Highlights from the first year of Odin observations

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\textsuperscript{\dagger} Deceased.
Abstract. Key Odin operational and instrumental features and highlights from our sub-millimetre and millimetre wave observations of H$_2$O, H$_2$O, NH$_3$, $^{15}$NH$_3$ and O$_2$ are presented, with some insights into accompanying Odin Letters in this A&A issue. We focus on new results where Odin’s high angular resolution, high frequency resolution, large spectrometer bandwidths, high sensitivity or/and frequency tuning capability are crucial: H$_2$O mapping of the Orion KL, W 3, DR 21, S 140 regions, and four comets; H$_2$O observations of Galactic Centre sources, of shock enhanced H$_2$O towards the SNR IC 443, and of the candidate infall source IRAS 16293-2422; H$_2$O detections in Orion KL and in comet Ikeya-Zhang; sub-mm detections of NH$_3$ in Orion KL (outflow, ambient cloud and bar) and O Ph, and very recently, of $^{15}$NH$_3$ in Orion KL. Simultaneous sensitive searches for the 119 GHz line of O$_2$ have resulted in very low abundance limits, which are difficult to accommodate in chemical models. We also demonstrate, by means of a quantitative comparison of Orion KL H$_2$O results, that the Odin and SWAS observational data sets are very consistently calibrated.

Key words. radio lines: ISM – ISM: molecules – space vehicles: instruments – submillimeter – techniques: spectroscopic – telescopes

1. Introduction

With the recognition that ion-molecule reactions play a major role in interstellar chemistry (Herbst & Klemperer 1973; Watson 1973; Dalgarno & Black 1976) it was proposed that H$_2$O and O$_2$, together with CO, O and C, must be important coolants of interstellar clouds, governing the cloud evolution towards formation of protostellar cores (Goldsmith & Langer 1978). This was still the current understanding at the time of the Odin (and SWAS, see below) feasibility and Phase A studies more than ten years ago (cf. Hjalmarson 1997). However, during the extended instrument planning phase, compelling evidence appeared from ground-based searches for O$^{18}$O (Maréchal et al. 1997, and references therein) that the O$_2$ abundance must be considerably lower than theoretically expected. To address this O$_2$ search sensitivity problem, a low-noise HEMT preamplifier at 119 GHz was added to the already rather complicated, cryo-cooled sub-millimetre radiometer system (Frisk et al. 2003, in this issue). The Odin O$_2$ search sensitivity in this way was improved by more than an order of magnitude.

NASA’s Submillimetre Wave Astronomy Satellite (SWAS) was launched in December 1998, and since then has performed extensive, simultaneous observations of the ortho-H$_2$O (1$_{10}$$-1_{01}$), $^{13}$CO (5 – 4) and CI (P$_{1}$ – P$_{0}$) lines, using an off-axis Cassegrain telescope of size 54 × 68 cm (Melnick et al. 2000a, and subsequent papers in that Apl issue). The observed H$_2$O abundances in quiescent dense interstellar clouds are estimated to be at least an order of magnitude lower than expected from pure gas-phase chemistry, most likely because of efficient molecular adsorption onto dust grain surfaces (e.g., Bergin et al. 2000). A large interstellar water abundance in terms of icy grain mantles has indeed been discovered by ESA’s Infrared Space Observatory (ISO, cf. van Dishoeck & Blake 1998; Ehrenfreud & Charnley 2000). Simultaneous deep SWAS searches for the O$_2$ (3$_3$ – 1$_2$) transition at 487 GHz have led to low abundance limits in a number of sources (Goldsmith et al. 2000), and recently also a very tentative O$_2$ detection in O Ph, presumably in an outflow wing (Goldsmith et al. 2002). The low O$_2$ abundance limits pose serious problems to chemistry models (e.g., Bergin et al. 2000).

