

Long-term photometry of the near-contact detached binary GR Tauri[★]

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Abstract. We present photometric observations of the eclipsing binary GR Tau made in 1985, 1993, 1998, 2000 and 2001. New times of minima and ephemerides are given. Based on a statistical analysis of the times of minima obtained by photoelectric and CCD photometry, the orbital period of the system is found to be decreasing with a rate of $\Delta P/P = 3.27 \times 10^{-10}$ during the past decades. An assumed hot spot on the trailing hemisphere of the secondary enables us to reconstruct the peculiar asymmetry of the distorted light curves. The photometric solution with the Wilson-Devinney code reveals that GR Tau is a near-contact detached system with the two components almost filling their respective critical Roche lobes. The photometric mass ratio turns out to be $q = 0.22$, and the masses and radii are derived for the components as $M_1 = 1.42 \pm 0.18 M_\odot$, $M_2 = 0.31 \pm 0.04 M_\odot$, $R_1 = 1.47 \pm 0.19 R_\odot$ and $R_2 = 0.70 \pm 0.09 R_\odot$. The authors discuss the possible evolutionary status of the component stars and the probable nature of the assumed hot spot.

Key words. stars: binaries: eclipsing – stars: individual: GR Tau

1. Introduction

GR Tau (=BD +20°0685) was discovered to be an eclipsing binary by Strohmeier et al. (1957) from their photographic survey for variable stars. The subsequent observations of Strohmeier et al. (1963) gave the first photographic light curve and a period of 0.474022 days for the system. Gotz & Wenzel (1962) assigned a spectral type of A9 for the primary component. Yamasaki et al. (1984) published the first photoelectric and spectroscopic observations of the binary. The spectral type of GR Tau was suggested to be about A5. They obtained a small mass function of $0.0101 M_\odot$ and a revised orbital period of 0.4298525 days. From the analysis of the light curves, they determined the first photometric solution of the system. Their solutions showed that GR Tau is a near-contact system, but their photometric mass ratio of $q = m_2/m_1 = 0.32$ leads to a mass of the primary component unacceptably small for an A5V star.

Lázaro et al. (1995) observed the system with visual and infrared photometry, and obtained *B*, *V*, *J* and *K* light curves. The near-contact configuration of GR Tau was confirmed by their Roche model photometric solution. From the analysis of the colours of GR Tau, however, they showed that the primary component's spectral type could be about A9V, instead of A5V as classified by Yamasaki et al. (1984). The orbital period was refined to be 0.42985060 days.

The asymmetry of the light curves of GR Tau has been noted by most of the previous studies since Yamasaki et al. (1984). It is also found that the light curve of the system

changes frequently with time as it presents always an asymmetry with Max I brighter than Max II. It may be due to the so-called O'Connell effect, as is widely found among the V1010 Oph type binaries. However, no detailed analysis in modeling the asymmetry as well as its variations has yet been made.

Because of the peculiar asymmetric feature of the light curve and the near-contact configuration, GR Tau has been listed as one of the programming targets for our long-term photometric observations. We have observed the system in five runs since 1985. In this paper, we report the long-term photometric observations made in the past 17 years. The light curves are analyzed using the W–D code (Wilson & Devinney 1971; Wilson 1979, 1990). The asymmetry and the changes of the light curves are discussed on the basis of a model with an assumed hot spot possibly caused by mass transfer between the component stars.

2. Observations and data reduction

The star GR Tau was observed in five runs made in the observing seasons of 1985, 1993, 1998, 2000 and 2001. The observations were all made at the Xinglong station of Beijing Astronomical Observatory. Table 1 gives relevant information on the observations.

The 1985 and 1993 observations were both carried out with the 60 cm reflector. The star BD +20°0684 was employed as the comparison, and BD +19°0642, as the check star. The differences in magnitude between the variable and comparison stars were corrected for differential extinction and transformed into the UB_v system. The *BV* light curves obtained in 1993 have been reported by Fang et al. (1994).

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[★] Tables A.1 to A.3 are only available in electronic at the CDS via anonymous ftp to cdsarc.u-strasbg.fr (130.79.125.5) or via <http://cdsweb.u-strasbg.fr/cgi-bin/qcat?J/A+A/395/587>

Table 1. Information on the observations of GR Tau.

Year	Telescope	Photometer	Filter	Comparison star
1985	60 cm	single channel	<i>B, V</i>	BD +20°0684
1993	60 cm	single channel	<i>B, V</i>	BD +20°0684
1998	85 cm	single channel	<i>B, V</i>	BD +19°0642
2000	85 cm	CCD	<i>V</i>	GSC 1258–171
2001	85 cm	three channel	<i>B, V</i>	GSC 1258–171

The observations in 1998 were made on the 85 cm Cassegrain telescope. In this run, the comparison star was BD +19°0642 and the check star was BD +20°0684. Data reduction was done in the same way as the 1985 and 1993 observations.

