A radio monitoring survey of ultra-luminous X-ray sources

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Abstract. We present the results of a radio monitoring campaign to search for radio emission from nearby ultra-luminous X-ray sources (ULXs). These sources are bright off-nuclear X-ray point sources with luminosities exceeding $L_X > 10^{39}$ erg s$^{-1}$. A well-defined sample of the 9 nearest ULXs has been monitored eight times over 5 months with the Very Large Array in A and B configuration. Our limiting sensitivity is $\mu$Jy (4$\sigma$) for radio flares and $\approx 60$ $\mu$Jy for continuous emission. In M82 two ULXs seem to have coincident compact radio sources, which are probably supernova remnants. No continuous or flaring radio emission has been detected from any other ULX. Thus, ULXs do not generally emit steady-state radio emission above radio powers of $1.5 \times 10^{27}$ W/Hz. The non-detections of the continuous emission are consistent with beamed or unbeamed radio emission from accreting black holes of $\lesssim 10^7 M_\odot$ based on the radio/X-ray correlation. Other published radio detections (M82, NGC 5408) are also discussed in this context. Both detections are significantly above our detection limit. If ULXs have flaring radio emission above $4 \times 10^{27}$ W/Hz we can give an upper limit on the duty cycle of the flares of 6%. This upper limit is in agreement with the observed number of flares in Galactic radio transients. Additionally we present a yet unreported radio double structure in the nearby low-luminosity AGN NGC 4736.

Key words. black hole physics – X-rays: binaries – galaxies: active

1. Introduction

Ultra-luminous X-ray sources are among the most widely discussed objects, but their true nature remains unknown. These objects have X-ray luminosities $L_X \sim 10^{39–42}$ erg s$^{-1}$ which seems too bright for normal (stellar-mass) black hole X-ray binaries (BHXBs), but far dimmer than normal active galactic nuclei (AGN). The first hints of these intermediate-luminosity X-ray point sources were found in the 1980s (Fabbiano 1989; Colbert et al. 1995). Observations with ROSAT, and subsequent high resolution X-ray satellites, confirmed these findings and showed that these intriguing sources are often not in the centers of the galaxies (Colbert & Mushotzky 1999; Roberts & Warwick 2000). More recent surveys (see e.g., Colbert & Ptak 2002; Ptak & Colbert 2004) find approximately one ULX in every five galaxies, confirming that ULXs are a common phenomenon.

One reason why ULXs are so intensely studied is their potential connection with intermediate mass black holes (IMBHs). Under the assumption that they emit isotropically, the Eddington limit gives a lower bound for the mass of the central black hole as large as a few hundred solar masses for nominal Eddington ratios of $\sim 0.1–1$. The existence of IMBHs would be interesting, as they could be the “missing link” between stellar-mass black holes and the supermassive black holes in the center of galaxies (see e.g., Ebisuzaki et al. 2001). Such objects would have important implications for cosmology and arise in many theories on the collapse of primordial stars.

While there is evidence that a few of these objects are indeed IMBHs (e.g., Strohmayer & Mushotzky 2003; Cropper et al. 2004), the creation and feeding mechanisms for black holes of this mass are still not well understood. The ULX phenomenon seems to be connected to star formation, see e.g., the Cartwheel Galaxy (Gao et al. 2003). This galaxy would need hundreds to thousands of IMBHs feeding from a yet unknown non-stellar mass reservoir. Alternatively, the ULX phenomenon may be connected to high mass XRBs (King 2004). Grimm et al. (2002) and Gilfanov (2004) show the existence of a universal X-ray luminosity function that may extend up to ULX luminosities by taking the star formation rate into account. There seems to be a dependence of the ULX abundance on the galaxy type, since dwarf galaxies host many of the nearby ULXs.

There is also an ongoing discussion whether the inner accretion disk temperatures of ULXs are too hot for IMBHs, since many have $kT_{\text{in}} \sim 1$ keV, which is more similar to those found in normal XRBs, see e.g., Colbert & Mushotzky (1999); Mizuno et al. (1999); Makishima et al. (2000). However,
recent XMM observations find lower inner disk temperatures for some ULXs, which is more compatible with IMBHs than earlier results from ASCA observations (see e.g. Miller et al. 2003, 2004).

Besides the explanation that ULXs are accreting IMBHs, other stellar-mass ULX models have been proposed: anisotropic X-ray emission (“mild beaming”) (King et al. 2001), super-Eddington accretion flows (Abramowicz et al. 1988; King 2004) and microblazars (Mirabel & Rodríguez 1999; Körding et al. 2002). In the microblazar model, the “ultra-luminous” X-ray emission is created by a relativistically boosted X-ray jet. While the jet is intrinsically weaker than the accretion disk, Doppler boosting can lead to the observed luminosities. If the jets also emit in the radio regime like XRBs in the hard state (cf., Fender 2001) or the very high state (see e.g., GRS 1915, Rodriguez et al. 1995), the radio emission is also boosted and may be observable. The jet/disk model (Falcke & Biermann 1995) assumes that the accretion flow and the jet form a symbiotic system; if there is a disk, there is always a jet in some form. Flat spectrum radio emission originates from the self-absorbed part of the jet, while one observes a steep synchrotron power law in the X-rays (Markoff et al. 2001). Thus, one possibility to distinguish between the microblazar model and the other possibilities is the presence of compact, luminous radio cores at the positions of ULXs.

