Optical and EUV observations of solar flare kernels

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ABSTRACT

We present high-resolution spectral observations, covering the entire optical region (3800–9000 Å), of a solar flare observed during a multi-wavelength campaign. The flare, recorded on 2002 January 11, was a medium solar flare event (GOES class C7.5). The spectral observations were carried out using the Hamilton echelle spectrograph on the coudé auxiliary telescope at Lick Observatory and with the Coronal Diagnostic Spectrometer (CDS) on board SoHO. The high signal-to-noise optical spectra are analysed using the same techniques as we applied to stellar flare data. Hα images obtained at Big Bear Solar Observatory (BBSO), plus magnetograms obtained with the Michelson Doppler Imager (MDI) on board SoHO and Transition Region And Coronal Explorer (TRACE) 1600 Å were used in the flare analysis. We observe stellar-like behaviour in the main solar chromospheric activity indicators, which show either filling-in or emission during the flare. We find that the Balmer and Ca ii lines show asymmetric profiles, with red-shifted wings and blue-shifted cores. This behaviour could be explained by material expanding. During the flare, the Mg i and Fe i lines show a filling-in of the line profile indicating that the flare affected the lower atmosphere. There is some evidence for pre-flare heating as seen in Fe xix 592 Å. Furthermore, O ν 629 Å shows an increase in flux some 10 min. before the coronal lines, perhaps indicating particle beam heating in the initial stages of the flare. We have also determined the main physical parameters at flare maximum. The electron densities and electron temperatures found for the flare imply that the Balmer emitting plasma originates in the chromosphere. The physical parameters obtained for the modelled flare are consistent with previously derived values for solar flares.

Key words. Sun: activity – Sun: photosphere – Sun: flares – Sun: chromosphere – line: formation – line: profiles

1. Introduction

Flares are events where a large amount of energy is released in a short interval of time, radiating at almost all frequencies in the electromagnetic spectrum. Flares (both solar and stellar) are believed to result when reconnection in the corona rapidly releases magnetic energy stored in the coronal fields (see reviews by Mirzoyan 1984; Haisch et al. 1991; Garcia Alvarez 2000). Many types of cool stars produce flares (Pettersen 1989), sometimes at levels several orders of magnitude more energetic than their solar counterparts. The exact mechanism(s) leading to the energy release and subsequent excitation of various emission features (X-ray, microwave, UV, and optical lines and continua) remains poorly understood. Solar flares are generally observed to have an impulsive phase marked by hard X-ray bursts (Dennis 1985) and microwave radio emission on time scales of seconds to minutes, indicating the presence of accelerated non-thermal particles (Lin 1974; Foukal 1990). In some cases, optical continuum radiation is also observed (Neidig 1989). The impulsive phase is usually followed by a gradual phase, lasting minutes to hours. The impulsive phase shows continuing enhancement of chromospheric and coronal emission followed by a slow decay during the gradual phase (Svestka 1976; Tandberg-Hanssen & Emslie 1988; Dennis & Schwartz 1989).

The Sun, due to its proximity, allows both high spatial and spectral resolution observations of flares. Satellites such as SoHO, TRACE and recently RHESSI, are continuously monitoring high-energy radiation, and magnetograms of the surface are routinely available. Ground-based sites, such as Big Bear Solar Observatory (BBSO) and Sacramento Peak Observatory routinely monitor Hα, Ca ii H and K, and take magnetograms. Simultaneous observations in different wavelength ranges (e.g. optical, X-rays, etc.) make it possible for the same flare to be studied at all levels in the solar atmosphere.

Despite the huge number of solar flare observations only a few experiments have studied the full optical spectrum during solar flares (Acampa et al. 1982; Donati-Falchi et al. 1984; Mauas 1990; Mauas et al. 1990). Johns-Krull et al. (1997)
reported full wavelength optical spectra with high spectral resolution and good time resolution for a solar flare. This type of observation allows one to apply the same techniques used in the analysis of stellar flares, which provides an important link for comparing studies of solar and stellar flares. The lack of full optical range observations for solar flares is primarily due to the design of solar spectrographs which allow very high resolution line profiles but usually only small wavelength coverage. The use of cross-dispersed echelle spectra with both good spectral resolution and good time resolution has been restricted almost totally to stellar objects. Johns-Krull et al. (1997) however modified the Hamilton spectrograph at Lick Observatory to obtain solar flare observations in 1993. This allowed them to observe the spectrum of a flare at high dispersion throughout the entire optical range. Unfortunately, there was little in the way of simultaneous satellite or other ground based data taken for this flare, which limited the range of the solar atmosphere which could be studied.