The 2000 observations were carried out on the 85 cm reflector with a red-sensitive Thomson TH7882 576×384 CCD photometer (Wei et al. 1990; Zhou et al. 2001) which provides a field of view of about 12' × 8'. Five stars in the field around GR Tau were selected as reference stars. We finally chose the star GSC 1258–171 as the comparison star through a try during the data reduction. Sky extinction was not taken into account in view of the close proximity of the program stars. The exposure times were 30 s. A single *V* filter was employed.

The 2001 observations were made on the 85 cm telescope with a 3-channel photometer. On the first two nights (Oct. 28 and 29, 2001), we obtained a completely covered light curve with *V* filter. The observations in the *B* passband made on Oct. 31 and Nov. 1 have not recorded a complete light curve but did cover two primary eclipses. The comparison star was the same as that used for the 2000 CCD observations. The exposure times ranged from 10 to 15 s depending on the weather conditions.

All measurements of the 1985, 1993 and 1998 photoelectric observations are given in the appendix. For the 2000 CCD and 2001 fast photometry observations we list only the normal points in the appendix in the form of phase vs. relative magnitude, as the amount of individual data is too large.

3. The period of GR Tau

By using the K–W method (Kwee & van Woerden 1956), 13 new times of minima were derived from our data. They are given in Table 2 along with previously published times of photoelectric and CCD minima collected since Yamasaki et al. (1984). These minimum times with a time-base of about 30 years enable us to discuss the long-term variations of the orbital period of GR Tau. Analyzing with the least square method, the following new ephemerides are derived:

$$\text{Min.}I = \text{HJD}2446415.0219(3) + 0^{\text{d}}42985101(7) \times E$$

$$\text{Min.}I = \text{HJD}2446415.0220(4) + 0^{\text{d}}42985171(6) \times E - 7.03(1.12) \times 10^{-11} \times E^2$$

As shown in Fig. 1 the quadratic ephemeris gives much better fitting to the observations than the linear ephemeris. The O–C values calculated from the second ephemeris above for each minimum time are given in Table 2. It is obvious that the period of GR Tau appears to show appreciable decreasing trend in the

Table 2. Times of photoelectric and CCD minima of GR Tau.

HJD	<i>E</i>	O–C	source
2 400 000+			
44544.3075	–4352.0	–.0011	Yamasaki & Okazaki (1984)
44573.1074	–4285.0	–.0020	Yamasaki & Okazaki (1984)
44578.2643	–4273.0	–.0013	Yamasaki & Okazaki (1984)
44579.1252	–4271.0	.0002	Yamasaki & Okazaki (1984)
46414.1619	–2.0	–.0022	Present study
46415.0209	.0	–.0025	Present study
46438.6622	55.0	–.0019	Mullis & Faulkner (1991)
47821.4944	3272.0	–.0056	Wunder et al. (1992)
47827.5148	3286.0	–.0008	Hanzl (1990)
47849.4363	3337.0	–.0012	Hanzl (1990)
47881.6752	3412.0	–.0006	Mullis & Faulkner (1991)
47889.4089	3430.0	–.0052	Hanzl (1990)
47945.2937	3560.0	–.0032	Hanzl (1990)
48288.3122	4358.0	–.0042	Agerer (1991)
48619.2973	5128.0	–.0074	Hanzl (1995)
48936.5294	5866.0	–.0019	Lázaro et al. (1995)
48938.6757	5871.0	–.0040	Lázaro et al. (1995)
48952.4314	5903.0	–.0022	Lázaro et al. (1995)
48953.2915	5905.0	–.0015	Lázaro et al. (1995)
48953.5019	5905.5	–.0059	Lázaro et al. (1995)
48954.3659	5907.5	–.0013	Lázaro et al. (1995)
49334.1394	6791.0	–.0012	Present study
49368.3190	6870.5	.0026	Agerer (1994)
49368.3199	6870.5	.0035	Agerer (1994)
49734.3334	7722.0	–.0025	Agerer (1996)
49760.3440	7782.5	.0042	Agerer (1996)
49760.3490	7782.5	.0092	Agerer (1996)
50013.5295	8371.5	.0061	Agerer et al. (1999)
50422.5188	9323.0	–.0046	Agerer & Hübscher (1998)
51139.5145	10 991.0	–.0011	Agerer et al. (1999)
51173.0380	11 069.0	–.0050	Present study
51176.0496	11 076.0	.0002	Present study
51198.4049	11 128.0	–.0014	Agerer et al. (1999)
51486.3991	11 798.0	–.0032	Agerer et al. (2001)
51571.0834	11 995.0	.0014	Present study
51573.0238	11 999.5	.0082	Present study
51574.0894	12 002.0	–.0044	Present study
52211.1277	13 484.0	–.0051	Present study
52211.3454	13 484.5	–.0023	Present study
52214.1353	13 491.0	–.0053	Present study
52220.1553	13 505.0	–.0049	Present study

past decades. The rate of the period decrease turns out to be $\Delta P/P = 3.27 \times 10^{-10}$ (~ 0.01 s yr^{–1}).

