High-frequency radio observations of the Kühr sample and the epoch-dependent luminosity function of flat-spectrum quasars*

R. Ricci2, I. Prandoni3, C. Gruppioni4, R. J. Sault2, and G. De Zotti5

1 SISSA/ISAS, via Beirut 2–4, 34014 Trieste, Italy
2 Australia Telescope National Facility, CSIRO, PO Box 76, Epping, NSW 2121, Australia
e-mail: Roberto.Ricci@csiro.au
3 INAF, Istituto di Radioastronomia, via Gobetti 101, 40129 Bologna, Italy
4 INAF, Osservatorio Astronomico di Bologna, via Ranzani 1, 40126 Bologna, Italy
5 INAF, Osservatorio Astronomico di Padova, Vicolo dell’Osservatorio 5, 35122 Padova, Italy
e-mail: dezotti@pd.astro.it

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ABSTRACT

We discuss our ATCA 18.5 and 22 GHz flux density measurements of Southern extragalactic sources in the complete 5 GHz sample of Kühr et al. (1981, A&AS, 45, 367). The high frequency (5–18.5 GHz) spectral indices of steep-spectrum sources for which we have 18.5 GHz data (66% of the complete sample) are systematically steeper than the low frequency (2.7–5 GHz) ones, with median $\alpha_{18.5}^5 = 0.76$, median $\alpha_{18.5}^{25} = 1.18$ ($S \propto \nu^{-\alpha}$), and median steepening $\Delta \alpha = 0.32$, and there is evidence of an anti-correlation of $\Delta \alpha_{18.5}\nu$ with luminosity. The completeness of 18.5 GHz data is much higher (89%) for flat-spectrum sources (mostly quasars), which also exhibit a spectral steepening: median $\alpha_{18.5}^5 = -0.14$, median $\alpha_{18.5}^{25} = 0.16$ ($S \propto \nu^{-\alpha}$), and median $\Delta \alpha = 0.19$. Taking advantage of the almost complete redshift information on flat-spectrum quasars, we have estimated their 5 GHz luminosity function in several redshift bins. The results confirm that their radio luminosity density peaks at $z_{\text{peak}} \approx 2.5$ but do not provide evidence for deviations from pure luminosity evolution as hinted at by other data sets. A comparison of our 22 GHz flux densities with WMAP $K$-band data for flat-spectrum sources suggests that WMAP flux densities may be low by a median factor of $\approx 1.2$. The extrapolations of 5 GHz counts and luminosity functions of flat-spectrum radio quasars using the observed distribution of the 5–18.5 GHz spectral indices match those derived directly from WMAP data, indicating that the high frequency WMAP survey does not detect any large population of FSRQs with anomalous spectra.

Key words: radio continuum: general – galaxies: nuclei – quasars: general – galaxies: luminosity function, mass function

1. Introduction

High radio frequency (10–100 GHz) sky surveys have started to become feasible only very recently. The first all-sky surveys at high radio frequencies have been provided by the WMAP satellite (Bennett et al. 2003). In particular, the complete sample of extragalactic sources drawn from its $K$-band catalog ($S \geq 1.25\, Jy$, $|b| \geq 10^\circ$; see De Zotti et al. 2005) comprises 155 extragalactic sources (plus a planetary nebula), over an area of 10.4 sr.

Another survey which will give a significant contribution to the knowledge of the $>10$ GHz extragalactic sky is the ongoing all-sky 20 GHz survey which is being undertaken with the Australia Telescope Compact Array (ATCA) in the Southern emisphere (a pilot survey covers about 1200 deg$^2$ to 100 mJy, Ricci et al. 2004b).

An example of a deeper survey is the 9th Cambridge survey carried out at 15 GHz with the Ryle Telescope (Waldram et al. 2003), which covers 520 deg$^2$ to $S = 25$ mJy, and reaches deeper flux densities on smaller areas.

Such surveys are of fundamental importance to directly provide information about the extragalactic populations dominating the sky at $>10$ GHz frequencies, and allow to test extrapolated models based mainly on lower frequency selected samples.

In this paper we present 18.5 and 22 GHz observations of the Southern extragalactic sources of the complete 5 GHz sample of Kühr et al. (1981), with the aim of better modeling the high frequency properties of radio galaxies and quasars, and check whether radio sources detected by WMAP in the $K$-band significantly differ from the ones dominating at lower frequencies.

