The WSRT wide-field H I survey

I. The background galaxy sample[★]

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Received 4 April 2003 / Accepted 19 May 2003

Abstract. We have used the Westerbork array to carry out an unbiased wide-field survey for H I emission features, achieving an RMS sensitivity of about 18 mJy/Beam at a velocity resolution of 17 km s⁻¹ over 1800 deg² and between $-1000 < V_{\text{Hel}} < +6500$ km s⁻¹. The primary data consists of auto-correlation spectra with an effective angular resolution of 49' FWHM, although cross-correlation data were also acquired. The survey region is centered approximately on the position of Messier 31 and is Nyquist-sampled over $60 \times 30^\circ$ in RA $\times$ Dec. More than 100 distinct features are detected at high significance in each of the two velocity regimes (negative and positive LGSr velocities). In this paper we present the results for our H I detections of external galaxies at positive LGSr velocity. We detect 155 external galaxies in excess of 8σ in integrated H I flux density. Plausible optical associations are found within a 30' search radius for all but one of our H I detections in DSS images, although several are not previously cataloged or do not have published red-shift determinations. Our detection without a DSS association is at low galactic latitude. Twenty-three of our objects are detected in H I for the first time. We classify almost half of our detections as “confused”, since one or more companions is cataloged within a radius of 30' and a velocity interval of 400 km s⁻¹. We identify a handful of instances of significant positional offsets exceeding 10 kpc of unconfused optical galaxies with the associated H I centroid, possibly indicative of severe tidal distortions or uncataloged gas-rich companions. A possible trend is found for an excess of detected H I flux in unconfused galaxies within our large survey beam relative to that detected previously in smaller telescope beams, both as function of increasing distance and increasing gas mass. This may be an indication for a diffuse gaseous component on 100 kpc scales in the environment of massive galaxies or a population of uncataloged low mass companions. We use our galaxy sample to estimate the H I mass function from our survey volume. Good agreement is found with the HIPASS BGC results, but only after explicit correction for galaxy density variations with distance.

Key words. galaxies: distances and redshifts – galaxies: evolution – galaxies: formation – galaxies: fundamental parameters – galaxies: luminosity function, mass function

1. Introduction

Unbiased wide-field surveys are an indispensable means for determining the physical content of our extended environment. The SDSS (Sloan Digital Sky Survey, York et al. 2000) is a prime example of the way in which such work is providing new insights at optical wavelengths into both the nearby and distant universe. At radio frequencies there have been a number of wide-field surveys for both continuum sources and, to a

lesser extent, emitters and absorbers in specific spectral lines. The HIPASS survey (Barnes et al. 2001) marks an important milestone in achieving high sensitivity to H I emission over more than half of the sky (given its ongoing extension of the Declination coverage from 0 to +25°). HIPASS has provided a deep inventory of both negative and positive LGSr (Local Group Standard of Rest) velocity H I emission features. The negative velocity features (Putman et al. 2002) are primarily associated in some way with the Galaxy and other Local Group objects, while the positive velocity features are primarily associated with moderately nearby (<100 Mpc) external galaxies (e.g. Kilborn et al. 2002). A northern hemisphere counterpart to the HIPASS survey is now underway in the form of HIJASS (H I Jodrell All Sky Survey, Lang et al. 2003).

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* Table 1 is only available in electronic form at the CDS via anonymous ftp to [cdsarc.u-strasbg.fr](ftp://cdsarc.u-strasbg.fr) (130.79.128.5) or via <http://cdsweb.u-strasbg.fr/cgi-bin/qcat?J/A+A/406/829> and Fig. 3 is only available in electronic form at <http://www.edpsciences.org>

An interesting component of the negative velocity H I sky are the so-called compact high velocity clouds, CHVCs (Braun & Burton 1999), which are isolated in position and velocity from the more extended high velocity H I complexes down to column densities below about $N_{\text{HI}} = 1.5 \times 10^{18} \text{ cm}^{-2}$. The suggestion has been made that these objects may be the most distant component of the high velocity cloud phenomenon, perhaps extending to 100's of kpc from their host galaxies. A critical prediction of this scenario (De Heij et al. 2002) is that a large population of faint CHVCs should be detected in the vicinity of M 31 (at declination $+40^\circ$) if enough sensitivity were available. While current observational data are consistent with this scenario, they are severely limited by the modest point source sensitivity available at northern declinations (within the Leiden/Dwingeloo Survey, Hartmann & Burton 1997) which is almost an order of magnitude poorer than that of HIPASS in the south.

We have undertaken a moderately sensitive large-area H I survey both to test for the predicted population of faint CHVCs near M 31 as well as to carry out an unbiased search for H I emission associated with background galaxies. We have achieved an RMS sensitivity of about 18 mJy/Beam at a velocity resolution of 17 km s^{-1} over 1800 deg^2 and between $-1000 < V_{\text{Hel}} < +6500 \text{ km s}^{-1}$. The corresponding RMS column density sensitivity for emission filling the $3000 \times 2800 \text{ arcsec}$ effective beam area is about $4 \times 10^{16} \text{ cm}^{-2}$ over 17 km s^{-1} . For comparison, the HIPASS survey has achieved an RMS of about 14 mJy/Beam at a velocity resolution of 18 km s^{-1} , yielding a slightly superior flux sensitivity. On the other hand, the column density sensitivity for emission filling our larger beam exceeds that of HIPASS by almost an order of magnitude. Since the linear FWHM diameter of our survey beam varies from about 10 kpc at a distance of 0.7 Mpc to more than 1 Mpc at 75 Mpc, it is only at Local Group distances that the condition of beam filling is likely to be achieved. Compared to the Leiden/Dwingeloo Survey, we achieve an order of magnitude improvement in both flux density and brightness sensitivity. We detect more than 100 distinct features at high significance in each of the two velocity regimes (negative and positive LGSr velocities). In this paper we will describe the survey observations and data reduction procedures in Sect. 2, followed by a presentation of the results for our H I detections of external galaxies in Sect. 3 and closing with a brief discussion of these results in Sect. 4. Our results at negative LGSr velocities will be presented in a companion paper.

2. Observations and data reduction

2.1. Survey strategy

Our survey area was defined to have an extent of 60×30 true degrees oriented in $\alpha_{2000} \times \delta_{2000}$ and centered on $(\alpha_{2000}, \delta_{2000}) = (10^\circ, 35^\circ)$, about 5° south of the M 31 nuclear position. Data were acquired in a drift-scan mode, whereby the 25 m telescopes of the WSRT array were kept stationary at a specified start position and the sky drifted past at the earth-rotation rate. Each telescope beamwidth is about 35 arcmin FWHM at an observing frequency of 1410 MHz. The fourteen telescopes of the array were split into two sub-arrays of

seven telescopes each. The two sub-arrays were pointed at declinations offset from one another by 15 arcmin, in order to achieve Nyquist-sampled declination coverage of the survey area in half the time that would otherwise be required. The recorded data were averaged over 60 s, corresponding to an angular drift of about 15 arcmin of right ascension, to yield Nyquist-sampling in the scan direction of the telescope beam.

Although the primary objective of the survey was acquisition of auto-correlation data, it was also desirable to acquire cross-correlation data simultaneously for the two sub-arrays of seven telescopes which observed the same set of positions on the sky. To this end, electronic tracking was employed during each 60 s integration directed at the sequence of central positions that was sweeping through the telescope beam at the earth-rotation rate. The two sub-arrays were each composed of six telescopes with short relative spacings (between 36 and 144 m) and a seventh telescope at a larger separation (of about 1.5 km). The duration of the drift-scan observations varied with declination from about 4.3 to 6.2 hours. A typical observing sequence consisted of a standard observation of a primary calibration source (3C48 or 3C286), a dual sub-array drift-scan observation and in some cases a second dual sub-array drift-scan observation. Each such session provided the survey data for a strip of either 60×0.5 or 60×1 true degrees. Thirty of the “double” sessions, lasting some 320 hours, could in principle provide the complete survey coverage.

In practise, the observations were distributed over some 52 sessions in the period 2002/09/04 to 2002/11/16. An effort was made to acquire the drift-scan data only after local sunset and before local sunrise to minimize solar interference. This was largely successful, with only a few hours of data showing the effects of the sun above the horizon. An effort was also made to insure that the drift-scan data was only acquired at moderately high elevations, both to eliminate the possibility of inter-telescope shadowing and to optimize the system temperature. Essentially all observations were done at elevations between 45 and 85 deg, for which the system temperature variations are observed to be less than 1 K, corresponding to about 3%. Repeat coverage of a number of scans was obtained in cases where instrumental failure or severe interference led to a significant increase in the noise level.