Galactic XRBs in the very high state are probably the nearest cousins to ULXs. Of these, GRS 1915+105 can be seen as the prototype. GRS 1915+105 shows bright radio flares of 1.5 Jy while the quiescent value is approximately 130 mJy (Rodriguez et al. 1995). An other highly variable XRB is Cyg X-3. While this source is also found in the low and high state, it shows major radio flares with a flux increase of a factor 10–100 on a time scale of a day which lasts for days or weeks (Ogley et al. 2001). Therefore, it has to be expected that ULXs also could be radio transients. These relativistically boosted radio flares could be detectable with the VLA. In addition, these observations can be used to estimate the number of radio flares occurring in a galaxy. Such estimates are important for the design of new digital radio telescopes such as LOFAR or the SKA, as they could easily search for transients given their multi-beam capabilities.

In this paper we describe a systematic search for steady-state or flaring radio emission from a well-defined sample of ULXs. In Sect. 2 we present our observation scheme and in Sect. 3 we show our results. Theoretical implications are discussed in Sect. 4. Our conclusions are given in Sect. 5. We discuss the low luminosity AGN NGC 4736 in the Appendix.

### 2. Observations

#### 2.1. Sample selection

ULXs are defined as having observed 0.3–8.0 keV luminosities $L_X > 10^{39}$ erg s$^{-1}$. We selected nine ULXs from a larger list of ~200 ULXs, which were found in the Chandra ACIS data publicly available on 4 Feb. 2003. X-ray luminosities were estimated from ACIS count rates using a power-law model with photon index $\Gamma = 1.7$ and the Galactic absorption column of the corresponding galaxy (calculated using the FTOOLS NH program). See Table 1 for $L_X$ values for the nine ULXs in our sample.

X-ray point source lists, used to create the larger ULX catalog, were created using the XASSIST automatic X-ray analysis program (Pak & Griffiths 2003, URL: www.xassist.org). Our nine sample objects are the nearest nine ULXs $(D < 5.5$ Mpc) in this list, with $\text{Dec} > -20^\circ$, excluding M 81 X-1. The latter source was excluded because it has already been studied with the VLA (more than 4 h of integration time, no published detections). Nevertheless, they represent a distance-limited sample that is unbiased in X-ray luminosity.

The distance limit was mainly used to reduce the sample size. A low distance also helps to give low limits on the total emitted radio power of the sources. This selection results in nine ULXs in seven fields of views (FOVs): M 33, NGC 2403, M 82 (2 ULXs), NGC 4736 (2 bright point sources, one may be the nucleus), NGC 520 and two fields in NGC 5457 which is also known as M 101. The positions of the ULXs are shown in Table 1. All host galaxies are spirals except M 82, which is classified as dwarf/irregular (see e.g., de Vaucouleurs et al. 1991).

#### 2.2. Observing scheme

The goal of these observations is to search for radio flares and continuous emission from ULXs using the VLA. The seven selected fields were observed eight times each. The observing runs were approximately four hours long and occurred between June and October 2003. This results in 34 min per source and epoch including the calibration scans. After the last epoch the data of all epochs were combined to give deep observations to search for continuous emission. To avoid confusion the observations were obtained in the VLA’s highest resolution configurations (A or B configuration).

The optimal receivers at the VLA to search for weak, flat spectrum point sources are the X-band receivers (8.49 GHz). They are more sensitive than the C and U band receivers, as long as the spectrum is flatter than a spectral index $\alpha$ ($S_{\nu} \propto \nu^{-\alpha}$) of 0.4. We expect that the ULXs have a flat or even inverted spectrum which could be detectable if they are beamed. Thus,
we observed at 8.49 GHz using the maximal bandwidth of 50 MHz. Due to the distribution of the right ascension of our sources we had to observe them in two groups (for the groups and the observing dates see Table 2).

As we are searching for very faint radio emission, we obtained phase-referenced observations. For all sources VLA phase-calibrators could be found within 10°, typically within 5° of the target source. The cycle time between calibrator and source was approximately 7 min. The phase correction between subsequent calibrator scans was typically ≤40′′ on the longest baselines, while the phase corrections on the shorter baselines are typically ≤10′′. Therefore, the phase calibrated images should only be degraded by a few percent due to phase errors. Amplitude calibration has been done using either 3C 286 or 3C 48. The phase-calibrated data has then been imaged with natural weighting to achieve the maximal sensitivity. In each epoch a source was observed for approximately 21 min, excluding phase and amplitude calibrator observations.

For M 82, we have also reanalyzed archival A-configuration data presented in Kronberg & Sramek (1985), as they have reported a radio flare of 7.07 mJy at C-band (6 cm). Intriguingly, the position of the flare is near the brightest ULX in this galaxy.

### 3. Observational results

#### 3.1. Error limits

The root mean square (rms) flux values of the observed sources have been calculated from natural weighted maps. For fields with detected sources the largest possible rectangular field excluding the source has been used. The theoretical rms value for a naturally weighted 21 min observation with 27 antennas is, according to the VLA Observational Status Summary (Taylor et al. 2004), 32 µJy. The rms flux values of the individual epochs and fields are shown in Table 3. In a single epoch the typical observed rms is about 35 µJy (except for M 82, see Sect. 3.2.1). If one combines all epochs, the rms goes down to ≤15 µJy, while the theoretical limit is 11 µJy. To be able to interpret the radio images we have to assess when a peak in the map is significant. Chandra has a point spread function (PSF) of approximately one arc-second. Wrobel et al. (2001) report that the astrometric error for phase referenced VLA observations (A-configuration) compared to the International Celestial Reference Frame as established by VLBI (Ma et al. 1998) is less than 10 mas for the X-band. The astrometric uncertainty of the VLA is therefore negligible compared to the Chandra position errors.

