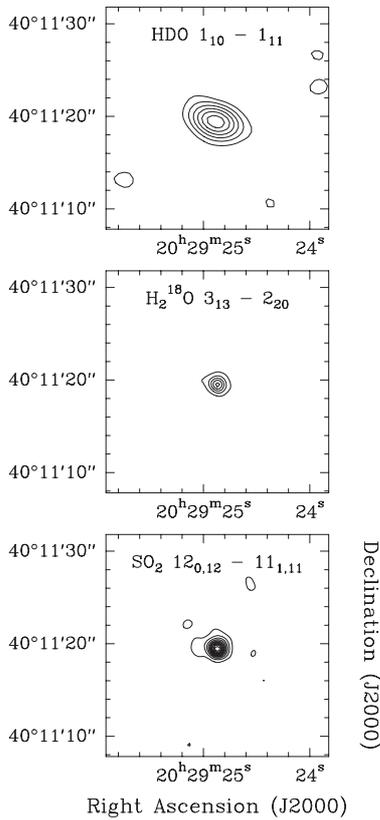
For the “dust” source VLA3, we calculate the mass from the observed 205 GHz flux density, assuming a dust temperature of 100 K, a standard gas to dust ratio of 100 and a mass opacity of 0.9 cm² per gram of dust (Ossenkopf & Henning 1994; Henning et al. 2000). The result is 0.8 M_{\odot} , which is inversely proportional to the assumed temperature. Interestingly, the spectral index of VLA3 indicates a value of the dust opacity index of $\beta \approx 1$, which is smaller than the “usual” value of ≈ 2 and suggests grain growth. This process is thought to occur in circumstellar disks, which is not inconsistent with the observed elongated shape of VLA3. The very compact 43 GHz emission from ionized gas seen by Van der Tak & Menten (2005) may then originate in the ionized inner part of the disk, a disk atmosphere, or a disk wind (Hollenbach et al. 1994; Lugo et al. 2004).

3.5. Interferometer observations of H_2O , HDO and SO_2 line emission

Figure 7 shows the maps of the HDO , $H_2^{18}O$ and SO_2 line emission observed with Bure. The beam sizes are the same as those of the continuum maps at that frequency. The rms noise levels of the line maps are 9 mJy/beam for HDO , 23 mJy/beam for $H_2^{18}O$, and 50 mJy/beam for SO_2 . In the case of SO_2 , the noise is limited by dynamic range problems. The line maps show compact emission, coincident with the “dust peak” of AFGL 2591. Columns 2–6 of Table 9 list the position, size and shape of the emission, obtained by fitting 2D Gaussians to the u, v data. Figure 8 shows spectra of the line emission, taken at the peak positions of the images. Columns 7–9 of Table 9 list the central positions, widths, and peak strengths of the lines, obtained by fitting Gaussian profiles to the spectra at the image maxima. The central velocities of the lines are consistent with the values measured at the 30-m telescope (Table 4). The width of the SO_2 line is consistent with the single-dish value, while the $H_2^{18}O$ line is 23% broader and the HDO line 50% broader. Within the errors, all of the single-dish flux is recovered by the interferometer: apparently, both telescopes trace the same gas, which has a compact ($\lesssim 1''$) distribution. Given the constraints on the source geometry from the comparison of single-dish and

Table 9. Positions, deconvolved sizes, velocities, and strengths of emission lines detected toward AFGL 2591 with the interferometer. Numbers in brackets denote uncertainties in units of the last decimal.

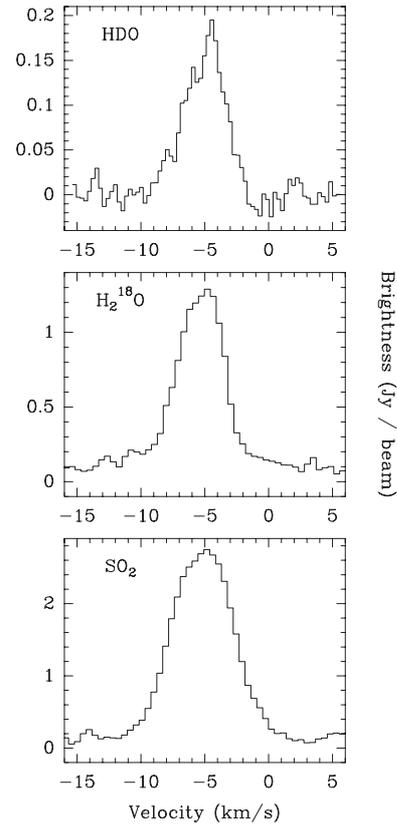
Line	RA (J2000) hh mm ss	Dec (J2000) dd mm ss	Major axis arcsec	Minor axis arcsec	Pos. angle deg	V_{LSR} km s ⁻¹	ΔV km s ⁻¹	Peak T_B K
HDO	20:29:24.872(11)	40:11:19.48(10)	≤ 1	≤ 1	...	-4.92(5)	3.63(10)	2.05(5)
H ₂ ¹⁸ O	20:29:24.8691(6)	40:11:19.493(8)	0.81(2)	0.71(3)	-47(12)	-5.29(3)	4.12(7)	24.26(36)
SO ₂	20:29:24.8688(4)	40:11:19.363(6)	1.00(2)	0.91(2)	-28(9)	-5.16(2)	5.44(5)	53.93(40)