Data was acquired in two 20 MHz IF bands centered at 1416 and 1398 MHz. The 18 MHz spacing of the two bands was chosen to provide a contiguous velocity coverage at a uniform nominal sensitivity. All auto- and cross-correlations were recorded for both linear polarizations in 512 uniformly weighted spectral channels across each 20 MHz band. A hanning smoothing was applied after the fact to minimize the spectral side-lobes of interference, yielding a spectral resolution of 78.125 kHz, corresponding to about 16.6 and 16.8 km s^{-1} in the two bands.

2.2. Data reduction

The drift-scan data for each sub-array was inspected and flagged in Classic AIPS using the SPFLG utility. Any questionable features appearing in the (time, frequency) display of each

auto-correlation baseline were critically compared amongst the 14 independent estimates (7 telescopes and 2 polarizations) that were available. Any features which were not reproduced in the other simultaneous spectra (from telescopes separated by as much as 2 km) were flagged. This allowed quite effective discrimination against interference.

Absolute flux calibration of both the auto- and cross-correlation data was provided by the observed mean cross-correlation coefficient measured for the standard calibration sources (3C48 or 3C286) of known flux density. The measured ratio of flux density to correlation coefficient averaged over all 14 telescopes and 2 polarizations was 300 ± 10 Jy/Beam. Although there are variations (typically less than about 10%) amongst the 28 independent receiver systems, the average gain of the system (at this frequency) has remained constant at the quoted value over a period of several years to better than 5 percent. The calibrator observation of each observing session was used both to determine the average gain value appropriate for the auto-correlations as well as providing phase and gain solutions appropriate for the calibration of the cross-correlation drift-scan data in that session.

Two different methods were employed to generate data-cubes of the auto-correlation data. The first method employed a local robust average of a 30 min sliding window to estimate the band-pass as a function of time and a 850 km s^{-1} sliding window to estimate the continuum level as a function of frequency. Only those values between the first and third quartiles were included in these averages, making them moderately robust to outliers, including HI emission features, in the data. This method could be applied blindly and produced the most uniform noise characteristics in the resulting cube. As such, it was well-suited for the automated detection of faint sources. However, moderately bright sources of HI emission that were extended either spatially or in velocity produced localized negative artifacts. The best results under these circumstances were obtained with the more complicated procedure outlined below: 1) a quadratic baseline in time was fit to the entire drift-scan (of 4 to 6 hour duration) and divided out after masking out any localized regions of emission, 2) a constant offset was determined and divided out of each frequency spectrum after masking out any regions of line emission (including the extended emission from the Galaxy), 3) a quadratic baseline in frequency was fit and subtracted from the lower half (the first 10 MHz) of each frequency spectrum after masking out any regions of line emission, 4) an alternative baseline solution (primarily for the upper 10 MHz of each band) was derived from a boxcar smoothed spectrum with 6 MHz box-width and subtracted from each frequency spectrum after masking out any regions of line emission, 5) a cubic baseline in time was fit and subtracted out of each drift-scan after masking out any localized regions of emission, 6) all unflagged data for each position and frequency were averaged, 7) the entire process (steps (1) through (6) above) was repeated using an updated mask to isolate regions of significant emission from the baseline determination.

The rationale for each step noted above was the following: step (1) was intended to compensate for long timescale

variations in the basic bandpass shape, step (2) for the possible contributions of bright continuum sources to the system temperature as function of time, steps (3) and (4) for residual corrections to the bandpass shape on short timescales (where it was found empirically that the two different methods gave somewhat superior results in the two halves of the band) and step (5) for residual corrections to the bandpass shape on intermediate timescales.

Although the survey strategy was designed to provide the most stable possible bandpass response, this proved to be somewhat disappointing. During the course of the observing campaign it was established that there were systematic variations in the bandpass shape at the level of about 1:1000 which were closely correlated with small variations in ambient temperature in the vicinity of the IF system electronics on timescales of 0.5 to 5 hours. Steps were taken to stabilize the airflow and ambient temperature which led to a substantial improvement in baseline stability. Even so, the remaining fluctuations are still a limiting factor to the final auto-correlation sensitivity, at least in the upper one third of each 20 MHz band, where they were most severe.

The drift-scan data were resampled in frequency to convert from the fixed geocentric frequencies of each observing date to heliocentric radial velocities at each observed position. A Gaussian smoothing with 1800 true arcsec FWHM was applied in RA before combining the 120 drift-scans into a data-cube for each 20 MHz band. A spatial resampling was then carried out to convert from the rectangular (α, δ) of the acquired data to a projected (α, δ) geometry. Finally, a spatial convolution with a 1800×900 arcsec FWHM Gaussian with PA = 0° was applied to introduce the desired degree of spatial correlation in the result.

The average primary beam of the WSRT array at an observing frequency of 1400 MHz has been determined with a holographic measurement (and can be found depicted at <http://www.astron.nl/wsrt/wsrtGuide/WSRT21BEAM.PS>). The central positive lobe has an area of 1535 arcmin^2 and is moderately well-fit by a Gaussian with 2235 arcsec FWHM. Peak sidelobe levels are about 0.1%. The same sequence of spatial smoothings described above, first a $900''$ box-car in RA to simulate the drift-scan, then a $1800''$ Gaussian in RA and then a $900 \times 1800''$ Gaussian in (α, δ) was applied to a digital representation of the primary beam. This allowed estimation of the effective beam size for the survey, which can be approximated by a Gaussian of 3020×2810 arcsec at PA = 90° , as well as the beam dilution factor needed to preserve the absolute flux scale in units of Jy/Beam.

The resulting RMS fluctuation level as function of heliocentric velocity is shown in Fig. 1 as determined from a fit to the peak of the histogram in each velocity channel. The effective velocity coverage extends from about $-1000 < V_{\text{Hel}} < 6500 \text{ km s}^{-1}$. An increased noise level is seen for $-200 < V_{\text{Hel}} < 150 \text{ km s}^{-1}$ due to intermediate velocity emission features associated with the Galaxy. Slightly elevated noise levels are also seen for $1700 < V_{\text{Hel}} < 2600 \text{ km s}^{-1}$ and $5500 < V_{\text{Hel}} < 6500 \text{ km s}^{-1}$ due to the problems of spectral baseline stability in the upper third of each 20 MHz band as noted above. A small velocity range near 1500 km s^{-1} also has a slightly higher noise due to a significant degree of data

Fig. 1. Variation of RMS fluctuation level with heliocentric velocity over the survey velocity coverage.

flagging at this velocity that was made necessary by recurring interference.

3. Results

An overview of the survey sky coverage, together with an indication of some of the galaxy detections, is given in Fig. 2. Peak H I brightness at a velocity resolution of 42 km s^{-1} FWHM is shown in the figure for the velocity interval $250 < V_{\text{Hel}} < 1500 \text{ km s}^{-1}$. The first contour in the figure corresponds to approximately 4σ at this velocity resolution. The total solid angle observed is 1800 deg^2 . The RMS fluctuation level varies slightly with recession velocity (as shown in Fig. 1). The positions of bright continuum sources (brighter than a few Jy) display residual fluctuations in excess of the nominal noise level. The resulting noise distribution is not entirely Gaussian, and for this reason we consider it necessary to adopt a conservative cut-off in our blind extraction of reliable H I detections. Rosenberg & Schneider (2002) have shown from their extensive simulations involving insertion of artificial sources into surveys of this type that an asymptotic completeness level is reached at a signal-to-noise ratio of about 8 (in terms of integrated signal strength relative to the error in the integral), while below this signal-to-noise ratio the completeness drops dramatically.

