Blank field sources in the ROSAT HRI Brera multiscale wavelet catalog

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ABSTRACT

The search for Blank Field Sources (BFS), i.e. X-ray sources without optical counterparts, paves the way to identification of unusual objects in the X-ray sky. Here we present four BFS detected in the Brera Multiscale Wavelet catalog of ROSAT HRI observations. This sample was selected on the basis of source brightness, distance from possible counterparts at other wavelengths, point-like shape and a good estimate of the X-ray flux \(f_X\). The observed \(f_X\) and the limiting magnitude of the optical catalogs fix a lower limit for our BFS on \(f_X/f_{\text{opt}} \lesssim 40\). This value puts them well beyond the 90\% threshold for usual source classes once the HRI energy band and proper spectral shape are taken into account, leaving room for speculation on their nature. Three BFS also show evidence of transient behaviour.

Key words. X-rays: general – stars: neutron – X-rays: galaxies: clusters – galaxies: AGN – X-rays: galaxies: quasars: general

1. Introduction

It is well known that one can discriminate among classes of X-ray sources, thereby evaluating the ratio between their fluxes in X-ray and optical bands (e.g. Maccacaro et al. 1988; Motch et al. 1998; Zickgraf et al. 2003), especially if some spectral information is supplemented. Blank Field Sources (BFS), i.e. X-ray sources with no optical counterparts, stand somehow apart from the bulk of X-ray zoology in that they represent rarer, less studied objects (e.g. Cagnoni et al. 2002). In particular, known X-ray sources are concentrated at not too high values of \(f_X/f_{\text{opt}}\); e.g., stars have \(f_X/f_{\text{opt}} \lesssim 1\), galaxies and AGN reach \(f_X/f_{\text{opt}}\) of a few tens. At higher values of \(f_X/f_{\text{opt}}\), peculiar populations start to show up: some examples are isolated neutron stars, type 2 quasars, extreme BL Lacs, or high-redshift clusters of galaxies.

The search for BFS aims at objects with high \(f_X/f_{\text{opt}}\), where the observed X-ray flux and the upper limit from the absence of an optical detection establish a lower limit on the \(f_X/f_{\text{opt}}\). A careful selection of bright BFS is therefore a powerful tool for finding population leftovers from usual source classes. The status of BFS is somehow a transitory condition that expresses our lack of knowledge of the real nature of these sources. New X-ray observations can provide temporal and spectral information on BFS, while deeper optical/near IR pointings can either reveal the counterpart or establish a much stronger lower limit on the \(f_X/f_{\text{opt}}\). With the aid of this information, BFS can be identified and classified properly.

In this paper we present a sample of luminous BFS from the ROSAT HRI Brera Multiscale Wavelet catalog (BMW–HRI: Panzera et al. 2003; Campana et al. 1999; Lazzati et al. 1999), with \(f_X/f_{\text{opt}} \gtrsim 40\), along with archival X-ray spectral data from other observations (when available). This is the first study to exploit all ROSAT HRI data for a BFS search (see Musso et al. 1998 for the first search on a much smaller dataset). A key factor in the identification process is the uncertainty of the X-ray position (see Cagnoni et al. 2002; Rutledge et al. 2003); in this respect, using HRI rather than ROSAT PSPC images for searching BFS is superior.

The paper is organized as follows: in Sect. 2 we briefly describe the BMW–HRI catalog and its cross-identification program with optical, infrared, and radio catalogs. In Sect. 3 we introduce the method of sample selection, while the final sample of BFS is presented in Sect. 4. In Sect. 5 we discuss the possible nature of BFS and future search plans are mentioned in Sect. 6.

\textsuperscript{1} Actually this value is \(f_{0.5-2\text{kev}}/f_{\text{BJ}}\), see Sect. 2.
2. The BMW–HRI catalog

2.1. Catalog description

The BMW–HRI catalog (Panzera et al. 2003) has been derived from an analysis of the ROSAT HRI data set with a source detection algorithm based on the wavelet transform (Lazzati et al. 1999; Campana et al. 1999).

The up-to-date version of the catalog contains 28,998 entries, down to a count rate of $\sim 10^{-4}$ cts s$^{-1}$. The total sky coverage is 732 deg$^2$ ($\sim$1.8% of the sky). For each entry name, position, count rate, extension, and relative errors are provided, along with derived parameters like flux and Galactic column density and ancillary information about the pointing. Furthermore, results of cross-correlations with GSC2, 2MASS, IRAS, and FIRST catalogs are reported (see Sect. 2.2 and Table 1). In the following, some of the basic characteristics of the catalog are discussed in order to permit understanding of the BFS selection method. Further details can be found in Panzera et al. (2003).

Extension. The wavelet transform is particularly powerful in dealing with extended sources. In practice, each detected source is characterized by its extension, along with the more usual parameters (count rate and position). Extended sources, especially when observed at large off-axis angles, are often broken in various point-like sources by other detection algorithms. This risk is minimized by the wavelet transform.