**Fig. 7.** Interferometric maps of line emission toward AFGL 2591 at the central velocity. For HDO (*top*), first contour and contour step are 30 mJy/beam. For H₂¹⁸O (*middle*) and SO₂ (*bottom*), first contour is 0.15 Jy/beam and contour step is 0.3 Jy/beam.

infrared column density estimates (Sect. 3.3), we conclude that the H₂¹⁸O and HDO emitting regions have sizes of $\sim 1''$.

3.6. Velocity structure of the line emission

Channel maps of H₂¹⁸O, HDO and SO₂ do not show clear changes of the emitting structure with velocity, but the signal to noise ratio of the Bure line data is high enough to locate the emission peak to a fraction of a synthesized beam width. Therefore, to study the velocity structure of the compact molecular gas, we have determined the position of the emission peak in each channel by a fit to the (u, v) data, and Fig. 9 shows the result. The H₂¹⁸O and SO₂ lines show a clear velocity gradient, where the redshifted gas is located in the South-West and the blueshifted gas in the North-East. The HDO shows the same trend, but not as clearly due to the lower angular resolution. In each case, the gradient is smooth, which suggests that the

**Fig. 8.** Interferometric spectra of line emission toward AFGL 2591 taken at the image maxima.

emission traces a physically coherent structure. Since the central position of the line emission coincides with that of the dust continuum peak, and its central velocity with that of the large-scale molecular envelope, the most plausible origins of the velocity gradient are outflowing motions in a bipolar cavity, or rotation in a circumstellar disk.

One central prediction of disk accretion models of low-mass star formation is that the outflow axis is perpendicular to the disk plane. The orientation of the structure seen in Fig. 9 does not agree with that of the large-scale outflow which is known to emanate from AFGL 2591. The position angle of the velocity gradient is 39° in H₂¹⁸O and 67° in SO₂ (measured East from North). The value for HDO is 13° , but this number is uncertain due to the large scatter in the data points, and not inconsistent with the value for H₂¹⁸O. In contrast, single-dish CO 3–2 and HCO⁺ 4–3 mapping (Hasegawa & Mitchell 1995b) shows an outflow of size $90 \times 20''$, embedded in an arcminute-scale outflow seen before in lower- J CO lines (e.g., Lada et al. 1984). The East-West orientation agrees with the

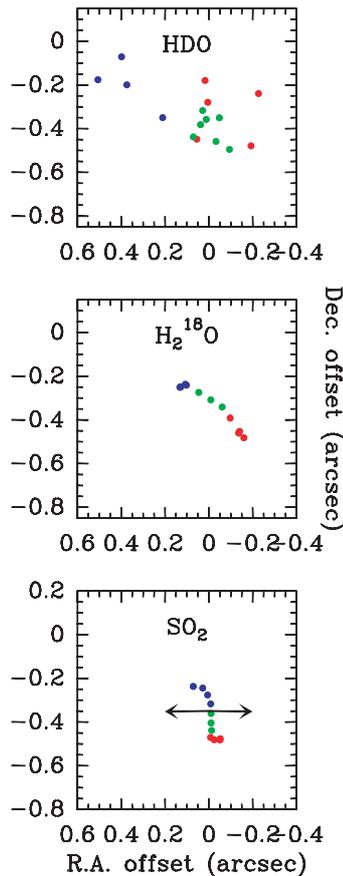


Fig. 9. Position of the emission peak in AFGL 2591 in each channel, as offset from the phase center of the Bure observations, for HDO (*top*), H_2^{18}O (*middle*) and SO_2 (*bottom*). Colour coding corresponds to velocity offset: red = redshifted ($-V_{\text{LSR}} = 2.5\text{--}4\text{ km s}^{-1}$), green = line center ($-V_{\text{LSR}} = 4.5\text{--}6\text{ km s}^{-1}$), blue = blueshifted ($-V_{\text{LSR}} = 6.5\text{--}8\text{ km s}^{-1}$). The arrow in the bottom panel indicates the orientation of the large-scale CO outflow.

positions of Herbig-Haro objects and spots of shock-excited H_2 emission (Tamura & Yamashita 1992; Poetzel et al. 1992). On smaller scales, VLBI observations of H_2O masers by Trinidad et al. (2003) show a shell-like structure which they interpret as an outflow cavity. The spots are spread over $0.02''$, in an elongated structure oriented about 20° West from North.

