

Light element abundances in the young open clusters NGC 3293, NGC 4755 and NGC 6231: Tracers of stellar evolution^{*}

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Abstract. The abundances of He (LTE), C, N, and O (NLTE) were derived for 21 B stars in three young open clusters. Almost all the stars show subsolar CNO abundances. However, the mean oxygen abundance for each programme cluster appears to be in marginal agreement with the most recent revisions of the solar value. After consideration of the CN abundances in this sample, there is no clear evidence of internal mixing. Only three stars among the non-supergiants seem to show a nitrogen enhancement. Two of them have a fairly low projected equatorial velocity (admittedly, they may be rapid rotators seen pole-on); the third one is a definite fast rotator. In the lower gravity stars (three stars in this sample with $\log g < 3.0$) some kind of mixing has apparently occurred. The supergiants do not differ significantly from the other programme stars in their respective helium contents. The mean helium abundance for each cluster is close to the standard value, $(\text{He}/\text{H}) \approx 11.0$.

Key words. open clusters and associations: individual: NGC 3293, NGC 4755, NGC 6231 – stars: abundances – stars: early type

1. Introduction

Several aspects of the intermediate stages of the evolution of massive stars are still poorly understood. In particular, little is known about the possible occurrence of internal mixing in these stars. In massive stars, this process may manifest itself observationally through the appearance at the stellar surface of material that has been mixed in the interior, hence of products of the CNO-cycle. The presence of such material can be diagnosed from determinations of the abundances of carbon, nitrogen and oxygen in early B stars near the end of their main-sequence lifetimes and slightly evolved from the main-sequence, in young open clusters. Constraints on the mixing models can be derived through investigation of possible dependencies of these abundances on stellar parameters.

Lyubimkov (1984, 1989a, 1989b) claimed that relationships exist between the abundances of He, C and N and stellar mass and age. More precisely, the abundances

of He and N increase with the stellar age, while carbon abundances decrease. The rate at which helium and nitrogen are enhanced is proportional to the stellar mass. The sum of C and N abundances remains constant, as expected for CN-processed material. On the other hand, Gies & Lambert (1992) carried out a comprehensive study of B supergiant and non-supergiant stars in order to check Lyubimkov's conclusions. They did not find clear evidence that nitrogen abundances progressively increase with stellar age. Nevertheless, they found evidence for the presence of CN-cycled material in a few main-sequence stars. The atmospheres of all supergiant stars appeared to be enriched in nitrogen.

Andrievsky et al. (1999) derived the abundances of carbon and nitrogen in hot main-sequence stars using NLTE calculations. For almost all stars the C and N abundances appeared to be subsolar, but with a solar C/N ratio, while for three stars a carbon deficiency was detected together with a nitrogen enhancement. This could be the signature of effective mixing in these three stars.

Recently, a thorough theoretical study on rotationally induced internal mixing and changes in the surface abundances of massive stars was performed by

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^{*} Based on spectra collected at the European Southern Observatory, La Silla, Chile, as part of programme 47.7-045.

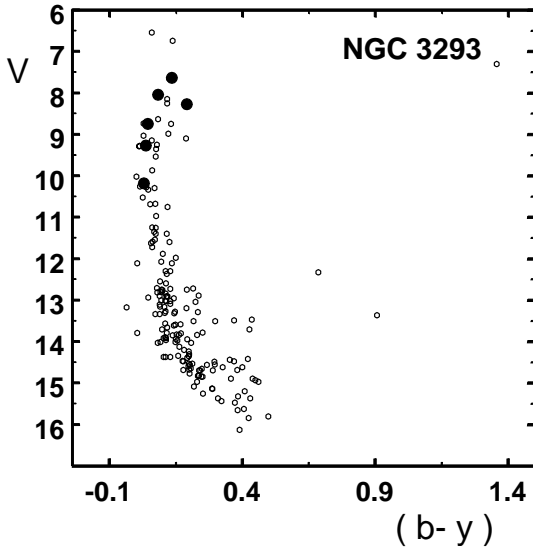


Fig. 1. Colour-magnitude diagram (CMD) for NGC 3293. Programme stars are shown as *filled black circles*. The photometric data are from Balona (1994).

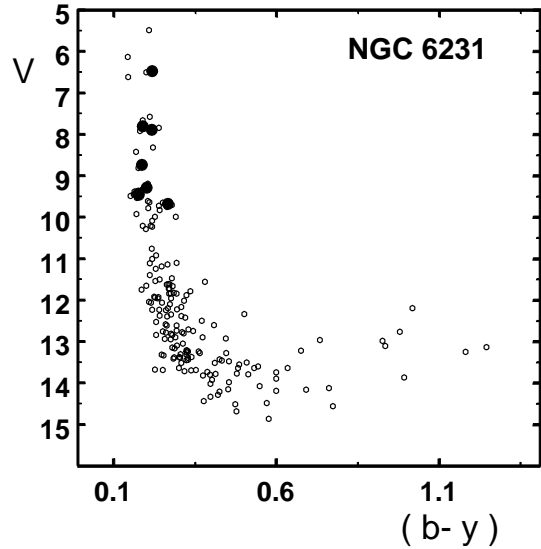


Fig. 3. CMD for NGC 6231. The photometric data are from Balona & Laney (1995).