On 20 February 2001, the Odin sub-millimetre wave spectroscopy satellite for astronomy and aeronomy research, equipped with an offset Gregorian telescope of diameter 110 cm, was launched by a Start-1 rocket from Svobodny in far-eastern Russia (Nordh et al. 2003, in this issue; von Schéele 2002). In this Letter we report some highlights from the first year of Odin observations and set the scene for a number of accompanying A&A Letters. We will here refrain from referring to the wealth of relevant recent papers on interstellar chemistry and cloud cooling and instead refer the reader to our accompanying Odin A&A Letters, and to Bergin et al. (2000), and other SWAS papers in that ApJ issue.

In Sect. 2 we summarise the Odin operation and data processing, available receiver and spectral line combinations, observation modes, and pointing and antenna characteristics. Section 3 contains a presentation and discussion of some early Odin highlights, with references to accompanying Odin A&A Letters and to some planned, subsequent papers.

2. Operational and observational characteristics

2.1. Operation, scheduling, and data processing

The satellite operation is described in a recent paper by Lundin (2002): “Finding the balance between autonomy on-board versus man-triggered actions from ground”. The scheduling time-lines of astronomical targets/receiver combinations – based on priorities agreed upon in the Odin astronomy team – are created by the Astronomy Mission Scientist, Å. Hjalmarson, and Aa. Sandqvist (Stockholm Observatory), using scheduling software developed by K. Volk (University of Calgary) and detailed command software developed and implemented by H.-G. Florén (Stockholm Observatory). The pipeline millimetre and sub-millimetre wave data processing and calibration are performed at the Odin Data Centre at Onsala Space Observatory. These processes as well as the data flow and storage, are described by Olberg et al. (2003), in this issue.

2.2. Receiver combinations, important molecular lines, and observation modes

Table 1 summarises the salient features of what is presented in more detail by Frisk et al. (2003) and Olberg et al. (2003), both in this issue. Odin accommodates two pairs of frequency tuneable, SSB (single sideband) tuned, sub-millimetre wave Schottky mixer receivers, together with a fixed-tuned HEMT receiver at 119 GHz, also operated in its SSB mode. Any four, three, or two (depending upon the satellite power available) of these front-end receivers can be combined with any of our three spectrometers: a broad band (1.1 GHz) acousto-optical spectrometer (AOS), and two flexible (bandwidth: 100–800 MHz;...
Table 1. Salient features of Odin.

<table>
<thead>
<tr>
<th>Feature</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Antenna size</td>
<td>110 cm</td>
</tr>
<tr>
<td>Beam size at 119/550 GHz</td>
<td>9.5°/2.1° (126°)</td>
</tr>
<tr>
<td>Main beam efficiency</td>
<td>90%</td>
</tr>
<tr>
<td>Pointing uncertainty</td>
<td>&lt;10&quot; (RMS)</td>
</tr>
<tr>
<td>Submm tuning range</td>
<td>486–504, 541–581 GHz</td>
</tr>
<tr>
<td>Submm noise temperature</td>
<td>3300 K (SSB)</td>
</tr>
<tr>
<td>HEMT centre frequency</td>
<td>118.75 GHz</td>
</tr>
<tr>
<td>2.5 mm noise temperature</td>
<td>600 K (SSB)</td>
</tr>
<tr>
<td>AOS bandwidth/resolution</td>
<td>1050 MHz / 1 MHz</td>
</tr>
<tr>
<td>AC1/2 bandwidth/resolution</td>
<td>100–800 MHz / 0.25–2 MHz</td>
</tr>
</tbody>
</table>

Table 2. Some low energy spectral lines observable by Odin.