On the other hand, if the line emission seen with Bure is due to a disk, the inclination is probably close to face-on. For example, the Gaussian fits presented in Table 9 imply axis ratios of 0.90 ± 0.02 for H_2^{18}O and SO_2 , corresponding to an inclination of $(26 \pm 3)^\circ$. A face-on disk is also consistent with the much higher outflow velocities seen in CO mid-IR absorption than in mm emission. The outflow is thus directed almost along the line of sight.

The total magnitude of the velocity gradient is 4.6 km s^{-1} over an offset of $0.3''$. Adopting a distance of 1 kpc and assuming the above inclination angle, the implied rotation period is $\sim 1000\text{ yr}$. The central star has a mass of $\approx 16 M_\odot$, based on the luminosity of the region (Van der Tak & Menten 2005). For this orbital period and stellar mass, Kepler’s third law implies an orbital radius of 250 AU, which is within a factor of 2 from the measured value. This agreement does not prove that the line

emission originates in a rotating disk, but it does indicate that the gas motion is controlled by stellar gravity. Higher resolution observations are necessary to resolve the velocity field of the compact line emission.

4. Radiative transfer analysis

Dividing the single-dish column densities in Table 7 by the $N(\text{H}_2)$ values based on (sub-)millimeter observations of dust continuum and C^{17}O lines in $\approx 15''$ beams (Table 1) leads to beam-averaged abundances of $2 \times 10^{-10}\text{--}1 \times 10^{-9}$ for HDO and $2 \times 10^{-7}\text{--}2 \times 10^{-6}$ for H_2O . On the other hand, the interferometer data of AFGL 2591 indicate a H_2O column density of $3.4 \times 10^{19}\text{ cm}^{-2}$, assuming optically thin emission with $T_{\text{ex}} = 100\text{ K}$, while the 205 GHz continuum indicates $N(\text{H}_2) = 2.1 \times 10^{24}\text{ cm}^{-2}$ for a dust temperature of 100 K. The $\text{H}_2\text{O}/\text{H}_2$ column density ratio of 1.6×10^{-5} is $\approx 10 \times$ higher than estimated from the single-dish data. To resolve this discrepancy, we have run radiative transfer models of the line emission.

4.1. Model description

The column density ratios evaluated in the previous section are crude estimates of the molecular abundances because they do not take excitation gradients along the line of sight into account. In protostellar envelopes and disks, where density and temperature vary over orders of magnitude, these gradients are very significant. Furthermore, column densities and their ratios do not give insight into the location of the gas along the line of sight. To estimate more accurate abundances, the line emission of H_2^{18}O and HDO in our sources has been modeled with the Monte Carlo radiative transfer program of Hogerheijde & van der Tak (2000)⁵. The application of this program to H_2O has been explicitly tested in a dedicated benchmark campaign (Van der Tak et al. 2005)⁶. Spectroscopic and collisional input for the modeling comes from the molecular database by Schöier et al. (2005)⁷. Besides collisional excitation, dust radiation is taken into account using grain properties from Ossenkopf & Henning (1994), Model 5. Radial profiles of the temperature and density of the envelopes of our sources were determined by Van der Tak et al. (2000b) based on single-dish maps of dust continuum and molecular line emission at (sub-)millimeter wavelengths.

4.2. Constant-abundance models

We first assume that H_2O and HDO are distributed evenly over the protostellar envelopes. Table 10 reports the results of these models. Stringent convergence criteria had to be applied, as the upper (emitting) level of the H_2^{18}O line has a relative population of only $\sim 10^{-7}$ in the outer parts of the envelopes. Nevertheless, Fig. 10 shows that the calculations are well converged. We calculate an optical depth for the H_2^{18}O line of $\sim 10^{-2}$. Maser

⁵ <http://www.mpifr-bonn.mpg.de/staff/fvandertak/ratran/>

⁶ <http://www.mpifr-bonn.mpg.de/staff/fvandertak/H2O/>

⁷ <http://www.strw.leidenuniv.nl/~moldata/>

Table 10. Envelope-averaged abundances of HDO and H₂O relative to H₂, derived from radiative transfer models.