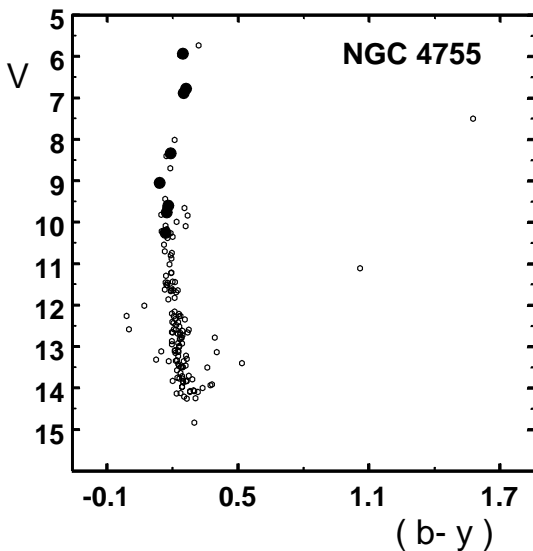


Fig. 2. CMD for NGC 4755. The photometric data are from Balona & Koen (1994).

Heger & Langer (2000). According to these authors rotational mixing is responsible for abundance anomalies found in stars with masses $10\text{--}20 M_{\odot}$.

It should be noted that the CNO abundance anomalies, similar to those produced by rotational mixing, can also be caused by the mass transfer in close binary systems when the atmosphere of the accreting star is contaminated with the companion's material processed in the incomplete CNO cycle of hydrogen burning.

In this work, we present NLTE abundance analyses of a sample of hot stars in three young open clusters: NGC 3293, NGC 4755 and NGC 6231. In Sect. 2, the observations are described. The methods used for the analysis of the data are presented in Sect. 3. In Sects. 4, 5 and 6, the adopted stellar parameters, the derived elemental

abundances, and their uncertainties are reported. Results are discussed in Sect. 7.

2. Observations

This study is based on observations carried out with the 3.6 m telescope of the European Southern Observatory at La Silla, Chile. Observations were carried out in May 1991 using the ESO Cassegrain Echelle Spectrograph CASPEC with the 52 lines mm^{-1} échelle grating and the blue-optimized cross-disperser grating. The detector was CCD #16, a UV-flooded Tektronix with 512×512 pixels of $27 \mu\text{m}$. The spectra cover the range $3900\text{--}5100 \text{ \AA}$ at a resolving power of 18 000. The reductions have been performed using the MIDAS echelle package. Details of the observations are given in Table 1. Column 3 gives the heliocentric Julian date of mid-observation, and Col. 4, the total exposure time. The S/N ratio per pixel given in the last column of the table has been computed from the measured rms in the continuum of the fully reduced spectrum.

The sample stars all appear in Mermilliod's (1999) lists of cluster members. This should ensure that, to the best of our present knowledge, all do indeed belong to their respective cluster. For NGC 4755, membership probabilities based on proper motions have actually been published by King (1981). This probability is above 90% for 6 of the 8 stars studied here, while it is 54% for NGC 4755-2 and 58% for NGC 4755-18. The programme stars have been selected so as to sample the widest possible range of masses at the top of the cluster sequence, taking into account the limitations set by the efficiency of the telescope-instrument combination. Their locations in the colour-magnitude diagrams of the studied clusters are shown in Figs. 1 to 3.

Table 1. Log of observations.

Star	<i>V</i>	HJD	Exp. (min)	<i>S/N</i>
NGC 3293				
8	8.27	2448399.548	30	230
13	10.18	2448400.530	105	270
16	8.74	2448399.509	42	240
18	9.27	2448400.588	45	260
20	8.04	2448400.465	12	180
22	7.63	2448399.471	18	100
NGC 4755				
2	5.93	2448398.565	12	180
3	6.78	2448398.590	18	190
5	8.33	2448398.622	36	260
6	9.05	2448399.607	80	190
7	9.76	2448398.716	160	130
18	9.60	2448399.694	105	180
106	6.88	2448400.714	9	220
414	10.25	2448400.670	120	170
NGC 6231				
161	7.89	2448400.871	18	200
232	9.67	2448399.793	120	200
286	9.43	2448399.870	75	220
287	9.28	2448398.817	90	200
292	8.73	2448398.873	45	180
293	6.47	2448398.891	6	210
323	7.80	2448398.909	24	170

3. Lines selected for analysis and calculations

The great majority of the programme stars have reasonably sharp spectral lines that lend themselves to direct equivalent width measurements. From lines present in overlapping orders, we estimated the upper limit of the uncertainty in the measured equivalent widths to be approximately 10%. The list of lines used in the analysis and their equivalent widths are gathered in Tables A.1, A.2 and A.3.

To represent the atmospheric structure of the studied stars, we used plane-parallel, LTE models interpolated from the grid of Kurucz (1992). The latter have been computed for three values of the microturbulence: 2, 4, and 8 km s⁻¹. The model corresponding to the value of the microturbulence closest to the value derived from the present abundance study (see Sect. 4.2) was adopted whenever the microturbulence did not exceed 8 km s⁻¹. For stars with larger microturbulence, grid models calculated with the maximum value of this parameter, 8 km s⁻¹, were used. Some models (high temperature/low gravity) are not available in the grid, so they were computed using the ATLAS9 code.

3.1. LTE determination of the helium abundance

Helium abundances were determined through LTE analysis of He I lines. For this purpose, we employed the WIDTH9 code of Kurucz. The oscillator strengths and damping constants were taken from CD-ROM 23 (Kurucz 1994).