Detection threshold and spurious sources. Each BMW–HRI source is also characterized by a detection probability value (peak significance). This value is the probability, in units of Gaussian $\sigma$, that the source is not spurious (i.e. a background fluctuation). This probability is the result of numerical simulations of random fields (without sources), repeated for different wavelet transform scales; it takes the number of background peaks mismatched as sources by the detection algorithm into account, for a given signal to noise ratio threshold in wavelet space. The mean number of spurious sources expected over an image is kept constant ($0.4$ spurious detections per field), varying the detection threshold. That is, the detection threshold depends on the mean background count value on the image and on the scale at which the search is performed (see Lazzati et al. 1999). The mean threshold value in the overall catalog is $\sim 4.2\sigma$.

Errors. In theory, the error for each parameter found with the wavelet transform algorithm can be estimated from the covariance matrix. In practice, such errors are reliable only when the number of source and background counts are sufficiently high ($\geq 2 \times 10^2$ counts per pixel). Otherwise, the distribution of the wavelet transform coefficients becomes Poissonian rather than Gaussian (in the lowest scales). If this is the case, a better estimate of the errors can be given by means of basic statistics (such as the error on the number of counts $N = \sqrt{N}$; for more details see Lazzati et al. 1999). The errors reported here and in the catalog are the maximum between the covariance matrix and the statistical estimates by Lazzati et al. (1999).

Boresight correction uncertainty. An additional source of errors on the absolute position determination of HRI sources is the uncertainty on the boresight correction. If the alignment between the node and the telescope optical axis is perfect, on-axis images are exactly in the center of the instrument. However, in the real situation there is always a finite misalignment that needs to be corrected. Unfortunately, uncertainties in the aspect solution, which describes the orientation of the telescope as a function of time, and errors in the alignment between the star trackers and the ROSAT X-ray Telescope introduce uncertainties in the boresight correction of variable size for each observation. In practice, the systematic offset between accurately known optical positions and X-ray positions can be used to evaluate the extent of boresight correction uncertainties. For the ROSAT HRI, the offset can be as large as 10″, even if in most cases it will be much less (David et al. 1998). The 10″ value is conservatively assumed as a fiducial value for the boresight offset. Usually, the statistical errors in the position from the detection algorithm, calculated as in the previous section, are much less than 10″ and therefore can be neglected in a first approximation.

Multiple detections. As each pointing is treated separately, not all of the catalog entries correspond to independent sources. An estimate of the overall number of independent sources can be given compressing the catalog in a 10° radius, and again the fiducial error value is indicated by the boresight uncertainty. This procedure selects only one source for each 10° cone radius, on the basis of an autocorrelation of the position. With this procedure, 20,433 sources are left. Obviously, the compression leads to loss of sources truly close to each other.

2.2. Multiwavelength catalog cross-correlations

In the following, we will largely use the term off-band to indicate a wavelength passband that is different from X-rays. The limiting flux of a given catalog in a given band will be generically denoted $f_{\text{lim,band}}$. For each of the BMW–HRI entries, cross-correlations with GSC2 (McLean et al. 2005), 2MASS Second Data Release (Kleinmann et al. 1994), IRAS Point Source Catalog (Beichman et al. 1988) and FIRST Survey Catalog (White et al. 1997) were performed by

<table>
<thead>
<tr>
<th>Catalog</th>
<th>Survey</th>
<th>Band</th>
<th>Coverage</th>
<th>Depth</th>
</tr>
</thead>
<tbody>
<tr>
<td>GSC2 (prel.)</td>
<td>POSSII</td>
<td>$BJ$</td>
<td>Dec: $&gt; 0$</td>
<td>22.5</td>
</tr>
<tr>
<td>2MASS SDR (2000)</td>
<td>$\ast$</td>
<td>$J$</td>
<td>50% sky</td>
<td>15.8</td>
</tr>
<tr>
<td>2MASS SDR (1989)</td>
<td>$H$</td>
<td>50% sky</td>
<td>15.1</td>
<td></td>
</tr>
<tr>
<td>IRAS PSC (1989)</td>
<td>$K$</td>
<td>50% sky</td>
<td>14.3</td>
<td></td>
</tr>
<tr>
<td>FIRST SC (2001)</td>
<td>FIRST 1.4 GHz</td>
<td>20% sky</td>
<td>$\sim 1$ mJy</td>
<td></td>
</tr>
</tbody>
</table>

Work in progress.
Panzer et al. (2003). Off-band catalog properties are summarized in Table 1. The adopted radius for the cross-correlation is 10′′ (see Sect. 2.1), assumed as positional X-ray uncertainty. In fact, the positional uncertainties for GSC2, 2MASS SDR and FIRST catalogs are <0.5′′ (3σ, McLean et al. 2005), <0.5″ (1σ, see the 2MASS documentation for an extensive discussion), <0.5′′ (90%, McMahon et al. 2001), respectively, so that they can be safely neglected for cross-identification purposes. An exception is the case of IRAS PSC, for which the positional accuracy varies with source size, brightness, and spectral shape; and it is different in different directions, but is usually better than 20′′ (see Beichman et al. 1988), so this last value has been used as cross-correlation radius with this catalog. Note that: a) when two or more entries in the correlating catalog are found, only the brightest source parameters are reported (but in any case the number of found sources is given), b) the sky coverage of the off-band catalogs is usually not complete (again see Table 1). In particular, the preliminary, unpublished version of the GSC2 (GSC2.3) catalog used in Panzer et al. (2003) lacks coverage of the zones where a bright source caused an overexposure of the Schmidt plates. For the 2MASS and FIRST, more complete catalogs are now available (2MASS All Sky Data Release, FIRST Survey Catalog 03Apr11 Version). Our final list was checked with them.