<table>
<thead>
<tr>
<th>Species</th>
<th>Transition</th>
<th>Frequency (MHz)</th>
<th>( \Delta \nu ) ( \times 10^{3} )</th>
<th>( n_{	ext{crit}} ) ( \times 10^{6} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>H(_{2})O</td>
<td>1(<em>{0} - 1(</em>{1}))</td>
<td>556 936.00 27</td>
<td>3\times10^6 ( ^c )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>H(_{2}^{18})O</td>
<td>1(<em>{0} - 1(</em>{1}))</td>
<td>547 676.44 27</td>
<td>3\times10^6 ( ^c )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>H(_{2}^{17})O</td>
<td>1(<em>{0} - 1(</em>{1}))</td>
<td>552 021.02 26</td>
<td>3\times10^6 ( ^c )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>NH(_{3})</td>
<td>0(<em>{1} - 0(</em>{0}))</td>
<td>572 498.07 29</td>
<td>4\times10^7 ( ^c )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>(^{13})NH(_{3})</td>
<td>0(<em>{1} - 0(</em>{0}))</td>
<td>572 112.79 29</td>
<td>6\times10^7 ( ^c )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>O(_{2})</td>
<td>3(<em>{3} - 2(</em>{2}))</td>
<td>487 249.38 26</td>
<td>10(^6) ( ^c )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>O(_{2})</td>
<td>1(<em>{1} - 1(</em>{0}))</td>
<td>118 750.34 6</td>
<td>6( ^b )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>CI</td>
<td>(^{3})P(<em>{1} - ^{3})P(</em>{0})</td>
<td>492 160.70 24</td>
<td>10(^6) ( ^c )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>CO</td>
<td>5 – 4</td>
<td>576 267.93 83</td>
<td>2\times10^5 ( ^c )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>(^{13})CO</td>
<td>5 – 4</td>
<td>550 926.30 79</td>
<td>2\times10^5 ( ^c )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>C(_{2})O</td>
<td>5 – 4</td>
<td>548 830.98 79</td>
<td>2\times10^5 ( ^c )</td>
<td>( ^c )</td>
</tr>
<tr>
<td>CS</td>
<td>10 – 9</td>
<td>489 751.05 129</td>
<td>5\times10^7 ( ^c )</td>
<td>( ^c )</td>
</tr>
</tbody>
</table>

\( ^a \) Upper state energy.
\( ^b \) Cloud density needed for excitation.
\( ^c \) Will be considerably reduced because of radiative trapping.

channel spacing: 0.125–1 MHz) hybrid auto-correlators (AC1 and AC2).

In Table 2 we have collected some relevant parameters for important molecules, the lower energy lines of which appear in the Odin receiver bands.

The tuning capability and easy selection of receiver combinations has allowed us to perform simultaneous observations of H\(_{2}\)O (at 557 GHz), or H\(_{2}^{18}\)O (at 548 GHz) with two different mixers – i.e., to observe the same line with two different velocity resolutions and to increase the sensitivity – together with a deep search for O\(_{2}\) at 119 GHz. In other observations we have selected H\(_{2}\)O (or H\(_{2}^{18}\)O), and NH\(_{3}\) (at 572 GHz), together with O\(_{2}\) at 119 GHz.

The Odin receiver set-up (Frisk et al. 2003), together with the design of the Attitude Control System (ACS, Jacobsson et al. 2002a,b) allows two complementary observation schemes: sky-switching and position-switching. By sky-switching we mean Dicke-switching via a mirror alternating between the telescope signal and one of two (selectable) sky beams of size 4.4′′, separated from each other by 28° and offset by 42° from the main telescope beam, always allowing a reference measurement toward cold, signal-free sky. Sky-switching allows the aforementioned full flexibility of receiver combinations. However, the intrinsic asymmetry of this switching mode leads to a power level offset and a stable low-level ripple and standing wave pattern, which must be removed by subtraction of a reference off-source spectrum. The necessity of spending time to get such an off-source spectrum means that this observation mode is most efficient in mapping programs.