3.2. Non-LTE determination of the CNO abundances

An updated version of Carlsson's (1986) NLTE code MULTI was used for the determination of elemental abundances of carbon, nitrogen and oxygen. The modifications of the code, atomic models and atomic level characteristics are given in detail in Korotin et al. (1999a, 1999b), Andrievsky et al. (1999), and Korotin et al. (1999c), as summarized below:

1. The atomic model of carbon used in the calculations includes 71 levels of C II, 12 levels of C III, and the ground level of the C IV. The radiative transitions between the first 38 levels of C II, three C III levels, and the ground level of C IV were considered.
2. The model of the nitrogen atom consists of 109 levels: 3 ground levels of N I, 93 levels of N II with $L \leq 4$ and $n \leq 6$, 12 levels of N III, and the ground state of N IV. Within the system of the nitrogen atom levels we considered the radiative transitions between the first 43 levels of N II, 5 N III levels, and the ground level of N IV.
3. The oxygen atomic model comprises 141 levels: 3 levels of O I, 132 levels of O II with $L \leq 5$ and $n \leq 8$, 5 levels of O III, and the ground state of O IV. In this system, we considered the radiative transitions between the first 49 levels of O II and the ground level of O III.

Our version of the MULTI code allows one to analyse the blends of lines of the same element. This option was used in cases where no single carbon lines are available in the spectrum (e.g., for NGC 4755-5).

4. Stellar parameters

4.1. Effective temperatures and gravities of the programme stars

Values of the effective temperature and gravity of the programme stars were derived from Strömgren photometry using the code written by T. T. Moon (based on the grid published in Moon & Dworetzky 1985) as modified by Napiwotzki (1994). Complete information on the photometric data for the programme stars can be found in the SIMBAD database, or the WEBDA database (Mermilliod 1999). As can be seen in SIMBAD, many of the studied stars are suspected to be variable. To minimize the uncertainty introduced in temperature and gravity determinations by photometric variability, mean indices were determined using all the Strömgren colour index measurements available in the literature. In this averaging process, weights were assigned to the indices coming from different sources on the basis of the number of measurements of the considered index in that reference. For a few stars, data from various authors show rather significant differences. The resulting uncertainties in the derived atmospheric parameters will be further discussed in Sect. 6. The adopted values of the mean photometric indices, and the resulting temperatures and gravities are given in Table 2.

Note that there exist indications that some of our program stars may be binary systems. For example, in

Table 2. Parameters of programme stars.

Star	DM ^a	Spectral type	(<i>b</i> - <i>y</i>)	<i>m</i> ₁	<i>c</i> ₁	β	<i>T</i> _{eff} (K)	log <i>g</i>	ξ_t (km s ⁻¹)	<i>v</i> sin <i>i</i> (km s ⁻¹)
NGC 3293 ^b										
8	-57° 3540	B0.5 III	0.192	-0.018	-0.019	2.570	26 000	3.60	13	31
13	-57° 3522	B2 V	0.030	0.069	0.050	2.620	24 350	4.15	3	24
16	-57° 3500	B0.5 V	0.045	0.047	0.015	2.592	25 000	3.75	12	33
18	-57° 3524	B1 V	0.038	0.052	0.038	2.606	24 500	3.95	6	28
20	-57° 3523	B1 III	0.084	0.018	-0.023	2.575	25 800	3.60	15	115
22	-57° 3506	B1 II	0.135	0.012	-0.021	2.564	24 900	3.30	15	55
NGC 4755 ^c										
2	-59° 4555	B3 Ia	0.248	-0.030	0.130	2.551	15 750	2.00	18	47
3	-59° 4566	B2 Ib	0.263	-0.054	0.125	2.568	17 100	2.40	18	47
5	-59° 4552	B1 III	0.192	-0.024	0.027	2.588	25 500	3.75	16	85
6	-59° 4564	B2 III	0.142	0.003	0.062	2.605	24 400	3.90	8	95
7	-59° 4528	B0.5 V	0.174	-0.007	0.048	2.601	25 100	3.90	6	37
18	-59° 4553	B2 IV	0.182	0.006	0.133	2.598	21 500	3.40	0	30
106	-59° 4543	B1.5 Ib	0.251	-0.055	0.025	2.559	21 350	2.80	18	47
414	-59° 4535	B1 V	0.169	0.022	0.139	2.632	22 350	4.10	2	67
NGC 6231 ^d										
161	-41° 11050	O8.5 III	0.217	-0.031	-0.032	2.585	30 000	4.20	20	41
232	521	B0.5 V	0.267	-0.046	0.055	2.629	26 470	4.60	8	45
286	378	B0.5 V	0.176	0.009	0.073	2.627	24 850	4.35	12	42
287	403	B1 V	0.201	0.004	0.089	2.635	24 300	4.45	2	8
292	-41° 11038	O9.5 V	0.187	-0.016	-0.046	2.596	31 700	4.60	20	83
293	-41° 11036	O9 Iab p	0.219	-0.025	-0.097	2.552	32 400	4.00	20	42
323 ^e	862	O9 III	0.190	-0.024	-0.068	2.587	33 370	4.60	20	52

^a For those stars of NGC 6231 that do not have a DM number, the WEBDA number (Mermilliod 1999) is given as second identifier.

^b Star numbering from Turner et al. (1980); MK types from Feast (1958), except for star 16 (Morgan et al. 1955).

^c Star numbering from Arp & van Sant (1958); MK types from Schild (1970), except for stars 2 (Walborn 1972), 7 and 414 (Feast 1963).

^d Star numbering from Seggewiss (1968); MK types from Levato & Malaroda (1980).