3. Sample selection

3.1. Catalog mask

The first sample selection was made on the basis of the following criteria:

a. blank field sources, i.e. without any cross identification in the other databases (see Sect. 2.2 and Table 1): 6061 catalog entries;
b. point-like sources (significance of the extension = 0): 4955 catalog entries;
c. bright sources, i.e. with \( f_X \geq 2.7 \times 10^{-13} \text{ erg s}^{-1} \text{ cm}^{-2} \): 275 catalog entries.

The reported numbers are catalog entries, i.e. no compression was applied to eliminate multiple detections (see Sect. 3.2 below). The \( f_X \) we used was computed from the observed source counts, in the 0.5–2 keV range, considering only channels 2–9 (e.g. David et al. 1998). The assumed spectral shape is a power law with photon index \( \Gamma = 2 \) (i.e. the same of the Crab pulsar and nebula) and column density \( 5 \times 10^{19} \text{ atoms cm}^{-2} \) (parameter # 48 in the BMW-HRI catalog). The chosen model for the count rate-flux conversion uses constant negligible column density, instead of the integrated Galactic one in the source direction. While the latter is commonly adopted in extragalactic surveys, it would introduce a larger error on low-latitude galactic BFS, which are expected to be low-luminosity, nearby objects (see Sect. 5). Furthermore, the chosen \( f_X \) derived with this choice is lower than the one with Galactic \( N_H \), so the obtained \( f_X/f_{\text{offband}} \) limit is lower. We chose to select sources on the basis of \( f_X \) and not of, for example, \( f_X/f_{\text{BJ}} \). This leads naturally to different \( f_X/f_{\text{offband}} \) ratios for sources in different parts of the sky, due to the variable depths of the surveys available to us. However, in such a way the selection is based on a directly observed quantity, the number of counts, not on the basis of a non-detection, and it provides a more homogeneous sample for future observations. We carefully evaluated the \( f_X/f_{\text{offband}} \) keeping into account the different depths and photometric band used in the different catalogs. The obtained lower limit for the \( f_X/f_{\text{offband}} \) from the GSC2 catalog are \( f_X/f_{\text{BY}} > 37 \) (Northern sky), \( f_X/f_{\text{BY}} > 59 \) (Southern sky), \( f_X/f_{\text{RF}} > 43 \) (Northern sky), \( f_X/f_{\text{RF}} > 129 \) (Southern sky). A detailed description of method and used values can be found in Chieregato (2005).

In contrast with the majority of similar studies, our sources are blank at various wavelengths (see again Table 1). We also note that, despite selecting point-like sources, we cannot exclude that the real nature of some source is extended (i.e. only the brightest peak has been detected with HRI).

3.2. Multiple detection removal

The removal of multiple detections from the sample could already be done using automatic compression criteria (position based), as in Sect. 2.1. However, given the limited entity of our sample of BFS candidates (275 elements), we preferred to remove multiple detections by hand to avoid, when possible, the loss of sources truly close to each other.

We used positional coincidence as the main criterion for identifying multiple detections. Another useful piece of information was the flux level. Unfortunately, positional coincidence can be weakened by the boresight uncertainty, while obviously the flux criterion is not fulfilled for highly variable sources. In practice multiple detections were reasonably identified in most of the cases. After multiple detections removal, we are left with 226 catalog entries.

3.3. Total counts cut

We made a further screening on the number of total counts, rejecting sources with less than 25 photons (on the original detection, i.e. before applying vignetting and Point Spread Function loss corrections). This was done in order to strengthen the derived source parameters, obviously, due to the statistics increase. After the total counts cut, we are left with 201 sources.

3.4. Are the BFS sources real?

We re-evaluated the number of spurious sources expected in this 201 subsample, again on the basis of simulations of random fields with the above masks. In particular, the flux cut is very effective in the removal of spurious sources. In fact, the expected number of spurious sources is ~2 in the subsample of 201 (see Moretti et al. 2004). Our BFS are, therefore, likely not to be fake detections.

3.5. Visual inspection: when a blank is not a blank

We inspected Digitized Sky Survey 2 (DSS2; McLean et al. 2000), Super Cosmos Sky (SSS; Hambly et al. 2001), 2MASS, FIRST, Sydney University Molonglo Sky Survey (SUMSS;
Bock et al. 1999), NRAO VLA Sky Survey (NVSS; Condon et al. 1998), and Westerbork Northern Sky Survey (WENSS; Rengelink et al. 1997) images for each of the 201 sources left. Furthermore, we cross-checked source positions with Simbad, NED, and with the VIZIER catalogs (in particular the USNO-B1 catalog, Monet et al. 2003). Single sources and fields already known and well studied or too complex have been excluded from the sample (e.g. Magellanic Clouds, M31, Orion Nebula). We also excluded sources positionally coincident with bright stars (holes in the GSC2 coverage), or at the periphery of optically extended emission (more than 10″ from the centroid and so escaping the automatic cross-identification). We found, and excluded, some cases in which our X-ray sources coincided with reliable USNOB1.0 entries (the criteria for the GSC2 detection are somewhat more conservative), or with counterparts at other wavelengths (IRAS Faint Sources Catalog, Moshir et al. 1990). We prefer not to be too severe at this stage, deferring strict cuts after the boresight correction. 72 sources survived this phase.