In this paper we present multi-wavelength observations of a solar flare observed at 20:07 UT on 11 January 2002, although the full observing campaign included the interval between 9–15 January 2002. Three flares were observed during the campaign, in this paper only one solar flare will be reported, a GOES C7.5 event, which had the most simultaneous coverage. The objective of this campaign was to monitor for flare events with the same sort of spectral coverage as is normally obtained in stellar events. Observing time was allocated for this project at Lick Observatory, Big Bear Solar Observatory, the CDS and MDI instruments onboard the SoHO, plus TRACE. Observations and data analysis are described in Sect. 2; while in Sect. 3 we discuss the flare evolution as seen by the different instruments. In Sect. 4 we analyse the behaviour of the chromospheric and photospheric lines during the flare, and the chromospheric and coronal flare energy budget. In Sect. 5 we analyse the main chromospheric physical parameters of the emitting region during the flare with the conclusions given in Sect. 6.

2. Instrumentation and observational material

2.1. Lick observations

The Hamilton echelle spectrograph (Vogt 1987) and the telescopes that feed the spectrograph, were not designed for solar observations. As a consequence of this, some unusual steps were needed to observe solar flares. We have followed the same observational setup, with some minor modifications, used by Johns-Krull et al. (1997) in order to obtain the present data. The 0.6 m Coude Auxiliary Telescope (CAT) was used to feed the spectrograph. The primary mirror was covered with a mask that stepped the aperture down to a circle of approximately 20 cm diameter. The CAT itself is fixed to the wall of the Observatory and is fed by a siderostat. This system feeds a white-light image to the spectrograph slit, which, for this study was 5.5″ wide by 10.0″ long. In order to cut down on the flux delivered to the spectrograph CCD, we inserted a BK38 filter behind the spectrometer slit. To further cut down on the flux, a piece of mylar was inserted in the filter holder along with the BK38 glass. Optical wavelength coverage was from 3800 to 9000 Å. This combination provides a spectral resolution of $R \sim 45,000$ and a temporal resolution of ~2 min. To improve the contrast of active regions on the Sun (compared to the image delivered to the spectrograph slit) and to enable the identification of solar flares, a 0.6 Å Daystar Hα filter was placed in front of the slit viewing guide camera. The siderostat could then be used to point the CAT to any feature visible in the Hα image, in our case, flare kernels.

The spectra were reduced with customized IDL software using a standard procedure discussed by Valenti (1994). This includes tracing each spectral order on the chip, flat-fielding and subtracting scattered light measured between orders. Use of the combined mylar and BK38 filter behind the slit resulted in excess fringing which is not fully removed in the standard flat-fielding procedure. In order to better remove this, we divide an extracted spectrum of a quiet region taken with the filter mounted behind the slit by an extracted spectrum of a quiet region taken without this filter in place. This results in a spectrum of the excess fringing which was then divided into each exposure taken with the combined mylar and BK38 to remove the fringes. Visual inspection indicates that residual fringing is less than 2%. In addition, the results shown below using this data are differential in nature (flare minus quiet, both observed with the combined mylar and BK38 filter), making this fringing less of a concern. Finally, wavelength calibration was accomplished by fitting a two-dimensional polynomial to nλ (where n is the echelle order number) as a function of thorium-line location in a lamp spectrum.

At BBSO, the active region NOAA 9773, where the flare occurred, was observed throughout the campaign. Figure 1 shows Hα images of the active region, taken at BBSO from the flare onset until the end of the gradual phase. Several eruptive centers can be observed during the flare onset. The flaring kernel site observed at Lick Observatory is marked in each image with the lower left of the 2 rectangles.

2.2. TRACE observations

The TRACE (Handy et al. 1999) observations used in this study were obtained from 20:02 UT until 20:32 UT on 02 January 11. The observations were taken in the 1600 Å passband which contains the C I lines, Fe II lines and UV continuum with a significant contribution from both the C IV 1548 and 1550 Å lines. The passband width is 275 Å. The exposure time of the images varies from 0.15 to 0.43 s. The TRACE angular pixel size is 0.5″ which gives a 1″ spatial resolution. The data were flat field corrected, normalized for the exposure and pointing corrected for the corresponding channel. The images (Fig. 2) are used to study the UV emission which during the flare is probably dominated by the C IV doublet emission. The C IV formation temperature is 100 000 K.

2.3. MDI observations

The MDI (Scherrer et al. 1995) observations were obtained at 20:02 and 20:07 UT in a full-disk mode. The resulting images have 34′ × 34′ field-of-view (FOV). The angular pixel size
Fig. 1. BBSO images of the active region NOAA 9773 that produced the 2002 January 11 solar flare. They show the evolution in time of the flare in the Hα line. In each image the position of the flaring knot is marked with a rectangle that represents the size of the optical slit. The CDS area of integration is also indicated (thicker line rectangle). The arrow shows the position of a disappearing filament.

Fig. 2. The evolution of the 2002 January 11 solar flare as seen in the TRACE 1600 Å images. In each image, the position of the Lick slit is shown with a thin line rectangle. The CDS area of integration is also indicated (thick line rectangle).

is $1.98''$ pixel$^{-1}$. We used Level-1.5 calibrated magnetograms for this study. The magnetograms were additionally corrected for a flux saturation for magnetic flux above 1200 G (for more details see Green et al. 2003, and the citation therein). The images were rotated to the time 20:32 UT (as all the other SoHO and TRACE images used in this paper). They were also converted in order to correspond to the TRACE view point (Earth view).