Yamasaki et al. (1984) concluded that the orbital period of GR Tau was unlikely to have any obvious long-term variation based on a statistical study of the times of minima observed over 15 years before 1980. The large scattering of their statistical data caused by the uncertainties in visual and photographic observations does not exclude the possibility of long-term changes of the orbital period as indicated by recent photoelectric and CCD photometry.

4. The light curves and photometric analysis

4.1. The asymmetry and variations of the light curves

Using the above quadratic ephemeris, phases are computed for all the observations and the corresponding light curves are displayed in Figs. 3–7. It is clear that all the light curves obtained in the five observing seasons show marked asymmetry

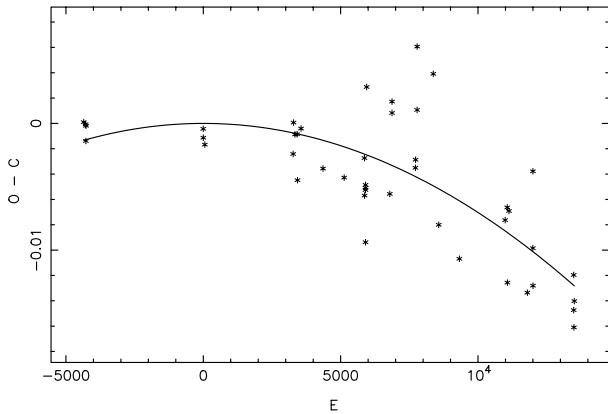


Fig. 1. The O–C diagram of the observed times photoelectric and CCD minima. The solid line represents the quadratic fitting curve.

Table 3. Main feature parameters of the V band light curves obtained in the five epochs.

Epoch	1985	1993	1998	2000	2001
Max. I–Min. I	–0.51	–0.50	–0.53	–0.53	–0.50
Max. I–Min. II	–0.22	–0.20	...	–0.23	–0.23
Max. II–Min. I	–0.44	–0.45	–0.44	–0.45	–0.43
Max. II–Min. II	–0.15	–0.15	...	–0.15	–0.16
Max. I–Max. II	–0.07	–0.05	–0.09	–0.08	–0.07

with Max. I brighter than Max. II. The general features are closely similar to those of the light curves previously obtained by Yamasaki et al. (1984). In Table 3 we list the main feature parameters of V-band light curves in the five different epochs. Since the secondary minima are not completely covered in the 1998 light curves, the corresponding values of Max. I–Min. II and Max. II–Min. II remain unknown.

The asymmetry of the light curve of GR Tau seems to be fixed in phase and not to migrate with phase like the photometric waves exhibited by the RS CVn stars. But the amplitude of asymmetry changed with time. From Table 3 we find that the amplitude Max. I–Min. I varied obviously with time, while the value of Max. II–Min. I seems to remain nearly constant. Thus the difference of brightness between Max. I and Max. II appears to depend primarily on the former.

4.2. The hot spot assumption

The peculiar asymmetry of the light curve seems to be something like the case in the stream-distorted light curves observed frequently in Algol-type systems with diminution of light in the phase interval 0.5–1.0, but for GR Tau there is no evidence of mass transfer from the secondary to the primary and, moreover, the near-contact configuration revealed by Yamasaki et al. (1984) and Lázaro et al. (1995) implies that there will be no space between the stellar surface and Roche lobe to allow the formation of an optically thick stream layer that would affect the brightness of the system.

On the other hand, the photometric asymmetry of GR Tau appears unlikely to be caused by the dark spot activity as is the

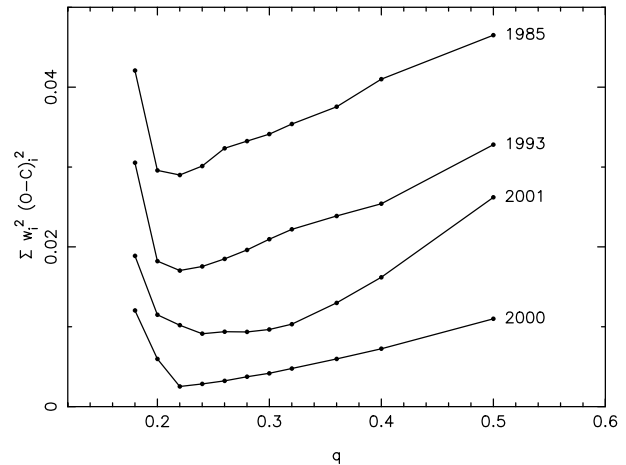


Fig. 2. The diagram of $\sum w_i^2(O-C)_i^2$ vs. q for the determination of solutions of the 1985, 1993, 2000 and 2001 light curves of GR Tau.

case in RS CVn systems. The primary component of GR Tau has a spectral type of A, which is much earlier than that of active stars. The secondary component is ascribed to be a late K or M dwarf star, but it contributes only a few percent to the total brightness (Yamasaki et al. 1984; Lázaro et al. 1995). In any dark spot model assumed for the system the photometric analysis will always produce an unrealistically low spot temperature.