^e Star 12 of Struve (1944)

a recent study, Garcia & Mermilliod (2001) report that NGC 6231-161, 232, 286 and 287 are definite spectroscopic binaries; stars 161 and 232 probably have double-lined spectra. The same authors also conclude that the binarity percentage among the OB stars of NGC 6231 is very high, more than 82%. Binarity can cause some ambiguity in photometric determinations of the stellar parameters, if the contribution of the secondary to the observed combined flux is significant (as is the case in particular for SB2s).

Another potential source of error in the estimation of the stellar parameters from Strömgren photometry is the possible presence of emission in the H β line, which could affect the β -index. For instance, we note that star NGC 4755-6 is reported as Be in the WEBDA database. It is quite possible that other targets could also be unrecognized Be stars. Therefore, the profile of H β in all the program stars was checked for appearance of emission features. The line seems to be purely in absorption in all stars, with the exception of NGC 6231-161, where it shows a small redshifted emission component.

In all cases where a sufficient number of O II and N II lines could be analyzed, the photometric *T*_{eff} determination was refined spectroscopically, by requiring that the abundances derived from individual O II and N II lines do not depend on their lower level excitation potentials. As an example, in Fig. 4 we show such a plot for O II lines of the star NGC 3293-8.

4.2. Microturbulent velocities

The microturbulent velocities ξ_t for the programme stars were determined by requiring NLTE oxygen and nitrogen abundances derived from individual O II and N II lines to be independent from their equivalent widths. The number of O II lines measured in the spectra is sufficient for a reliable determination of ξ_t in all studied stars except for NGC 6231-161, 293, and 323. For the latter, the value of ξ_t estimated for NGC 6231-292 is adopted, based on the similarity of their spectral types. For some stars, the number of N II lines is also sufficient to derive a value of the microturbulence. In all cases, the results obtained from

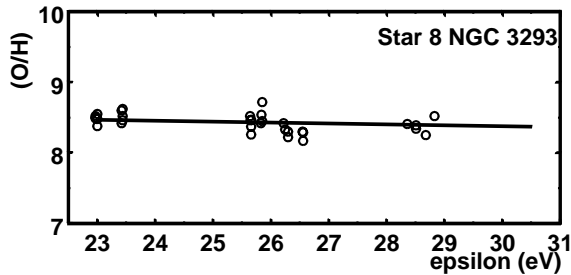


Fig. 4. Oxygen abundance from individual O II lines vs. their excitation potentials.

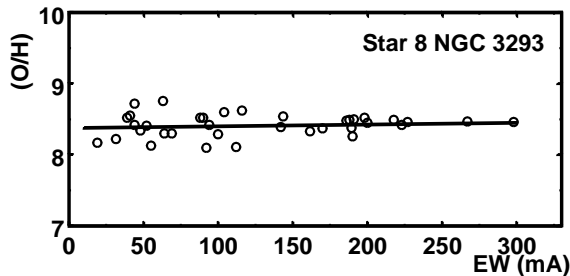


Fig. 5. Oxygen abundance from individual lines vs. their equivalent widths: determination of the microturbulent velocity.

consideration of the N II and O II lines are mutually consistent. In Fig. 5 we show an example of microturbulence velocity determination. A visible dependence between the abundance results from individual O II/ N II lines on their equivalent widths appear when the microturbulent velocity is varied within 10–15% of the adopted value.

Results of the NLTE ξ_t determinations are given in Table 2. As one can see, for some stars, the derived microturbulent velocities are rather high. Let us briefly discuss the high values found for B supergiants and O stars in the sample.

In several studies, it was found that some additional significant nonthermal broadening is required to explain observed line widths in the spectra of early B supergiant stars. For example, values in the range of 20–35 km s⁻¹ were reported by Lamers (1972), Van Helden (1972), Lennon & Dufton (1986), Gies & Lambert (1992), Smartt et al. (1997), McErlean et al. (1998), Vrancken et al. (2000). Such velocities are of the order of magnitude of, or higher than, the thermal velocity of the hydrogen atoms in the atmospheres of these stars. For main-sequence early B stars, high microturbulent velocities were also derived by e.g. Gies & Lambert (1992).

It is known that the atmospheric structure of hot and luminous stars is significantly complicated by the presence of large stellar winds. These winds are characterized by strong high-velocity turbulence (shock-wave perturbed winds). As shown by Groenewegen & Lamers (1989), such a turbulence in A-B-O supergiant stars and O dwarfs can be described by a ξ_{turb} parameter which equals approximately 10% of the wind terminal velocity. On average, O stars (both main-sequence and supergiants) have higher terminal wind velocities, therefore a stronger turbulence can be expected for them (even for A supergiants velocities

of about 20 km s⁻¹ are found; see McCarthy et al. 1997). As wind turbulence and a significant velocity gradient in the atmosphere (caused by the wind) can, to some extent, be imprinted on the profiles of the strong photospheric lines in the spectra of hot luminous stars, rather high microturbulent parameters are derived (see Kudritzki 1992; Lamers & Achmad 1994; and discussion in McErlean et al. 1998).

Admittedly, one could artificially decrease the microturbulent velocities of the programme stars by excluding the strong lines (say, with equivalent widths of about 300 mÅ). Nevertheless, this is not justified since these lines are appropriate for NLTE analysis: thus they are included in the present study. However, in agreement with Dufton (1972) and Smartt et al. (1997), we do not attach a great physical meaning to the determined microturbulent velocities, and we consider the ξ_t value as an empirical fitting parameter to avoid any systematic trend in the [O/H] or [N/H] vs. equivalent width relations.