Candidate H I detections in our combined data-cubes were determined by two different methods. In the first instance, the SAD source finding algorithm within Classic AIPS was used to extract all local peaks in excess of 3 times the local RMS level in datacubes having Nyquist-sampled velocity resolutions of 2, 5, 10, 20 and 50 times the basic velocity channel separation of 8.4 km s^{-1} . A reduced list of emission candidates having a detected peak in excess of 5σ in at least two different velocity channels or two different velocity smoothings was extracted for further analysis. The aim of this procedure was to reliably recover significant detections spanning a wide range in observed linewidth. A complimentary list of candidates was determined from visual inspection of subsequent channel maps as well as subsequent position-velocity projections using the KVIEW display program (Gooch 1995) for the data-cubes at velocity resolutions of 2, 5, 10, 20 and 50 times the basic velocity channel

separation of 8.4 km s^{-1} . Candidate features were rejected if their response in our data-cubes before spatial smoothing were inconsistent with the telescope beam response in either α or δ , implying the time variable signature of interference. The properties of some 500 candidate detections were then estimated in detail. In particular, the integrated line strength was determined for each candidate by extracting the single spectrum from our spatially smoothed data-cubes with the highest total flux density. The underlying assumption is that essentially all galaxies will be unresolved with our effective FWHM beamsize of $3020 \times 2810 \text{ arcsec}$, corresponding to $73 \times 68 \text{ kpc}$ at the nearest galaxy distance of 5 Mpc. This assumption was tested by comparing the peak with the integrated flux of a Gaussian fit to an image of integrated H I for each source. While some sources show clear signs of confusion from nearby companions (which will be discussed in detail below), there were no instances of our having significantly resolved single galaxies with our survey beam. The associated error in flux density was determined over a velocity interval of $1.5 \times W_{20}$ (where W_{20} is the velocity width of the emission profile at 20% of the peak intensity) together with the actual RMS fluctuation level associated with this velocity interval.

Only the 155 candidates with an integrated flux density exceeding 8 times the associated error were retained. Spectra of each detection are shown in Fig. 3 for the single spatial pixel with the maximum integrated H I signal. The source centroid positions were determined from either a Gaussian or a parabolic fit to the peak in images of integrated H I line strength over the full velocity extent of each detection. The positional accuracy is dependent on the signal-to-noise ratio, and is expected to be roughly $\text{FWHM}/(s/n)$, implying about 3 arcmin RMS in both α and δ for the lowest significance detections.

3.1. Galaxy properties

The properties of our H I selected detections are summarized in Table 1. In addition to the position of the H I centroid and integrated flux density, F_{HI} in units of Jy-km s^{-1} , we tabulate the heliocentric recession velocity, V_{Hel} , and velocity width at 20% of the peak brightness, W_{20} . The integrated flux density has been converted to an H I mass by first calculating the recession velocity in the Local Group Standard of Rest frame, $V_{\text{LGSR}} = V_{\text{Hel}} + 300 \sin l \cos b$, then assuming a Hubble constant, $H_0 = 75 \text{ km s}^{-1} \text{ Mpc}^{-1}$ to derive the distance, $D = V_{\text{LGSR}}/H_0$ in Mpc, and finally using $M_{\text{HI}} = 2.356 \times 10^5 D^2 F_{\text{HI}}$ under the simplifying assumption of negligible H I opacity.

Seven of our detections have recession velocities that lie near the transition in our velocity coverage from the lower to the upper band of 20 MHz width, at $V_{\text{Hel}} \sim 2800 \text{ km s}^{-1}$. While our velocity coverage is complete across this transition, the spectral baseline extent is severely impaired for such objects. Consequently, there is a large systematic uncertainty in the integrated flux density, recession velocity and velocity width of these detections. Two of our detections (UGC 11864 and UGCA 020) lie at the edge of our nominal sky coverage, so that they are not properly sampled and also have large

Fig. 2. Illustration of the survey sky coverage and galaxy detections between $250 < V_{\text{Hel}} < 1500 \text{ km s}^{-1}$. The peak brightness of H I emission at a velocity resolution of 42 km s^{-1} FWHM is shown in this velocity interval with a linear greyscale between 5 and 100 mJy/Beam and contours drawn at (1, 2, 5, 10, 20 and 50) times 45 mJy/Beam.

uncertainties in their properties. We indicate these large uncertainties by a “:” suffix in Table 1 and do not use our measured properties for these sources in the subsequent analysis.

The spatial distribution of our H I detections derived from the distance calculated as above is shown in Fig. 4, both for the entire depth of our survey (about 80 Mpc) and out to 20 Mpc. A moderate galaxy concentration out to about 15 Mpc is followed by an apparent void over much of our surveyed field out to about 45 Mpc, which in turn is followed by a substantial increase in detected number density out to about 80 Mpc. A galaxy filament along the eastern edge of our coverage connects the nearby and more distant concentrations.

3.2. Optical ID's

Cataloged counterparts of our H I detections were sought in the NASA/IPAC Extragalactic Database (NED) on 2003/03/01 within an extended error circle of 30 arcmin radius. This search radius was chosen since it corresponds to the radius of the first null in the primary beam of the WSRT telescopes. Only objects within this radius can contribute significantly to our detected H I fluxes. Identifications with cataloged galaxies having published red-shifts was possible in most cases. The

NED ID's of our H I detections are listed in Table 1 together with the angular offset of the NED position from that of the H I centroid. Some objects deserving special attention are noted below.

J2202+4838, corresponding to $(l, b) = (96.5, -5.4)$, has no cataloged optical counterpart nor candidate galaxy visible in the DSS. Given the low galactic latitude of this line-of-sight this is perhaps not too surprising.

J2339+2509 appears to be associated with CGCG 476-100, although no previous red-shift is available for this galaxy.

J0136+4759 appears to be associated with an uncataloged LSB galaxy at $(\alpha, \delta) = (01:36:40, +48:03:40)$, lying very near a bright foreground star.

J0239+3015 is very likely associated with the NED galaxy [VR94] 0236.9+3003 with tabulated photometry by Vennik & Richter (1994), but without a previous red-shift determination.

J0302+4852 appears to be associated with the cataloged source HFLZOA G144.00-08.53 which has no previous red-shift determination.

In the final column of Table 1 we give an indication of known and possible companions of our detections. We list the number (truncated at a maximum value of nine) of NED

Fig. 3. H I spectra of detected galaxies. The complete set of panels for all detected galaxies is only available in the electronic version of the paper at <http://www.epdsciences.org>. The catalog designation is indicated in the upper left of each panel, while the likely optical ID is indicated in the upper right.

galaxies within a 30 arcmin search radius of the primary optical ID which have (a) a known red-shift within 400 km s^{-1} of the primary ID, (b) a known red-shift between 400 and 1000 km s^{-1} of the primary ID, and (c) an unknown redshift. These three categories of possible companion galaxies have been used to define a confusion index relevant to our survey made up of the

three digits “abc”. All sources with a confusion index of 100 or greater are indicated by a “c” entry in the Note column of Table 1. A total of 85 of our 155 detections are classified as “unconfused” by this criterion. The number of possible companions in the third category considered (no known red-shift) deserves some further comment. Although some of our

Fig. 4. Wedge diagrams for the H I detected galaxies in our survey. All detected galaxies are depicted on the left, while only those within 20 Mpc are shown on the right.

Fig. 5. Comparison of the expected RMS positional error with the observed positional offsets of our H I detections from the primary optical counterpart. The dotted lines correspond 1, 2.5, 5 and 10σ . Dashed lines are drawn at an offset of 3 arcmin and 10 kpc, which may distinguish cases of internal asymmetries from confusion with gas-rich companions. Galaxies with no cataloged companions within 30 arcmin and 400 km s^{-1} are plotted as filled circles, while those with such companions are plotted as open circles.

detected galaxies have a large number of objects (as many as 66) in this category, it is often merely an indication that the general field has received intensive study, usually directed at a distant background galaxy cluster. While this category remains an ambiguous and non-uniform measure of possible companions, it still gives some indication of what is known of the galaxy environment.

We compare the anticipated centroid error of our H I detections (from $\sqrt{2} \text{ HWHM}/(s/n)$) with the angular and linear offsets of the NED ID's in Fig. 5. The dotted lines in the figure correspond to 1, 2.5, 5 and 10 times the estimated centroid error. Nominally unconfused galaxies (those with no known

companions within 30 arcmin and 400 km s^{-1}) are plotted as filled circles in the figure, while galaxies with known nearby companions are plotted as the open circles. The mean observed offset for the unconfused galaxies is 4.8 arcmin and 66 kpc, although the majority of these have a low significance. A small concentration of significant centroid offsets is seen below about 3 arcmin and 10 kpc. These cases are indicated by an “o” entry in the Note column of Table 1. This component may be due to asymmetries in the H I distribution of individual objects, since it corresponds to sub-galactic dimensions. Most of the large observed offsets occur in cases of galaxies with cataloged companions. In addition, there are a small number of instances

of larger angular offsets of high significance in apparently unconfused galaxies. Positional offsets larger than 10 kpc and 5σ are seen in 6 instances, and larger than 10 kpc and 10σ in two. These cases have been indicated in the Note column of Table 1 by a “+” symbol entry for offsets larger than 10 kpc and 5σ and a pair of “+” symbols for offsets larger than 10 kpc and 10σ . This component of offsets is suggestive of either severe asymmetries in the gas distribution of single galaxies or nearby uncataloged gas-rich companions within the telescope beam.

