

Diffuse galactic sub-MeV γ -ray excess from high-energy electrons

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Received 22 July 2005 / Accepted 11 September 2005

ABSTRACT

Recent INTEGRAL observations have confirmed the existence of diffuse sub-MeV radiation from the inner Galaxy. We propose a new explanation for the origin of this diffuse component in terms of electrostatic bremsstrahlung radiation from high-energy electrons. This electrostatic bremsstrahlung interpretation requires a very flat power law energy spectrum of high-energy electrons in the inner Galaxy which is reasonable if many point sources with significant spectral dispersion contribute to the injection of high-energy electrons in the Galaxy. The electrostatic bremsstrahlung interpretation avoids the enormous source power problem that previous explanations by radiation from low-energy electrons have.

Key words. cosmic ray electrons – diffuse galactic gamma rays – radiation processes

1. Introduction

The OSSE (Kinzer et al. 1999) and COMPTEL (Strong et al. 1994) instruments on board of the *Compton Gamma Ray Observatory* established that the diffuse galactic gamma ray continuum emission extends down to photon energies below 100 keV. The truly diffuse origin of this radiation has been proven by correlated SIGMA measurements to estimate the galactic point source contribution and the analysis of the GINGA measurements of the galactic ridge emission at much lower energies (Yamasaki et al. 1997). Recent INTEGRAL observations (Strong et al. 2003; Lebrun et al. 2004; Strong et al. 2004, 2005) indicated that besides the positronium continuum, dominating the emission between 200 and 400 keV, a second diffuse galactic continuum emission component is present with a flat power law photon number spectral index of -1.7 from 20 keV up to 1 MeV.

In the past this sub-MeV gamma-ray emission has been interpreted as nonthermal electron bremsstrahlung radiation from low-energy electrons (Skibo & Ramaty 1993; Skibo et al. 1996; Schlickeiser 1997; Dogiel et al. 2002) in the interstellar medium. As the main energy loss processes of such low-energy electrons are Coulomb and ionization interactions, the electrons nonthermal bremsstrahlung efficiency is small, so that this interpretation requires large amounts of low-energy interstellar electrons. Attributing these to cosmic ray electron sources, Skibo & Ramaty (1993) and Skibo et al. (1996) estimate that integrated over the whole Galaxy a source power of $\sim(1 - 10) \times 10^{42}$ erg s^{-1} in these low-energy (<10 MeV) electrons is required to maintain these electrons against the severe

Coulomb and ionization losses. This power exceeds the power supplied to the nuclear cosmic ray component by at least two orders an order of magnitude, and represents a serious problem for the general problem of the origin of cosmic rays. It is the purpose of this work to provide an alternative interpretation of the Compton observatory and INTEGRAL sub-MeV measurements. We show that the sub-MeV diffuse gamma-ray emission can be explained by the electrostatic bremsstrahlung radiation from high-energy electrons which is inverse Compton scattering of electrostatic plasma waves into transverse electromagnetic photons. This electrostatic bremsstrahlung interpretation avoids the enormous source power problem of the radiating electrons, and relies on the existence of interstellar electrostatic plasma waves generated by the interaction of supernova outflows with the ambient interstellar medium.

2. Synchrotron and electrostatic bremsstrahlung from high-energy cosmic-ray electrons

2.1. High energy cosmic ray electrons

According to current understanding high-energy ($\gamma_e > 10^4$) galactic cosmic ray electrons (negatrons and positrons) consist of two distinct populations: primary and secondary electrons. Primary cosmic ray electrons are accelerated in the same sources as high-energy cosmic ray hadrons whereas secondary electrons result from inelastic hadron-hadron collisions of cosmic ray hadrons in the interstellar medium. The equilibrium power law energy spectrum of secondary electrons then is very steep ($\propto \gamma_e^{-3.7}$) because the secondaries production

spectrum ($\propto \gamma^{-2.7}$) reflecting the cosmic ray hadron equilibrium spectrum is steepened by radiative synchrotron and inverse Compton losses. Because of this steep spectrum secondary electrons cannot produce the sub-MeV gamma-ray emission (Mörsberger & Schlickeiser 1997).

The interstellar cosmic ray primary electron spectrum results from the competition between injection from many discrete sources as shell-type supernova remnants, diffusive propagation in the interstellar medium and radiative losses. A number of authors (Brecher & Burbidge 1972; Pohl & Esposito 1998; Strong et al. 2000; Büsching et al. 2001; Büsching et al. 2005) have demonstrated that in this case the energy spectra of primary electrons and hadrons vary significantly over the galaxy and may be much harder than locally observed. The reasons are twofold:

- 1) the superposition of many power laws with dispersion in individual spectral index values does not produce a simple power law (Brecher & Burbidge 1972) at all energies. At high energies it flattens due to the dominance of those point sources with the flattest spectral index;
- 2) at high energies the radiative loss times of electrons are extremely short. As the bulk streaming speed of cosmic rays is limited to sub-Alfvénic values they penetrate not very far from their sources.