2.4. CDS observations

The Coronal Diagnostic Spectrometer (CDS) (Harrison et al. 1995) onboard SoHO provided EUV spectral coverage of one of the flare kernels of the flare (see Figs. 1 and 2). Rasters of $20'' \times 202''$ made out of five adjacent 6 s exposures, using the $4'' \times 240''$ slit and pixel scale of $3.36''$ in the North-South direction, were obtained for four spectral windows.
centered at \( \text{Fe} \ xix \ 592.2 \ \text{Å} \), \( \text{Fe} \ xvi \ 360.8 \ \text{Å} \), \( \text{Mg} \ x \ 624.9 \ \text{Å} \), and \( \text{O} \ v \ 629.7 \ \text{Å} \). The time needed to build every raster was 84 s which resulted in an average cadence of 106 s between the data points. The position of the CDS slit over a kernel is noted in Figs. 1 and 2. CDS pointing is always given in advance to the observations while the Lick pointing was done during the flare onset.

Standard reduction was applied to the CDS data correcting for bias, flat-field, cosmic rays, and instrumental effects such as horizontal shifts due to the rotation of the scan mirror and rotation and tilt in the spectrum due to the misalignment between grating and detector, and grating and slit. Standard software was used in the data reduction and line fitting. An area of \( 8.1'' \times 6.7'' \) was tracked within the rasters from 19:43 UT to 21:00 UT and its emission integrated producing a lightcurve for each of the cited spectral lines (see Fig. 4 later).

2.5. Instrumental offset corrections

The pointing offset of the instruments used in the present study was corrected using the TRACE images as a reference. First, all the images used in this study were rotated to the time 20:32 UT. Then the BBSO \( \text{H} \alpha \) and TRACE 1600 Å images were coaligned using the sunspots as a reference point. A 20'' Solar_\( Y \) offset was found. The next step was to coalign TRACE and SOHO instruments, using MDI images. Both instruments recorded white light images which were used to determine the offset via cross-correlation. That offset was also used to compare TRACE and CDS. Finally, the entire TRACE dataset was visually inspected and compared for reassurance.

2.6. GOES

In addition, we used the GOES (1–8 Å) and (0.5–4 Å) pass-band data in order to compare the coronal and chromospheric behaviour during the flare and obtain a rough estimate of the energy released in this region.

3. The flare evolution

On 2002 January 11 we observed the flare event discussed here. The flare onset was first detected by the BBSO instrument at around 19:55 UT. The first optical spectrum recorded at Lick Observatory was at flare maximum at 20:07 UT. The flare took place in active region NOAA 9773 and was classified as C7.5. This implies, based on the soft X-ray (SXR) emission monitored by the GOES satellite, an emitted flux of \( 7.5 \times 10^{-6} \ \text{Wm}^{-2} \) at the Earth or \( 5.3 \times 10^{24} \ \text{ergs}^{-1} \) at the solar surface. This was therefore a medium size flare.

Active region NOAA 9773, N14W27, was classified following the Mt Wilson magnetic classification introduced by Hale et al. (1919) as a \( \beta\gamma\delta/\beta\gamma\delta \), which implies a bipolar group with an irregular distribution of polarities. The magnetic flux density in the active region ranged from \( -2000 \) to \( 3000 \ \text{G} \) at 19:02, 19:32 and 20:02 UT while the magnetogram at 20:32 UT gives values from \( -1480 \) to \( 1230 \ \text{G} \) (see Fig. 3). The total unsigned magnetic flux derived from the measured MDI longitudinal magnetic field decreased by about 30% from 20:02 to 20:32 UT. Note that this is the flux measured only in the FOV shown in Fig. 3.

The \( \text{H} \alpha \) images consist of several bright kernels during the impulsive phase and bright post-flare loops during the gradual phase. A filament was present several arcsecs north of the flare eruption from at least 18:02 UT and was still present at the start of the flare at 20:02 UT. By 20:32 UT it had disappeared. Unfortunately, the time resolution does not permit us to comment on its relation to the observed flare activity. The higher time resolution data are the TRACE observations.

All the optical spectral observations reported in this work were made on the marked kernel (see Fig. 2). The quiescent solar spectra were taken near the active region, but almost one
Fig. 4. Evolution of the flare as seen by GOES (top panel), CDS: Fe xix 592 Å, Fe xvi 360 Å, Mg x 624 Å and O v 629 Å (middle panel) and Lick: Hα, Ca II K and He I D3 lines (bottom panel).

day after the flare. This allows us to minimize the presence of small-scale structures other than the flare itself. However, a shorter interval between the flare and quiet spectra would have been desirable.