In Fig. 6 we present an atlas of images taken from the second generation digital sky survey of the Space Telescope Science Institute for all of the H I detections with somewhat ambiguous optical ID’s. Those fields are depicted which have a significant angular offset (more than 5σ) from the optical ID, or which have no prior red-shift determination. In each case a red 30×30 arcmin field centered approximately on the H I centroid was extracted.

3.3. Previous H I detections

Previous measurements of the H I content of our detections were available within NED and LEDA (and the references tabulated there) in 132 of 155 cases. We plot our flux densities against the LEDA values in Fig. 7. The H I fluxes tabulated in LEDA correspond to weighted averages of all previously published values. Those cases marked in the last column of Table 1 as being possibly confused in our telescope beam (having one or more cataloged companion galaxies within 30 arcmin radius and 400 km s^{-1}) are plotted separately as the open circles in Fig. 7. The distribution is consistent with essentially the same absolute flux scale for the isolated galaxies. A linear regression solution (fit to the linear fluxes rather than their logarithm) is overlaid on the data in Fig. 7 and has a slope of 0.992, corresponding to a mean flux-scale discrepancy of less than 0.8%. The confused galaxies of our sample show both a larger scatter and a systematic trend for an excess H I flux detected in our larger telescope beam.

3.4. H I in galaxy environments

An important difference between the flux measurements reported here and those in the literature is the large effective beam size of our survey. Indeed, compared to the $3.1 \times 3.7'$ FWHM beam of the upgraded Arecibo telescope, our $46 \times 49'$ beam has a 200 times greater solid angle. The linear FWHM diameter of our survey beam varies from about 70 kpc at the nearest galaxy distance of 5 Mpc to more than 1 Mpc at the furthest galaxy distance of 75 Mpc.

Given our larger beam area it is interesting to search for any systematic increase in the H I we detect relative to what has been detected previously in a smaller beam. In the first instance we plot the ratio of our survey H I flux relative to that tabulated by LEDA as function of distance and H I mass in Fig. 8. Only those galaxies for which the flux ratio had a signal-to-noise greater than 5 are plotted, after taking account of the uncertainties in both our value and that tabulated by LEDA.

The data-points are relatively few in number and quite noisy. Essentially no excess flux is seen as a function of distance (the distribution has a correlation coefficient, $r = 0.080$ and Student’s $t = 0.269$, corresponding to a probability of significance, $P = 60\%$), while a weak trend of excess detected H I flux may be present as function of H I mass ($r = 0.195$, $t = 0.946$, $P = 85\%$). The linear regression solutions with equal weights given to all points are overlaid in both cases.

The LEDA data has been obtained from a wide variety of sources with a corresponding variety in both beam size and calibration strategy. To eliminate these variables from the flux comparison we also plot the ratio of our survey H I flux relative to that measured previously with the Arecibo telescope as function of distance and H I mass in Fig. 9. The Arecibo data is taken from the Pisces-Perseus supercluster survey (Giovanelli & Haynes 1985; Giovanelli et al. 1986; Giovanelli & Haynes 1989; Wegner et al. 1993; Giovanelli & Haynes 1993). We have only plotted the data for apparently isolated galaxies; those with no cataloged companions lying within a radius of 30 arcmin and 400 km s^{-1} as indicated in the last column of Table 1. The flux ratio calculated from the *observed* Arecibo H I flux is plotted as the filled circles with error bars. A corrected Arecibo H I flux is also listed in these references, in which approximate corrections are applied for telescope pointing errors, a model of the finite angular extent of the target galaxies and the likely effect of H I self-absorption. Since the correction for H I self-absorption should affect both our measurements to an equal degree, we undo this correction to the Arecibo fluxes before calculating the ratio. We plot the *corrected* Arecibo to WSRT flux ratio as the dotted open circles in the figure. The error bars in the corrected flux ratio do not take account of the uncertainties in the correction procedure.

While the data-points are even fewer in number, they may suggest a weak systematic excess of detected H I flux within our larger survey beam. For the flux ratio as a function of distance the distribution is characterized by $r = 0.120$, $t = 0.404$ and $P = 65\%$, while as a function of mass $r = 0.183$, $t = 0.619$ and $P = 73\%$. The corrected flux ratio still shows an excess for the majority of sources (with $r = 0.162$, $t = 0.475$ and $P = 68\%$ as a function of distance and $r = 0.187$, $t = 0.819$ and $P = 78\%$ as a function of mass), although the incidence of several corrected flux ratios significantly less than unity casts some doubt on the general reliability of the correction procedure.

3.5. New H I detections

Twenty-three of the objects listed in Table 1 have been detected in H I for the first time in our survey. We comment briefly on each of these objects below.

J2202+4838, at $(l, b) = (96.5, -5.4)$ is the only object in our 8σ sample with no apparent optical counterpart within a 30 arcmin search radius in the second generation DSS images. The predicted extinction in this direction is moderate but not extreme, $A_B = 1.32 \text{ mag}$ (Schlegel et al. 1998).

UGC 11923, classified merely as type “S”, has received relatively little study, no doubt due in part to its position

Fig. 6. Atlas of ambiguous optical ID's. All H I detections are depicted which have a significant angular offset (more than 5σ) from the optical ID, or which have no prior red-shift determination. Countours at 80, 90 and 97% of the peak in integrated H I are overlaid on a 30×30 arcmin red frame from the second generation DSS. A square-root transfer function is used for the optical image.

Fig. 7. Comparison of our survey integrated H I fluxes with previous determinations in the literature. The objects indicated by open circles have cataloged companions within 30 arcmin and 400 km s^{-1} and are classified as confused. The linear regression solution excluding these objects is plotted as a dotted line and has a slope within 0.8% of unity.

$(l, b) = (94.9, -9.3)$ and consequently relatively high extinction, $A_B = 1.74 \text{ mag}$ (Schlegel et al. 1998). This is a moderately gas-rich system, with $\log(M_{\text{HI}}) = 10.06$.

UGC 11929 is a little-studied S0 galaxy with IRAS fluxes in the 60 and $100 \mu\text{m}$ bands of 1.3 and 2.7 Jy.

KKR 71 is a nearby irregular system that has received little study.

CGCG 514-098 is a moderately distant unclassified galaxy with IRAS fluxes in the 60 and $100 \mu\text{m}$ bands of 1.3 and 3.3 Jy.

CGCG 476-100 is an unclassified galaxy with no previous red-shift determination.

UGC 64 is a nearby low mass system with a peculiar optical morphology, possibly suggestive of tidal interaction (Vorontsov-Velyaminov 1977).

And IV was only recently recognized (Ferguson et al. 2000) as a low mass dwarf in the background (at 5 to 8 Mpc distance) of M 31, rather than being closely associated with M 31 as had been thought previously. In fact, together with UGC 12894, UGC 64 and UGC 288, And IV forms a nearby filament of low-mass galaxies.

CGCG 501-016 is an unclassified galaxy that was host to SN 1995am with IRAS fluxes in the 60 and $100 \mu\text{m}$ bands of 0.2 and $<0.84 \text{ Jy}$.

CGCG 502-039 is an unclassified galaxy with IRAS fluxes in the 60 and $100 \mu\text{m}$ bands of 1.1 and 1.74 Jy.

CGCG 521-039 is an unclassified galaxy which has received little study.

J0136+4759 appears to be associated with a previously uncataloged LSB galaxy at $(\alpha_{2000}, \delta_{2000}) = (01:36:40, +48:03:40)$ as illustrated in Fig. 6. The optical galaxy was presumably not

recognized previously due to its close proximity to a moderately bright foreground star.

NGC 661 is classified as an E+, with the UGC noting a diffuse companion at 3.8 arcmin offset in PA 261 east of north. Chamaraux et al. (1987) report a non-detection ($<0.57 \text{ Jy-km s}^{-1}$) at the optical position and red-shift measured in one Arecibo beam (3.9 arcmin FWHM), suggesting that our detection, with its 9.4 arcmin offset, may be due to either a companion of NGC 661 or tidal debris.