We know the positions of the ULXs with the accuracy of Chandra. During the first five epochs the VLA was in the A configuration and our beam width was ≈240 mas. The PSF of Chandra, which has a full width half maximum of around one arcsecond, is therefore covered by 17 beams. The probability that we detect a random peak in 17 independent beams is given by the error function. Under the assumption that these beams are independent, the chance that there is a random 3σ peak at the position of a given ULX is 4%, the chance of a 4σ peak 0.1%. We can therefore only accept peaks at the positions of the ULXs if they exceed 4σ. To search for other radio flares in the whole map we have to increase the 4σ limit. We typically map 50′′ by 50′′, which corresponds to roughly 14 000 “independent” beams, so there will usually be one 4σ peak in the map. Therefore point sources without known positions are not likely to be real unless they have a signal-to-noise higher than ≈5σ.

#### 3.2. Detections

##### 3.2.1. M 82

In the starbursting galaxy M 82 many compact radio sources (CRSs) have been found (see e.g., Kronberg & Wilkinson 1975). Many of them have been identified as supernova remnants (SNRs) by direct imaging (e.g., Pedlar et al. 1999) or by spectral observations (see Allen & Kronberg 1998). M 82 also shows strong diffuse radio emission, making it more difficult to achieve the optimal noise limit for the short A-configuration snapshot observations (zero-spacing, dynamic range limitations). The longer baselines are less contaminated by this diffuse emission. Thus, the best rms values have been obtained by only using baselines exceeding 300 kλ for the A and early AnB config observations. This omits approximately half of the baselines, increasing the theoretical rms by a factor of ≈√2 to 44 µJy. The rms values for a single epoch still do not reach this value but are slightly higher: ≈70 µJy. For the 6th to the 8th epoch smaller cut-offs had to be utilized while the array configuration approached B-configuration. In Table 4 we give the X-ray and radio positions. While the X-ray positions are named with capital letters, the corresponding radio detections

### Table 2. Observing dates of the different epochs. The sources are observed in two groups. The first group consists of M 33, NGC 2403 and M 82, while NGC 4736, NGC 5204 and the two FOVs in NGC 5457 belonging to the second group.

<table>
<thead>
<tr>
<th>Epoch</th>
<th>Group 1</th>
<th>Group 2</th>
<th>Configuration</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>2 June</td>
<td>4 June</td>
<td>A</td>
</tr>
<tr>
<td>2</td>
<td>3 July</td>
<td>3 July</td>
<td>A</td>
</tr>
<tr>
<td>3</td>
<td>28 July</td>
<td>5 July</td>
<td>A</td>
</tr>
<tr>
<td>4</td>
<td>18 Aug.</td>
<td>17 Aug.</td>
<td>A</td>
</tr>
<tr>
<td>5</td>
<td>7 Sep.</td>
<td>30 Aug.</td>
<td>A or AnB</td>
</tr>
<tr>
<td>6</td>
<td>15 Sep.</td>
<td>15 Sep.</td>
<td>AnB</td>
</tr>
<tr>
<td>7</td>
<td>03 Oct.</td>
<td>07 Oct.</td>
<td>AnB</td>
</tr>
<tr>
<td>8</td>
<td>16 Oct.</td>
<td>18 Oct.</td>
<td>B</td>
</tr>
</tbody>
</table>

### Table 3. rms flux values for the different FOVs. Besides the rms of each epoch we give the rms value for map of the combined dataset. All values are given in µJy. For a discussion on M 82 see Sect. 3.2.1.

<table>
<thead>
<tr>
<th>Epoch</th>
<th>1</th>
<th>2</th>
<th>3</th>
<th>4</th>
<th>5</th>
<th>6</th>
<th>7</th>
<th>8</th>
<th>Comb.</th>
</tr>
</thead>
<tbody>
<tr>
<td>M 33</td>
<td>37</td>
<td>36</td>
<td>37</td>
<td>60</td>
<td>41</td>
<td>71</td>
<td>54</td>
<td>17</td>
<td></td>
</tr>
<tr>
<td>2403</td>
<td>31</td>
<td>36</td>
<td>32</td>
<td>31</td>
<td>59</td>
<td>38</td>
<td>45</td>
<td>37</td>
<td>13</td>
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<td>32</td>
<td>31</td>
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<td>35</td>
<td>39</td>
<td>77</td>
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<td>14</td>
</tr>
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<td>33</td>
<td>35</td>
<td>75</td>
<td>34</td>
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<td>37</td>
<td>51</td>
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<td>13</td>
</tr>
<tr>
<td>N4736</td>
<td>45</td>
<td>47</td>
<td>49</td>
<td>55</td>
<td>45</td>
<td>61</td>
<td>63</td>
<td>62</td>
<td>21</td>
</tr>
</tbody>
</table>

Reference Frame as established by VLBI (Ma et al. 1998) is less than 10 mas for the X-band. The astrometric uncertainty of the VLA is therefore negligible compared to the Chandra position errors.

We know the positions of the ULXs with the accuracy of Chandra. During the first five epochs the VLA was in the A configuration and our beam width was ≈240 mas. The PSF of Chandra, which has a full width half maximum of around one arcsecond, is therefore covered by 17 beams. The probability that we detect a random peak in 17 independent beams is given by the error function. Under the assumption that these beams are independent, the chance that there is a random 3σ peak at the position of a given ULX is 4%, the chance of a 4σ peak 0.1%. We can therefore only accept peaks at the positions of the ULXs if they exceed 4σ. To search for other radio flares in the whole map we have to increase the 4σ limit. We typically map 50′′ by 50′′, which corresponds to roughly 14 000 “independent” beams, so there will usually be one 4σ peak in the map. Therefore point sources without known positions are not likely to be real unless they have a signal-to-noise higher than ≈5σ.
are named as indexed small letters. The X-ray sources D and F are the well-known ULXs (Matsumoto et al. 2001).