3.6. Cheshire cat: elimination of fluctuations

We put aside those primary sample sources which were not detected in a longer HRI or ROSAT PSPC pointing (the PSPC had the same passband but a substantially larger effective area than the HRI). These could be variable sources detected at peak flux, but they are also candidates for false detections. No source was inserted in this subclass on a non-detection basis with other instruments, because of the different passband; i.e., given the absence of spectral information on our sources, the non-detection could be due to a peculiar spectral shape. For some sources the suspicion of a false detection is strengthened by the number of non-detections or by the respective length of the pointings. However, there are cases for which source parameters like detection probability and number of counts seem to indicate a real source. Furthermore, as discussed in Sect. 3.4, we expected only ~2 spurious detection in 201 sources. Of the 72 sources, 16 were placed in this transient candidate subsample, while for the other 56 there was no evidence of a transient nature.

3.7. Pin-pointing the sources to the sky: boresight correction

We performed some astrometry on the remaining fields including those of transient candidates, in order to get rid of the 10″ fiducial boresight uncertainty and therefore to fully exploit the HRI angular resolution capability. Furthermore, the error given by the detection algorithm (see Sect. 2.1), while not being entirely statistical, can as a first approximation be treated as random, so that the usual Gaussian relations can be used. In contrast, the offset given by the boresight uncertainty varies in a random way in the ensemble of all the pointings, but is systematic in nature for all the sources in the single pointing, weakening the rejection of distant optical associations to BFS.

For each pointing, we matched X-ray sources to optical counterparts (optical positions of known X-ray emitters), if any, or to optical catalog sources distant less than 10″. We excluded X-ray extended sources and sources with no optical catalog entry in 10″. Ambiguous cases, i.e. with two or more optical sources present, were treated individually using the distance and the optical luminosity as criteria for the identification and then checked a posteriori. Even if only another X-ray source was present, the shift for this source was applied to the pointing, although obviously in these cases the corrected positions and uncertainties have to be taken cum grano salis. Generally, the new error is less than 10″; even if it remains around 10″, the new position should be free of systematic boresight uncertainty. Note that performing the boresight correction only after the selection of Sect. 3.1 can bring to the loss of BFS. However, this allows us to deal with a limited number of sources and boresight corrections.

4. The data

The elimination of sources with an off-band counterpart in a 4σ radius from the boresight corrected position left us with our final sample of four sources (three of them transient candidates). This 4σ limit assures us that starting with 1340 sources (i.e. sources outside the inner 3′ in the BMW–HRI catalog and obeying to the cuts above unless lacking of counterparts), only 0.08 of them would lie outside the error region by chance.

In summary, our final BMW–HRI Blank Fields Sources are:

- Blank: no counterpart at other wavelengths in a 4σ radius from the boresight corrected position.
- Bright: \( f_X \geq 2.7 \times 10^{-13} \text{ erg s}^{-1} \text{ cm}^{-2} \). Note that from the GSC2 magnitude limit, this is equivalent to \( f_X / f_{B1} \), respectively, \( \geq 37 \) and \( \geq 59 \) for Northern and Southern sources.
- Point-like: extension significance is 0.
- Well-detected: total (uncorrected) counts \( \geq 25 \).

After this selection we end up with one persistent BFS, plus three transient BFS. We report the new positions and error radii (as obtained from the boresight correction procedure), the flux, the detection probability (i.e. the probability of detecting a background fluctuation as a source), the total number of counts, the Galactic coordinates, the integrated Galactic column density, the distance from the nearest off-band association (in terms of error radii), the lower limit on the \( f_X / f_{B1} \), and the upper limit on the radio emission (see Table 2). Note that all but one (1BMW200739.8-484819) of the sources have a radio flux upper limit though with different depth.

4.1. Properties of BFS

The final error radius estimate for our sources is quite different from case to case, ranging from \(~3"\) to \(~7"\); the latter large error is due to the large off-axis angle and to the low number of counts (see Fig. 1). In fact, all our sources are quite offset from the center of the pointing (\(>15"\)). This depends on the selection procedure used and needs some comment. First, we concentrate on the consequences of the requirement of absence of offband counterparts in a 10″ radius. This strongly affects the overall distribution in detector coordinates. In fact, the fiducial 10″ radius for the initial correlation, chosen to match the
Table 2. BFS parameters.