The alternative Odin switching method is true position-switching, where the entire satellite, as rapidly as possible, is used to nod the telescope beam between the target position and a selected signal-free off position. Position-switching was considered as an important option during the design of the Odin ACS (Jacobsson et al. 2002a,b). Therefore the satellite light is sufficiently rapid that the in-orbit measured on-source observing efficiency is as high as 40% for off-positions closer than 1° from the targets. The position-switching mode has been used in observations where very broad lines are expected, e.g., in observations of H\(_{2}\)O towards Galactic Centre targets and in extragalactic H\(_{2}\)O searches. The baseline structure achieved is very smooth across a 1.1 GHz band and the remaining ripple and standing wave pattern is very low, often invisible, in this symmetric switching, even after very long integrations (cf. Sandqvist et al. 2003, in this issue). Therefore the position-switching mode has also been used in our first searches for new spectral lines/molecules – simultaneous searches in the Orion KL region for PH (at 553.4 GHz) and for \(^{13}\)NH\(_{3}\) (at 572.1 GHz). The only drawback in using position-switching is a loss of flexibility in our choice of receiver combinations, as we can then simultaneously use only one pair of sub-mm mixers, or the other pair of sub-mm mixers together with the 119 GHz receiver (cf. Table 1 of Olberg et al. 2003).

2.3. Antenna characteristics and pointing capability

Odin is equipped with an offset Gregorian telescope of diameter 110 cm. The RMS surface accuracy of the main mirror is <8 μm, i.e. <\( \lambda/70 \) at 557 GHz, and the system is (for aeronomy reasons) designed to have a side-lobe level well below ~25 dB. From Jupiter mapping at 557 GHz we estimate a FWHM beam size of 126′′ (2.1′) and a main beam efficiency of about 87 ± 6%. The corresponding theoretically calculated values are 125′′ and 88% at 557 GHz, and 9.5′ and 92% at 119 GHz, respectively (for further details, see Frisk et al. 2003).

The Odin attitude control system and attitude reconstruction capability are described in two recent papers by Jacobsson et al.: “The high-performing attitude control system of the scientific satellite Odin” (2002a), and “Star tracker/gyro calibration and attitude reconstruction for the scientific satellite Odin – in flight results” (2002b). The astronomical targets used to observationally delineate the pointing stability are Jupiter, the strong H\(_{2}\)O emission from several comets (Leach eux et al. 2003, in this issue), and the equally strong Orion KL H\(_{2}\)O outflow emission (Olofsson et al. 2003, in this issue). Frisk et al. (2003) show that Odin’s long-term pointing repeatability and the short-term attitude reconstruction accuracy are very good, and at least within <10′′ (RMS). Up to 05 October 2002, all Odin observations were performed with a known antenna beam.
offset of about 30′′ (counted from the expected beam location in the star-tracker reference coordinate system). For Odin’s beam size of 126′′ such a pointing offset would cause an intensity decrease of at most 15% (for a point source). Our correction attempt, implemented on 05 October 2002, has reduced the current pointing offset to be below 10′′.

3. Results and discussion

We here present some highlights from the first year of Odin observations, where Odin’s higher angular resolution, higher frequency resolution, larger spectrometer bandwidths, higher sensitivity or/and frequency tuning capability (compared to SWAS, cf. Hjalmarson 2003) are crucial.

3.1. Selected Odin H\(_2\)O and H\(_{18}\)O highlights

Figure 1 shows a number of Odin ortho-H\(_2\)O (1\(_{10}^-1\_01\)) spectra at 557 GHz – observed in the sky-switching mode and using the broad band (BW = 1.1 GHz), lower resolution (1 MHz) AOS spectrometer. We have in the different panels also indicated simultaneously used auto-correlation spectrometer bandwidths (allowing higher velocity resolution). The observational results are discussed below.

3.1.1. Orion KL mapping

Figure 1a displays the high velocity H\(_2\)O outflow/shock emission from the Orion KL molecular cloud core, where the water abundance appears to be as high as 5 \(\times\) 10\(^{-4}\) and essentially all oxygen is locked-up in gas phase H\(_2\)O (ISO: Harwit et al. 1998; Wright et al. 2000; SWAS: Melnick et al. 2000b; Odin mapping: Olofsson et al. 2003, in this issue). Figure 1b shows the narrow H\(_2\)O line observed 2′ south of the outflow centre, where the water abundance is estimated to be several orders of magnitude lower. In an accompanying A&A Letter, Olofsson et al. present the first results from our Odin H\(_2\)O mapping at 1′ spacing of the Orion KL molecular cloud core, together with Odin’s detections of H\(_{18}\)O at 548 GHz at the cloud centre (outflow and quiescent gas) and from quiescent gas 2′ south of the outflow. The mapping was performed simultaneously in two mixers, one connected to the broad band AOS and the other to a higher velocity resolution auto-correlation spectrometer. Our Orion KL H\(_2\)O and H\(_{18}\)O outflow/shock emission spectra will be discussed further in Sect. 3.5.