Our VLA observation was centered on ULX D of Table 4. Thus, the bright X-ray point source A is 18.8′′ away from the phase tracking center. As we are observing in continuum mode, our sensitivity is already reduced for this ULX due to bandwidth smearing. The distance corresponds to 72 beam widths, thus, the peak flux will be reduced by 5% (Taylor et al. 2004). We have not detected any radio emission from source A.

The brightest CRS in the FOV ($e_2, 41.95 +57.5$ in the notation of Muxlow et al. 1994, abbreviated to MPW) is in good agreement with the position of a bright X-ray point source (E). This source is probably a SNR (e.g., Allen & Kronberg 1998) that may be immersed in a molecular cloud (Pedlar et al. 1999). As the radio and X-ray positions coincide, the radio and X-ray maps are correctly aligned.

Near the ULX D we have detected a CRS ($d_1$) with a continuum flux of 1.8 mJy in all eight epochs observed in this campaign. This source is also detected by MPW with 1.8 mJy as $41.3+59.6$ and is identified as a SNR (Allen & Kronberg 1998). As expected, the flux of the source is stable during the 8 epochs within 1σ of the rms. It has a distance of 0.77′′ to the ULX, in agreement with the Chandra position. We did not detect the flaring source of Kronberg & Sramek (1985) in our VLA observations.

For the ULX F we find two possible radio counterparts as shown in Table 4. The brighter one is also given in earlier papers as $42.21 +59.0$ while the weaker one is barely visible in the maps of MPW but is not given in their table. Both sources are not variable over our observing period.

Besides the CRS near the reported positions of the ULXs we have not found any radio flare in M 82. However, Kronberg & Sramek (1985) found a 7.07 mJy flare near the brightest ULX (D) in M 82 ($2.3 \times 10^{19}$ W/Hz). We have reanalyzed their data for direct comparison with our results. The position of the M 82-flare is on the 0.5′′ circle for ULX D, its offset (0.77′′) is consistent with the ULX X-ray position.
The radio source is 1.4′′ away from the Chandra position of the ULX, further away than the typical position accuracy of Chandra, but it is still possible that it is connected to the ULX. However, besides this peak there is another similar maximum another arcsecond away, and both maxima are found on a side-lobe. Thus, it is possible that the source could be an instrumental artifact. Besides this, we do not detect significant (>4σ) radio emission from the ULXs in the individual epochs.

### 4. Theoretical interpretations

#### 4.1. Flares

The nearest known Galactic cousins of ULXs are the highly accreting black hole XRBs. XRBs in the low hard or the high state are highly variable in the radio and X-ray regime (cf., van der Klis 1989). The radio flux can vary by a few hundred percent, for example Cyg X-1 has a typical 15 GHz radio flux of approximately 10 mJy and reaches up to 30 mJy (Pottschmidt et al. 2003). Stronger accreting objects like GRS 1915+105 or Cyg X-3 show violent radio flares of a factor 10 to 100 greater than the normal fluxes. During these flares GRS 1915+105 reaches 1.5 Jy (Rodriguez et al. 1995) and Cyg X-3 goes up to more than 15 Jy (Watanabe et al. 1994), while the “steady” state emission is around 100 mJy for both sources. The brighter radio bursts of Cyg X-3 may be due to a higher Doppler factor, as the inclination of this source is only 14° (Mioduszewski et al. 2001). GRS 1915+105 on the other hand is seen nearly edge on (angle to the line of sight: 66°, Rodriguez et al. 1995).

One possibility to estimate the luminosities of radio flares that one could expect from beamed emission of ULXs is by taking the transients GRS 1915+105 and Cyg X-3 as an example. The average distance of our ULXs is approximately 4.6 Mpc. GRS 1915+105 has a distance of 11 kpc (Fender et al. 1999b) while Cyg X-3 is 9 kpc away (Predehl et al. 2000). Without any relativistic beaming these sources would have a flux of 8.5 μJy and 57 μJy, which is below our detection limit. However, already a mild beaming factor of 20 for GRS 1915 or an additional factor of 3 for Cyg X-3 would bring the flares into our detection limit (0.15 mJy). As GRS 1915 is seen with an inclination angle of 66°, its Doppler factor will be around one. Thus, a beaming factor of 20 can be reached with a moderate Lorentz factor of Γ ≈ 3 (Lind & Blandford 1985), if the jet
points roughly at the observer (inclination angle < 15°). The Doppler factor of Cyg X-3 is uncertain, as the observed luminosity is probably already beamed. But an additional factor of 3 is easy to obtain if the jet is pointing directly towards the observer. If the ULXs are flaring in radio we should be able to detect the beamed radio flares.

We have not found a single significant flare for all our sources in the monitoring campaign. In the individual epochs our 4σ sensitivity is on average 0.15 mJy.