The lower panel in Fig. 4 shows the evolution of the Hα, He I D3 and Ca II K lines. All three lines have a similar behaviour, although He I is more impulsive-like. The photospheric lines (Mg I triplet) and metallic lines (Fe I) have their maximum enhancement roughly at the same time as the other chromospheric lines (see Table 1). The decline in the Ca II K line is delayed with respect to the Balmer lines.

As the CDS lines were taken slightly off-set from the optical flare kernel, direct comparison with the optical data is difficult. O V 629 produces a peak at around 19:57 UT, some 10 min before the coronal lines peak. Such behaviour has been seen many times from SMM data, where the transition region lines are seen to correlate to within 1 s of hard X-ray bursts of at least 50 keV (Cheng et al. 1981; Woodgate et al. 1983). CDS has also observed before such a precursor brightening with no coronal counterpart (Brosius 2003). It has been interpreted in terms of particle beams which deposit much more energy in the chromosphere and transition region than in the corona (see Brosius 2003; Brosius & Phillips 2004, for a further discussion), which would explain the absence of a coronal response. O V produces another two peaks, one at 20:07 UT and a further one at 20:29 UT. Mg x and Fe XVI also peak at 20:07 UT with a further increase in flux starting at 20:20 UT, giving a broad peak at 20:25 UT. The Fe XVI initially starts to increase in flux at 20:02 UT. The hottest line observed, Fe XIX shows a very different behaviour; first showing a peak at 20:02 UT and a second larger increase at 20:15 UT. That peak is simultaneous to the one observed in GOES data (20:16 UT) which starts rising at 20:04 UT. It is unclear whether the first peak seen in the Fe XIX line corresponds to the pre-flare heating as noted by other authors, e.g. Harra et al. (2001). The problem here is that CDS is observing the activity at a single flare kernel, while GOES observes the total flare integrated flux increase. A blend with a cool line does not seem to be the reason, because the precursor brightening does not show up in the Fe XIX lightcurve. The behaviour after the GOES and Fe XIX peak at around 20:16 UT can be better reconciled if we think about hot loops seen in Fe XIX which slowly cool down passing through different temperature regimes and being observed first in Fe XVI and Mg X, and later in O V, spanning around 20 min. The sharp increase in the Mg X lightcurve 20 min after the flare maximum (20:07 UT) is an artifact consequence of the solar rotation tracking and the dimensions of the CDS pixels. It shows up also in Fe XIX. That is not the case in the rest of the instances where solar rotation is corrected.

4. Chromospheric activity indicators

The wavelength range covered by our observations (3800–9000 Å), allows us to probe the behaviour of different chromospheric lines during the solar flare, and compare them with their stellar counterparts. Our solar observations are relative-intensity, so we have to transform them into fluxes. In order to make this transformation, we have first corrected for limb darkening (Pierce & Waddell 1961) at the position of the flare (20° from disk center) and then calibrated our observed continuum fluxes by reference to standard solar continuum fluxes from Allen (1976). We combined those values with the excess equivalent width (EW) of the Balmer lines shown in Table 1.

In order to obtain the excess EW, we first subtract the quiet solar spectrum from the flare spectra. This is the normal procedure followed in stellar flare observations where many of the lines are in emission even when quiescent. Then, we measure the EW with respect to a local continuum interpolated from a linear fit to line-free nearby continuum portions included within the spectral interval. In this analysis we consider positive and negative values of the EW as referring to lines in absorption and emission, respectively. Table 1 also shows the excess EW for other chromospheric activity indicators such as the Ca II lines, Mg I triplet and the He I D3 line. The fluxes for the Balmer lines (Hα to He) and Ca II are reported in Table 2.

4.1. The Balmer lines

Figure 5 shows the flare plus quiet Sun line profile, the quiet Sun profile, and the subtracted profile for the main chromospheric activity indicators observed during flare maximum, at 20:07 UT. We note that all the Balmer lines show a filling-in during flare maximum, which decreases as the flare evolves (see Table 1). We note that the peak intensity in Hα is greater...
Table 1. Flare observations and excess equivalent widths in angstroms for the activity chromospheric indicators observed during the solar flare at 20:07 UT on 11 January 2002. The sign of the EW follows the usual convention (negative values for emission lines).

<table>
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<th>Hβ</th>
<th>Hγ</th>
<th>Hδ</th>
<th>Hε</th>
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<tr>
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<td>-209</td>
<td>-165</td>
<td>-133</td>
<td>-108</td>
</tr>
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<td>-157</td>
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<td>-176</td>
<td>-137</td>
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<tr>
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Table 2. Fluxes in units of $10^6$ erg s$^{-1}$ cm$^{-2}$ at the solar surface for the Balmer, Ca II and He I lines observed during the solar flare at 20:07 UT on 11 January 2002.

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<th>Hγ</th>
<th>Hδ</th>
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<td>1.02</td>
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than in Hβ and Hγ. In the stellar case, the strongest Balmer line is usually Hγ or Hδ as reported for a flare on AD Leo (García-Alvarez et al. 2003). In solar flares, Hβ and Hγ are sometimes observed with a larger flux than Hα as reported by Zirin et al. (1982).