3.1.2. Mapping of H\(_2\)O in W3, DR21, and S140

Odin has observed H\(_2\)O (1\(_{10}^-1\_01\)) emission in a considerable number of interstellar clouds and is mapping the emission in some selected regions. Our mapping of the W3 molecular cloud at 1′ spacing is presented in this A&A issue by Wilson et al. (2003), and an illuminating Odin/SWAS comparison of the H\(_2\)O emission towards W 3 (IRS 5) is provided in Sect. 3.6. Our DR 21 H\(_2\)O mapping (of a 5′ × 10′ area, using 1′ spacing), together with the results of a recent simultaneous search for H\(_{18}\)O and NH\(_3\) (1\(_0 \rightarrow 0_0\)), will be presented in a future paper. Another highlight may well be the detection of very narrow (4 km s\(^{-1}\)) and twice as strong H\(_2\)O emission towards the S 140 bright rim (PDR), and also at a somewhat different velocity (all compared to the SWAS spectrum, of width 5.7 km s\(^{-1}\), observed at the cloud core 1′ away; Snell et al. 2000). The result of the Odin S 140 strip maps at 1′ spacing across this favourably located PDR will appear in a future paper.

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\(^{1}\) The intensity scale in all figures is not corrected for a main beam efficiency of 0.9.
3.1.3. $\text{H}_2\text{O}$ and $\text{H}^{18}\text{O}$ observations of comets

Figure 1c displays the strong, very narrow (in the broad band AOS) $\text{H}_2\text{O}$ line at 557 GHz observed from comet C/2001 A2 (LINEAR). Odin has observed and mapped the water production in several comets at high spectral resolution (80 m s$^{-1}$), as discussed by Lecacheux et al. (2003, in this issue). In the most recent bright comet, 153P/2002 C1 (Ikeya-Zhang), weak $\text{H}^{18}\text{O}$ (1$_{0}-1_{0})$ emission was also detected (at an intensity level about 90 times lower than that of $\text{H}_2\text{O}$). The $\text{H}_2\text{O}/\text{H}^{18}\text{O}$ integrated intensity ratio is consistent with an isotope ratio $^{18}\text{O}/\text{H}_2\text{O} = 450 \pm 50$ and verifies the estimated very high $\text{H}_2\text{O}$ production rate, $26 \times 10^{28}$ molecules s$^{-1}$, where we rely on a rather elaborate excitation/radiative transfer model (see Lecacheux et al. 2003). The comet observations were performed in the Odin sky-switching mode. Two mixers were tuned to $\text{H}_2\text{O}$, or $\text{H}^{18}\text{O}$, both connected to high-resolution autocorrelation spectrometers (their selected bandwidth, 100 MHz, is indicated in Fig. 1e), and one of them also to the AOS.

3.1.4. $\text{H}_2\text{O}$ emission and absorption in Sgr A clouds

Figure 1d shows the intricate mixture of $\text{H}_2\text{O}$ emission and absorption observed by Odin in the direction of the Galactic Centre, Sgr A*. Interestingly, a very similar spectral shape has been observed in the HCO$^+$ (J = 1 - 0) line (Linke et al. 1981). This Odin result, and even more intriguing data from nearby positions, where the sub-millimetre wave continuum is stronger, is discussed in an accompanying A&A Letter by Sandqvist et al. (2003). These beautiful Sgr A results, paired with the very flat broad-band baselines resulting from Odin’s position switching observations, have encouraged us to begin deep searches for $\text{H}_2\text{O}$ emission in some nearby starburst galaxies, where shock enhanced $\text{H}_2\text{O}$ emission could result from a large number of outflows caused by extensive star formation.