The distance of M 33 is around 0.84 Mpc (Freedman et al. 1991). Besides this nearby galaxy, all other observed galaxies have distances in the range from 3.6 Mpc to 5.4 Mpc. Thus, the detection limits of flares will only vary by a factor of two for those galaxies. For the average distance of 4.6 Mpc the upper limit on the radio power \( S_\nu \) of flares is \( 3.8 \times 10^{17} \) W/Hz. This corresponds to a 5 GHz radio luminosity (\( \nu S_\nu \)) of \( 1.9 \times 10^{34} \) erg s\(^{-1}\) assuming a flat spectrum. For M 33, however, the limit is reduced to a radio power of \( 1.1 \times 10^{16} \) W/Hz and a radio luminosity of \( 5.7 \times 10^{32} \) erg s\(^{-1}\).

The non-detection may be due to the unknown time scales of the radio flares in ULXs. We observed our sources once or twice a month. If the time scale of a radio flare is only a day, we are strongly under-sampling the radio light-curve. The time scale of the boosted flares is unclear and depends on the physical process creating the flare. If the flare is created inside the jet, e.g., similar to the shock in jet models used for blazars (see e.g., Marscher & Gear 1985), the observed time scale of the flare will be reduced by the Lorentz factor. For \( \Gamma \approx 5 \) the time scale could be as short as a few hours. However, it will be extremely bright, even a few mJys are possible. On the other hand, if the flare is created by enhanced injection of material into the jet by the disk for an extended time, the observed time scale will be the same as the intrinsic one.

The fact that we have not found a single flare in all our epochs can be translated to an upper limit for the probability that an average ULXs of our sample has a radio flare brighter than \( 3.8 \times 10^{17} \) W/Hz, which should be detectable. If we assume that the probability that an ULXs is flaring at a given time is similar for all observed ULXs, and one flare is uncorrelated with earlier or later flares, the flares should be Poisson distributed. The probability distribution describing how many events we detect given the probability of a detection is

\[
P(\lambda, n) = \frac{\exp(-\lambda) \lambda^n}{n!},
\]

where \( \lambda \) denotes the expectation value of the distribution and \( n \) is the number of events.

For this study we have to exclude M 82 as the rms in the M 82 maps is much higher than in the other fields. The six fields have been observed eight times which yields 48 samples. Let \( \rho \) denote the upper bound of the duty cycle of the ULX. Thus, it is an upper limit for the probability that one ULX flares at a given moment. \( \rho \) has to be chosen such that we should have almost certainly detected at least one event if the probability that a given ULX flares is \( \rho \). Here \( \rho \) will be chosen such that we should detect one or more flares with a probability of 95%. This leads to

\[
0.95 = \sum_{n=1}^{\infty} P(48\rho, n) = 1 - \exp^{-48\rho}.
\]

This results in \( \rho = 0.06 \), i.e. the duty cycle of radio flares exceeding our detection threshold in ULXs is < 6%. The expectation value for the number of flares is 2.9, our detection of no flares is in agreement with this value at a 5% level.

If we know the time scale of a typical radio flare in an ULX we can convert the upper limit of the flaring probability to an upper limit of the number of flares a ULX can have in a year. Let \( \Delta t_{\text{flare}} \) denote the average length of a flare. As the radio flares of GRS 1915 have a time scale of days, we will use \( \Delta t_{\text{flare}} \approx 2 \) days as a reference. Thus, the upper limit of the probability to detect a flare in a single observation corresponds to an upper limit of \( 10^{16} \) radio flares per year. This upper limit is still higher than the number of bright flares in GRS 1915+105.

Besides looking for radio emission from ULXs we can use these observations to derive upper limits on the amount of radio transients occurring in a galaxy. Our galaxies have an average distance of around 4.6 Mpc and we mapped an area of 51′′, giving an observed area of 1.3 kpc\(^2\). We are searching for radio flares from unknown sources, therefore, it is also not known what quantity correlates with the number of flares. While for ULXs the star formation rate might be appropriate, be the number of stars or black holes could be more appropriate for another class of flaring objects. As we are observing a few small fields in different galaxies, it is very hard to derive the exact number of observed solar masses or even the observed star formation rate. As most of our observed galaxies (besides M 82, which we exclude) are non-starbursting galaxies, the star formation rate per solar mass will be of the same order of magnitude. In order to get a rough estimate of how much mass we have actually observed, we assume a similar mass density as in our Galactic neighborhood. In our Galactic neighborhood the mass density of the disk is of the order of 200 \( M_\odot \) per pc\(^2\). Therefore one field of view contains a mass of approximately \( 3 \times 10^8 M_\odot \). Thus, we expect fewer than 0.2 flares in a single observation of \( 10^6 M_\odot \).

4.2. Background sources

We have detected continuous radio sources only inside NGC 5457 (excluding M 82). In all other fields there were no background sources visible. However, we only map an area of 51′′× 51′′. Fomalont et al. (2002) find that the number density of sources above a flux density \( S \) is approximately

\[
N = 0.099 \pm 0.01 \left( \frac{S}{40 \mu Jy} \right)^{-1.11 \pm 0.13} \text{arcmin}^{-2}.
\]  

Our limiting sensitivity (4σ) is approximately 56 \( \mu Jy \). Thus, this predicts that on average there should be 0.049 ± 0.007 sources in a given observed 51′′× 51′′ map. Our findings are in agreement with this result. The probability that a background source is found coincident with a ULX, within the position accuracy of Chandra (1 arcsec), is below 2 \times 10^{-5}. 

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**References:**

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N = 0.099 \pm 0.01 \left( \frac{S}{40 \mu Jy} \right)^{-1.11 \pm 0.13} \text{arcmin}^{-2}.
\]  

Our limiting sensitivity (4σ) is approximately 56 \( \mu Jy \). Thus, this predicts that on average there should be 0.049 ± 0.007 sources in a given observed 51′′× 51′′ map. Our findings are in agreement with this result. The probability that a background source is found coincident with a ULX, within the position accuracy of Chandra (1 arcsec), is below 2 \times 10^{-5}. 