Source	HDO 10 ⁻⁸	H ₂ O 10 ⁻⁵
W3 IRS5	1	0.8
W33A	3	1.0
AFGL 2591	2	6.0
NGC 7538 IRS1	3	0.8

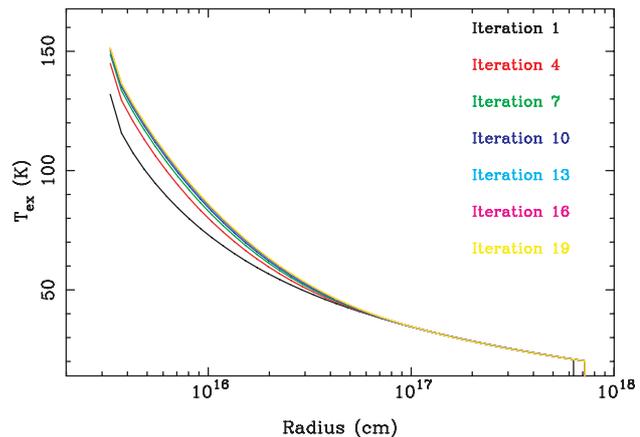
action, as often observed in the H₂O 183 GHz line, does not occur at H₂¹⁸O abundances of $\sim 10^{-8}$. The derived abundances are factors of ~ 30 higher than the line-of-sight averages from Sect. 3.3, due to the strong excitation gradients in these sources.

4.3. Jump models

Although the constant-abundance models in the previous section fit the strength of the H₂¹⁸O line and the average strength of the HDO lines well, they have several shortcomings. First, they predict *FWHM* sizes of the H₂¹⁸O and HDO emission in AFGL 2591 of 4–5'', significantly larger than the observed $\lesssim 1''$. Second, the fit residuals of HDO are correlated with energy level, in the sense that low-excitation lines tend to be overproduced and high-excitation lines underproduced. Both effects suggest that H₂O and HDO are not distributed evenly throughout the sources, but have enhanced abundances in the warm inner envelopes. Single-dish spectra and interferometer maps of H₂CO and CH₃OH (Van der Tak et al. 2000a) and of SO and SO₂ (Van der Tak et al. 2003) show the same effect.

Assuming that the water is produced by evaporation of icy grain mantles, we have run “jump” models for H₂¹⁸O in AFGL 2591. The parameters of this model are the abundance in the warm inner region and the temperature at which the ice evaporates. Laboratory studies indicate that the evaporation temperature lies in the 90–110 K range, depending on ice composition and structure (Pontoppidan et al. 2003; Fraser et al. 2005). With only one transition of H₂¹⁸O observed, it is not possible to constrain both parameters simultaneously, so we have initially fixed the boundary of the inner region at the 100 K point. The water abundance is assumed negligible outside this radius, which for AFGL 2591 lies at 2000 AU. The upper level of the observed transition is too high to set useful limits on the water abundance in the outer envelope. We find that the observed line flux and source size are reproduced for H₂O/H₂ = 1.4×10^{-4} which represents a major fraction of the available oxygen.

Alternatively, the ice mantles may evaporate at a somewhat different temperature. We have run a model with the H₂O abundance in the warm gas fixed at 2×10^{-4} , and varied the radius of the inner region to match the data. The best-fit model of this kind has the ice evaporating at $T = 115$ K. We consider both H₂O jump models plausible; multi-line observations

**Fig. 10.** Excitation temperature of the H₂O 3₁₃–2₂₀ line as a function of radius for the AFGL 2591 model. The curves for the last few iteration stages in the Monte Carlo calculation coincide, which shows that the simulation has converged.

of H₂¹⁸O are needed to rule out either model. This result is consistent with modeling of the SWAS data of AFGL 2591 (Boonman et al. 2003) which indicates an evaporation temperature of 90–110 K.

The present data do not constrain the abundance of H₂O/H₂ in the outer envelope very well. Boonman et al. (2003) derive an upper limit of $\sim 10^{-8}$ from a combined analysis of ISO-SWS, ISO-LWS and SWAS data which cover a range of energy levels. However, such a low H₂O abundance would imply an HDO/H₂O ratio of \sim unity in the outer envelope, which is implausibly high for this type of object. Our best-fit model to the SWAS data has H₂O/H₂ $\sim 10^{-6}$ in the outer envelope, inconsistent with the results by Boonman et al, but uncertain because based on only one transition. Clearly, *Herschel*-HIFI data are needed to settle this issue.

For HDO, “jump” models were run for each source except W3 IRS5, where too few lines were observed to constrain such models. In these models, the jump occurs at the fixed location of $T = 100$ K, and the HDO abundance inside and outside this radius are allowed to vary independently. The “jump” models reproduce all the single-dish line fluxes to within 50%, which is about the expected error margin. Table 11 reports the results of these models. For AFGL 2591, the size of the 81 GHz line emission measured with Bure acts as an extra constraint. The best-fit constant-abundance model predicts an emitting region of $\approx 5''$ *FWHM*, whereas the jump model predicts $\approx 2''$, consistent with the measured value, so that this latter model is favoured.