<table>
<thead>
<tr>
<th>Name</th>
<th>RA J2000</th>
<th>Dec J2000</th>
<th>bII</th>
<th>Err. rad.</th>
<th>Flux erg cm$^{-2}$ s$^{-1}$</th>
<th>Prob. $\sigma$</th>
<th>Cts</th>
<th>$N_H$ 10$^{20}$ cm$^{-2}$</th>
<th>$f_X/f_{opt}$</th>
<th>Close ass. (sigma)</th>
</tr>
</thead>
<tbody>
<tr>
<td>1BMW042142.4-571541</td>
<td>04 21 42.6</td>
<td>-57 15 39</td>
<td>-42.54</td>
<td>2.9</td>
<td>6.5x10$^{-13}$</td>
<td>14.0</td>
<td>742</td>
<td>1.8</td>
<td>&gt;141</td>
<td>5.3</td>
</tr>
<tr>
<td>1BMW135703.0+181122</td>
<td>13 57 02.9</td>
<td>+18 11 21</td>
<td>72.45</td>
<td>7.1</td>
<td>3.5x10$^{-13}$</td>
<td>4.2</td>
<td>112</td>
<td>2.1</td>
<td>&gt;47</td>
<td>6.3</td>
</tr>
<tr>
<td>1BMW200739.8-484819</td>
<td>20 07 39.8</td>
<td>-48 48 19</td>
<td>-32.31</td>
<td>5.8</td>
<td>3.0x10$^{-13}$</td>
<td>4.3</td>
<td>55</td>
<td>5.1</td>
<td>&gt;65</td>
<td>5.2</td>
</tr>
<tr>
<td>1BMW043306.8+155307</td>
<td>04 33 06.8</td>
<td>+15 53 07</td>
<td>-21.16</td>
<td>6.5</td>
<td>2.9x10$^{-13}$</td>
<td>4.2</td>
<td>34</td>
<td>17.0</td>
<td>&gt;40</td>
<td>4.9</td>
</tr>
</tbody>
</table>

Fig. 1. BMW–HRI position error: the statistical errors on position of BMW–HRI sources for the different number of counts as a function of the off-axis angle. Error bars are at 2$\sigma$. The change in the slope of the curves corresponds to angular resolution loss at the change of image rebinning. An offset of 0.2' and 0.4' has been applied to points in the 50 counts bin and 100 counts bin, respectively, for visualization purposes.

The worst case boresight error, is huge in comparison to the statistical positional error for the innermost sources, but it corresponds to the intrinsic error for the outermost sources (see Fig. 1). Sources in the innermost part are preferentially associated with optical counterparts, which may not be the true ones, since they have been searched for over a region much larger than their error boxes. The overall effect is to deplete the distribution from inner sources, enhancing the proportion of outer sources. This bias is present independently of the real size of the boresight correction error. Second, there is a trend overestimating outer source counts (especially at low values), since it is favorable to detect them on top of a positive background fluctuation (Eddington bias, see Hasinger et al. 1993; Moretti et al. 2002). This again adds on the enhancement of outer sources. The first bias leads to the loss of BFS, but it does not affect our selected sources. The effect of the second bias on the estimated flux can be evaluated a posteriori from the duration of the pointings and the number of counts. Very conservative assumptions give a factor of ~2 as the maximum overestimation.

4.2. 1BMW042142.4-571541

This source is the brightest of our sample and the one with the highest number of detected photons. The good statistics involve a small detection algorithm error radius, so the position determination is quite accurate (see Fig. 2). The boresight correction was performed with the known position of the quasar HE 0419-5657 and the nucleus of NGC 1547 as well as with eleven other sources matched to optical catalogs. The source is close (~15'') but actually more than 5$\sigma$ to a bright star (BJ = 13.31, $F = 11.61$ in the GSC2.2), which makes the lower limit on the $f_X/f_{opt}$ less reliable, in particular in the red band. There are two fake USNO B1.0 sources along the saturation spikes of the bright star (see Fig. 2). The source was observed with Einstein (1E 0420.7-5723) and with the ROSAT PSPC (1WGAJ0421.7-5716) and was also detected as a RASS Bright Source (1RXS J042144.0-571601). The WGA computed flux is $8.8 \times 10^{-13}$ erg cm$^{-2}$ s$^{-1}$, consistent with the RASS count rate. Taking the caveats on the count rate to flux conversion into account, this could mean that there is no evidence for long term flux variations.

We re-extracted the archival PSPC observation of this source (sequence rp700034m00). The source was very close to the PSPC rib; despite this, we attempted a rough spectral analysis. We extracted source photons from a circle of radius 200'' centered on source position. We used a circle of radius 240'' as background region, centered in an empty region south of the source. The collected counts were 321. We rebinned channels by a variable factor in order to have at least 20 photons in each bin. Channels 1–11 and 136–256 were ignored. Figure 3 shows the resulting spectrum, which we then
Table 3. Spectral fits for PSPC observation of 1BMW042142.4-571541.

<table>
<thead>
<tr>
<th>Model</th>
<th>$N_H$ $10^{20}$ cm$^{-2}$</th>
<th>$kT/T$ keV</th>
<th>Flux (0.1–2 keV) $10^{-13}$ cgs</th>
<th>$x^2_{\text{red}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>B.body</td>
<td>0.0$^{+0.3}$</td>
<td>0.19$^{+0.02}_{-0.03}$</td>
<td>5.6 1.4</td>
<td></td>
</tr>
<tr>
<td>B.body 1.8(frozen)</td>
<td>0.15$^{+0.04}_{-0.03}$</td>
<td>6.9 1.9</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Pow.-law</td>
<td>1.8$^{+0.8}_{-0.7}$</td>
<td>1.8$^{+0.2}_{-0.3}$</td>
<td>6.5 1.3</td>
<td></td>
</tr>
<tr>
<td>Pow.-law 1.8(frozen)</td>
<td>1.8$^{+0.2}_{-0.3}$</td>
<td>6.5 1.2</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Bremss. 1.3$^{+1.0}_{-0.9}$</td>
<td>1.46$^{+0.17}_{-0.14}$</td>
<td>6.2 1.3</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Bremss. 1.8(frozen)</td>
<td>1.03$^{+0.25}_{-0.29}$</td>
<td>6.2 1.2</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