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4.3. Steady state emission

In M 82 two reported Chandra ULXs positions are in agreement with the positions of CRSs that are probably radio SNR. For all other galaxies we have not detected any continuous emission near the ULXs. The upper flux limit for the observed sources is ≈60 μJy (4σ). For the steady-state emission ULXs should lie on the Radio/X-ray correlation, which was first established for XRBs (Corbel et al. 2000, 2003; Gallo et al. 2003) and has been extended to AGN (Merloni et al. 2003; Falcke et al. 2004). Simple jet scaling (see e.g., Falcke et al. 2004) predicts that the X-ray and radio emissions are correlated and scale according to:

\[ L_X \propto L_R^{1.38} M_e^{0.81} \]  

(2)

where \( L_X \) and \( L_R \) denote the X-ray and radio luminosities. The exact value of the exponents depends on the spectral indices of the Radio and X-ray emission. In Fig. 3 we plot the 3–9 keV X-ray luminosity against the 5 GHz radio luminosity. The radio luminosity is always calculated assuming a flat spectrum, i.e., it is defined as \( L_{5\text{GHz}} = 5 \text{GHz} \cdot S_{5\text{GHz}} \) where \( S_{5\text{GHz}} \) denotes the radio power at 5 GHz. The Chandra X-ray luminosities are corrected for the different energy band assuming a photon index of 1.6. Besides a sample of low hard state XRB (GX 339-4, V404 Cyg, XTE J1118+480, 4U 1543-47, GS 1354-64) we show the two XRBs used to extrapolate to the expected radio flare luminosities: GRS 1915 and Cyg X-3. The data-points have been taken from Gallo et al. (2003). The brightest 1.5 Jy flares of GRS 1915 are not included in the plot. Additionally we present a sample of AGN consisting of low luminosity AGN, FR-I radio galaxies (Fanaroff & Riley 1974) and radio and X-ray selected BL Lac objects. This figure is similar to Fig. 3 in Falcke et al. (2004). Furthermore, we show the theoretical radio/X-ray scaling for black holes of different masses.

A strict radio/X-ray correlation will only be valid for the steady-state emission of low hard state, jet-dominated sources. For high state objects the radio emission will be quenched (see, e.g., Fender et al. 1999a; Gallo et al. 2003). For one possible correlation extending to the very high state see Maccarone et al. (2003). Furthermore, we cannot expect that the correlation holds for the flares of the transient sources. Those flaring sources are only included for comparison.

The empirical correlation can be used to compare the upper limits (4σ) of this radio monitoring campaign with the expectations. The radio flux of high or very high state objects may be reduced compared to the low/hard state objects. The correlation gives an upper bound for the expected radio flux. Only the upper limit for M 33 is close to the value expected from the correlation for stellar mass black holes. All other upper limits are far above the expected radio flux from microblazars with stellar mass black holes. Our non-detections are consistent with either beamed or non-beamed radio emission from stellar or intermediate mass black holes that are \( \leq 10^3 M_\odot \).

Table 6. The radio power of our sample for continuous emission and flares. The radio detections in M 82 are probably SNRs while in NGC 4736 we probably detect the nucleus of the galaxy. For all other sources we present upper limits for continuous emission over all epochs and flares. The distances are taken from Tully (1988).

<table>
<thead>
<tr>
<th>Name</th>
<th>Dist.</th>
<th>( L_X ) [10^{39} \text{ erg s}^{-1}]</th>
<th>( L_{R,\text{flare}} ) [10^{33} \text{ erg s}^{-1}]</th>
<th>( L_{R,\text{cont}} ) [10^{33} \text{ erg s}^{-1}]</th>
</tr>
</thead>
<tbody>
<tr>
<td>M 33</td>
<td>0.7</td>
<td>1</td>
<td>&lt;0.43</td>
<td>&lt;0.2</td>
</tr>
<tr>
<td>NGC 2403</td>
<td>4.2</td>
<td>2</td>
<td>&lt;13</td>
<td>&lt;5.4</td>
</tr>
<tr>
<td>M 82-D</td>
<td>5.2</td>
<td>3</td>
<td>&lt;45</td>
<td>290</td>
</tr>
<tr>
<td>M 82-F</td>
<td>5.2</td>
<td>6</td>
<td>&lt;45</td>
<td>307</td>
</tr>
<tr>
<td>NGC 4736-1</td>
<td>4.3</td>
<td>1</td>
<td>&lt;21</td>
<td>177</td>
</tr>
<tr>
<td>NGC 4736-2</td>
<td>4.3</td>
<td>1</td>
<td>&lt;21</td>
<td>&lt;9.2</td>
</tr>
<tr>
<td>NGC 5204</td>
<td>4.8</td>
<td>2</td>
<td>&lt;18</td>
<td>&lt;7.1</td>
</tr>
<tr>
<td>NGC 5457-1</td>
<td>5.4</td>
<td>2</td>
<td>&lt;23</td>
<td>&lt;9.1</td>
</tr>
<tr>
<td>NGC 5457-2</td>
<td>5.4</td>
<td>1</td>
<td>&lt;23</td>
<td>&lt;9.1</td>
</tr>
</tbody>
</table>