3.1.5. Shock enhanced $\text{H}_2\text{O}$ towards IC 443 G

Figure 1e shows the result of the IC 443 supernova shock impinging upon the nearby molecular cloud region (e.g., White et al. 1987; Ziurys et al. 1989), in terms of $\text{H}_2\text{O}$ shock chemistry and shock excitation. The Odin $\text{H}_2\text{O}$ signal is about a factor of 2 stronger than the corresponding SWAS line, which indicates an Odin/SWAS beam filling ratio of this amount and a source size similar to the Odin 2.1’ beam (if we assume a Gaussian brightness distribution, see Sect. 3.4). An analysis of Odin’s observations of shock enhanced $\text{H}_2\text{O}$ emission will appear future paper.

3.2. First detections of the NH$_3$ (1$_{0}-0_{0}$) line

Simultaneously with observations of $\text{H}_2\text{O}$, or $\text{H}^{18}\text{O}$, and searches for $\text{O}_2$, we also tuned one of our receivers to the ortho-NH$_3$ (1$_{0}-0_{0}$) line at 572 GHz. Early results from these observations are i) the detection of quiescent cloud as well as outflow NH$_3$ (1$_{0}-0_{0}$) emission in the Orion KL cloud core, quiescent cloud emission 2’ south, and the first detection of NH$_3$ emission from the Orion Bar (Larsson et al. 2003, in this issue), and ii) the first detection of the NH$_3$ (1$_{0}-0_{0}$) line in a dark cloud (Liseau et al. 2003, in this issue). In view of these Odin NH$_3$ results, we have used the above mentioned receiver combinations more often in recent observations. Since the $\text{H}_2\text{O}$ (1$_{0}-1_{0}$) and NH$_3$ (1$_{0}-0_{0}$) transitions have similar excitation requirements (cf. Table 2), while the two molecules have rather different chemical behaviour, this approach could also provide new constraints for our interpretation of the $\text{H}_2\text{O}$ and NH$_3$ data.

3.3. Searches for new molecules / spectral scans

Very recently we have taken the first steps in the spectral scan / molecular search area, made possible by Odin’s tuneable sub-millimetre receivers. We have carried out simultaneous searches in the Orion KL region for PH (at 553.4 GHz) and for the 1$_{0}-0_{0}$ transition of $^{15}\text{NH}_3$ (at 572.1 GHz). A preliminary reduction of the data set appears to show a detection of the $^{15}\text{NH}_3$ (1$_{0}-0_{0}$) transition. This result will be important for our NH$_3$ radiative transfer modelling, and ultimately also for our $\text{H}_2\text{O}$ (1$_{0}-1_{0}$) analysis.

3.4. Deep Odin searches for molecular oxygen

Simultaneously with our sub-millimetre wave observations of $\text{H}_2\text{O}$, $\text{H}^{18}\text{O}$, and NH$_3$, we have been searching for molecular oxygen, normally using the dedicated 119 GHz $\text{O}_2$ (1$_{0}-1_{0}$) HEMT preamplifier channel, but in the case of the Galactic Centre sources also using the $\text{O}_2$ (3$_{3}-1_{3}$) transition at 487 GHz. In spite of some frustrating phase-locking problems, Odin’s sensitive 119 GHz searches have led to very low 3σ $\text{O}_2$ abundance limits, especially in cold dark clouds, as reported in this A&A issue by Pagani et al. (2003). Odin here has improved upon the SWAS upper limits (Goldsmith et al. 2000) by factors of 20 and 40 in L 134 N and TMC-1, and by factors of a few in some star forming regions. Such low $\text{O}_2$ abundance limits are very difficult to accommodate in current chemical models (cf. Pagani et al. 2003; Goldsmith et al. 2002).