[ZBS97] A31 was first detected in the AHIS survey (Zwaan et al. 1997), where it is noted as having $\log(h^{-2} M_{\text{HI}}) = 8.97$. Our detection centroid is offset by about 10 arcmin and appears to be significantly more massive.

NGC 780 is an unclassified galaxy, with apparent stellar plumes extending in several directions from the main galaxy body, suggesting a recent merger remnant.

CGCG 538-034 is classified as S0 and has IRAS fluxes in the 60 and $100 \mu\text{m}$ bands of 3.0 and 3.0 Jy.

UGC 01830 is classified as an SB0/a and has IRAS fluxes in the 60 and $100 \mu\text{m}$ bands of 1.2 and 3.1 Jy. The UGC notes: “Very compact core, double ring halo.”

IC 1815 is classified as an SB0, and has a diffuse stellar halo.

J0239+3015 appears to be associated with the NED object [VR94] 0236.9+3003 with tabulated photometry by Vennik & Richter (1994), but without a previous red-shift determination.

J0240+4221 = IRAS 02371+4223 is an unclassified galaxy with IRAS fluxes in the 60 and $100 \mu\text{m}$ bands of 0.7 and 1.3 Jy.

UGC 2172 represents the lowest integrated H I flux detection of our sample with only $3.1 \pm 0.35 \text{ Jy-km s}^{-1}$, corresponding to $\log(M_{\text{HI}}) = 7.85$. A previous unsuccessful search for H I in this galaxy by Schneider et al. (1992) was directed at the incorrect velocity interval, since an optical red-shift only became available in 1999.

NGC 1161 is classified as an S0 galaxy and has previously been searched for H I emission by Haynes et al. (1990 using the Green Bank 300 ft telescope down to an RMS sensitivity of 3.5 mJy/Beam over 5.5 km s^{-1} velocity channels. The extreme positional offset of our detection (13.5 arcmin corresponding to 100 kpc) coupled with the GB300’ non-detection suggest that we are likely detecting either tidally stripped gas at large radii or a gas-rich companion rather than NGC 1161 itself.

HFLZOA G144.00-08.53 is an obscured system ($A_B = 2.05 \text{ mag}$ (Schlegel et al. 1998), classified as a dE, although the compact central concentration is surrounded with more diffuse stellar emission in the DSS. No previous red-shift is available for this source.

3.6. The H I mass function

An important application of blind H I surveys is the characterization of the general population of neutral gas-rich objects without the inevitable bias associated with the study of an optically-selected sample. The 155 H I detections which follow from our 8σ limit on integrated H I flux, form a relatively small, but moderately complete sample with which

Fig. 8. Plots of our survey integrated H I flux relative to those measured previously as function of distance (left) and H I mass (right) for apparently isolated galaxies. Only those galaxies with no cataloged companions within a radius of 30 arcmin and 400 km s^{-1} are plotted. Linear regression solutions are overlaid.

Fig. 9. Plots of our survey integrated H I flux relative to that measured with the Arecibo telescope as function of distance (left) and H I mass (right) for apparently isolated galaxies. Our survey beam probes a 200 times larger solid angle that varies in linear diameter from about 70 kpc at 5 Mpc to more than 1 Mpc at 75 Mpc. Only those galaxies with no cataloged companions within a radius of 30 arcmin and 400 km s^{-1} are plotted. The *observed* flux ratios are plotted as filled circles with error bars, while dotted open circles are used for flux ratios that have been corrected in an approximate way for telescope pointing errors and estimated source extent as seen with Arecibo. Linear regression solutions are overlaid.

to characterize the population. Indeed, extensive simulations with synthetic sources in a comparable survey led Rosenberg & Schneider (2002) to conclude that above an effective signal to noise ratio of 8, their sample of H I selected objects was essentially complete. Since our sample is limited in flux density, which scales as the inverse square of distance for a given H I mass, it is clear that we sample a very different survey volume at low H I mass relative to high. For example, an object with $\log(M_{\text{HI}}) = 7.5$ and $W_{20} = 40 \text{ km s}^{-1}$ can only

be detected out to $D = 5.4 \text{ Mpc}$, while Galactic H I emission extends out to about $V_{\text{Hel}} = 200 \text{ km s}^{-1}$, corresponding to $V_{\text{LGSR}} \sim 350 \text{ km s}^{-1}$ and $D = 4.7 \text{ Mpc}$, leaving a possible detection volume of only some 10 Mpc^3 . On the other hand an object with $\log(M_{\text{HI}}) = 10$ and $W_{20} = 300 \text{ km s}^{-1}$ can be detected beyond the edge of our survey volume at about $D = 88 \text{ Mpc}$, corresponding to more than 10^5 Mpc^3 . When comparing our detections in different H I mass bins it becomes important to

Fig. 10. Optical luminosity distributions over the survey region in the indicated intervals of $\log(D)$. The solid curves are Schechter functions with $M_B^* - 5 \log_{10} h = -19.66$, $\alpha = -1.21$ and the indicated density relative to $\Phi_* = 1.61 \times 10^{-2} h^3$. The “standard” optical luminosity function parameters are taken from Norberg et al. (2002).

consider whether the average space density of galaxies is actually uniform within our survey volume.

The wedge diagrams of our HI detections shown in Fig. 4 already gave some indication for non-uniformity of the space density along the line-of-sight, despite the rather substantial solid-angle of our survey. A moderate galaxy density is seen between 5 and 15 Mpc, followed by an apparent void and subsequently another enhancement. Since this is difficult to quantify directly on the basis of our limited number of detections, we have instead considered the space density of optically cataloged galaxies as a function of recession velocity. We extracted from the LEDA database all galaxies of known LGSR recession velocity (greater than $+250 \text{ km s}^{-1}$) and integrated B -band magnitude within the spatial and velocity boundaries of our survey region. A total of 1774 galaxies are cataloged within LEDA within our survey boundaries. We adopt an approximate completeness limit of $m_{BT} = 14.9$, based on the turn-over in

the cumulative distribution function. Constraining the selection to $m_{BT} < 14.9$, we retain about $N_{\text{Gal}} = 530$ optically cataloged galaxies in our survey volume. After sorting these galaxies by distance, they have been divided into 25 overlapping sub-samples with sample populations varying linearly from a minimum of $N_{\text{Pop}} = 40$ galaxies at the nearest distances to $N'_{\text{Pop}} = 100$ galaxies at the maximum distance. Adjoining sub-sample populations share more than half of their membership to insure sufficient sampling of the density variation with distance. (This has been accomplished by choosing the start index, $i = 1 \dots N_{\text{Gal}}$, of sub-sample, $j = 1 \dots 25$, in the ordered galaxy list using the prescription $i = (N_{\text{Gal}} - N'_{\text{Pop}})^{((j-1)/24)^{0.8}}$.) We plot the optical B -band luminosity distributions of these 25 sub-samples in Fig. 10. The relevant intervals of $\log(D)$ are indicated at the top of each panel. Also plotted is a “standard” luminosity function taken from Norberg et al. (2002),

$$d\Phi(M)/dM = \Phi_* \ln(2.512)(L/L_*)^{\alpha+1} \exp(-L/L_*),$$

Fig. 11. The derived galaxy over-density in our survey volume as function of distance (top) and the distribution of detected HI masses with distance (bottom). The over-density is derived from fits to the luminosity distribution of optical galaxy sub-samples as shown in Fig. 10. The dotted line in the top panel indicates the same measure of optical galaxy over-density derived for the HIPASS survey volume ($\delta < 0$), while the dashed line is that for the northern hemisphere ($\delta > 0$).

with L_* corresponding to $M_B^* - 5 \log_{10} h = -19.67$, $\alpha = -1.21$ and the indicated density relative to $\Phi_* = 1.61 \times 10^{-2} h^3$, which they derive from more than 10^5 galaxies in the 2dF red-shift survey. As can be seen in the panels of the figure, the luminosity distributions are moderately complete, with only the occasional down-turn in the lowest luminosity bin. All of these distributions can be reasonably well-described by the same ‘standard’ luminosity function, where the only permitted variable in a χ^2 minimization was the galaxy space density relative to the global average value of $\Phi_* = 1.61 \times 10^{-2} h^3$ found by Norberg et al. Values of $\Delta\chi^2 = 1$, corresponding to 1σ errors in the over-density, varied from about 20% at the smallest distances to 10% at the largest distances. These appear to be realistic error estimates under the assumptions that the density is a smoothly varying distribution in our survey volume and that the shape of the luminosity function is not also a function of distance. The large degree of overlap of the sub-samples in adjoining distance intervals does not have an adverse impact on the error estimate, but simply insures sufficient sampling of changes in the density with distance.