In Sect. 2 we briefly present the 18.5 and 22 GHz flux density measurements. The main high frequency properties of Southern Kühr sources are analysed in Sect. 3. In Sect. 4

* Table 1 is only available in electronic form at the CDS via anonymous ftp to cdsarc.u-strasbg.fr (130.79.128.5) or via http://cdsweb.u-strasbg.fr/cgi-bin/qcat?J/A+A/445/465
we estimate the 5 GHz luminosity function (LF) of FSRQs at different cosmic epochs and compare it with predictions of a recent model by De Zotti et al. (2005). The observed 5–18.5 GHz spectral index distribution is then exploited to extrapolate the LF to 22.8 GHz, the central frequency of the WMAP K-band. The extrapolated LF is compared with a direct estimate from the WMAP data. In Sect. 5 we summarize our main conclusions.

Throughout this paper we adopt a flat $\Lambda$ cosmology with $\Omega_M = 0.7$ and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$.

2. The data

In March 2002 we have carried out total intensity and linear polarization measurements at 18.5 GHz of 249 of the 258 southern ($\delta < 0^\circ$) extragalactic sources in the 5 GHz all-sky 1 Jy sample (Kühr et al. 1981; Stickel et al. 1994) with the Australia Telescope Compact Array (ATCA), using the prototype high-frequency receivers mounted on antennas CA02, CA03, and CA04. The array configuration was quite compact, giving a HPBW = 15′.6. Full details on observations and data reduction are provided by Ricci et al. (2004a, hereafter referred to as Paper I), who carried out an analysis of linear polarization data. Many of these sources had already been targeted with the ATCA as part of a search for potential 22 GHz calibrators. The remaining 22 GHz total intensity measurements were taken with ATCA in January 2001 using the three antennas of the 750C array configuration providing a HPBW = 7′.0.

The 18.5 and 22 GHz flux densities are listed in Table 1, where, for each observed source, we also report the 5 GHz flux density, the spectral index between 2.7 and 5 GHz (both from Kühr et al. 1981), the source type (from Stickel et al. 1994) and the source redshift (mostly from Stickel et al. 1994, with some additional redshifts from the NED database).

Both 18.5 and 22 GHz data were reduced as discussed in Paper I. In particular, flux densities are derived using non-imaging model-fitting techniques, which assume a point source model. We notice that the 22 GHz data have larger uncertainties, are less homogeneous and, having better angular resolution, can suffer more from resolution effects. Therefore, our analysis mostly relies on 18.5 GHz data.

Whenever a source at 18.5 GHz is poorly modeled as a point source (see Paper I for details), such source is labeled “R” (for “resolved”) in the last column of Table 1.

Small differences with the 18.5 GHz flux densities, and associated errors, reported by Ricci et al. (Paper I), are due to a new analysis of the measurements. The 18.5 GHz errors now include the flux calibration uncertainty of 5%. Calibration errors, estimated at 10%, are also included in the 22-GHz data.

Following Stickel et al. (1994), we have classified as flat-spectrum, sources with 2.7–5 GHz spectral index $\alpha_{2.7} < 0.5$ ($S_\nu \propto \nu^{-\alpha}$); sources with larger values of $\alpha_{2.7}$ are classified as steep-spectrum. Of the 249 sources observed, 146 are flat-spectrum and 103 are steep-spectrum.

In the following analysis we exclude “resolved” (“R”) sources. This means 130 out of the 146 flat-spectrum sources in the original complete sample (89%) and 68 of the 103 steep-spectrum sources (66%) are retained. Special emphasis will be given to Flat-Spectrum Radio Quasars (FSRQs), which constitute the population which is best represented in our sample.

3. Data analysis

Since the high frequency information on steep-spectrum sources is seriously incomplete (see above), and, as a consequence, any conclusion on their spectral properties may be biased, we only note that they show evidence of a spectral steepening. The median $\alpha_{5.7} = 0.76$ ($S_\nu \propto \nu^{-\alpha}$) (the mean is 0.81 ± 0.02), while the median $\alpha_{18.5} = 1.18$ (mean 1.28 ± 0.06), and the median steepening is $\Delta \alpha = 0.32$ (mean 0.47 ± 0.06). The steepening is anti-correlated with luminosity (see Fig. 1) or redshift (the Spearman’s rank correlation coefficient yields a probability of no correlation of $3 \times 10^{-4}$). This translates into an anti-correlation of $\alpha_{18.5}$ with luminosity (or z), possibly reversing the positive correlation with luminosity of the low frequency spectral index (see, e.g., Dunlop & Peacock 1990).