4.3. Projected rotational velocities

In order to estimate the projected rotational velocities for the programme stars we used the SYNSPEC code (Hubeny et al. 1994). Determination of $v \sin i$ values in the studied stars was carried out by fitting the synthetic and observed profiles of the line He I 4437 Å. This line was used because:

1. it is a single, unblended line;
2. it is not very strong in any of the stars, and its profile is not expected to be sensitive to NLTE effects;
3. it is also quite insensitive to the microturbulent velocity (see, e.g., McErlean et al. 1998).

For all the stars the fit was performed using the code SYNSPEC. In Fig. 6 we show some examples of the profile fitting. Estimated projected rotational velocities are presented in Table 2. The estimated error in the $v \sin i$ determinations is about 20% of the adopted value. Recently, Garcia & Mermilliod (2001) presented projected rotational velocities for stars in NGC 6231. Our results for the stars of this study are in marginal agreement with their determinations, but in average slightly smaller.

5. Elemental abundances

Elemental abundances were determined using all available CNO lines with equivalent widths up to 300 mÅ. The results are presented in Table 3. Some examples of observed and NLTE synthetic profiles for selected CNO lines are shown in Figs. 7–9. A comparison between NLTE and LTE abundances (which were also calculated, but not presented in this paper), shows that for the majority of the programme stars the differences between NLTE and LTE are typically about 0.1–0.2 dex, although in a few cases these differences exceed 0.5 dex.

Helium abundances were determined from the line He I 4437 Å. This absorption line is not strong in any

Table 3. NLTE abundances of CNO elements and LTE abundance of He.

Star	(C/H)	σ	(N/H)	σ	(O/H)	σ	(He/H)	σ
NGC 3293								
8	8.22	0.11	7.78	0.15	8.41	0.15	11.20	
13	8.38	0.10	7.81	0.13	8.65	0.15	11.12	
16	8.10	0.14	7.71	0.16	8.42	0.15	11.13	
18	8.16	0.15	7.69	0.16	8.65	0.15	11.07	
20	8.15	–	7.92	0.12	8.52	0.17	11.25	
22	8.18	–	7.68	0.13	8.34	0.15	11.25	
Mean	8.20	0.10	7.77	0.10	8.50	0.13	11.15	0.07
NGC 4755								
2	7.84	0.07	7.90	0.17	8.10	0.12	11.40	
3	7.70	0.02	7.97	0.12	8.14	0.16	11.25	
5	8.30	–	7.94	0.05	8.36	0.18	11.15	
6	8.65	–	8.04	0.32	8.60	0.16	11.15	
7	8.11	0.08	8.13	0.19	8.42	0.28	11.15	
18	8.01	0.18	7.46	0.15	8.30	0.25	10.95	
106	8.00	0.13	8.09	0.16	8.12	0.18	11.30	
414	7.95	–	7.83	0.15	8.86	0.17	11.10	
Mean	8.07	0.29	7.92	0.21	8.36	0.26	11.18	0.14
NGC 6231								
161	8.30	0.15	7.95	–	7.76	0.13	10.50	
232	8.48	0.17	7.98	0.29	8.63	0.18	11.00	
286	8.43	0.17	7.79	0.30	8.30	0.21	10.80	
287	8.13	0.25	7.75	0.17	8.92	0.16	10.73	
292	8.09	–	7.80	–	8.48	0.15	11.03	
293	–	–	–	–	8.20	0.04	10.63	
323	–	–	–	–	7.81	0.02	10.95	
Mean	8.29	0.17	7.85	0.10	8.30	0.42	10.81	0.20
Sun ^a	8.55		7.97		8.87		10.99	

^a Solar abundances are from Grevesse et al. (1996).

of the programme stars. It is not seen in the spectrum of NGC 6231-323, where two other helium lines were used: 4009 Å and 4143 Å.

One may question if LTE analysis of helium lines yields reliable abundances of this element. This question has been addressed in the literature. The most comprehensive NLTE studies of helium lines are those by Auer & Mihalas (1973) and Dimitrov & Kubat (1988). In both studies, the authors investigated departures from LTE (specifically the behaviour of “*b*” factors) for helium lines in a wide spectral region. They concluded that over a wide range of effective temperatures and gravities, NLTE corrections are negligible for the blue-violet spectral region, but that they become very significant in the red-infrared. In particular, for the line HeI 4437 Å, the difference between NLTE and LTE equivalent widths reaches only 5%. Thus, it is comparable with the error of our equivalent width measurements, and it cannot significantly affect the helium

abundances that we infer. Accordingly, the latter appear reliable.

6. Uncertainties

As mentioned above, the photometric indices for a given star reported by different authors can, in some cases, differ rather markedly. This may be partly caused by differences in detector sensitivity (see remarks in Balona 1994). But most likely, actual stellar variability is a significant contribution. As a matter of fact, almost all our programme stars are known to be variable (for example, stars 16, 18, 20 in NGC 3293, stars 2, 6, 7, 18, 106 in NGC 4755, stars 161, 232, 282, 292, 293 in NGC 6231). Even if a star is not reported as a variable, its photometric variability can be suspected from the existence of rather large radial velocity variations (SIMBAD database).