3.5. $\text{H}_2\text{O}$ and $\text{H}^{18}\text{O}$ in Orion KL

The physical and chemical conditions in the Orion KL core region are discussed by Olofsson et al. (2003), who provide a number of relevant references. The $\text{H}_2\text{O}$ spectrum observed by Odin’s 2.1’ beam towards Orion KL is displayed in Fig. 2. No appreciable (narrow) ambient cloud emission, nor semi-narrow “hot core” emission, is apparent in this spectrum (for details see Olofsson et al.). The $\text{H}_2\text{O}$ emission predominantly emanates in the well-known outflow/shock region, and the emission from the “blueshifted” (with respect to the ambient cloud, “ridge”, velocity of 8–10 km s$^{-1}$) part of the flow is considerably weaker than the “redshifted” part. The abrupt $\text{H}_2\text{O}$ brightness decrease at velocities blueward of about 10 km s$^{-1}$ is not an Odin instrument artifact, since this surprising feature is verified in our independently observed AOS spectrum (Fig. 1a). Absorption by low excitation $\text{H}_2\text{O}$ foreground gas (in the molecular cloud envelope) at ambient cloud velocities can affect the line shape. In addition, self-absorption in the optically thick blueward
outflow is a possibility. In fact, this effect has been observed by Odin in expanding cometary H2O coma (Lecacheux et al. 2003). Narrow, deep self-absorption features are indeed very common in molecular cloud H2O (10−100) spectra observed by SWAS (e.g. Ashby et al. 2000) and Odin (see Sects. 3.1 and 3.7), but are not apparent in our Orion ambient cloud spectra.

In Fig. 2 we also show the corresponding SWAS water spectrum (from Melnick et al. 2000b), observed with a beam size of 3.3′×4.5′. The SWAS spectrum here has been scaled by a factor of 2.4 to match the outflow line-wings observed by Odin. The very good match of the H2O line-wings as observed by the two telescopes (after scaling the SWAS spectrum by a factor 2.4) can be used to roughly estimate the size of the outflow/shock region. If we assume that the distribution has a Gaussian shape, the Odin/SWAS beam filling ratio R can be expressed as

\[ R = 15.876 + \theta^2 \]

where \( \theta \) is the source size expressed in arcsec. We have used an Odin (Gaussian) beam area of (126 arcsec)² and a SWAS beam area of 3.3′×4.5′×3600 arcsec². For a very small source, R approaches 3.37. Since the previously used scale factor can be interpreted as the Odin/SWAS beam filling ratio, Eq. (1) leads to an H2O outflow FWHP size of circa 110′ for a Gaussian brightness distribution. A very similar size is estimated from our Odin H2O mapping (cf. Olofsson et al. 2003). Moreover, our preliminary deconvolution of the H2O “position-velocity cube” mapped by Odin indicates a rather extended lower intensity H2O outflow/shock emission region.

In addition, the larger SWAS beam appears to “pick up” ambient cloud H2O emission not visible in the Odin spectrum.

This result is verified by a convolution of the Odin’s H2O map of the Orion KL core region (Fig. 1 of Olofsson et al.) to the angular resolution of SWAS. The resulting “SWAS convolved” Odin map centre spectrum is shown in Fig. 3, together with the H2O spectrum observed by SWAS. The two almost identical spectra also demonstrate the very consistent calibrations of the Odin and SWAS H2O data sets.

The H18O spectrum observed by Odin towards Orion KL is displayed in Fig. 4, along with the corresponding SWAS spectrum (Melnick et al. 2000b). The SWAS H18O spectrum here has been scaled by a factor 3.3 to match the red wing of the Odin spectrum. This value is close to the theoretical beam filling ratio for a very small source, 3.37. Using Eq. (1) we estimate that the size of the red H18O outflow (moving away from us) must be ≤50″. We also should point out that the SO2 (286,23−285,23) line visible near −65 km s⁻¹ also scales by a factor of ≈3.3, indicating a similarly small source, which is expected for this high energy line (\( E_u = 475 \) K). As in the case of H2O the blueward (counted from 10 km s⁻¹) H18O outflow emission is weaker than the redward one.