Since low-luminosity (and low z) sources are mostly galaxies, while quasars are high-luminosity (and high z) sources (see Fig. 1), we have larger steepenings for galaxies (median $\Delta \alpha = 0.34$) than for quasars (median $\Delta \alpha = 0.20$).

A high-frequency steepening is also observed for flat-spectrum sources, which are mostly FSRQs. The median values are $\alpha_{2.7} = -0.14$ ($S_\nu \propto \nu^{-\alpha}$), $\alpha_{18.5} = 0.16$, $\Delta \alpha = 0.19$, while the mean values are, respectively, $-0.13 \pm 0.03$, $0.11 \pm 0.04$, and $0.24 \pm 0.05$. The best-fit Gaussian distribution of 2.7–5 GHz spectral indices has a mean of $-0.10$ and a dispersion of 0.30, while that of 5–18.5 GHz spectral indices has mean 0.13 and dispersion 0.44 (Fig. 2). The large value for the dispersion of the distribution of $\alpha_{18.5}$ may be partly due to measurement errors and to uncertainties in the source classification; we will come back to this in Sect. 4.2. The Spearman’s rank correlation statistic detects a positive correlation of $\alpha_{18.5}$ with luminosity (or z), although with a lower significance than for steep-spectrum sources (probability of no correlation $2 \times 10^{-3}$).
3.1. Comparison with WMAP K-band flux measurements

It is interesting to compare our 18.5 GHz flux density measurements of flat-spectrum unresolved sources of the Kühr sample with those of WMAP in the K-band (22.8 GHz). A cross-correlation with the WMAP catalog (Bennett et al. 2003) yields 59 matches within angular separations of 11 arcmin (the Bennett et al. 2003, criterion for WMAP source cross-identifications). For the un-matched flat-spectrum sources in our sample we adopt an upper flux density limit in the K-band of \( S_K = 1.25 \text{ Jy} \), i.e. the completeness limit of the WMAP survey, as derived by De Zotti et al. (2005; see also Argüeso et al. 2003). We notice, however, that Bennett et al. (2003) give a less conservative estimate of \( S_K = 0.75 \text{ Jy} \). Using the Kaplan-Meier estimator (routine KMESTM in the software package ASURV Rev 1.2 (Isobe & Feigelson 1990)) we get an estimated median value of \( \log(S_{18.5 \text{ GHz}} / S_K) = 0.098 \) (corresponding to \( S_{18.5 \text{ GHz}} / S_K = 1.25 \)). This includes WMAP flux density upper limits by implementing the survival analysis methods presented in Feigelson & Nelson (1985) and Isobe et al. (1986). The mean value of \( \log(S_{18.5 \text{ GHz}} / S_K) = 0.126 \) is less reliable because it is significantly affected by outliers (probably variable sources).

The median \( \alpha_{18.5} \approx 0.16 \) would imply a significantly lower value \( \log(S_{18.5 \text{ GHz}} / S_K) = \log(S_{22.8 \text{ GHz}}) = \log(S_{18.5 \text{ GHz}}) - 0.16 \log(22.8/18.5) \), or \( \log(S_{18.5 \text{ GHz}} / S_K) = 0.015 \) and using the actually observed distribution of \( \alpha_{18.5} \) for flat-spectrum sources (see Fig. 2) we obtain an even lower mean value \( \log(S_{22.8 \text{ GHz}} / S_K) = 0.0097 \approx 0.0099 \), suggesting that the WMAP K-band flux densities are systematically lower than the ATCA flux densities by a factor of \( \approx 1.2 \). Using our measurements at the closer frequency of 22 GHz (which has, however, a larger calibration uncertainty) we find a median value \( \log(S_{22 \text{ GHz}} / S_K) = 0.057 \) (mean \( \log(S_{22 \text{ GHz}} / S_K) = 0.101 \pm 0.025 \)), confirming the indication of a small, but non-negligible offset.