4.2. The Ca II H and K lines and the Ca II infra-red triplet

Enhanced emission cores in the Ca II H and K lines are the primary optical indicators of chromospheric activity. These lines are sensitive to changes in the electron density and temperature in the chromosphere. Figure 5 shows the Ca II H and K lines turning into emission during the flare. We also note that the Ca II H and K lines show central reversal. This is a common feature in solar plage and network features (Bocchialini & Gouttebroze 1996) and in stellar emission lines on active red dwarfs.

The second feature we observe in the Ca II H and K lines, is that the line shape is not symmetrical, with two peaks of different intensity, as if the reversal were not at the center of the line. This may be the signature of changes in the nature of the original source and/or in the state of the atmosphere. It may suggest that the rest frame of the line differs between the source region and the overlying atmosphere, with relative motions of a few km s$^{-1}$ along the line of sight. Apart from these features, we also note that, during the flare, the excess...
emission in the Ca II H line is, on average, around 20% larger than the Ca II K line (see Table 1).

The Ca II infra-red triplet (IRT) \( \lambda8498, \lambda8542, \) and \( \lambda8662 \) lines are another important stellar chromospheric activity indicator (Linsky et al. 1979; Dempsey et al. 1993; Gunn & Doyle 1997). These lines are formed in the lower chromosphere of the Sun (Vernazza et al. 1981). Foing et al. (1989) found that the Ca II IRT lines correlate well with the Ca II H and K peaks. Unlike the bluer Ca II lines, the Ca II IRT lies in a region with a well-defined continuum which makes the analysis simpler. Shine & Linsky (1972, 1974) reported that in the quiet solar atmosphere these lines are in absorption, while they show core brightening and eventually develop self-reversed emission when observed in bright plages. These authors concluded that the emission resulted from a steepening of the temperature gradient in the chromosphere. Figure 5 shows the Ca II IRT lines turning into emission during the flare. They keep this state well into the gradual phase. The Ca II \( \lambda8498 \) line is the strongest of all the IRT lines (see Table 1). It shows an excess emission around 39% and 43% larger than the Ca II \( \lambda8542 \) and \( \lambda8662 \) lines respectively. The excess emission of the Ca II \( \lambda8498 \) line is roughly the same as for Ca II H.

The decline in all the Ca II lines is delayed with respect to the Balmer lines. The rapid increase in the continuum and Balmer line fluxes observed during flares (García-Alvarez et al. 2002; Houdebine 2003) is consistent with an increase of the local electron density due to increased photo-ionization of Hydrogen (Houdebine et al. 1991). However, increases in the Ca II photo-ionization continua would deplete the Ca II ground state population. The increase in electron density creates larger collisional excitation rates that could balance the previous effect. These two effects would force a different behaviour of the Ca II lines compare to the Balmer lines as we obtained from the observations.

4.3. The HeCa I D3 line

All the changes observed in the spectral lines during the flare produce either filling-in or outright emission. He I D3 has been previously detected in absorption in active stars where it was attributed to plages or coronal radiation (Huemoerder 1986). It has also been seen in emission during stellar flares (Montes et al. 1997, 1999; Oliveira & Foing 1999). Landman (1981) studied high resolution spectra of the He I D3 line in solar plage regions. In the Sun, it is in emission in the strongest flares (Tandberg-Hanssen 1967) and appears in absorption in weak to medium flares and in eruptive prominences (Zirin 1988). Zirin (1975) reported that the mechanism for the appearance of the He I D3 line in absorption involves ionization by coronal X-rays and subsequent recombination. The He I D3 line is harder to excite than the Balmer lines, due to its higher excitation potential. Andretta & Giampapa (1995) suggested, on the basis of their NLTE calculations that some of the He I spectral features in the Sun and late-type stars could be linked to regions of enhanced UV and X-ray emission.

Johns-Krull et al. (1997) found a very distinct behaviour for this line, however, their analysis was for a much larger solar flare. Andretta & Giampapa (1995) found that this line can be formed in two different atmospheric regions: in the upper chromosphere at temperatures of 7000–8000 K and in a plateau above 20000 K. They conclude that in weaker chromospheres, formation in the upper chromosphere dominates.

In the present flare, we observe He I D3 in absorption during the flare, while it is undetectable in the quiescent state. Variability in its EW is clearly observed during the flare in overall agreement with the variability observed in other lines, although it is more impulsive-like.