In order to model the asymmetric light curves of GR Tau, a natural alternative is to try the hot spot assumption. Shaw (1990) found that such asymmetry of light curves with Max. I brighter than Max. II appears to occur regularly in the V1010 Oph-type near-contact systems. From the IUE satellite observations of V1010 Oph, Shaw & Guinan (1990) found that there is an additional thermal source responsible for the Lyman α emission from the trailing hemisphere of the secondary component. They suggested further that the photometric asymmetry of V1010 Oph may also be caused by the additional thermal source, a phenomenon which could be typical for the V1010 Oph-type stars. For a similar, near-contact binary V361 Lyr, Zhai & Fang (1995) used a hot-spot model and gave a thorough treatment of its striking asymmetric light curves. In the case of GR Tau, we find rapid decrease of the system orbital period. This implies there could be mass transfer from the hot primary onto the hemisphere of the cool secondary component causing extra light emitted from the later type star. Thus a hot-spot model could well be the most preferable one to describe the light-curve asymmetry of GR Tau.

By assuming the existence of a hot spot on the trailing hemisphere of the secondary, the W–D synthetic light-curve program is employed to carry out the photometric analysis of the observed light curves.

4.3. Photometric solutions

The 1992 version of the Wilson–Devinney code (Wilson 1990) was employed for our light curve analysis. All the individual light curves were combined into normal points and were used to compute the solutions. In deriving the photometric solution, we made an un-spotted synthesis for the light curves of 1985, 1993, 2000 and 2001 at the first stage. Since 1998 light curves

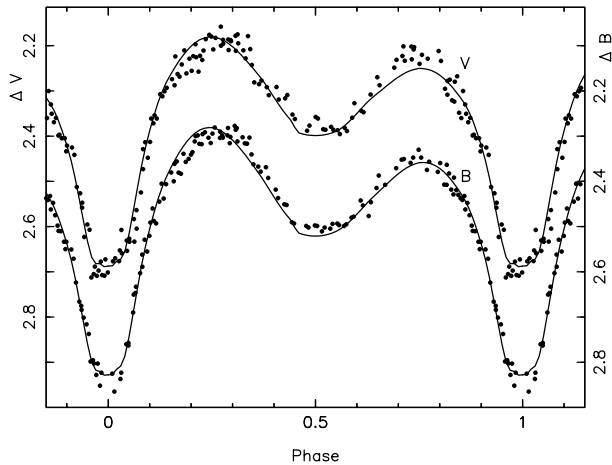


Fig. 3. The *BV* light curves of GR Tau obtained in 1985. Dots present the observations and the solid lines the theoretical light curves.

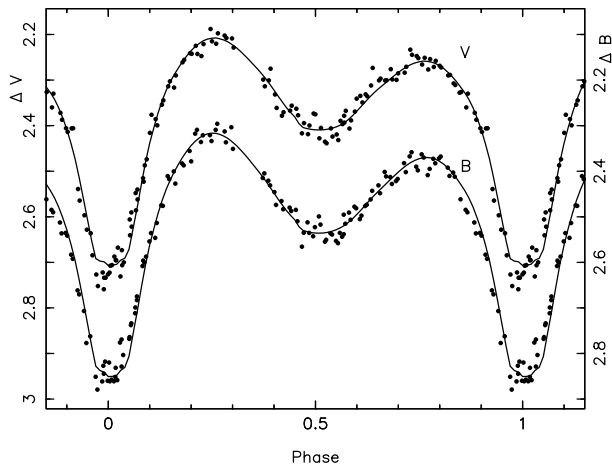


Fig. 4. The same as Fig. 3, but for the year 1993.

do not completely cover the secondary eclipse, they are not included in the test calculation. The following assumptions are made for the light-curve analysis: temperature of the primary component is set at $T_1 = 7500$ K according to Lázaro et al. (1995); gravity exponents and albedos are taken to be $g_1 = 1.0$ and $A_1 = 1.0$ for the primary (radiative envelope) and $g_2 = 0.32$ and $A_2 = 0.5$ for the secondary (convective envelope); the limb darkening coefficients, according to Al-Naimiy (1978), are assigned as $x_1 = 0.60$, $x_2 = 0.73$ in *V* and $x_1 = 0.79$, $x_2 = 0.92$ in *B*, respectively.

The adjustable parameters are the phase shift ϕ_0 , the orbital inclination i , the temperature T_2 , the luminosity L_1 and L_2 , the surface potentials Ω_1 and Ω_2 , and the mass ratio $q = m_2/m_1$. According to the analysis of Lázaro et al. (1995), the mass ratio of GR Tau should be around 0.22 from its small mass function $f(m) = 0.0101 M_\odot$. On the other hand Yamasaki et al. (1984) derived a mass ratio of 0.32. We just fix the mass ratio at a series of values of 0.18, 0.20, 0.22, 0.24, 0.26, 0.28, 0.30, 0.32, 0.34, 0.36, 0.40 and 0.50. All the test solutions started from Mode 2. In this way, the initial solutions are carried out at each assumed mass ratio for each light curve.