Our derived variation of galaxy density within our survey boundaries as function of distance is plotted in the upper panel of Fig. 11, with the 1σ error bars noted above. The local galaxy density (within about 10 Mpc) is slightly below the standard 2dF value. This plummets to some 20% of the average in the void near $D = 25$ Mpc, slowly climbs to an over-density centered at 70 Mpc and subsequently declines. The impact of such a variation of density can be judged in the lower panel of Fig. 11, where our detections are plotted as a function of HI mass and distance.

Fig. 12. Distribution of effective survey volume (bottom) and density-corrected luminosity function (top). The solid and dashed curves in the lower panel contrast the “discrete” and “integral” formulations of the effective survey volume as function of the limiting distance, D_{max} , under the simplifying assumption that most detections occur near the limiting distance. The symbols in the top panel give the observed (open square) and density-corrected luminosity functions (filled circles for the discrete-, and crosses for the integral formulations of V_{eff}) in our entire survey volume. The solid and dashed curves are Schechter function fits to the density-corrected data from the discrete and integral formulations.

The suggested method of introducing density corrections in optical luminosity functions or an HIMF (Saunders et al. 1990; Rosenberg & Schneider 2002) is the calculation of a so-called effective volume, $V_{\text{eff}} \propto \int_{D_{\text{min}}}^{D_{\text{max}}} \rho(D) D^2 dD$, to replace the physical volume $V_{\text{tot}} \propto \int_{D_{\text{min}}}^{D_{\text{max}}} D^2 dD$, where the integrals extend over the entire distance range over which each source could have been detected. We considered that a more straightforward method of achieving the desired result, of a “uniform density” survey volume, might be simply $V'_{\text{eff}} \propto \rho(D) \int_{D_{\text{min}}}^{D_{\text{max}}} D^2 dD$, where the density has been taken out of the integral and is only evaluated at the distance of the detection in question. Although somewhat challenging to compare these two formulations of V_{eff} directly, this can be done approximately by considering that the largest number of detections of any distribution which is rising at its faint end will be near the limiting distance, D_{max} . With this simplifying assumption we compare these two formulations in the lower panel of Fig. 12, where the ratio of effective to physical volume is plotted as a function of D_{max} . The “discrete” formulation of V'_{eff} is plotted as the solid line and exactly traces the distribution of over-density plotted in Fig. 11 after normalization with the physical volume. The “integral” formulation of V_{eff} is plotted as the dashed line in Fig. 12. This formulation shows large systematic departures from the discrete one, with peaks and troughs shifted to higher distances and having decreased amplitude.

Fig. 13. Distribution of detected H I masses (bottom) and the derived HIMF (top). The open squares are derived from a straightforward application of $1/V_{\text{tot}}$, while the filled circles have been corrected for the variation of galaxy density with distance. The dashed and solid curves are the best-fitting Schechter functions with $\log(M_*) = 10.2$, $\Theta_* = 9.5 \times 10^{-4}$ and $\alpha = -1.48$ for $1/V_{\text{tot}}$, and $\log(M_*) = 9.85$, $\Theta_* = 55. \times 10^{-4}$ and $\alpha = -1.28$ after the density correction. The dotted curve represents the reciprocal of the limiting survey volume as function of mass.

The density-corrected optical luminosity functions which follow from the discrete and integral formulations are compared in the upper panel of Fig. 12. The open squares in this figure indicate the accumulated $(1/V_{\text{tot}})$ luminosity function over our entire survey volume with no density correction, while the filled circles and crosses are the density-corrected luminosity functions. Both forms of density correction give a substantial improvement in recovering the template luminosity function of Norberg et al. (2002). The best χ^2 fits (constrained only to have the Norberg et al. power-law of -1.21) are overlaid in the figure as the solid and dashed curves. The other Norberg et al. Schechter function parameters are recovered to within the 1σ errors in both cases, although the fit residuals are significantly higher in the case of the integral formulation, leading to a higher value of the reduced $\chi^2 = 33.3$, compared to $\chi^2 = 12.5$ for the discrete formulation. In view of the much lower reduced χ^2 of the discrete density-correction method we have chosen to utilize this approach in correcting our HIMF.

We are now in a position to determine the HIMF from our galaxy detections. The histogram of galaxy detections as function of H I mass is shown in the bottom panel of Fig. 13, while the corresponding mass functions are shown in the top panel of the figure. The mass function was calculated by accumulating each galaxy detection divided by the maximum volume out to which that object would still have satisfied our detection criterion of an 8σ integrated flux density, normalized as usual to a binwidth of one dex in H I mass. We have explicitly taken account of the variation of survey sensitivity with recession velocity (as illustrated in Fig. 1) in the calculation of the limiting survey volume for each detected galaxy. This corresponds to the classical $(1/V_{\text{tot}})$ method developed by Schmidt (1968).

Open squares are used in the figure to indicate the $(1/V_{\text{tot}})$ datapoints, while the filled circles have also been corrected for the variation of density with distance in our survey region (as plotted in Fig. 11). Based on our comparison of the “discrete” and “integral” formulations of density-correction discussed above, we chose to apply the method of discrete density correction, in which $V'_{\text{eff}} \propto \rho(D) \int_{D_{\text{min}}}^{D_{\text{max}}} D^2 dD$. We have also considered density correction of the HIMF with the integral formulation of V_{eff} and find similar results, but with a reduced χ^2 value of the best-fitting Schechter function almost three times as large ($\chi^2_{\text{min}} = 15$ compared to 6.3).

We have included all 155 of our blind H I detections in the mass function, although in those cases where our flux density measurement of the primary optical ID had a large uncertainty, either due to possible confusion by nearby companions (which applies to 70 of our detections as noted in the last column of Table 1) or to a recession velocity near $V_{\text{Hel}} \sim 2800 \text{ km s}^{-1}$ (the transition in our velocity coverage from the lower to the upper band of 20 MHz width) we have used the LEDA flux values to calculate the H I mass rather than our own. In addition, we have considered all cataloged galaxies that could contribute to source confusion within our survey beam in the vicinity of our 70 confused detections. Those confusing galaxies that had tabulated LEDA fluxes sufficient to satisfy our 8σ limit on integrated H I flux, were also accumulated in the HIMF. This consideration led to an additional 14 nearby companion galaxies being incorporated into the mass function. The total number of galaxies contributing to our HIMF is 169.

The error bars of the datapoints in Fig. 13 are determined solely by the square-root of the H I mass-bin occupancy, which leads to errors that vary between about 20 and 100%. These

Fig. 14. Contours of χ^2 for pairs of our fit parameters are shown both before (top row) and after (bottom row) correction for galaxy density variations with distance. Contours are drawn at $\Delta\chi^2 = 1, 4$ and 9 corresponding to $1, 2$ and 3σ for one degree of freedom. The third parameter is kept fixed at the best-fitting value when plotting each pair.

errors dominate the error budget since they exceed the random errors associated with the density-correction by a factor of between 2 and 5.

The dotted line in Fig. 13 represents the approximate inverse search volume of our survey as function of mass, where we have assumed a relationship between H I mass and linewidth of the form: $W_{20} = 0.16 M_{\text{HI}}^{1/3} \text{ km s}^{-1}$, for M_{HI} in solar units. The best-fitting Schechter functions of the form:

$$\Theta(M) = \Theta_* \ln(10) (M_{\text{HI}}/M_*)^{\alpha+1} \exp(-M_{\text{HI}}/M_*),$$

are over-laid on the data points. The straightforward ($1/V_{\text{tot}}$) points are best-fit with $\log(M_*) = 10.15 \pm 0.1$, $\Theta_* = 9.5 \pm 3 \times 10^{-4}$ and $\alpha = -1.5 \pm 0.1$, indicated by the dashed line in the figure. The best-fit values after correction for galaxy density as given by the upper panel of Fig. 11 are $\log(M_*) = 9.85 \pm 0.07$, $\Theta_* = 55 \pm 15 \times 10^{-4}$ and $\alpha = -1.28 \pm 0.1$, indicated by the solid line in the figure. Contours of χ^2 for pairs of our fit parameters are shown in Fig. 14 both before and after correction for galaxy density variations with distance. The contours are drawn at $\Delta\chi^2 = 1, 4$ and 9 corresponding to $1, 2$ and 3σ for one degree of freedom. In each plot, the third parameter is kept fixed at the best-fitting value. From the χ^2 contours it is clear that the solutions for $\log(M_*)$ and α are well separated, while combinations involving the galaxy density become somewhat degenerate.