4.4. The photospheric lines

Strong neutral metal absorption lines that are formed in the lower chromosphere/temperature minimum region are very interesting for solar-stellar activity studies. The Mg I triplet consists of low excitation lines emitted when the low chromosphere is heated, usually at foot-points of flux loops passing
through the flare. It is a good diagnostic of photospheric activity (Basri et al. 1989). A filling-in of these lines has been found in some RS CVn systems (Gunn & Doyle 1997; Montes & Ramsey 1998). This is observed in the last panel of Fig. 5. Together with the filling-in of Fe I lines, which are formed in the upper photosphere. This suggests that the flare affected the lower chromosphere. We found in the quiet Sun spectra no excess equivalent width for Mg I or Fe I. Figure 6 shows a time series of the analysed photospheric lines during the flare. Only a small filling-in was found for the Na I doublet. This is surprising since both neutral metals, Mg I and Na I, form in the lower chromosphere and, taking into account that the ionization energies are 7.6 and 5.1 eV respectively, one would expect the Na I Doublet to show a stronger filling-in than the Mg I Triplet. A similar behaviour was found by Gunn & Doyle (1997) in the RS CVn system ER Vul.

4.5. Line asymmetries

It is well known that the profiles of some strong chromospheric emission lines in the spectra of solar flares have asymmetries (see e.g. Svestka 1976, and references therein). It is important to measure these asymmetries during flares as they may allow us to measure the velocities of the emitting plasma. Blue and red asymmetry can be found in the same flare, depending on the phase. High temporal resolution is required in order to perform this kind of study. Several authors (i.e. Zarro et al. 1988; Canfield et al. 1990a; Falchi et al. 1992) have found red asymmetries and have interpreted this excess emission in the red wing as due to downward-moving chromospheric plasma. Canfield & Gayley (1987) tried to match the line profiles using a model that included chromospheric condensation with moderate success. Schmieder et al. (1987) reported blue-shifted Hα absorption during the gradual phase of three major two-ribbon flares, which was interpreted by the authors as due to upward chromospheric motion. Canfield et al. (1990b) found two cases of small regions with blue-shifted Hα, which occurred only early in the impulsive phase. Evidences of blue asymmetries have also been reported during solar flares (Heinzel et al. 1994). Red asymmetries have often been interpreted as the result of chromospheric downward condensations (Canfield et al. 1990a). In fact, recent line profile calculations (Heinzel et al. 1994; Ding & Fang 1997) show that the chromospheric downward condensations can explain both blue and red asymmetries.

In Fig. 5, we note that all the Balmer and Ca II lines have asymmetric profiles, at flare maximum, with red-shifted wings and blue-shifted cores. The line cores are shifted by about $-10$ km s$^{-1}$ from rest. The lack of blue shifted cores in the Ca II IRT$\lambda8498$, the weakest of the Ca II IRT, could be telling us something about height discrimination. Note that the photospheric lines during the quiet state and flare state match (see Fig. 5), suggesting that these asymmetries are unlikely to be caused by a shift in the detector. Additionally, at the end of the gradual phase the profiles become symmetric as they were during the quiet state. This behaviour in the line profiles could be explained by material expanding. The blue-shift in line center could be produced by the material closest to the observer moving upwards. The red-shift in the wings would then arise from material located further down in the solar atmosphere, moving downwards. We note that the Ca II H and K lines are the only ones that show self-reversals. The central emission itself is due to the chromospheric outward temperature rise, and the central dip is a non-LTE or scattering effect due to the dilution of radiation near the surface and the corresponding decrease in excitation. The Balmer lines do not have the central emission due to their lesser dependence on temperature. Detailed analysis and modelling of the variations observed in the chromospheric line profiles during the solar flare will appear in a subsequent paper (García-Alvarez et al. 2005).

4.6. The chromospheric and coronal flare energy budget

In order to estimate the energy released during the solar flare, we use the line fluxes given in Table 2. We then multiply those
Table 3. The energy encompassed by the slit and the total optical energy emitted during the solar flare at 20:07 UT on 11 January 2002.

<table>
<thead>
<tr>
<th>Line</th>
<th>Energy in slit (10^{27} erg)</th>
<th>Total energy (10^{28} erg)</th>
</tr>
</thead>
<tbody>
<tr>
<td>H\alpha</td>
<td>2.27</td>
<td>1.79</td>
</tr>
<tr>
<td>H\beta</td>
<td>1.12</td>
<td>0.88</td>
</tr>
<tr>
<td>H\gamma</td>
<td>0.84</td>
<td>0.66</td>
</tr>
<tr>
<td>H\delta</td>
<td>0.62</td>
<td>0.49</td>
</tr>
<tr>
<td>H\epsilon</td>
<td>0.53</td>
<td>0.42</td>
</tr>
<tr>
<td>Ca\textsc{ii} K</td>
<td>1.07</td>
<td>0.84</td>
</tr>
<tr>
<td>Ca\textsc{ii} H</td>
<td>1.52</td>
<td>1.19</td>
</tr>
<tr>
<td>Ca\textsc{ii} 8498 Å</td>
<td>1.29</td>
<td>1.04</td>
</tr>
<tr>
<td>Ca\textsc{ii} 8542 Å</td>
<td>0.95</td>
<td>0.73</td>
</tr>
<tr>
<td>Ca\textsc{ii} 8662 Å</td>
<td>0.89</td>
<td>0.70</td>
</tr>
<tr>
<td>Total</td>
<td>11.08</td>
<td>8.72</td>
</tr>
<tr>
<td>GOES (1–8 Å)</td>
<td>–</td>
<td>3.27</td>
</tr>
<tr>
<td>GOES (0.5–4 Å)</td>
<td>–</td>
<td>0.23</td>
</tr>
</tbody>
</table>