The influence of errors in the adopted atmospheric parameters on the derived abundances is illustrated for the

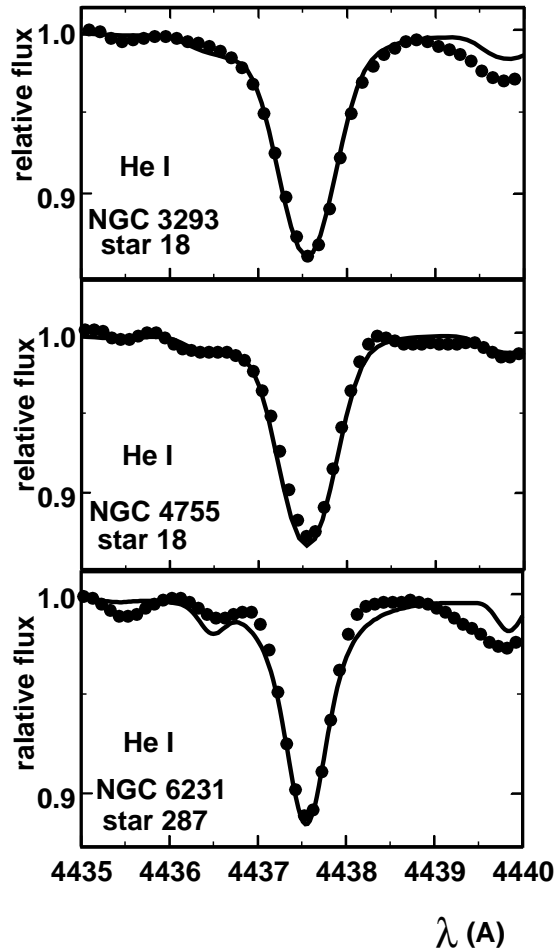


Fig. 6. Observed and calculated profiles of the line He I 4437 Å for some programme stars.

case of star NGC 3293-16. Based on published photometric measurements of this star, we obtain: $(b - y) = 0.045 \pm 0.012$, $m_1 = 0.047 \pm 0.022$, $c_1 = 0.015 \pm 0.016$, and $\beta = 2.592 \pm 0.002$. In the simplest case (all the indices are either simultaneously increased, or simultaneously decreased), we derive the following parameter uncertainties: $T_{\text{eff}} \approx 24\,400\text{ K} - 25\,800\text{ K}$, $\log g \approx 3.7 - 3.8$. With both the temperature and the gravity at the lower limit of their uncertainty range, the derived NLTE abundances are:

$$\begin{aligned} (\text{C}/\text{H}) &= 8.06 \pm 0.16, \\ (\text{N}/\text{H}) &= 7.67 \pm 0.16, \\ (\text{O}/\text{H}) &= 8.44 \pm 0.14. \end{aligned}$$

With the highest temperature and highest gravity, the abundances are:

$$\begin{aligned} (\text{C}/\text{H}) &= 8.15 \pm 0.13, \\ (\text{N}/\text{H}) &= 7.77 \pm 0.17, \\ (\text{O}/\text{H}) &= 8.40 \pm 0.15. \end{aligned}$$

Note that with the adopted parameters for this star ($T_{\text{eff}} = 25\,000\text{ K}$, $\log g = 3.75$), we get:

$$\begin{aligned} (\text{C}/\text{H}) &= 8.10 \pm 0.15, \\ (\text{N}/\text{H}) &= 7.71 \pm 0.16, \\ (\text{O}/\text{H}) &= 8.42 \pm 0.15. \end{aligned}$$

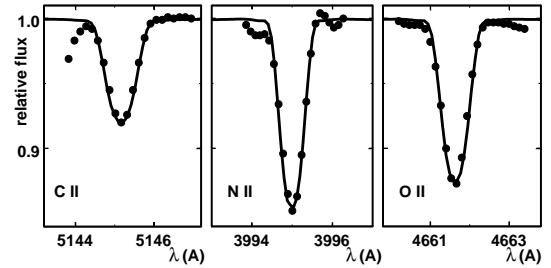


Fig. 7. Observed and calculated profiles of selected CNO lines for star NGC 3293-13.

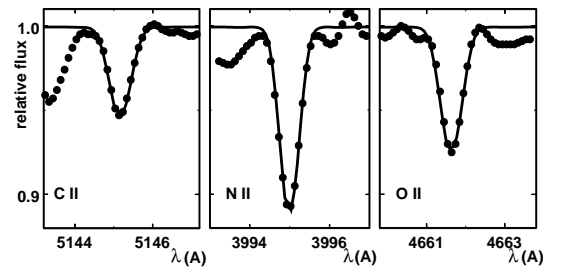


Fig. 8. Same as Fig. 7 but for star NGC 4755-18.

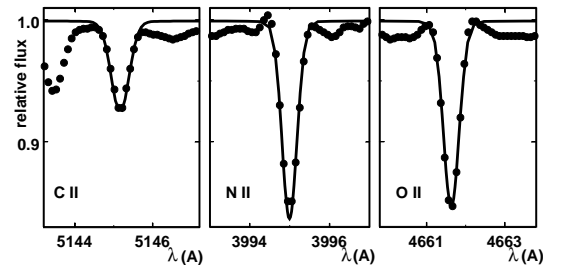


Fig. 9. Same as Fig. 7 but for star NGC 6231-287.