After that, the hot-spot assumption mentioned above is employed to perform a further analysis for all the observations

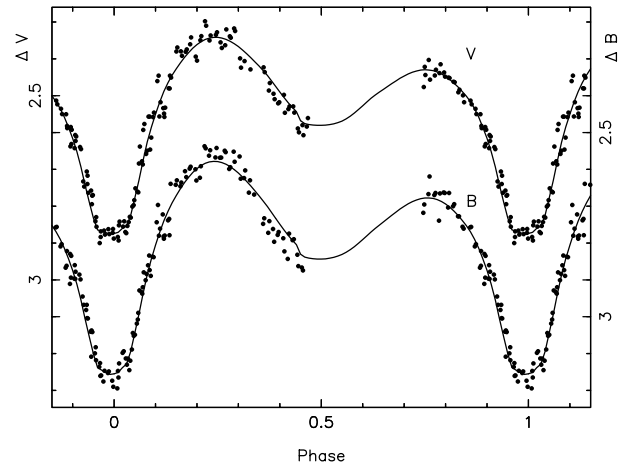


Fig. 5. The *BV* light curves of GR Tau in 1998.

of the light curves. The spot parameters are the spot temperature T_s (given as a fraction of the surrounding photospheric temperature), the spot radius (defined as half of the angular radius subtended by the spot at the stellar center), the spot latitude (0 – 180° , measured from the north to the south pole of star) and spot longitude (0 – 360° , measured for each star counterclockwise from the meridian at the substellar point to the other component). In our hot-spot model, the spot latitude is fixed at 90° . The preliminary spot longitude could be found approximately from the phase of the maximum brightness of the light hump caused by the hot spot on the secondary component. The other two spot parameters, T_s and r_s , are calculated by adjusting the theoretical light curves to fit approximately the observed distorted light curves. After several runs, the parameters of both the Roche-model and the spot model are corrected alternately and best solution is derived for each assumed mass ratio. By referring to the distributions of $\sum w^2(O-C)^2$ vs. q_i for each program light curve as shown in Fig. 2, it is found that the most probable solution would be at around $q = 0.22$ for the 1985, 1993 and 2000 light curves and at $q = 0.24$ for the 2001 light curve. The final convergent solution of the photometric mass ratio is derived to be $q = 0.215 \pm 0.003$ for the 1985, $q = 0.222 \pm 0.002$ for the 1993, $q = 0.219 \pm 0.001$ for the 2000, and $q = 0.245 \pm 0.002$ for the 2001 light curve, respectively. The solutions are given in Table 4. As for the 1998 light curves, no Roche solution could be found from the above analysis. Here we adopt the Roche parameters for the 1993 light curve to fit the 1998 light curves. Subsequently the parameters of the hot-spot model for the 1998 light curves are derived through synthesis of the residual distortions.

5. Discussion and conclusions

We have presented a long-term photometric study of GR Tau. The results of period analysis and light-curve synthesis allow us to draw the following conclusions.

1) GR Tau appears to be a near-contact detached binary with the two components each almost filling its Roche lobe. Such a near-contact configuration has already been reported by Yamasaki et al. (1984), but that is based on the photometric solution of part of their 1980 light curves. The solution of

Table 4. Photometric solutions of the 1985, 1993, 1998, 2000 and 2001 light curves.

Parameter	LC 1985	LC 1993	LC 1998	LC 2000	LC 2001
i (degrees)	87.68 ± 0.52	87.15 ± 0.73	.	87.42 ± 0.18	87.58 ± 0.26
$q = m_2/m_1$	0.215 ± 0.003	0.222 ± 0.002	.	0.219 ± 0.001	0.246 ± 0.002
T_1 (K)*	7500	7500	.	7500	7500
T_2 (K)	3682 ± 109	3624 ± 82	.	3641 ± 32	3650 ± 27
g_1^*	1.0	1.0	.	1.0	1.0
g_2^*	0.32	0.32	.	0.32	0.32
A_1^*	1.0	1.0	.	1.0	1.0
A_2^*	0.5	0.5	.	0.5	0.5
$x_1 B^*$	0.73	0.73	.		
$x_1 V^*$	0.60	0.60	.	0.60	0.60
$x_2 B^*$	0.92	0.92	.		
$x_2 V^*$	0.79	0.79	.	0.79	0.79
Ω_1	2.3050 ± 0.0122	2.3132 ± 0.0074	.	2.2912 ± 0.0040	2.3555 ± 0.0058
Ω_2	2.2930 ± 0.0080	2.3060 ± 0.0043	.	2.2848 ± 0.0030	2.4770 ± 0.0096
r_1 (pole)	0.4737 ± 0.0018	0.4731 ± 0.0011	.	0.4777 ± 0.0008	0.4691 ± 0.0011
r_1 (point)	0.5887 ± 0.0063	0.5918 ± 0.0041	.	0.6156 ± 0.0062	0.6071 ± 0.0083
r_1 (side)	0.5116 ± 0.0023	0.5108 ± 0.0015	.	0.5173 ± 0.0012	0.5062 ± 0.0015
r_1 (back)	0.5324 ± 0.0026	0.5322 ± 0.0016	.	0.5398 ± 0.0014	0.5295 ± 0.0017
r_2 (pole)	0.2319 ± 0.0106	0.2340 ± 0.0049	.	0.2380 ± 0.0036	0.2197 ± 0.0037
r_2 (point)	0.2988 ± 0.0392	0.3025 ± 0.0235	.	0.3282 ± 0.0296	0.2534 ± 0.0078
r_2 (side)	0.2405 ± 0.0125	0.2428 ± 0.0058	.	0.2475 ± 0.0043	0.2258 ± 0.0042
r_2 (back)	0.2686 ± 0.0210	0.2711 ± 0.0097	.	0.2791 ± 0.0075	0.2430 ± 0.0059
$L_1/(L_1 + L_2)B$	0.996 ± 0.003	0.997 ± 0.002	.		
$L_1/(L_1 + L_2)V$	0.992 ± 0.004	0.993 ± 0.002	.	0.994 ± 0.001	0.994 ± 0.001
Ω_{in}	2.2696	2.2866	.	2.2793	2.3436
Spot model					
co-latitude*	90.00	90.00	90.00	90.00	90.00
co-longitude	270.40 ± 0.87	272.10 ± 1.25	272.80 ± 1.57	271.82 ± 0.46	270.95 ± 0.42
r_{spot}	25.73 ± 0.48	25.64 ± 0.35	27.60 ± 0.44	27.48 ± 0.30	28.60 ± 0.28
T_{spot}/T_2	2.040 ± 0.027	1.928 ± 0.032	2.150 ± 0.045	2.038 ± 0.025	2.100 ± 0.028
$\sum w^2(O-C)^2$	0.0263	0.0167	.	0.0024	0.0090