4. The epoch-dependent luminosity function of FSRQs

4.1. The 5 GHz luminosity function

To estimate the epoch-dependent 5 GHz luminosity function of FSRQs we exploit the full Kühr et al. (1981) sample, including the Northern portion, not observed at 18.5 GHz. The redshift information is essentially complete: Stickel et al. (1994) list redshifts for 198 of the 214 FSRQs in that sample and a search in the NASA-IPAC Extragalactic Database (NED) yielded redshift measurements for 8 additional objects. In summary, the redshift completeness is 94% in the flux density bin \( 0 \leq \log S_5(\text{Jy}) < 0.2 \), 98% in the bin \( 0.2 \leq \log S_5(\text{Jy}) < 0.4 \) and 100% at higher flux densities.

We have estimated the 5 GHz luminosity function, averaged over intervals \( \Delta \log L = 0.5 \), in six redshift bins (0.01 < \( < 0.4 \), 0.4 < \( < 0.8 \), 0.8 < \( < 1.2 \), 1.2 < \( < 2 \), 2 < \( < 3 \), 3 < \( < 4 \)) using the classical \( 1/V_{\text{max}} \) method (Schmidt 1968):

\[
\Phi(L_5, z) = \frac{\sum_{L_i \in \{L_5 \pm \Delta L/2, L_5 + \Delta L/2\}} w_i}{V_{\text{max}, i}}
\]

(1)

where \( V_{\text{max}} \) is the maximum volume accessible to each source and \( w_i \) is the weight factor correcting for the redshift incompleteness. \( V_{\text{max}} \) is given by the integral of the volume element from the lower bound of the redshift bin to \( z_{\text{max}} = \min(z_{\text{up}}, z_{\text{lim}}) \), where \( z_{\text{up}} \) is the upper bound of the bin and \( z_{\text{lim}} \) is the redshift at which the source would have a flux density of 1 Jy (i.e. the lower limit of the Kühr sample). The weight factor \( w_i \) is set at \( w_i = 1/0.94 \) at and \( w_i = 1/0.98 \), respectively, for sources in the two lowest flux density bins, defined above, and at \( w_i = 1 \) for brighter sources. The K-correction was computed using the spectral index \( \alpha_{2.7} \) of each source (see Kühr et al. 1981). The area covered by the catalog is 9.81 sr. The 68% confidence intervals on the luminosity function were computed assuming Poisson statistics with an effective number of sources per bin given by:

\[
n_{\text{eff}} = \frac{1}{\int \sum_{i \in L_{\text{max}}}} \int \frac{1}{V_{\text{max}, i}} \frac{1}{i}
\]

(2)

In Fig. 3 the estimated luminosity functions (points) are compared with the model by De Zotti et al. (2005), averaged over the same luminosity and redshift bins (lines). The agreement is generally satisfactory. In spite of its simplicity (it assumes pure luminosity evolution), the model correctly reproduces the positive evolution up to \( z \sim 2.5 \) and the subsequent negative evolution.

4.2. Extrapolation to high frequencies

We now use the WMAP K-band catalogue to derive the 22.8 GHz luminosity function and compare it to the one extrapolated from 5 GHz. After having multiplied the WMAP fluxes by 1.2 (see Sect. 3), and using the 4.85 GHz flux densities from the GB6 (Gregory et al. 1996) or PMN (available at http://www.parkes.atnf.csiro.au/) catalogs, 143 of them have been classified as flat-spectrum based on the
Fig. 3. Comoving luminosity function of FSRQs in the Kühr et al. (1981) sample for different redshift bins (points), compared with the model by De Zotti et al. (2005) (lines).

Fig. 4. Differential counts of FSRQs in the complete WMAP K-band sample (see text), compared with the extrapolations of 5 GHz counts using the Gaussian representations of the 5−18.5 GHz spectral index distribution of Fig. 2, namely $\langle \alpha \rangle = 0.13$ and $\sigma = 0.30$ (solid line). The WMAP fluxes have been multiplied by 1.2 (see Sect. 3). The error bars show the 68% confidence intervals for Poisson statistics (Gehrels 1986).

With $\sigma = 0.30$, close to the dispersion of the 2.7−5 GHz spectral index distribution, the agreement is good (see solid line in Fig. 4), indicating that high frequency surveys do not detect many FSRQs with “anomalous” radio spectra.