For some of the studied stars, another important source of uncertainty comes from the fact that they belong to double-lined binary systems (see Sect. 4). In such systems, the observed spectrum results from the superposition of the individual spectra of the two components. Our observational material does not allow us to separate the contributions of each of them. If the lines of the components could be distinguished, their equivalent widths would be reduced due to the dilution of the continuum of one star by that of the other. In the present case, the lines appear single in all the stars observed, so that they must be contributed by both members of the binary. Accordingly, our analysis is self-consistent, in that the abundances are derived from the same combination of the contributions of the two components as the stellar atmospheric parameters. The “average” abundances that are determined in that way can be expected to represent reasonably well the composition of the two stars, provided that they are sufficiently similar, in particular in mass and evolutionary state. This seems rather likely for members of such early type SB2s – any significant mass difference between the two members of a pair would almost unavoidably imply such a difference of luminosity that the less massive one

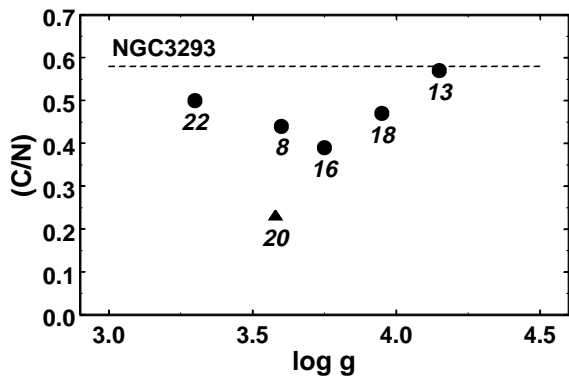


Fig. 10. (C/N) ratio vs. surface gravity for the stars of NGC 3293. *Circles* identify the stars with projected rotational velocity less than 55 km s^{-1} , *triangles* represent the stars with $v \sin i > 95 \text{ km s}^{-1}$. The solar (C/N) ratio is shown by the dashed line.

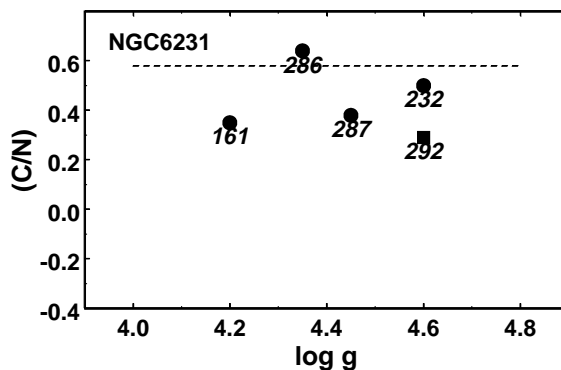


Fig. 12. (C/N) ratio vs. surface gravity for the stars of NGC 6231.

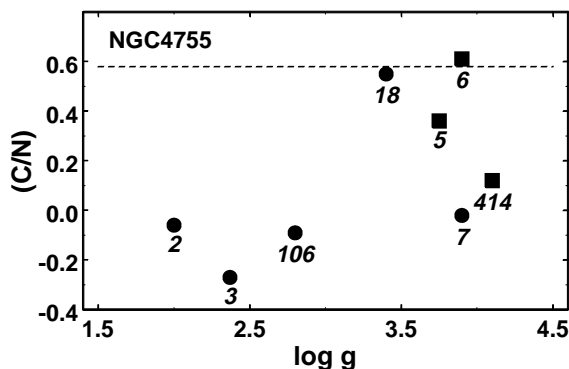


Fig. 11. (C/N) ratio vs. surface gravity for the stars of NGC 4755. *Squares* correspond to the stars with projected rotational velocity between 55 km s^{-1} and 95 km s^{-1} .

would not contribute significantly to the observed colours and spectrum. In other words, although they cannot be exactly quantified from the analyzed spectra, uncertainties of abundances determinations in the SB2 systems of the sample are expected to be small. However, their possible influence should be kept in mind in the interpretation of the results.

7. Discussion

7.1. Helium abundance

Ahumada & Lapasset (1995) give the ages of NGC 3293 ($\log A = 7.40$), NGC 4755 ($\log A = 6.85$), and NGC 6231 ($\log A = 6.50$). From Table 3 one can infer a mean helium abundance for NGC 3293 of $(\text{He}/\text{H}) = 11.15 \pm 0.07$; for NGC 4755, it is $(\text{He}/\text{H}) = 11.18 \pm 0.14$, and for NGC 6231, $(\text{He}/\text{H}) = 10.81 \pm 0.20$. The surface abundance of helium in the youngest of the three studied clusters, NGC 6231, may be marginally less than in the two older ones. Nevertheless, the difference is too small to derive any evolutionary conclusion. The supergiants of our sample do not differ significantly in their helium abundance from other programme stars, but they may show a hint of helium enhancement.

7.2. Carbon, nitrogen, and mixing

In Figs. 10 to 12, the ratio (C/N) is plotted against the surface gravity of the sample stars. If B-type main-sequence stars undergo large-scale internal mixing, the consequences should show through in the ratio of the surface abundances of carbon to nitrogen. At first sight, Figs. 10 to 12 may appear to support the view that CN-cycled material can appear at the surface of some B stars during their main-sequence lifetime. For instance, this is suggested by the scatter of the (C/N) values among the stars of NGC 6231, which all have large surface gravities. Nevertheless, closer inspection of Table 3 reveals that the main contribution to low (C/N) values comes from lower carbon abundances, while nitrogen remains nearly solar in almost all cases. Even for the two stars of NGC 4755 with the lowest gravities (stars 2 and 3), where mixing processes might be expected to modify the surface abundances, nitrogen appears to be solar. One possible interpretation of this puzzling result is to suppose that in some stars, present-day C and N surface abundances are really altered by CNO-cycle processing, but that the initial abundances of these elements were subsolar.