Fig. 3. Power law spectral fit of the PSPC observation of 1BMW0421.4-571541.

fitted with XSPEC (v.11.2). Due to the poor statistics, we considered a limited number of models, as shown in Table 3. The spectrum is definitely soft, without evidence of absorption excess. The best fit absorbed blackbody model tends to null absorption, with a 90% upper limit on column density of $\sim 3 \times 10^{20}$ atoms cm$^{-2}$. The power law fit gives a relatively better reduced $\chi^2$ (1.3 versus 1.4). Furthermore, the best column density value is quite close to the integrated Galactic value of $1.2 \times 10^{20}$ atoms cm$^{-2}$ (see Fig. 3). Freezing absorption at the Galactic value, the reduced $\chi^2$ does not change appreciably. The bremsstrahlung fit gives the best reduced $\chi^2$ (1.3) and the column density is $\sim 1.3 \times 10^{20}$ atoms cm$^{-2}$.

4.3. Transient BFS

In this section three sources are presented (see Table 2) that, while fulfilling all the other criteria for BFS, show additional strong evidence of transient behaviour. In fact, despite their brightness in the BMW–HRI catalog ($f_X \geq 2.7 \times 10^{-13}$ erg s$^{-1}$ cm$^{-2}$), for each of them there is at least another ROSAT HRI or PSPC pointing where they have not been detected (deeper than the HRI exposure that revealed them).

1BMW135703.0+181121. 1BMW135703.0+181121 was detected in a $\sim 5.5$ ks HRI pointing in Jan. 1998; we included it in the possible transient subsample since it was not detected in three longer PSPC pointings: $\sim 11.4$ ks, Jan. 1992; $\sim 8$ ks, Jul. 1990; $\sim 5.6$ ks, Jul. 1992; also it was not seen in a shorter $\sim 1.1$ ks HRI pointing in Jan. 1992. The high number of counts detected points toward a real, transient source. The boresight corrected position was obtained with three sources. There are no counterpart at other wavelength in a 45″ radius ($\sim 6.3\sigma$, see Fig. 4).

1BMW200739.8-484819. This source was detected in a $\sim 3.2$ ks pointing in Oct. 1996, while no corresponding source was observed in four other pointings of comparable length (2.6–3.6 ks) at distance of $\sim$ days, nor in two $\sim 11$ ks ROSAT PSPC pointings in Nov. and Apr. 1992. The closest optical catalog source is at more than 30″ (see Fig. 5) from the boresight corrected X-ray position (the boresight correction was performed with the other two X-ray sources detected in the Oct. 1996 HRI pointing). Two objects, A and B in Fig. 5, are clearly visible in the blue SSS image, at $\sim 16''$ ($\sim 2.7\sigma$) and $\sim 13''$ ($\sim 2.3\sigma$), respectively. Object B (BJ = 22) is reported to be extended, while object A (BJ = 22.3) is possibly extended too. Oddly, while nothing is apparent at that position in the red SSS image, a source can be seen both in the shallower ESOI RED image and in the SSS I image, at a position almost coincident with object B with RF = 20.2 (more than 2 mag brighter than in the Schmidt plate limit) and IN = 19.0. An obvious hypothesis could be the presence of...
5. Discussion: nature of the blank fields

In order to proceed further with our hypotheses on the BFS nature, we need to compare the $f_X/f_{B1}$ limits of BFS with existing classification schemes. We chose the reference scheme of the RASS-Hamburg optical identification program (Zickgraf et al. 2003), which uses similar X-ray (0.1–2.4 keV) and optical (Johnson B) bands. We converted the RASS-Hamburg limits (see Table 6 of Zickgraf et al. 2003) to BMW–HRI and GSC2 $f_X/f_{B1}$ (the deepest GSC2 band), keeping the different X-ray energy band and spectral shape used for flux evaluation in mind as well as the different optical band. The resulting classification scheme is shown in Table 4 (see Chieregato 2005 for details about the conversion). The $f_X/f_{B1}$ lower limits of BFS, $\simeq 37$ in the Northern sky and $\simeq 59$ and in the Southern sky (due to different GSC2 depth, see Table 1), are a factor of $\sim 2$ and $\sim 3$ beyond the limits of our reference values for different classes, respectively. This points to the possibility that BFS are unusual objects, left out of classifications, or at least very peculiar members of the more ordinary categories.

5.1. BFS as unusual sources

It is well known that some rare class of sources can reach extreme values of $f_X/f_{B1}$ (see Cagnoni et al. 2002). Here we will focus on our specific sample and discuss the possibility that our BFS belong to the following categories:

a. Isolated Neutron Stars (INS).
b. X-ray binaries.
c. High Redshift or Dark Clusters of Galaxies.
d. Type 2 Quasars.
e. Extreme BL Lac.