* Assumed.

Lázaro et al.'s (1995) resulted in a near-contact semi-detached configuration for the system with the secondary filling its Roche lobe. However, being taken from an average Roche model solution of their 1992 light curves, the distortion of the light curves, albeit slight, has not been taken into account. In our analysis all the observations of the light curves have been used to perform the photometric solution. By adopting a hot-spot model, our theoretical light curves give much better fittings to the observations of all the 1985, 1993, 2000 and 2001 light curves without all-phase systematic deviations. Furthermore, such a near-contact status of GR Tau is supported by the fact of orbital period decrease as evidenced by O–C analysis.

2) The W–D approach in the analysis of light curves with total eclipses can in general yield a reliable mass ratio for the binary systems under study. If we take the average values of the mass ratio $q = 0.22$ and orbital inclination $i = 87^\circ.42$ obtained from the photometric solutions for the four light curves, and combine them with the spectroscopic mass function $f(m) = 0.0101 M_\odot$ (Yamasaki et al. 1984), we can derive the absolute dimensions of GR Tau as:

$$\begin{aligned}
 M_1 &= 1.42 \pm 0.18 M_\odot & R_1 &= 1.47 \pm 0.19 R_\odot \\
 & & L_1 &= 6.14 \pm 1.70 L_\odot \\
 M_2 &= 0.31 \pm 0.04 M_\odot & R_2 &= 0.70 \pm 0.09 R_\odot \\
 & & L_2 &= 0.08 \pm 0.02 L_\odot.
 \end{aligned}$$

From their photometric solution, Yamasaki et al. (1984) deduced an unacceptably small mass and radius for the primary component. Our results have greatly reduced the discrepancy of the derived mass and radius compared with the spectral type. However, the new values of mass and radius deduced for the primary are still somewhat small when compared to its spectral type of A9V.

If the primary component of GR Tau is a normal main sequence star, for a spectral type A9V, its mass and radius should be about $M_1 = 1.68 M_\odot$ and $R_1 = 1.54 R_\odot$ according to the calibration from Schmidt-Kaler (1982). To match the deduced mass for the primary with its spectral type, an assumed mass ratio of $q = 0.206$ would be appropriate for the mass function of $f(m) = 0.0101 M_\odot$, or else, the mass function should be taken to be about 0.0120 with our photometric mass ratio $q = 0.22$. Since we know that the mass ratios derived from all the photometric solutions to date (Yamasaki et al. 1984; Lázaro et al. 1995 and present work) are all larger than 0.22, we suggest that the discrepancy of the deduced mass in relation to the spectral type of GR Tau mentioned above might be due to the slight underestimation of the mass function if the photometric mass ratio is more reliable. A new high-precision radial velocity study based on high S/N spectroscopy is especially desirable to settle the problem.