3.6. H2O in the W3 region

In Fig. 5 the Odin 557 GHz H2O spectrum observed towards W3 (IRS 5) is compared with the corresponding SWAS spectrum (unpublished data provided by Melnick et al.). Both spectra are double peaked, similar to CO (3−2) and (6−5) spectra, and almost lack emission around –38 km s⁻¹, which is the cloud core velocity as observed in the C18O (2−1), 13CO (6–5) and C34S (5–4) lines (Hasegawa et al. 1994; Tieftrunk et al. 1995). Hence the double peaked H2O spectrum can be attributed to strong self-absorption. While the peak intensity of the Odin spectrum is about twice as strong as that from SWAS, according to Eq. (1) indicating a Gaussian FWHP size of the emission region of about 150′′, it is also evident
that the “blueshifted” peak of the Odin spectrum is considerably weaker than the “redshifted” one, which is opposite to the situation in the SWAS spectrum. The reason for this difference was clear immediately after our Odin observations, since we had also included in the observations a $3 \times 3$ point reference map at $2\arcmin$ spacing, centred on W 3 (IRS 5). In the $2\arcmin$ west position, near W 3 (IRS 4), the Odin H$_2$O signal was in fact considerably stronger than at the map centre. While centred on W 3 (IRS 5), the larger SWAS beam is picking up H$_2$O emission from W 3 (IRS 4) as well. Since we here really could benefit from the smaller Odin $2\arcmin.1$ beam, a mapping of the W3 region at $1\arcmin$ spacing rapidly was scheduled. The results are presented in an accompanying Letter by Wilson et al. (2003).

3.7. H$_2$O in the IRAS 16293-2422 core

Figure 6 shows the H$_2$O spectrum observed by Odin centred on the the very nearby (160 pc), deeply embedded young stellar object IRAS 16293–2422. This IRAS source is a prominent outflow source (e.g. Walker et al. 1988; Castets et al. 2001; Lis et al. 2002). The self-reversal appearance of the spectral lines have been modelled by infall models (cf. Schöier et al. 2002). The H$_2$O spectrum, with its characteristic absorption dip, shows a similar asymmetry as the HCO$^+$ (3–2) spectrum observed by Lis et al. (2002) toward the E1 peak about 50$\arcsec$ east of the IRAS source. The HCO$^+$ (3–2) spectrum toward the IRAS source of Lis et al. (2002) shows the opposite asymmetry, i.e. the blue-shifted side is the strongest.

For comparison we also show, in Fig. 6, the SWAS H$_2$O spectrum toward IRAS 16293–2422 which has a very similar shape but is about a factor of 2 weaker. This would be expected if the source size is similar to our $2.1\arcmin$ beam.

In Fig. 6 we have also included a model H$_2$O spectrum as would be observed by a $2.1\arcmin$ antenna beam. This model spectrum is based on the infall model parameters of Schöier et al. (2002). The radiative transfer was solved by means of accelerated lambda iteration (ALI, e.g. Rybicki & Hummer 1991). We found that the infall source alone could not explain the deep and rather narrow absorption. In order to model the absorption dip we had to include a halo with properties very close to those determined in the parental cloud by the ammonia observations of Menten et al. (1987). By including a $15 M_\odot$ halo having an H$_2$ density of $10^4$ cm$^{-3}$, a kinetic temperature of 11 K, a turbulent velocity width of 0.3 km s$^{-1}$, we obtained the spectrum shown in Fig. 6. The ortho-H$_2$O abundance throughout the entire cloud was set to $2 \times 10^{-9}$. Apparently, the observed
H$_2$O spectrum is dominated by emission from the outflow and there are no obvious indications of infall. In fact, a central model source with a static velocity field combined with the halo works equally well (although the core ortho-H$_2$O abundance could be a magnitude higher because of more effective trapping).

It is worth stressing again that the observed H$_2$O spectrum resembles the HCO$^+$(3–2) spectrum toward the E1 peak, which is suggested to be an interaction region of the outflow and the dense ambient gas (Lis et al. 2002).

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References