In the following we assume that the line-of-sight averaged cosmic-ray electron differential number density in the inner Galaxy is given by

$$N(\gamma) = N_0 \gamma^{-s} \exp(-\gamma/\gamma_0) \quad (1)$$

where the spectral index s can be as small as 1.5 which is compatible with the smallest observed synchrotron radiation spectral index of shell-type supernova remnants in the Galaxy (Clark & Caswell 1976; Göbel et al. 1981).

2.2. Electrostatic bremsstrahlung

Electrostatic bremsstrahlung, i.e. the inverse Compton scattering of electrostatic plasma waves into transverse electromagnetic waves by relativistic electrons, is very similar to synchrotron radiation which is the inverse Compton scattering of the virtual photons of the magnetic field into transverse electromagnetic waves by relativistic electrons in the Weizsäcker-Williams approach. Therefore its polarisation and emission properties are similar to that of synchrotron radiation (Chiuderi & Veltri 1974; Windsor & Kellogg 1974). Recently, one of us (Schlickeiser 2003) has proposed electrostatic bremsstrahlung as an alternative (to synchrotron radiation) nonthermal radiation process in jets of active galactic nuclei. Electrostatic bremsstrahlung becomes competitive to synchrotron radiation if the energy density in electrostatic waves U_L is of the same order as the magnetic field energy density U_B . Our interpretation here indicates that this is indeed the case in the interstellar medium, as plenty of electrostatic plasma waves are generated by the interaction of the supernova remnant shock waves with the ambient interstellar medium. Modeling the interaction of the supernova

shock waves with its ambient medium kinetically as a cloud of electron-proton plasma moving through the electron-proton interstellar medium (Pohl et al. 2002; Schlickeiser et al. 2002) indicates that the resulting energy density of electrostatic waves is

$$U_L = \frac{2}{3} n_e m_p \langle V_s \rangle^2 f \quad (2)$$

where n_e denotes the interstellar electron density, $\langle V_s \rangle$ the average supernova shock wave speed and f the volume fraction of interstellar medium swept by supernova explosions. For the hot coronal phase ($f = 0.8$, $n_e = 10^{-3} \text{ cm}^{-3}$) and $\langle V_s \rangle = 300 \text{ km s}^{-1}$ Eq. (2) yields $U_L = 8.5 \times 10^{-13} \text{ erg cm}^{-3}$, which is of the same order as the galactic magnetic field energy density $U_B = 6.4 \times 10^{-13} \text{ erg cm}^{-3}$ for a 4 μG magnetic field strength.

2.3. Intensities of sub-MeV radiation from electrostatic bremsstrahlung and synchrotron radiation

In order to calculate the synchrotron and electrostatic bremsstrahlung radiation intensities we use the monochromatic approximations of the respective differential spectral powers (Schlickeiser 2003)

$$p_{\text{synch}}(\nu, \gamma) = 2c\sigma_T U_B \gamma^2 \delta(\nu - \nu_s \gamma^2), \quad (3)$$

$$p_{\text{eb}}(\nu, \gamma) = \frac{128}{27} c\sigma_T U_L \gamma^2 \delta(\nu - \nu_b \gamma^2), \quad (4)$$

where c denotes the speed of light, σ_T the Thomson cross section, $U_B = B^2/8\pi$ the magnetic field energy density and U_L the energy density in electrostatic plasma waves. $\nu_s = 16.8(B/4\mu\text{G}) \text{ Hz}$ and $\nu_b = 400(n_e/10^{-3} \text{ cm}^{-3})^{1/2} \text{ Hz}$ are the characteristic synchrotron and electrostatic bremsstrahlung frequencies, respectively.

We then obtain for spontaneous synchrotron emission coefficient

$$\begin{aligned} j_{\text{synch}}(\nu) &= \frac{1}{4\pi} \int_1^\infty d\gamma p_{\text{synch}}(\nu, \gamma) N(\gamma) \\ &= \frac{c\sigma_T U_B N_0}{4\pi \nu_s} (\nu/\nu_s)^{(1-s)/2} \exp\left(-\left(\frac{\nu}{\nu_s \gamma_0^2}\right)^{1/2}\right) \end{aligned} \quad (5)$$

for frequencies $\nu > \nu_s$. Likewise, the spontaneous electrostatic bremsstrahlung emission coefficient is