Another possibility is that, in such a near-contact binary system, the primary of GR Tau might be a low-mass and

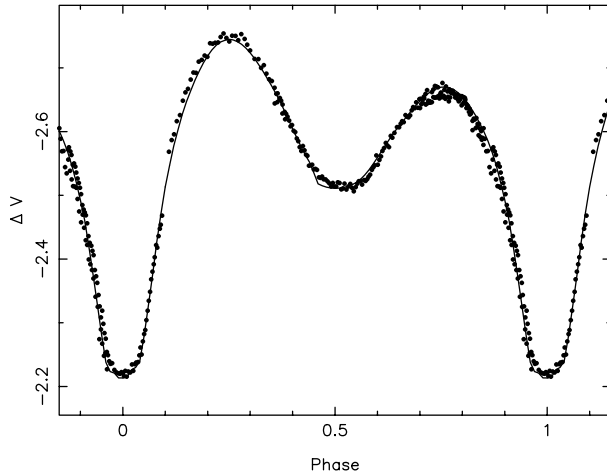


Fig. 6. The V light curve of GR Tau observed in 2000.

low-luminosity star as suggested by Yamasaki et al. (1984). In our results, the secondary has a temperature of about 3640 K, which corresponds to a spectral type of M1–2 if it is a dwarf star. Taking the values of absolute dimensions deduced for GR Tau, we plot the two components on the mass-radius diagram and find that the primary is still on the main sequence while the secondary is evidently an evolved star far away from the main sequence. The evolutionary status of GR Tau is something like that in other near-contact detached systems with the less massive secondary actually more evolved than the primary. It is suggested that the evolved secondary component should have been the more massive component but is now the less massive one after losing a considerable part of its original mass through Roche lobe overflow. And at present, the secondary could have just shrunken back from its Roche lobe making the system a near-contact detached system. The present massive primary component would have gained additional mass from the secondary. This could explain why it looks strange with respect to its spectral type.

3) The hot spot assumption gives a good description of the asymmetry of the light curves of GR Tau. The photometric solutions indicate that the hot spot remains at nearly the same position, with almost the same temperature and size in both epochs of 1985, 1993, 1998, 2000 and 2001. Taking the average value of $T_{\text{spot}}/T_2 = 2.05$ and $r_{\text{spot}} = 27^\circ.01$, the additional luminosity emitted by the hot spot can be calculated as:

$$L^* = S_{\text{hem}} \sigma T_2^4 S_{\text{spot}} ((T_{\text{spot}}/T_2)^4 - 1)$$

where $S_{\text{spot}} = 1 - \cos r_{\text{spot}}$ is the spot area as a fraction of the hemispheric area S_{hem} of the star.

The physical nature of the hot spot is still open for discussion. In view of the probable mass transfer in the system revealed by the period changes observed, we suggest that the hot spot on the secondary could be caused by the mass accreted from the primary component. Suppose that the accreting mass falls from the inner Lagrange point L1 directly onto the surface of the secondary, the luminosity caused by the accreting mass would be approximately represented by the following expressing according to Zhai & Fang (1995):

$$L_{\text{acc}} = GM_1 \frac{dm}{dt} (\Omega_2 - \Omega_{\text{in}}) / A$$

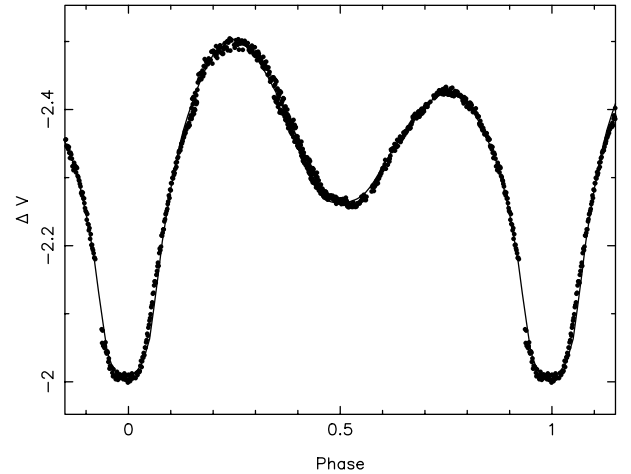


Fig. 7. The same as Fig. 6, but observed in 2001.

where G is the gravitation constant, $\frac{dm}{dt}$ is the mass-accretion rate, A is the distance between the two components, Ω_2 and Ω_{in} are the dimensionless potentials of the stellar surface and the critical Roche lobe, respectively. If $L^* = L_{\text{acc}}$, or we assume that the enhanced luminosity of the secondary component is entirely from the mass accretion, we find the mass-accretion rate to be $\frac{dm}{dt} = 1.86 \times 10^{-7} M_{\odot}/\text{yr}$, which corresponds to a period-change rate of $\Delta P/P = -1.34 \times 10^{-9}$.

The observed period-change rate is about $\Delta P/P = -3.27 \times 10^{-10}$, the absolute value of which is much less than what we predicted. This implies that the possible mass transfer we infer from our photometry is not strong enough to create the spot luminosity unless some additional energy, besides the potential energy which we have taken into account, has also been released by the mass flow. Usually the original velocity of the accreting mass around the Lagrange point is not equal to zero as we have assumed, thus the original kinetic energy of the mass flow could probably serve as part of the additional energy released. Another possible energy source might be the thermal energy brought about directly by the mass flow through the high temperature plasma wind from the primary component. A satisfactory model should be carefully constructed in order to give an comprehensive description of the mass accretion and the physical nature of the hot spot, but this goes beyond the scope of the present work.