Redshift measurements were found in the literature for 99 of the 110 objects in the WMAP FSRQ sample. We have estimated their luminosity function for the same luminosity and redshift bins as for the Kühr sample (see Fig. 3), except for the highest redshift bin, because we have only one WMAP source at $z > 3$. The K-corrections have been computed using the median 4.85−22.8 GHz spectral index $\alpha_{22}^{18.5} = 0.16$; this has been preferred to $\alpha_{4.85}^{22.8}$ of individual objects which are too liable to variability. In Fig. 5 these estimates (points) are compared with the extrapolations of the 5 GHz luminosity function as modeled by the De Zotti et al. (2005), which provides a good representation of the data (see lines), when we use the 5−18.5 GHz Gaussian spectral index distribution with $\langle \alpha \rangle = 0.13$ and $\sigma = 0.30$.

5. Conclusions

The ATCA 18.5 and 22 GHz flux density measurements of Southern extragalactic sources in the complete 5 GHz sample of Kühl et al. (1981) have been analyzed. The non-imaging technique used yielded reliable flux density measurements for 66% of steep-spectrum sources. Although this incompleteness requires caution in dealing with the data, some interesting indications emerge. First the high frequency (5−18.5 GHz) spectral indices are systematically steeper than the low frequency (2.7−5 GHz) ones, the median steepening being $\Delta \alpha = 0.32$. There is a hint of a larger steepening for sources classified as galaxies than for those classified as QSOs. Since QSOs have generally higher luminosities and higher redshifts than galaxies, this difference translates into an anti-correlation of $\Delta \alpha$ with...
the radio luminosity (or z) when we consider the full steep-spectrum sample. As a consequence, the positive correlation of the radio luminosity with the low frequency spectral index (see, e.g., Dunlop & Peacock 1990), seems to get reversed when we consider the high frequency spectral index. More complete data would however be necessary to test the reliability of this indication.

The information on 18.5 GHz flux densities of flat-spectrum sources (mostly FSRQs) has a much better completeness level (89%). A high-frequency steepening is again present, with median values $\alpha_{2.7}^5 = -0.14$, $\alpha_{18.5}^5 = 0.16$, $\Delta \alpha = 0.19$. Luminosities of these sources are positively correlated, though weakly, with both the low and the high frequency spectral indices.

In addition to having almost complete 18.5 GHz data, FSRQs have almost complete redshift information, that has been exploited to estimate their luminosity function in several redshift bins. The 5 GHz luminosity functions, computed using the full 1 Jy sample, while confirming that the radio luminosity density peaks at $z_{\text{peak}} = 2.5$ (Dunlop & Peacock 1990; Shaver et al. 1996, 1999) do not provide evidence for deviations from pure luminosity evolution, reported by other analyses (Hook et al. 1998; Vigotti et al. 2003), perhaps due to the poor statistics. In particular we notice that the derived epoch-dependent 5 GHz luminosity function is fairly well represented by the model by De Zotti et al. (2005).

We exploit the multi-frequency 2.7, 5 and 18.5 GHz information, for Southern Kühr FSRQs to better extrapolate the De Zotti et al. model to 22.8 GHz, i.e. the central frequency of the WMAP K-band survey. This allows us to properly take into account the spectral steepening of FSRQs between 5 and 18.5 GHz.

The extrapolated model is compared to the sample of FSRQs detected in the K-band by WMAP. To this respect, it is interesting to note that a direct comparison between ATCA 18.5/22 GHz and WMAP K-band flux density measurements for flat-spectrum sources hints towards a systematic under-estimation of WMAP fluxes by a factor $\sim 1.2$.

After correcting the WMAP K-band fluxes for the above-mentioned offset, we are able to get very good agreement between the differential counts of WMAP K-band FSRQs and the modeled extrapolation to 22.8 GHz. The same happens when we analyze the luminosity properties of the FSRQ WMAP sample. The K-band epoch-dependent luminosity function derived from the sample of WMAP FSRQs is fairly well reproduced by the modeled 5 GHz luminosity function, extrapolated to 22.8 GHz as described above.

Such results imply that the WMAP survey does not detect substantial numbers of FSRQs with anomalous spectra.

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