**Table 4.** $f_X/f_{B1}$ for various classes of sources, adapted to BMW–HRI fluxes and GSC2 BJ band.

<table>
<thead>
<tr>
<th>Class</th>
<th>Lower limit</th>
<th>Upper limit</th>
<th>$B-V$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Stars</td>
<td>$1.3 \times 10^{-5}$</td>
<td>0.6</td>
<td>$-0.5 \pm 2$</td>
</tr>
<tr>
<td>White Dwarfs</td>
<td>$1.2 \times 10^{-5}$</td>
<td>$9.3 \times 10^{-4}$</td>
<td>$-0.3 \pm 1.5$</td>
</tr>
<tr>
<td>Cat. var.</td>
<td>$1.8 \times 10^{-2}$</td>
<td>4.6</td>
<td>$-0.1 \pm 1$</td>
</tr>
<tr>
<td>Galaxies</td>
<td>$3.1 \times 10^{-2}$</td>
<td>11.1</td>
<td>$0.4 \pm 1.5$</td>
</tr>
<tr>
<td>Gal. clusters</td>
<td>0.2</td>
<td>16.9</td>
<td>$1.0 \pm 1.5$</td>
</tr>
<tr>
<td>AGN (w. BLLacs)</td>
<td>0.22</td>
<td>12.0</td>
<td>$-0.5 \pm 1$</td>
</tr>
</tbody>
</table>

Isolated Neutron Stars. Isolated Neutron Stars (INS; see Treves et al. 2000; Haberl 2003; 2004) are extreme BFS. The bona-fide $f_X/f_{B1}$ lower limit used to assess their nature is 1000, but for optically identified objects the real value can be as high as $10^4$. The X-ray spectrum is optimally fitted by a blackbody with $kT \sim 60$–100 eV. The optical emission lies a factor of a few over the Rayleigh-Jeans tail of the blackbody and no emission outside the optical/soft X-ray range has been found. The low column densities from the X-ray spectra and the parallax measured distance for RXJ 1856.5-3754 (Walter & Lattimer 2002) makes this source an intrinsically faint and near object, probably emitting from the neutron star surface.

Different kinds of isolated neutron stars can also reach high $f_X/f_{B1}$ values: Anomalous X-ray Pulsars (AXP; Israel et al. 2002), close-by radio pulsars with substantial cooling X-ray thermal emission (e.g. PSR 0656+14) and Geminga-like objects (Bignami & Caraveo 1996). AXPs are young, luminous, and distant objects that should be much closer to the Galactic plane than our BFS. The radio limit, though not strong, argue against close-by cooling radio pulsars. This does not apply to Geminga-like neutron stars, for which the radio emission is only marginal. Converting the X-ray flux of Geminga to the HRI pass-band and using $V = 25.6$ yields $f_X/f_V \sim 67$, making it consistent with our BFS candidates. In the cooling framework, the same physical mechanism is at the basis of INS emission and of the soft thermal component of young radio pulsars and Geminga-like objects. Based on this scenario, Popov et al. (2003, 2005) produced the log $N$–log $S$ of cooling neutron stars (all the three categories), allowing us to estimate that only $\sim 1$ unidentified INS is expected for our flux and sky coverage. If our steady BFS (1BMW042142.4-571541) is an INS, no optical counterpart could be revealed up to the $\sim 28$ mag; therefore, this possibility could be strengthened by the absence of optical sources in deeper observations, reaching the critical $f_X/f_{B1}$ lower limit of 1000, corresponding to magnitude $\sim 25$–26, together with a better position from Chandra. At the same time, the well-defined properties of the X-ray emission could be...
tested by observations with an instrument with sufficient soft response, like the XMM-Newton EPIC pn camera.

**X-ray binaries.** Some kind of X-ray binaries can reach high values of $f_X/f_{\text{opt}}$. In particular, BFS searches have already discovered Ultraluminous X-ray Sources (ULX; e.g. Cagnoni et al. 2003). However, the observed flux level of BFS is too low for a Galactic object, so none of our BFS can belong to nearby galaxies, excluding ULXs, supersoft sources, and bursting LMXBs. Therefore, we consider known classes of X-ray binaries unlikely (see, however, the discussion on possible transients).

**Clusters of Galaxies.** BFS research has already been effective for the discovery of clusters of galaxies since Cagnoni et al. (2001). Furthermore, several distant clusters have already been discovered in the BMW–HRI catalog as extended objects (Moretti et al. 2004). Distant clusters may escape this selection if only the X-ray peak is bright enough to be detected, resulting in a point source, while the rest is concealed in the background. This effect is strongly favoured at large off-axis angles (and consequent PSF degradation), like those at which BFS are detected. Therefore, if some BFS are galaxy clusters, we expect them to be at quite a high redshift. The $f_X/f_{\text{opt}}$ of clusters of galaxies is the highest among the X-ray emitting classes of Table 4. For normal galaxy clusters, the values reached are usually lower than those of BFS (and if only the peak of the X-ray emission is computed, the resulting $f_X/f_{\text{opt}}$ is further lowered). Despite this, there are two effects that could boost the $f_X/f_{\text{opt}}$. The first, and more effective, is the redshift of the optical emission. In fact, for a redshift of 0.75, the 4000 Å break of the cD galaxy (or of any early type galaxy of the cluster) would be redder than the GSC2 F band. Even a lower redshift could result in a sufficient enhancement of the $f_X/f_{\text{opt}}$. Second, the X-ray peak might not be coincident with the optical one, if the cluster is not virialized, or if there is no cD galaxy, like in Bautz-Morgan type III clusters, or if there is a central cooling flow. Optical-near IR deeper observations should reveal a substantial increase in galaxy counts in the proximity of BFS and, if there is any, the cD galaxy of the cluster. An even stronger test would be X-ray observations with Chandra or XMM-Newton, which should reveal an extended source.