$$\begin{aligned} j_{\text{eb}}(\nu) &= \frac{1}{4\pi} \int_1^\infty d\gamma p_{\text{eb}}(\nu, \gamma) N(\gamma) \\ &= \frac{16c\sigma_T U_L N_0}{27\pi \nu_b} (\nu/\nu_b)^{(1-s)/2} \exp\left(-\left(\frac{\nu}{\nu_b \gamma_0^2}\right)^{1/2}\right) \end{aligned} \quad (6)$$

for frequencies $\nu > \nu_b$. At high frequencies the emission is optically thin and the radiation intensities are given by the line-of-sight integrals $I_{\text{synch}}(\nu) = \int_0^\infty dl j_{\text{synch}}(\nu)$ and $I_{\text{eb}}(\nu) = \int_0^\infty dl j_{\text{eb}}(\nu)$. For ease of exposition we assume constant values of the galactic magnetic field strength, the electrostatic plasma wave energy density and the interstellar electron density so that the characteristic frequencies ν_s and ν_b will be constants.

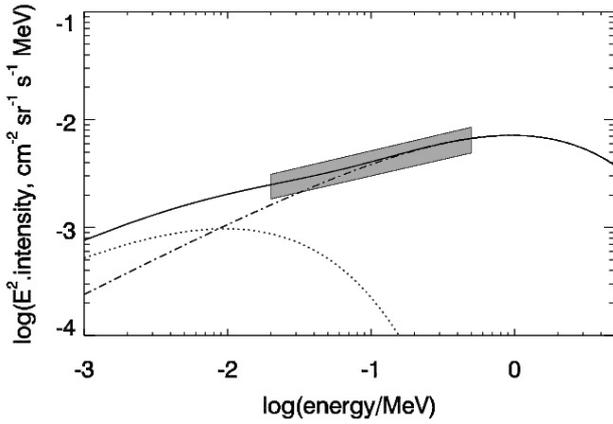


Fig. 1. Comparison of the best fit to the sum (solid line) of the synchrotron and electrostatic bremsstrahlung intensities with the observed diffuse emission (shaded area) from the Galactic disk corrected for the positronium contribution. The shaded area represents the power-law fit to the non-positronium component of the diffuse emission for the region $330^\circ < l < 30^\circ$, $-10^\circ < b < +10^\circ$ given in Strong et al. 2005. The best fit results for the parameter values $\nu_b/\nu_s \approx 100$, $s = 1.62$, $U_L/U_B = 3.0$, $\nu_b/\nu_s = \epsilon_b/\epsilon_s = 100$ and $\epsilon_s = 5$ keV. The dotted line shows the synchrotron flux, the dot-dashed line the electrostatic bremsstrahlung flux.

In Fig. 1 we compare the best fit of the normalised (at $h\nu = 500$ keV) sum of the synchrotron and electrostatic bremsstrahlung intensities with the observed diffuse emission from the inner Galaxy corrected for positronium contributions. The reduced $\chi^2_{\text{minp.d.f}} = 0.0015$ is obtained for the values $s = 1.62$, $U_L/U_B = 3.0$, $\nu_b/\nu_s = \epsilon_b/\epsilon_s = 100$ and $\epsilon_s = 5$ keV. During the χ^2 -test we keep the values of the ratio $\nu_b/\nu_s = 100$ and the high-energy synchrotron radiation cutoff $\epsilon_s = 5$ keV fixed, and only vary the electron spectral index s and the ratio of energy densities U_L/U_B . The contour intervals of the latter two parameters with different confidence level are shown in Fig. 2.

The value of the high-energy exponential cut-off of the radiating electrons γ_0 is determined by requiring that the continuum hard X-ray emission up to 10 keV is predominantly due to synchrotron radiation requiring the value $\epsilon_s = h\nu_s\gamma_0^2 = 5$ keV. Together with the frequency ν_s this implies the value $\gamma_0 = 2.7 \times 10^8 (B/4\mu\text{G})^{-1/2}$, i.e. maximum electron energies of $10^{14} (B/4\mu\text{G})^{-1/2}$ eV which compare favorably well with the maximum energy of shock accelerated electrons found in 14 individual young shell supernova remnants (Reynolds & Keohane 1999). This choice then also defines the high-energy cut-off of the electrostatic bremsstrahlung emission $\epsilon_b = h\nu_b\gamma_0^2$ at about 500 keV.

The relative X-ray (<10 keV) synchrotron to sub-MeV (<1 MeV) electrostatic bremsstrahlung intensities require a best fit electrostatic plasma wave energy density of $U_L = 3 U_B$ of about three times the magnetic field energy density. Figure 2 indicates that with 99 percent confidence the average galactic energy density in electrostatic waves is greater than 80 percent of the average galactic magnetic field energy density.