![Fig. 3](http://www.edpsciences.org/aa)
4.4. Other radio detections of ULXs

There are currently two radio detections mentioned in the literature, here we discuss their interpretation within the microblazar model. Kaaret et al. (2003) have detected radio emission from an ULX with an X-ray luminosity $1.1 \times 10^{36}$ erg s$^{-1}$ in NGC 5408. They find a 4.8-GHz radio flux of 0.26 ± 0.04 mJy, well above our detection limit (17σ for continuous emission). Assuming a distance of 4.8 Mpc (Karachentsev et al. 2002) this corresponds to a radio power $7.2 \times 10^{17}$ W/Hz, which is also above the upper limits for our sample of ULXs ($1.5 \times 10^{17}$ W/Hz). For the X-band they give an upper limit of 0.12 mJy (3σ). The spectral index is therefore larger than $\alpha > 1.0$. We would expect that the radio cores of microblazars have a flat to inverted spectrum, by far flatter than the observed steep spectrum. The radio source does not seem to be a transient as it was detected with a roughly similar flux in a subsequent observation (Kaaret private communication). As we have the radio and X-ray luminosity we can use the radio/X-ray correlation to compare the ULX in NGC 5408 with XRBs and AGN. The source seems to be too radio loud for a stellar mass black hole, as shown in Fig. 3. Under the assumption that the source should lie on the correlation, we can derive a mass for the object of $\sim 1000 M_\odot$. Otherwise, the observed emission could be explained by a flare, or by extended emission like a radio supernova. At least this object seems to be different to the ULX sample observed by us.

Kronberg & Sramek (1985) found a 7.07 mJy flare in M 82 ($L_R = 1.1 \times 10^{36}$ erg s$^{-1}$), which we also confirmed by reanalyzing their data. The flare position is within the errors of the brightest ULX in M 82 which has up to $9 \times 10^{36}$ erg s$^{-1}$. This flare is bright compared to GRS 1915, which we expect to show boosted 0.4 mJy flares. However, the value does not seem to be unrealistic for a major flare and the uncertainty of the flux estimate. Up to now no second event has been reported, even though M 82 has been frequently observed. Thus, these events seem to be extremely rare. This is consistent with our conclusion that the duty cycle of ULX flares in our sample is very low (<6%). For comparison we have also shown the flare in Fig. 3. However, the radio/X-ray correlation is only valid for steady-state emission.

Our limiting sensitivity (4σ) of the individual epochs is ≈0.15 mJy ($4 \times 10^{17}$ W/Hz for average distance, rms $\approx 36 \mu$Jy), with the exception of M 82. We have not found any significant radio flare in the sample. If they last only a few days, we are heavily under-sampling the light-curve. However, we can give an upper bound of 6% for the probability of finding a flare brighter than $4 \times 10^{17}$ W/Hz in a given observation of an ULXs. This translates to an upper bound of 11 $\mu$Jy flares per year with a time scale $\delta t_{\text{flare}}$. This is much more than observed for example in GRS 1915. Thus this monitoring campaign cannot rule out or support the idea of relativistically beamed flares in ULXs.

We have reconfirmed the 7 mJy flare near the brightest ULX in M 82 in archival data from Kronberg & Sramek (1985). This corresponds to a radio power of $2.3 \times 10^{19}$ W/Hz. Such a radio flare can be well understood in the context of relativistically beamed emission from the ULX. Another possible explanation is a SN similar to SN 1983n (Kronberg et al. 2000).

We have detected a yet unreported variable double source in NGC 4736. As the position of the stronger source coincides with the position of the nucleus of the LLAGN as reported in Nagar et al. (2002), we assume that this double source is probably connected to the LLAGN and not to the phenomenon of ULXs.

The search for continuous emission from ULXs was only successful for M 82 (limiting flux: $\approx 60 \mu$Jy, 4σ). Here two CRs (probably SNRs) could be associated with ULXs. Thus, ULXs do not generally emit steady-state radio emission above radio powers of $1.5 \times 10^{17}$ W/Hz. Using the radio/X-ray correlation we have shown that our non-detections are consistent with either beamed or non-beamed radio emission from stellar or intermediate mass black holes that are $\leq 10^5 M_\odot$. The search for a radio loud ULXs is therefore still open. Our results suggest that detections will be extremely rare for current telescopes. Nevertheless, the investigation of these intriguing sources should continue.

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Appendix A: The LLAGN in NGC 4736

As we have already mentioned in Sect. 3.2.2, we have detected a double source in NGC 4736 show in Fig. 2. One of these sources is stronger (2.2–1.2 mJy, source b) and variable while the other source seems slightly extended and is fairly stable in flux (1 mJy, source a). The positions are given in Table 5. The peak fluxes and the ratio of the two sources is shown in Table A.1; for a light-curve of the peak fluxes see Fig. A.1.

Due to the low fluxes of both sources it is not possible to self-calibrate the data itself. The fluxes can therefore be reduced by phase errors. They can furthermore be changed by amplitude calibration. However, the flux-ratio of the two sources should be independent of these effects. In the radio the variability can also be seen; it varies between 2 and 1.2 on a time scale of months as also shown in the table. We therefore conclude that at least the stronger source $b$ is variable. This indicates that source $b$ is a compact object, probably the nucleus of the low-luminosity AGN. It has the position that was previously published by Nagar et al. (2002). Source $a$ is not mentioned in earlier publications, but the Nagar et al. (2002) rms value was significantly higher (0.34 mJy). We have reanalyzed their data, and detect a 4σ peak at the position of
Radio flux and beam size of the double source in NGC 4736.

Table A.1. Radio flux and beam size of the double source in NGC 4736.