The “jump” models indicate that the HDO/H₂O ratio is $\approx 5 \times 10^{-4}$ in the inner region. This value corresponds to an enhancement over the interstellar D/H ratio by 100, and is consistent with the observational limits on the HDO/H₂O ratio in the solid state. We conclude that the bulk of the HDO and H₂O seen with the 30 m telescope and the Plateau de Bure interferometer is evaporated ice.

Table 11. Abundances of HDO and H₂O in the inner and outer envelopes of the sources, derived from radiative transfer models.

Source	HDO/H ₂		H ₂ O/H ₂		HDO/H ₂ O	
	10 ⁻⁹		10 ⁻⁴		10 ⁻⁴	
	inner	outer	inner	outer	inner	outer
W33A	200	10
AFGL 2591	100	4	1.4–2 ^a	10 ⁻² –10 ⁻⁴ ^b	5	40–4000
NGC 7538 IRS1	100	20

^a Values for ice evaporating at $T = 100$ K and $T = 115$ K.

^b From Boonman et al. (2003).

5. Discussion

5.1. Geometry of AFGL 2591: a massive circumstellar disk?

An important result from this study is that the H₂¹⁸O 203 GHz emission from AFGL 2591 is very compact (800 AU diameter) and coincident within the errors with the continuum emission from VLA3. The 1.3 mm continuum data indicate a mass of $0.8 M_{\odot}$, which is $\approx 5\%$ of the stellar mass of $\approx 16 M_{\odot}$ (Van der Tak & Menten 2005). One possible interpretation of these data is that we are observing the circumstellar disk surrounding the central star of AFGL 2591. The mass ratio seems plausible for a young protostar and given the observed axis ratios of the line and continuum emission (Tables 8 and 9), this putative disk is likely to be close to face-on (inclination 26–38 degrees). This inclination is compatible with the spectral energy distribution of the source in the near- to mid-infrared range (e.g., Whitney et al. 2003). However, there are certainly other possible interpretations of the data, at least as far as the line emission is concerned.

Another significant result from the interferometer observations is that neither H₂¹⁸O nor SO₂ is distributed in a spherically symmetric fashion around the central star. Figure 9 shows that the red and blue wings of the lines are offset from one another along a roughly NE–SW direction. The orientation seems to be somewhat different for water and SO₂. While one clearly needs much better angular resolution to interpret the data, we conclude that the velocity gradient in Fig. 9 is probably due to either disk rotation or the effect of the interaction of an outflow with the sides of a cavity (presumably directed toward us). SO₂ is a known outflow indicator (Schilke et al. 1997) and, more in general, is found associated with shocked gas. It seems likely that in a wide-angle wind inclined at a moderate angle to the line of sight, different outflow tracers will show different orientations in a representation such as that of Fig. 9. Such differences in orientation may particularly occur if the jet driving the outflow were to be precessing, which VLBI observations of H₂O 22 GHz indicate for AFGL 2591 (Trinidad et al. 2003). Although higher resolution imaging is required to settle this issue, we note that the column densities of H₂O, SO₂ and H₂ are well above the values in other massive molecular outflows (Beuther et al. 2002), making this scenario implausible.

Near-infrared speckle imaging of AFGL 2591 by Preibisch et al. (2003) supports our picture of the geometry of the source. These images (as well as older ones) show several loops of emission extending due West from the central source, with

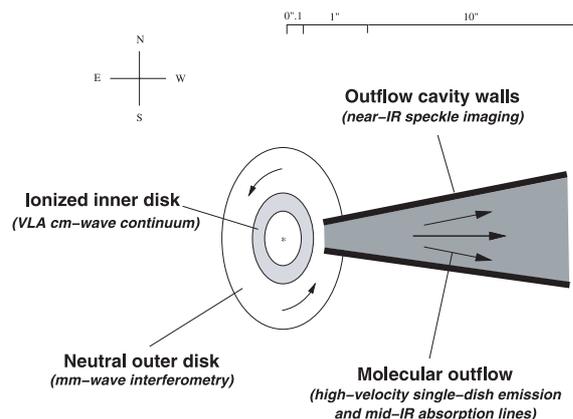


Fig. 11. Sketch of the inner part of the AFGL 2591 region, as projected on the sky, with the observational characteristics of each physical component indicated. At a distance of 1 kpc, 1'' corresponds to 1000 AU.

major axes of $\approx 10''$ and axis ratios of ≈ 3 . This emission presumably traces a limb-brightened cavity around the approaching outflow lobe. Since the outflow is expected to be perpendicular to the disk, the Western orientation of the loops is nicely consistent with the N–S orientation of the 205 GHz continuum emission. The near-infrared images thus indicate that the Western part of the disk is tilted away from us. The axis ratios of the loops are larger than that of the disk in the Bure images, as expected because the outflow is an intrinsically elongated structure, unlike the disk. The proposed outflow orientation is also consistent with the other observations mentioned in Sect. 3.6. In these data, the general direction of the outflow is East-West, with some diversity among tracers as expected in the proposed pole-on orientation of the system.