For comparison the HIMF derived by Zwaan et al. (2003) from the HIPASS Bright Galaxy Catalog, based on 1000 H I selected galaxies in the 2π steradians below $\delta = 0$, has $\log(M_*) = 9.79 \pm 0.06$, $\Theta_* = 86 \pm 21 \times 10^{-4}$ and $\alpha = -1.30 \pm 0.08$. Although completely at odds with our ($1/V_{\text{tot}}$) values, good

agreement is apparent between these values and our own after application of the galaxy density correction. Similar considerations apply to HIMF derived by Rosenberg & Schneider (2002) from the Arecibo Dual-Beam Survey (ADBS) who find $\log(M_*) = 9.88$, $\Theta_* = 58 \times 10^{-4}$ and $\alpha = -1.53$, although with somewhat poorer agreement in the faint end slope.

Although our choice of a minimum significance of 8σ in integrated H I flux is expected to result in a high degree of completeness in our sample (cf. Rosenberg & Schneider 2002), this can also be tested by evaluating the average value of V/V_{max} (Schmidt 1968). For a complete sample of a uniform density volume we expect $\langle V/V_{\text{max}} \rangle = 0.5$. In the absence of density corrections we actually find $\langle V/V_{\text{max}} \rangle = 0.35$ for our sample, while after discrete density correction we find $\langle V/V'_{\text{eff}} \rangle = 0.57$, where D_{max} is based on the local 8σ limit (which varies with distance as in Fig. 1) over a linewidth of $1.5 \times W_{20}$. If instead we define D_{max} by the local 8σ limit over a linewidth of $1.2 \times W_{20}$ we obtain $\langle V/V_{\text{max}} \rangle = 0.31$ and $\langle V/V'_{\text{eff}} \rangle = 0.50$. This may suggest that we have been somewhat too conservative in assessing significance based on the larger velocity intervals. However, a roll-off in completeness at the lowest significance levels can also result in an elevated expectation value of $\langle V/V_{\text{eff}} \rangle$ (Rosenberg & Schneider 2002). Our value of $\langle V/V'_{\text{eff}} \rangle = 0.57$ is comparable to that found by Rosenberg & Schneider (2002) for the ADBS $\langle V/V_{\text{eff}} \rangle = 0.60$.

4. Summary and discussion

Our unbiased H I survey of 1800 deg^2 in the northern sky has allowed recognition of a number of significant points

regarding the H I content, distribution and environment of nearby galaxies. From the analysis of some 500 candidate detections we have extracted a moderately complete sample of 155 galaxies (listed in Table 1 and illustrated in Fig. 3) with an integrated H I flux in excess of 8σ at distances between 5 and 80 Mpc. Seven of the detections occur so near the boundary of our two segments of velocity coverage (near $V_{\text{Hel}} = 2800 \text{ km s}^{-1}$), and two so near the edges of our spatial coverage, that their derived parameters are unreliable (although the detections themselves are secure). This leaves 146 detections with derived parameters of high quality. A plausible optical galaxy ID was found within a 30 arcmin search radius for all but one of the 8σ detections, although one object was previously uncataloged and three others had no previous red-shift determination. Twenty-three objects (or their uncataloged companions) are detected in H I for the first time.

We have characterized the environment of each detected galaxy by performing a search within NED for all cataloged objects within a radius of 30 arcmin (corresponding to a possible contribution within our telescope beam). These (potential) companions have been tabulated in three categories in Table 1, namely; (a) confused for objects within 400 km s^{-1} of the primary ID, (b) unconfused for objects offset by 400 to 1000 km s^{-1} from the primary ID, and (c) possibly confused for objects of unknown red-shift. It will remain difficult to assess the actual likelihood of association for objects in this last category, until red-shift determinations become available. For the moment we will regard only those objects with entries in category (a) as “confused” and all others as “unconfused”.

We determine agreement of our absolute flux scale to the weighted average of all previous determinations of the H I flux (as tabulated by LEDA) to better than 1% for our unconfused detections, as shown in Fig. 7. Confused objects show a systematic excess H I flux in our large survey beam.

4.1. Centroid offsets of gas and stars

Since our survey was not targeted at known galaxies, we have an independent determination of the position centroid for each detected object. The majority of apparent offsets between the gaseous and stellar distributions are consistent with the substantial uncertainties that follow from a large survey beam and only moderate signal-to-noise, as shown in Fig. 5. However, a number of significant centroid offsets (greater than 5σ) are detected in nominally unconfused galaxies which are indicated in Table 1 by entering the symbol “o”, “+” or “++” in the Note column. These have been divided somewhat arbitrarily into two categories, depending on whether the linear centroid offset is less than 10 kpc (category “o”) or greater than 10 kpc (categories “+” and “++”). The reasoning behind this division is that a 10 kpc limit may mark a plausible distinction between internal asymmetries of individual objects and the larger scales that are more likely to indicate external gaseous components.

Determining the cause of these significant centroid offsets requires higher resolution imaging. One extensive source of high resolution imaging is the WHISP survey (Kamphuis et al. 1996, <http://www.astro.rug.nl/~whisp>) which has

targeted some 200 UGC galaxies north of $\text{Dec} = 20^\circ$ with synthesis observations using the WSRT array. WHISP observations are currently available for only a small fraction of the 155 galaxies in our 8σ sample. The WHISP results for each galaxy are summarized in a web-accessible data overview consisting of a series of images of the integrated H I distribution and accompanying velocity field at three different angular resolutions, of about 15, 30 and 60 arcsec, together with a global H I profile, a major axis position-velocity plot and an optical reference image. The most relevant component of this overview for our purposes is the distribution of integrated H I at the lowest angular resolution of 60 arcsec, where the highest surface brightness sensitivity is reached. For the four of 11 instances of significant centroid offset noted in Table 1 that have already been imaged in the WHISP survey we comment briefly on what is seen in the 60 arcsec integrated H I image:

- NGC 7640 o: asymmetric with extensions;
- UGC 12732 o: possible companions;
- UGC 731 o: asymmetric;
- NGC 925 o: asymmetric with extensions.

All four cases that have been imaged with high resolution show large-scale asymmetries or possible uncataloged companions. Our tentative conclusion is that our measured centroid offsets are indeed indications of substantial asymmetries and the presence of possible uncataloged gas-rich companions in the immediate vicinity of the primary ID.

4.2. Uncataloged companions

The issue of uncataloged gas-rich companions is also addressed by the comparison of our survey flux densities with those measured previously for nominally unconfused galaxies in Figs. 8 and 9. Although the comparison with the heterogeneous LEDA data has a large degree of scatter, the comparison with the Arecibo data may indicate an excess of H I at large radii.

The Arecibo data were taken from the Pisces-Perseus supercluster survey (Giovanelli & Haynes 1985; Giovanelli et al. 1986; Giovanelli & Haynes 1989; Wegner et al. 1993; Giovanelli & Haynes 1993). We plot ratios of both the *observed* and the *corrected* H I flux density. In the latter case approximate corrections for telescope pointing errors and a model of the galaxy extent were applied to the Arecibo data. These data were only available for 20 of our unconfused detections, which vary in distance from about 7 to 70 Mpc. The Arecibo beam (3.3 arcmin FWHM at the time of those observations) has a linear dimension that varies from about 7 to 70 kpc over the distance range above, while our survey beam varies from about 100 kpc to 1 Mpc in diameter. Four of the unconfused galaxies have *observed* flux ratios, R , of unity within our 1σ errors, including the nearest object. All of the other galaxies have an excess detected H I flux in our larger survey beam which varies from about 10% to 300%. The comparison of *corrected* flux ratios is more ambiguous, since at least five of the data-points now have apparent flux ratios which are significantly less than unity. As indicated in Sect. 3 above, each of our survey detections was examined for evidence of source resolution effects

and our absolute flux scale is well-defined, making such apparent deficits in our detected H I flux difficult to understand. Given the approximate nature of the corrections applied to the observed Arecibo fluxes it seems likely that they may (on occasion) result in a degree of over-compensation for missed flux.