<table>
<thead>
<tr>
<th>Julian date</th>
<th>Peak 1 [mJy]</th>
<th>Integ. 1 [mJy]</th>
<th>Peak 2 [mJy]</th>
<th>Integ. 2 [mJy]</th>
<th>Ratio peak</th>
<th>Beam size [mas]</th>
</tr>
</thead>
<tbody>
<tr>
<td>59 524</td>
<td>2.25</td>
<td>2.1</td>
<td>1.13</td>
<td>1.4</td>
<td>2.0</td>
<td>251</td>
</tr>
<tr>
<td>59 540</td>
<td>1.45</td>
<td>1.7</td>
<td>1.23</td>
<td>1.8</td>
<td>1.2</td>
<td>279</td>
</tr>
<tr>
<td>59 562</td>
<td>1.36</td>
<td>1.4</td>
<td>1.04</td>
<td>1.3</td>
<td>1.3</td>
<td>261</td>
</tr>
<tr>
<td>59 585</td>
<td>1.30</td>
<td>1.3</td>
<td>0.81</td>
<td>1.1</td>
<td>1.6</td>
<td>268</td>
</tr>
<tr>
<td>59 598</td>
<td>1.91</td>
<td>2.0</td>
<td>1.10</td>
<td>1.7</td>
<td>1.7</td>
<td>260</td>
</tr>
<tr>
<td>59 614</td>
<td>1.50</td>
<td>1.6</td>
<td>1.04</td>
<td>1.3</td>
<td>1.4</td>
<td>414</td>
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<tr>
<td>59 636</td>
<td>1.72</td>
<td>1.44</td>
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<td>504</td>
</tr>
<tr>
<td>59 647</td>
<td>1.82</td>
<td>1.69</td>
<td></td>
<td></td>
<td>1.1</td>
<td>827</td>
</tr>
</tbody>
</table>

Fig. A.1. Light-curve of NGC 4736. Note that the VLA array configuration changed from A to B during the last epochs.

source a (1.4 mJy). Source a was therefore already visible in March 1998.

For the discussion of the other source we have to consider that the array configuration of the VLA changed between epochs five and eight from the A-configuration to B-configuration, increasing the beam size from 260 mas to 840 mas. Extended emission will therefore contribute more to later epochs than to the first. This explains the increasing peak fluxes of the source a from 1 mJy to 1.7 mJy in the last epochs. The effect that the peak flux increases with the beam-size can be seen if one uses a uvtaper of 400 kλ, which roughly doubles the beam-diameter. The peak flux increases from 1.1 to 1.3 mJy in epoch one. The extension of source a is also visible in the difference between the peak and the integrated flux. For source b these two values are usually similar, while for source a the integrated value is higher by approximately 30%.

It is puzzling that the integrated fluxes are more variable (1.1–1.8 mJy) than the peak fluxes. We checked that this effect is not due to phase errors; epoch one and two have phase jumps between successive calibrator scans of less than 40°, usually less than 10°. This would result in only slight (<6%) amplitude losses. One problem of the integrated flux is that the integration boxes are small, as both point sources are less than an arc-second away. Another possibility could be that source a is barely resolved, small changes in the beam size and the atmosphere could then lead to the observed variability. An extended (>200 mas) source which is variable on time scales of months at a distance of 4 Mpc is not possible, as its size would be larger than 5 pc. However, there could be a compact substructure in the second source which is varying.

During the last epoch we also obtained a C-band map of NGC 4736. As the VLA was already in the B-configuration, the two sources could not be separated. Both sources combined give a peak flux of 4.1 mJy and an integrated flux of 7.3 mJy.

To check the spectrum of the two sources we obtained U-band images during the last epoch. Source b has a peak and integrated flux of 1.7 mJy. The weaker source two has a peak flux of 1.1 mJy and 0.95 mJy integrated flux. The missing integrated flux is probably due to the small integration box around the peak needed to avoid source b. The diameter of the beam was 480 mas. We have reanalyzed the old data from Mar. 1998 to check for the second source. Both peaks are visible in the map with fluxes of 2.235 mJy and 1.403 mJy respectively. However, the rms was 0.34 mJy making accurate flux measurements impossible.

We cannot compare the U-band flux directly to the X-band flux, as the beam size of the X-band observations is about twice the U-band beam size. As the flux of the second source is roughly constant during the first epochs, we compare the U-band flux with the last epoch in the A-configuration and taper the data to get a beam-size of 480 mas. The flux of the second source increases to 1.41 mJy. This results in a spectral index between 15 and 8.4 GHz of α = 0.42.

To check whether these sources are compact we have observed NGC 4736 with the European VLBI Network. While source b was detected with a correlated flux of 0.86 mJy, source a was not detected. The rms in this observation was 0.1 mJy.

No other low-luminosity AGN (LLAGN) in 2 cm survey of Nagar et al. (2002) shows such a double structure. The nature of the second source is unknown. Maoz et al. (1995) report for this galaxy a second UV source 2.5" away from the nucleus (see also Maoz et al. 2005) and suggest that this galaxy could be in the final stage of a merger. We have not detected their second source. While the position angle of the UV source is ~3° the second radio source is at ~49°. We conclude that source a could either be a hotspot or a jet feature of the LLAGN. The spectral index is fairly flat. Thus, the hotspot has to be self-absorbed. If this source was observed at a greater distance and the two components unresolved, the spectrum would peak at a few GHz, as source b has a flat spectrum. This source could therefore be a close-up of a GHz peak source (GPS). An other explanation is that inside the extended emission there is a compact flat spectrum source below the detection limit (0.3 mJy) of
our EVN observations. If confirmed, this would be an important result as this could either be a second nucleus or the radio core of an ultra-luminous X-ray source.

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