It also is worth noting that the source-averaged $N(\text{H}_2\text{O})$ toward AFGL 2591, corrected to a 0.8 source size, is $5 \times 10^{19} \text{ cm}^{-2}$, two orders of magnitude higher than the value derived from ISO 6 μm data (Boonman & van Dishoeck 2003). This difference again suggests an asymmetrical distribution of material around the central source as, for example, in a disk geometry. Along the line of sight to the central source, one observes across the outflow cavity but column densities are much larger in the disk or perhaps (see above) toward the cavity edges where outflow and inner envelope interact. In this scenario, the high excitation temperatures derived from the ISO data indicate that the cavity is hot (several 100 K).

Figure 11 summarizes our combined interpretation of the new and the existing observations of the central region of

AFGL 2591. The circumstellar disk and the molecular outflow are embedded in a large-scale molecular envelope, observed as the low-velocity single-dish emission and mid-infrared absorption. For clarity, this large-scale envelope is not drawn here, but it is depicted in Fig. 11 of Van der Tak et al. (1999), along with additional large-scale features. Note that part of the cm-wave emission observed with the VLA may arise in the base of the outflow.

5.2. Chemistry of water

It is also of great interest that the “jump” models in Sect. 4.3 show that our observations are consistent with the idea that the H_2^{18}O emission traces gas with a temperature above 100 K. Moreover, the H_2O abundances derived from these models are equal to or a few times higher than those estimated for H_2O ice in these sources (Van Dishoeck 1998; Boonman et al. 2003). Thus, we believe that the $3_{13}\text{--}2_{20}$ transition of H_2^{18}O can be used to trace the behaviour of high temperature gas where water ice has evaporated. Since the interferometer recovers all single-dish H_2^{18}O line flux, we conclude that the H_2^{18}O emission appears to be an excellent tracer of the inner ~ 1000 AU of protostars. In our model of AFGL 2591, the mass inside the 100 K point is $0.2 M_\odot$, similar to the mass derived from the Bure continuum data.

The warm-up of the central region to 100 K is evidently sufficiently recent that gas-phase chemistry has not had time to modify the abundances substantially. The $\text{H}_2\text{O}/\text{H}_3\text{O}^+$ ratio in W3 IRS5 is close to the equilibrium value of 1000 (Phillips et al. 1992). For W33A, AFGL 2136, AFGL 2591, S140 IRS1 and NGC 7538 IRS9, our unpublished JCMT observations of the H_3O^+ 364 GHz line indicate upper limits of $T_{\text{MB}} < 0.13\text{--}0.21$ K. Because of the large Einstein A coefficient of the transition, its excitation temperature is probably significantly below that of H_2O and HDO. Assuming $T_{\text{ex}} = 25$ K and an ortho/para ratio of 2, our upper limits on $N(\text{H}_3\text{O}^+)$ are $(4\text{--}7) \times 10^{13} \text{ cm}^{-2}$. These numbers correspond to lower limits on the $\text{H}_2\text{O}/\text{H}_3\text{O}^+$ ratio of 2000 in W33A and 6000 in AFGL 2591. Thus, in these sources, gas-phase chemistry seems not to have had time to return the $\text{H}_2\text{O}/\text{H}_3\text{O}^+$ ratio to its equilibrium value since the evaporation of the grain mantles.

The time scale t_{ion} to reach chemical equilibrium between H_2O and H_3O^+ can be estimated by realizing that H_3O^+ is produced in reactions between H_2O and molecular ions, in particular HCO^+ , He^+ and H_3^+ . Destruction of H_3O^+ is by dissociative recombination, which for 25% re-forms water, but for 60% makes OH and 15% O (Jensen et al. 2000). We thus estimate t_{ion} as $(\alpha_L n_X)^{-1}$, where α_L is the Langevin reaction rate of $\sim 10^{-9} \text{ cm}^3 \text{ s}^{-1}$, and the concentration of molecular ions n_X is given by the balance of cosmic-ray ionization on the one hand and reactions with CO and O on the other. Models by Van der Tak & van Dishoeck (2000) indicate $n_X \sim 10^{-4} \text{ cm}^{-3}$, so that $t_{\text{ion}} \sim 3 \times 10^5 \text{ yr}$. We conclude that the evaporation of grain mantles in AFGL 2591 has taken place less than ~ 0.1 Myr ago. The time scale may be even shorter, given the mass loss rate of $\sim 10^{-4} M_\odot \text{ yr}^{-1}$ measured in the CO outflow (Hasegawa & Mitchell 1995b). Assuming that the $0.8 M_\odot$ disk