The straightforward conclusion that can be drawn from the *observed* flux ratio plots, it that in most cases, the H I distribution must be significantly more spatially extended than the Arecibo beam, even when this beam subtends 50–70 kpc.

High resolution imaging will be necessary to determine, on a case-by-case basis, where the excess detected flux actually resides. WHISP data (<http://www.astro.rug.nl/whisp>) are currently available for 7 of the 20 unconfused galaxies that have Arecibo flux measurements. Inspection of the 60 arcsec integrated H I image in the WHISP database yields the following assessment:

- NGC 7286: $R = 1.79$, possible extensions;
- UGC 12693: $R = 1.33$, possible extensions;
- UGC 12713: $R = 1.05$, asymmetric;
- UGC 1856: $R = 1.15$, nothing unusual;
- NGC 972: $R = 1.38$, asymmetric, extensions, possible companions;
- NGC 1012: $R = 1.21$, 10% of H I flux in uncataloged companion;
- NGC 1056: $R = 1.72$, extensions.

The two cases showing the largest excess H I flux in the WSRT survey beam are notable for having possible extensions at low surface brightness in the distribution of integrated H I in the WHISP data. Only in one of these seven cases, has a distinct uncataloged companion been detected in the WHISP imaging. Deeper imaging of a larger field-of-view will be needed to search for additional uncataloged companions to account for the excess detected H I flux.

An intriguing possibility for the location of the excess detected H I flux is that it resides in a relatively diffuse distribution subtending a few 100 kpc in the vicinity of the primary target. This is exactly the type of hypothetical distribution, in the environment of M 31, which motivated the negative velocity component of our wide-field H I survey. Such distributions were found (De Heij et al. 2002) to provide the best-fit to the spatial and kinematic distribution of the compact high-velocity cloud population in the vicinity of the Galaxy. The best-fitting models of this type consist of gas bound to low-mass dark-matter halos with a steep power-law ($\alpha = -1.7$) distribution in number as function of neutral gas mass and are concentrated around their major galaxy host in a Gaussian distribution with a spatial dispersion of 150–200 kpc. The total H I mass predicted in these Local Group models to survive (ram-pressure- and tidal-stripping) to the present day amounts to some $1.2 \times 10^9 M_\odot$. Compared to the $8 \times 10^9 M_\odot$ of H I in M 31 and the Galaxy, this corresponds to an excess H I mass of about 15% distributed on scales of a few hundred kpc. Only a handful of the rare, massive components might be identifiable as discrete objects, while the rest of the distribution would merely contribute to a diffuse enhancement of H I mass, centered on the systemic velocity of the host.

This possibility can and should be tested with a dedicated experiment.

4.3. Spatial variance of the HIMF

In Sect. 3 we derive an HIMF from our 8σ sample of background galaxies. Despite the fact that our survey region covered some 1800 deg^2 and therefore sampled a variety of environments along the line-of-sight, it is impossible to overcome the fact that in a flux-density limited sample, all low mass detections must of necessity be very nearby. In consequence, the power-law slope and normalization of the HIMF are strongly dependent on whether or not there are significant variations of the average galaxy density with distance. We have quantified such a variation by utilizing cataloged optical galaxies in our survey volume as shown in Figs. 10 and 11. The absolute density of optical galaxies as function of distance was determined by fitting for the normalization of the “standard” luminosity function determined by Norberg et al. (2002) in a sequence of heavily over-lapping sub-samples. Density variations with distance of more than an order of magnitude are derived within our survey volume. Accounting for this variation leads to very substantial changes in the best-fitting HIMF parameters, as seen in Fig. 13. Before correction we obtain $\log(M_*) = 10.15 \pm 0.1$, $\Theta_* = 9.5 \pm 3 \times 10^{-4}$ and $\alpha = -1.5 \pm 0.1$, and after correction $\log(M_*) = 9.85 \pm 0.07$, $\Theta_* = 55 \pm 15 \times 10^{-4}$ and $\alpha = -1.28 \pm 0.1$, where the error estimates come from the 2σ contours of $\Delta\chi^2$ in the fit parameters shown in Fig. 14.

Only after applying the density correction is statistical agreement realized with the HIPASS BGC values (Zwaan et al. 2003) $\log(M_*) = 9.79 \pm 0.06$, $\Theta_* = 86 \pm 21 \times 10^{-4}$ and $\alpha = -1.30 \pm 0.08$ and to a lesser extent with the Arecibo Dual-Beam Survey (ADBS) values (Rosenberg & Schneider 2002) $\log(M_*) = 9.88$, $\Theta_* = 58 \times 10^{-4}$ and $\alpha = -1.53$.

Schneider et al. (1998) and Rosenberg & Schneider (2002) also consider the variation of number density of optical galaxies in the ADBS survey region. However they determine only a relative, rather than an absolute density and do so on the basis of galaxy number counts rather than explicit fitting to the complete portion of the luminosity function. They conclude that density corrections have only a minor impact on the form and normalization of the HIMF in their sample, although their normalization does increase from 48 to $58 \times 10^{-4} (\text{Mpc}^{-3} \text{ dex}^{-1})$, when relative density corrections are applied. These authors also make use of the “integral” formulation of density correction discussed in Sect. 3 rather than the “discrete” formulation which we find leads to substantially reduced fit residuals.

The absolute normalization of the HIPASS BGC HIMF is not a trivial procedure (Zwaan et al. 2003). Since a maximum likelihood method has been employed, the shape of the HIMF should be well-determined, but only in terms of a relative density. The normalization is determined after the fact by carrying out integrations of the derived selection function. Since a reasonable range of detected masses in the HIPASS sample is only achieved inside of about 25 Mpc, the normalization is, of necessity, also tied to this distance range. To assess the impact of a possible variation of galaxy density with distance within the

HIPASS sample, we have also fit for the optical galaxy normalization using all LEDA galaxies at $\delta < 0$, just as in Sect. 3, for our own survey volume. Given the larger optical galaxy sample size, we defined 25 overlapping sub-samples with sample populations varying linearly from a minimum of 100 galaxies at the nearest distances to 1500 galaxies at the maximum distance. This distribution is plotted as the dotted line in Fig. 11. The southern hemisphere has a galaxy density which is equal to the Norberg et al. (2002) value between about 8 and 25 Mpc. At larger distances there appears to be a smooth decline to about 30% of this density by 80 Mpc. At distances smaller than about 8 Mpc there also appears to be a decline in galaxy density to about 50% of nominal. It seems that over the critical distance range of 10–25 Mpc, the galaxy density within the HIPASS survey volume is essentially the nominal one, suggesting that the HIPASS BGC value of $\Theta_* = 86 \pm 21 \times 10^{-4} \text{ Mpc}^{-3} \text{ dex}^{-1}$ should be quite reliable. The apparent down-turn in galaxy density below about 8 Mpc may have some consequence for the apparent shape of the HIPASS BGC HIMF below $\log(M) = 7.5$, since such systems could only be detected out to 8 Mpc in the HIPASS data.

For comparison we have also determined and plotted the optical galaxy density in the same way for the entire northern hemisphere in Fig. 11. The northern hemisphere distribution is quite different than the southern, with a moderate over-density (about 50%) inside of 15 Mpc, followed by a relative dearth of galaxies between 20 and 60 Mpc. It seems quite conceivable that the apparent discrepancies between the HIPASS BGC and the ADBS HIMF parameters may be a consequence of such large-scale differences in the galaxy distribution.

Acknowledgements. We acknowledge useful discussions and feedback on methods with L. Staveley-Smith and M. Zwaan. The Westerbork Synthesis Radio Telescope is operated by The Netherlands Foundation for Research in Astronomy under contract with the Netherlands Organization for Scientific Research. We have made use of the LEDA database (<http://leda.univ-lyon1.fr>). This research has made use of the NASA/IPAC Extragalactic Database (NED) which is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration. The Digitized Sky Surveys were produced at the Space Telescope Science Institute under U.S. Government grant NAG W-2166. The images of these surveys are based on photographic data obtained using the Oschin Schmidt Telescope on Palomar Mountain and the UK Schmidt Telescope. The plates were processed into the present compressed digital form with the permission of these institutions.

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Online Material

Fig. 3. H I spectra of detected galaxies. The complete set of panels for all detected galaxies is only available in the electronic version of the paper at <http://www.epdsciences.org>. The catalog designation is indicated in the upper left of each panel, while the likely optical ID is indicated in the upper right.

Fig. 3. continued.

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