values by the projected slit size on the solar surface (3.97 \times 10^{17} \text{ cm}^2) and integrate the value over time in order to find the energy emitted in each line. These values are shown in Table 3. This accounts for the energy encompassed by the slit during the solar flare. In order to estimate the total energy radiated during the flare, we have to multiply those values by a scaling factor. The scaling factor, equal to 7.9, is obtained by comparing the solar flare. In order to estimate the total energy radiated during the flare, we have to multiply those values by a scaling factor. The scaling factor, equal to 7.9, is obtained by comparing the observed area of the solar flare in the H\alpha images (Fig. 1) and the area of the projected slit. We have assumed that the strength of every line, during the flare, is uniform over the whole area. The total energy emitted for each line, during the flare, is then obtained by multiplying the slit energies by this scaling factor. These results are shown in Table 3. We note that these energies are probably overestimates.

In addition, we also calculate the energy radiated in both GOES pass-bands during the flare. Integrating the radiated power between the times of our optical observations (20:07 to 20:27 UT), we obtain a total energy of \sim 3.5 \times 10^{28} \text{ erg}. Several authors obtained energy budgets for stellar flares simultaneously observed in the optical and the X-ray range (i.e., Butler et al. 1986; Doyle et al. 1988; Hawley & Pettersen 1991). The H\gamma and Ca\textsc{ii} K lines release a similar amount of energy during those stellar flares, \sim 10^{29} \text{ erg}.

The total energy emitted during the present solar flare by the H\gamma and X-ray emission follows the linear logarithmic correlation found by Butler et al. (1986) for stellar flares and extended for the solar case (Butler 1993). The continuum energy release, observed in almost all stellar flares, can be even larger than the soft X-ray energy released. However, most of the flares observed on the Sun, including the one analysed in this work, do not show a detectable increase in the continuum emission, due to the relatively strong photospheric solar background.

The above values for the radiative losses are only for the upper chromosphere or the corona. (Bruner & McWhirter 1988) have shown that, for a variety of solar plasmas, the total radiated power output, excluding hydrogen line radiation, may be estimated even if one knows the intensity of only one spectral line, e.g. O \textsc{v} 629 Å. Using the same time limits as the optical data and an area of 2.79 \times 10^{17} \text{ cm}^2, we derive a total power of 1.95 \times 10^{29} \text{ erg} for the temperature interval 4.0 \leq \log T_e \leq 8.0. This is approximately a factor of two greater than the H Balmer radiative losses, and thus would suggest that adding in the additional H losses, e.g. the Lyman lines, that the total radiated losses in H would be comparable to that radiative in the other lines in the upper chromosphere, transition region and corona.

5. Balmer decrement line modelling

Balmer decrements (flux ratio of higher series members to H\gamma) have also been used to derive plasma densities and temperatures in flare star chromospheres (Kunkel 1970; Gershberg 1974; Katsova 1990). Jevremovi\v{c} et al. (1998) developed a procedure (Balmer Decrement Fitting Procedure, BDFP hereafter) to fit the Balmer decrement in order to determine various physical parameters in the flaring plasma. García-Alvarez et al. (2002), based on the BDFP, obtained the first detailed trace of physical parameters during a stellar flare. These authors also applied the BDFP to literature data from a solar flare (Johns-Krull et al. 1997) and from a a stellar flare (Katsova 1990) obtaining values for the physical parameters in good agreement with those reported in those respective works.

The BDFP is based on the solution of the radiative transfer equation, using the escape probability technique (Drake 1980; Drake & Ulrich 1980) and a simplified picture of the flaring plasma as a slab of hydrogen with an underlying thermal source of radiation which causes photoionization. This source, which represents the photosphere, can be exposed to additional heating during the flare. This procedure assumes that the Balmer lines are formed in the upper chromosphere. We note that this simplified picture leaves out some important processes (i.e. cooling from Mg\textsc{ii}, C\textsc{ii} and other metals). The BDFP minimizes the difference between the observed and calculated Balmer decrement using a multi-directional search algorithm (Torczon 1991, 1992). This allows us to find the best possible solution for the Balmer decrement in a four dimensional parameter space where the parameters are: electron temperature, electron density, optical depth in the Ly\alpha line, and the temperature of the underlying source (background temperature). The best solutions for the Balmer decrement allow us to calculate the effective thickness of the slab of hydrogen plasma with which we are able to determine the surface area of the emitting plasma from the calculated emission measure. Although the temperature of the underlying source could be a free parameter in our code, we have fixed a lower limit at 5500 K for numerical stability reasons.