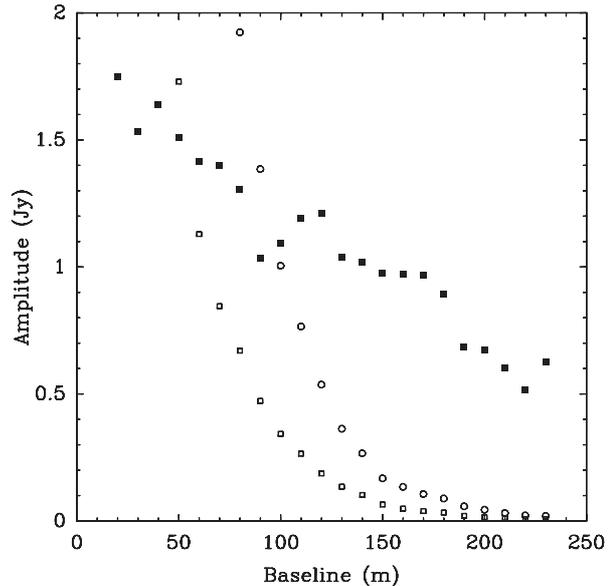


Fig. 12. Filled squares: visibility amplitude of H_2^{18}O line emission observed toward AFGL 2591 with the Plateau de Bure interferometer, integrated over velocity and binned. Superposed are model points for constant abundance (open squares) and for the “jump” model (open circles).

accretes at the same rate, the disk lifetime is only $\sim 10^4$ yr. This value is similar to the age estimate of 3×10^4 yr from multi-species chemical modeling of the envelope of AFGL 2591 (Doty et al. 2002).

It may be asked whether the Monte Carlo treatment of radiative transfer in Sect. 4 is useful given our uncertainty on the geometry. In fact, we think that our spherically symmetric model is adequate to determine the mass of warm (>100 K) gas necessary to explain the observations, and this mass is, for an essentially optically thin line, geometry independent. However, understanding line profiles as well as observations at still higher angular resolution will require a more sophisticated treatment such as axisymmetric modeling.

Figure 12 illustrates the limitations of our models by comparing them to the interferometer data in the uv plane. Although both models reproduce the total flux and the jump model also the source size as estimated through Gaussian fits, neither model reproduces the observations in detail. This discrepancy hints at the existence of additional geometrical structure which is not present in the model. The two models predict different source sizes, but the same overall emission shape, probably because the H_2^{18}O line is a tracer of warm gas.

Finally, our HDO data and the implied HDO/ H_2O abundance ratios are interesting in combination with the result from the “jump” models that the gas-phase water abundance is consistent with material which has recently evaporated off grains. Thus the observed HDO/ H_2O should reflect the ratio of these species in the solid state and indeed, our result is consistent with current limits on HDO ice.

5.3. Comparison with low-mass protostars

The observations of H_2^{18}O presented here have shown that in the studied high-mass protostars the water emission likely originates in the envelopes, and that the H_2O abundance jumps in the inner warm region where the grain mantles sublimate. A similar analysis of the ISO-LWS spectra of two low-mass protostars (Ceccarelli et al. 2000; Maret et al. 2002) shows that also in those cases, a jump in the water abundance at approximately the radius where the dust temperature reaches 100 K is needed to explain the observed far-infrared water line spectrum. However, the water abundance in the warm gas is strikingly different: $\sim 10^{-4}$ in the high-mass, and $\sim 3 \times 10^{-6}$ in the low-mass protostars respectively. Mid-infrared spectra of low-mass protostars give solid water abundances of $\sim 10^{-4}$ as for high-mass protostars (Pontoppidan et al. 2004). It is not obvious why the water abundance in the warm gas should be only a few percent of the ice abundance. As Sect. 4.3 discusses, the exact ice evaporation temperature depends somewhat on the ice composition and structure, but not enough to make a factor of 100 difference. Unless the water abundances in the studied low-mass protostellar envelopes are affected by the large (80'') ISO-LWS beam, it seems that the break-down of evaporated water ice is faster around low-mass than around high-mass stars. This conclusion is somewhat surprising since water is destroyed by ion-molecule chemistry and the few available data suggest that if anything, the ionization rate is higher around high-mass than around low-mass stars (Van der Tak 2005). In the future, *Herschel*-HIFI data will be helpful to settle this issue.