These high-energy primary electrons will also inverse Compton scatter any additional target photons in the inner

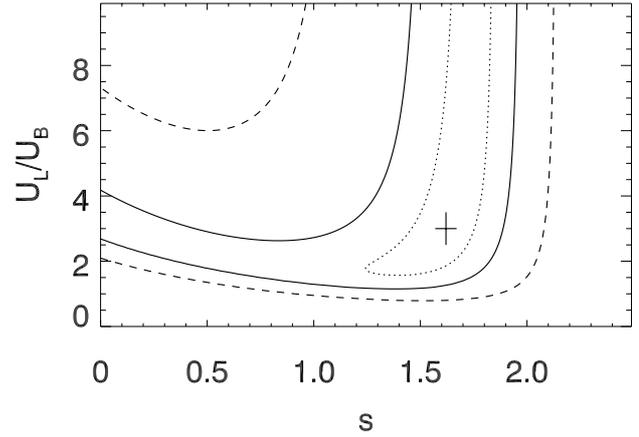


Fig. 2. 20 percent (0.25σ), 68 percent (1σ) and 99 percent (2.6σ) confidence intervals of the free parameters electron spectral index s and energy density ratio U_L/U_B .

Galaxy as UV, optical, infrared or microwave background photons. Denoting with U_a the energy density of any additional photon component, we find in the Thomson limit that the ratio of additional inverse Compton radiation to electrostatic bremsstrahlung at the high-energy cutoff

$$\frac{\nu I_{\text{ICa}}}{\nu I_{\text{eb}}} = \frac{U_a}{U_L} \quad (7)$$

is determined by the respective energy density ratio. For the 2.7 K microwave background photons we infer

$$\frac{\nu I_{2.7}(30 \text{ TeV})}{\nu I_{\text{eb}}(0.4 \text{ MeV})} = 3.3 \quad (8)$$

yielding at 30 TeV $\nu I_{2.7}(30 \text{ TeV}) \approx 0.003 \text{ MeV s}^{-1} \text{ cm}^{-3} \text{ rad}^{-1}$ from the inner Galaxy. At EGRET energies (10 GeV) this contribution amounts about

$$\begin{aligned} \nu I_{2.7}(10 \text{ GeV}) &= 0.003 \left(\frac{10 \text{ GeV}}{30 \text{ TeV}} \right)^{0.75} \\ &= 7.4 \times 10^{-6} \text{ MeV s}^{-1} \text{ cm}^{-3} \text{ rad}^{-1} \end{aligned} \quad (9)$$

which is three orders of magnitude smaller than the observed flux from the inner Galaxy (Schlickeiser et al. 1997). Even higher energy densities of optical and infrared photons in the inner part of the Galaxy would not contradict the flux observed by EGRET.

3. Discussion and summary

Recent INTEGRAL observations have confirmed the existence of diffuse sub-MeV radiation from the inner Galaxy. Here we have presented a new explanation for the origin of this diffuse component in terms of electrostatic bremsstrahlung radiation from high-energy electrons. This electrostatic bremsstrahlung interpretation requires a very flat power law energy spectrum of high-energy electrons in the inner Galaxy which is reasonable if many point sources with significant spectral dispersion contribute to the injection of high-energy electrons in the Galaxy. The electrostatic bremsstrahlung interpretation avoids the enormous source power problem that previous explanations

by radiation from low-energy electrons have, and relies on the existence of interstellar electrostatic plasma waves probably generated by the interaction of supernova outflows with the ambient interstellar medium. Our excellent best fit to the observed diffuse emission from the inner Galaxy indicates that the averaged galactic energy density of interstellar electrostatic plasma waves is greater than 80 percent of the averaged galactic magnetic field energy density. The maximum energy of the radiating electrons is about 10^{14} eV. Our proposed interpretation as radiation from high-energy electrons then implies a rather structured angular distribution of this radiation component, as the radiating high-energy electrons have small radiative lifetimes and therefore do not propagate very far away from their sources. For an average spatial diffusion coefficient of $D(E) \approx 10^{27}(E/\text{GeV})^{0.3} \text{ cm}^2 \text{ s}^{-1}$ and the radiative loss time t_s the mean distance travelled by an 10^{14} eV electron is about $\langle x \rangle = \sqrt{2Dt_s} \approx 30$ pc corresponding to 0.2 degrees if located near the Galactic center. Alternatively, by measuring $\langle x \rangle$ at different photon energies one may determine the value of the average spatial diffusion coefficient $D(E)$.

Acknowledgements. R. Schröder and R. Schlickeiser acknowledge support by the Deutsche Forschungsgemeinschaft through grant Schl 201/16-1.

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