4) All the light curves of GR Tau obtained by us in the five epochs are asymmetric. They are basically similar to the 1980 light curves published by Yamasaki et al. (1984). On the other hand, Hanzl (1990) and Lázaro et al. (1995) report nearly symmetric light curves of the system. The light curve of GR Tau seems to vary between asymmetric and symmetric forms on a time scale of several years. Comparing the symmetric light curves with the asymmetric ones, we notice that the asymmetric light curves always have deeper primary minima (Max. I–Min. I) than the symmetric light curves. It is suggested that the variation in the light curves of the system could be due to changes of the hot spot in the secondary component. Well-planned observational studies in the future are needed to understand better the true nature and behaviour of the binary system including the hot spot deduced by conventional theory.

Appendix

Table A. 4. Normal points of the 2000 light curve.

Phase	ΔV	Phase	ΔV	Phase	ΔV	Phase	ΔV	Phase	ΔV
0.0082	-2.2208	0.2203	-2.7435	0.4172	-2.5931	0.6182	-2.5919	0.8197	-2.6381
0.0200	-2.2286	0.2298	-2.7478	0.4317	-2.5786	0.6336	-2.6060	0.8351	-2.6165
0.0309	-2.2355	0.2436	-2.7484	0.4468	-2.5597	0.6495	-2.6236	0.8484	-2.5975
0.0463	-2.2675	0.2571	-2.7489	0.4621	-2.5480	0.6610	-2.6329	0.8618	-2.5808
0.0606	-2.3318	0.2713	-2.7469	0.4736	-2.5441	0.6723	-2.6365	0.8728	-2.5612
0.0718	-2.3866	0.2877	-2.7383	0.4855	-2.5354	0.6837	-2.6521	0.8865	-2.5388
0.0861	-2.4440	0.2984	-2.7303	0.5009	-2.5317	0.6990	-2.6605	0.9010	-2.4926
0.1117	-2.5775	0.3131	-2.7149	0.5152	-2.5301	0.7133	-2.6730	0.9140	-2.4509
0.1242	-2.6061	0.3280	-2.7024	0.5298	-2.5287	0.7256	-2.6712	0.9272	-2.4081
0.1378	-2.6373	0.3403	-2.6906	0.5406	-2.5271	0.7370	-2.6845	0.9407	-2.3410
0.1506	-2.6644	0.3498	-2.6804	0.5529	-2.5326	0.7518	-2.6801	0.9541	-2.2866
0.1664	-2.6888	0.3648	-2.6683	0.5672	-2.5375	0.7672	-2.6785	0.9644	-2.2521
0.1823	-2.7112	0.3806	-2.6492	0.5819	-2.5483	0.7821	-2.6742	0.9786	-2.2289
0.1961	-2.7208	0.3958	-2.6278	0.5963	-2.5655	0.7935	-2.6667	0.9959	-2.2205
0.2096	-2.7385	0.4067	-2.6099	0.6072	-2.5758	0.8043	-2.6648		

Table A. 5. Normal points of the 2001 V light curve.

Phase	ΔV	Phase	ΔV	Phase	ΔV	Phase	ΔV	Phase	ΔV
0.0061	-2.0066	0.2062	-2.4743	0.4064	-2.3365	0.6056	-2.3070	0.8072	-2.3947
0.0214	-2.0121	0.2199	-2.4774	0.4199	-2.3171	0.6199	-2.3247	0.8211	-2.3843
0.0330	-2.0326	0.2344	-2.4873	0.4333	-2.3010	0.6327	-2.3383	0.8330	-2.3666
0.0465	-2.0754	0.2464	-2.4856	0.4468	-2.2864	0.6461	-2.3495	0.8465	-2.3507
0.0599	-2.1339	0.2594	-2.4872	0.4603	-2.2731	0.6596	-2.3609	0.8600	-2.3288
0.0734	-2.1884	0.2728	-2.4869	0.4729	-2.2677	0.6731	-2.3771	0.8733	-2.3050
0.0869	-2.2400	0.2873	-2.4803	0.4878	-2.2634	0.6866	-2.3844	0.8861	-2.2781
0.1004	-2.2862	0.3002	-2.4724	0.4981	-2.2613	0.7000	-2.3922	0.8996	-2.2404
0.1138	-2.3219	0.3134	-2.4628	0.5136	-2.2537	0.7102	-2.3998	0.9130	-2.1938
0.1273	-2.3526	0.3278	-2.4457	0.5279	-2.2492	0.7260	-2.4109	0.9210	-2.1714
0.1401	-2.3755	0.3404	-2.4274	0.5410	-2.2506	0.7400	-2.4159	0.9412	-2.0598
0.1534	-2.3990	0.3535	-2.4122	0.5533	-2.2591	0.7533	-2.4162	0.9533	-2.0322
0.1662	-2.4235	0.3663	-2.3964	0.5613	-2.2595	0.7668	-2.4147	0.9668	-2.0114
0.1793	-2.4497	0.3798	-2.3754	0.5802	-2.2774	0.7802	-2.4119	0.9798	-2.0075
0.1934	-2.4649	0.3932	-2.3586	0.5938	-2.2926	0.7937	-2.4029	0.9933	-2.0053

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