We apply the BDFP to the C 7.5 solar flare observed during our solar campaign. We normalised the flare-only fluxes for H\alpha, H\beta, H\gamma, H\delta and H\epsilon to H\gamma for three intervals around flare maximum.

In Fig. 7 we plot the observed and the fitted Balmer decrements for the solar flare. We note from the plots that, the slope of the Balmer decrement is shallower for the flare maximum.
than for the other two, as seen in stellar flares. We applied the BDFP to the Balmer decrements observed around flare maximum. The electron temperatures we obtain are around 8000 K for the first five Balmer decrements observed during the flare. The electron densities range from $5 \times 10^{15}$ to $7 \times 10^{15} \text{ cm}^{-3}$, roughly the same as García-Alvarez et al. (2002) obtained for the solar flare observed by Johns-Krull et al. (1997). The electron densities found for the solar flare imply that the Balmer emitting plasma originates somewhere in the chromosphere, where the electron density range from $\sim 10^{15} \text{ cm}^{-3}$ to $\sim 10^{12} \text{ cm}^{-3}$ (Gabriel & Mason 1982). The effective thickness of the slab of hydrogen plasma around flare maximum is $\sim 100 \text{ km}$. The value for Log $\tau(\text{Ly}-\alpha)$ are around 4.5, which means that the plasma is optically thick in this line.

We have used the best fit solutions for the Balmer decrement of the solar flare combined with the derived flux for the H$\beta$ line, to calculate the surface area of the emitting plasma. The value we obtain is $4.9 \times 10^{18} \text{ cm}^2$. This value is roughly the same as that inferred from the BBSO image obtained during flare maximum ($3 \times 10^{18} \text{ cm}^2$). Comparing the values for the physical parameters obtained here (C 7.5 flare) with those obtained by García-Alvarez et al. (2002) for the M 7.7 flare observed by Johns-Krull et al. (1997), we find slightly larger values for all the model parameters in the stronger flares.

6. Summary

We have presented a multi-wavelength analysis of a solar flare observed at 20:07 UT on 11 January 2002 which took place in the active region NOAA 9773. The main results of this work are:

1. The flare shows several eruptive centers in both the H$\alpha$ and Ca II K images. The H$\alpha$ and Ca II K images consist of two components, a bright kernel during the impulsive phase and bright loops during the gradual phase.

2. The main chromospheric activity indicators show either filling-in (Balmer lines) or emission (Ca II lines) during the flare. Photospheric lines such as the Mg I triplet and Fe I lines also show filling-in during the event. This may indicate that the flare reaches the lower chromosphere. The He I $D_2$ line appears in absorption during the flare, being undetectable during the quiet state.

3. We observe that all the Balmer and Ca II lines show asymmetric profiles, with red-shifted wings and blue-shifted cores. This behaviour could be explained by material expanding. The blue-shift in the line center could be produced by the material closest to the observer moving upwards. The red-shift in the wings would then arise from material located further down in the solar atmosphere, moving downwards.

4. We observe that the Ca II lines decline more slowly than the Balmer lines, as is observed in M dwarf flares. The photospheric lines (Mg I Triplet) and metallic lines (Fe I) have their maximum enhancement roughly at the same time as the chromospheric lines, which suggests that the energy supply to the photosphere would be the same. The temporal behaviour suggests that all the chromospheric lines were formed in the same atmospheric region.

5. In the EUV lines, observing a close but different kernel to the Lick pointing, we detect a precursor brightening in the transition region O V line peaking 10 min before the flare maximum. After the flare, the evolution shows a typical pattern of hot loops cooling down from $8 \times 10^6 \text{ K}$ (Fe XIX line) to transition region (250 000 K) temperatures in an interval of around 20 min.

6. Although the estimates for the energy released in the optical lines and X-rays are crude, a total flare energy in excess of $10^{30} \text{ erg}$ was derived. The optical flare showed no detectable increase in the continuum emission, as happens in most of the flares observed on the Sun, probably due to the relatively strong photospheric solar background.

7. We have determined the main chromospheric physical parameters for the flare maximum. We found electron densities of $\sim 10^{14} \text{ cm}^{-3}$ and electron temperatures of $\sim 8000 \text{ K}$ around flare maximum implying that the Balmer emitting plasma originates somewhere in the chromosphere, rather than in the photosphere. The modelling method also yields an estimate of the flare emitting area ($4.9 \times 10^{18} \text{ cm}^{-2}$) which is roughly the same that the observed value ($3 \times 10^{18} \text{ cm}^{-2}$).

There is an obvious need for more solar flare observations of this kind to be compared with stellar flares, particularly if one can obtain simultaneous coverage with UV and X-ray satellites. Furthermore, on the Sun, due to its proximity, we can study flares in a very detailed way in all the wavelength ranges (optical, X-ray, etc.). So, observing the Sun in a similar way to which we observe the stars will, hopefully, give a better understanding of the unresolved processes causing the observed stellar activity.

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