**BL Lacs.** BL Lacs stand somehow apart from the bulk of AGNs in that they can reach high $f_X/f_{\text{opt}}$ values. The unified spectral energy distribution of BL Lacs (Fossati et al. 1998; Ghisellini et al. 1998) can be described by a two-peak emission, in which the lower energy peak (radio to X-ray) is commonly ascribed to synchrotron emission, while the high energy peak (X-ray to $\gamma$-ray) is probably due to inverse Compton emission. In this scenario, BL Lacs are distinguishable on the basis of the shifts of the peaks, creating the so-called blazar sequence. By tuning the position of the peaks, it is possible to produce the old categories of X-ray selected and radio selected BL Lacs, the former being those capable of high $f_X/f_{\text{opt}}$ values.

In order to check the viability of BL Lacs hypothesis for BFS, we consider the different $\alpha_{\text{opt}}$ (the power law index connecting the X-ray flux to the optical flux) and $\alpha_{\gamma}$ (the power law index connecting the X-ray flux to the radio flux) for Einstein Slew Survey BL Lacs (Perlman et al. 1996) and the upper limit for BFS. Despite the poor radio limits, the locus occupied by BFS is partially superimposed on the most extreme Slew BL Lac, with the brightest BFS, 1BM042142.4-571541, exhibiting an $\alpha_{\text{opt}}$ limit substantially lower. If these objects were indeed BL Lacs, they would probably stretch the sequence in the $\alpha_{\text{opt}}-\alpha_{\gamma}$ plane, representing a still undiscovered extreme population. The identification of the optical and possibly of the radio counterpart would be very important for testing the BL Lac hypothesis. In this case, further X-ray observations should show a power law spectrum with photon, without evidence of absorption excess.

**Type 2 Quasars.** Evidences for a substantial population of type 2 Quasars, as requested by unification models (Urry & Padovani 1995; Comastri et al. 1995), has been continuously increasing in the recent past. In particular, the resolution of the largest fraction of the cosmic X-ray background in discrete sources (Hasinger et al. 1998; Mushotzky et al. 2000; Campana et al. 2001; Giacconi et al. 2002; Moretti et al. 2003; Brandt & Hasinger 2005) has led to infer the existence of an adequate number of highly absorbed and luminous AGNs, responsible for the hard part of the background. Up to now, while only a few type 2 quasars have been firmly identified (e.g. Stern et al. 2002; Norman et al. 2003), many candidates are being produced either by deep fields, pencil beam searches (Chandra Deep Field North, Barger et al. 2003; Chandra Deep Field South, Szokoly et al. 2003; Lockman Hole, Mainieri et al. 2002), or by dedicated shallower surveys, with a substantial sky coverage (Hellas2XMM, Baldi et al. 2002; ChaMP, Green et al. 2003; SEXSI, Harrison et al. 2003; BMW-Chandra, Romano et al. 2005). The most striking characteristic of type 2 quasars is the severe absorption excess. In particular, this makes them Extremely Red Object, with $R-K > 5$. Therefore, if some BFS are type 2 quasars, we expect them to have a relatively bright near-IR counterpart. Unfortunately, the near-IR limit of the 2MASS (Table 1) is not sufficient to confirm or rule out this possibility. Deeper near-IR observations are therefore mandatory to test the viability of this scenario. Future X-ray spectral information could also be extremely useful, confirming or excluding the need for absorption excess to fit the spectra. A number of sources capable of very high $f_X/f_{\text{opt}}$ ratios were recently discovered in Chandra Deep Fields, the so-called Extreme X-ray Objects (EXO; Koekemoer et al. 2004). The nature of these sources is still not clear, but they share the extremely red colour with type 2 quasars. However, these sources have much fainter X-ray fluxes than BFS ($\leq 4 \times 10^{-15}$ erg cm$^{-2}$ s$^{-1}$), so they are likely to belong to different populations.

The nature of BFS reported here is probably not unique. In order to accommodate the possible transients in one of the above classes, a key factor is variability. The most variable candidates, BL Lacs, can vary up to a factor of $\sim 100$ (e.g. Mkn 501; Pian et al. 1998), enough to explain the non-detections. This makes BL Lacs appealing candidates for possible transients. The variability of type 2 quasars, instead, seems insufficient. Other known transient source classes usually can not produce such high $f_X/f_{\text{opt}}$ values. An alternative explanation could be that of a new kind of transient, capable of great flux variability but reaching much lower peak maximum luminosity. According to this idea, the $f_X/f_{\text{opt}}$ would be greatly boosted by the non-simultaneity of X-ray and optical observations.
6. Follow-up projects

We wish to stress that the interpretation of BFS requires future multi-wavelength observations. Deeper optical and near-IR observations should allow us to find the optical counterpart (except possibly for INS), and future X-ray observations should give spectral and temporal insights. We have already started a program of optical follow-up. In the future, more will be known about these mysterious and fascinating lacunae in the X-ray/optical sky.

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