HDO emission, practically at all the same frequencies of the present work, has been observed towards the low-mass protostar IRAS 16293-2422 (Parise et al. 2005). Similarly to the present work, the authors analyzed the observations with an abundance jump model. They found that the abundance of HDO has a jump from a value $\leq 10^{-9}$ in the outer envelope, where the dust temperature is lower than 100 K, to $\sim 10^{-7}$ in the inner envelope. The abundance found in the region of sublimated ices is therefore very similar in high- and low-mass protostars. However, the HDO/ H_2O ratio is substantially different. In high-mass protostars it is at most 3×10^{-3} , whereas in the low-mass protostar IRAS 16293-2422 it is ~ 0.03 , namely ten times larger and close to the observational limit on solid HDO (Parise et al. 2003). This difference is as expected from recent measurements of extreme molecular deuteration around low-mass protostars, where doubly and triply deuterated molecules have been detected with fractionations of a few percent (for references, see Van der Tak 2005). The degree of deuteration in high-mass protostellar envelopes is much lower, starting from the failure to detect H_2D^+ in massive protostars (Pagani et al. 1992; Stark et al. 1999). Possibly, for high-mass protostars, the very cold and dense pre-collapse phase where the CO freezes out onto the grain mantles lasts only a short time. The present measurement of HDO/ H_2O in high mass protostars, compared with the value found in IRAS 16293-2422, confirms this hypothesis: high-mass protostars show indeed a lower degree of water deuteration.

The water ice on grain surfaces is laid down in relatively low density molecular cloud gas (10^3 – 10^4 cm^{-3}) where

H atoms are as abundant as O atoms and thus there is a relatively high probability of H_2O forming subsequent to O sticking to a grain. After formation, the ice presumably stays frozen until the dust is heated up by protostellar radiation. It is the temperature of this ‘‘primordial’’ low-density molecular cloud gas that counts for the chemistry of HDO/ H_2O . The present observations suggest that this temperature is higher in the Giant Molecular Clouds producing high-mass stars than in low-mass star-forming regions such as Taurus and Ophiuchus.

6. Conclusion and outlook

The chemical composition of star-forming matter depends on both temperature and time. For over a decade, people have tried to use the time dependence to estimate chemical ages of star-forming regions. These estimates, the so-called ‘‘chemical clocks’’ remain unreliable, probably due to uncertainties in the initial conditions as well as in the physical structure (cf. Van der Tak 2005). This paper has shown that the temperature dependence may instead be used to apply ‘‘chemical filters’’. In particular, we have used the H_2O molecule to image the material at $T > 100$ K, filtering out the surrounding cooler material. The success of this filter lies in the evaporation of icy grain mantles which enhance the H_2O gas-phase abundance by two orders of magnitude at $T > 100$ K.

By using this chemical filter, we have shown that the dust and water inside 2000 AU from the central star in the AFGL 2591 region is asymmetrically distributed. Most likely, the observations trace a circumstellar disk of diameter 800 AU which rotates at about Keplerian speed. The result of the SO_2 observations is qualitatively the same as for H_2^{18}O , but differs in the details. In this case, the ‘‘chemical filter’’ may be a bit leaky because the SO_2 abundance in the large-scale envelope is non-negligible, and because SO_2 is also abundant in the bipolar outflow of AFGL 2591. Furthermore, comparison of our results with those for the prototypical low-mass system IRAS 16293 shows that the H_2O abundance in the warm gas around young high-mass stars is much higher, but the HDO/ H_2O ratio much lower than around low-mass protostars.

In the future, observations with the PACS camera and the HIFI spectrometer onboard the *Herschel* space observatory will further refine this picture (Walmsley & van der Tak 2005). Observations of multiple H_2O and H_2^{18}O lines will constrain the excitation and ortho/para ratio of water much better than is possible from ground-based data. The high sensitivity will allow us to measure the water abundance in each chemical zone (disk/outflow/envelope), not only around high-mass protostars, but also around lower-mass objects. Given our estimated outflow contribution of 10–20% to the H_2^{18}O and HDO lines in HIFI-sized beams, the contribution from outflows to H_2^{16}O spectra from HIFI will probably be much larger, which should be considered in the planning of the HIFI observations. New, deep searches for H_3O^+ in high-mass protostellar envelopes with APEX, HIFI and ALMA would also be useful. Observations of the H_2^{18}O and SO_2 lines in AFGL 2591 on longer baselines are necessary to resolve the velocity field and test other possibilities such as a binary system, as seen in the W3 (H_2O) source by Wyrowski et al. (1999).

Acknowledgements. We thank the staffs of the IRAM 30 m, JCMT, and Plateau de Bure telescopes for assisting with the observations, especially Jan Martin Winters and Jérôme Pety at IRAM Grenoble. The JCMT data were taken in collaboration with Ewine van Dishoeck and Annemieke Boonman at Leiden Observatory. Holger Müller at the University of Cologne kindly provided H₂O term values to spectroscopic accuracy.

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