The problem of subsolar metallicities in B stars has been widely debated recently. This feature of the chemical composition of B stars has been found in virtually all studies (see, e.g., Gies & Lambert 1992; Cunha & Lambert 1994; Daflon et al. 1999; Andrievsky et al. 1999; Rolleston et al. 2000). Subsolar abundances were reported for C-N-O elements, as well as for heavier species. Hou et al. (2000) conclude that at present there is no satisfactory explanation for the supermetallicity of the Sun relative to young objects (specifically B stars) in the solar neighbourhood. Luck et al. (2000) proposed that conclusions about subsolar metallicities may suffer from ignoring effects such as atomic diffusion in atmospheres of B main-sequence stars. Another explanation was recently proposed by Snow (2000). Such processes as sedimentation and ambipolar diffusion during stellar formation could lead to some decrease in the abundances of some elements.

Possibly, the long-standing problem of subsolar C and N abundances in B stars may eventually find an even simpler solution, as recently happened with oxygen after

revision of its abundance in the Sun. A recent work of Holweger (2001) based on optical and infra-red lines placed the solar oxygen abundance at $(O/H) = 8.74 \pm 0.08$, that is, significantly lower than accepted before. An even lower oxygen abundance in the Sun was recently deduced by Prieto et al. (2001). By applying a 3-D time-dependent hydrodynamical model of the solar atmosphere, those authors re-analyzed the forbidden line $[O\ I] 6300 \text{ \AA}$, finding that the oxygen abundance is 8.69 ± 0.05 . In fact, the mean oxygen abundances that we derive in the studied clusters are now, within errors, in marginal agreement with the latter value.

As one can see from Figs. 10 to 12, the main-sequence stars (i.e., those having gravities higher than, say, 3.75 dex) are mostly near the line corresponding to atmospheric C-N abundances that have not yet been altered through mixing with CN-processed material (stars 7 and 414 in NGC 4755 are two obvious exceptions). Among the group of stars with $\log g > 3.75$, there are stars with nearly solar abundances of carbon and nitrogen, as well as stars with decreased abundances of both carbon and nitrogen. If the net result of the competing effects of upward radiative forces and downward gravitational settling is similar for carbon and nitrogen ions, selective radiative diffusion could, in principle, lead to a comparable dilution of C and N in the atmospheres of B stars. This process can be effective only if the rotational velocity is not too high (but an upper limit has not been estimated for B stars). For example, star NGC 4755-6, with $v \sin i = 95 \text{ km s}^{-1}$, has solar C and N abundances, while star NGC 3293-22, with a more moderate rotational velocity ($v \sin i = 55 \text{ km s}^{-1}$) shows both decreased carbon and nitrogen. The latter case may indicate that rotational velocities as high as 50–100 km s^{-1} do not effectively prevent radiative diffusion in the atmosphere.

Several stellar evolutionary grids have been computed with particular attention to the evolution of the surface abundances (e.g., Maeder & Meynet 1988; Schaller et al. 1992). In these tables, in the models with masses 9–12 M_{\odot} (which correspond to the main-sequence stars of the present sample), one can note that the first change of the surface abundances of carbon and nitrogen happens when the model reaches the red giant region. Contrary to this prediction of standard stellar evolution, in this study, CN-processed material seems to appear at the stellar surface for hot supergiants (stars 2, 3, 106 in NGC 4755). These stars could have already passed the red giant stage. A possibility to create a N-rich main-sequence star through rotationally induced mixing was considered by Maeder (1987). The author succeeded in showing that this effect can modify the surface abundances of a 25 M_{\odot} rapidly rotating star. Whether this should be the case for hot supergiant stars with lower masses (and rather small projected rotation velocities) is not clear.

Very recently Heger & Langer (2000) performed a theoretical study considering the evolution of surface abundances for rotating massive stars. The authors found that the abundance ratios (C/N), (C/O) and (N/O)

significantly depend upon the rotational velocity. For example, after approximately 10 Myr (the mean age of our sample of cluster) surface [C/N] ratio should be about -0.3 dex if the star rotates with equatorial velocity 200 km s^{-1} . From Figs. 10–12 one can see that the great majority of our non-supergiant stars show a similar value of (C/N), but we should note that they seem to rotate slowly for their spectral types (although some of them can be pole-on rotators), and for rotational velocities of about 100 km s^{-1} a negligible modification on the surface CNO abundances is expected (see Heger & Langer’s Fig. 9). Also, according to Howarth & Smith (2001) the only very rapid rotators are nitrogen enhanced.

On the other hand, Heger & Langer (2000) predict that rotation should decrease surface (C/O) ratio. In fact, the typical (C/O) value for our program stars is about the same as for the Sun (but with significant deviations, both positive and negative, for some stars), therefore the comparison between our determinations and the theoretical predictions does not seem to support the supposition that surface abundances in massive young B stars are rotationally altered.

8. Conclusion

The C-N abundances determined for the sample of stars considered in this study do not show clear evidence for the occurrence of internal mixing for most main-sequence stars. Only three non-supergiant stars, NGC 4755-7, NGC 4755-414, and NGC 3293-20 seem to show an enhanced nitrogen abundance together with decreased carbon. Two of them do not have a large projected rotational velocity (but admittedly, they might be rapid rotators seen almost pole-on). In the lower gravity stars (stars 2, 3 and 106 in NGC 4755), CNO-processed material appears to be present at the surface, suggesting that some kind of mixing has occurred. This conclusion is based on the assumption that initial abundances of carbon and nitrogen in these stars were subsolar. The problem of subsolar abundances itself is not yet resolved and deserves further attention. If the hypothesis of initial subsolar abundances is not valid, then one needs to identify a mechanism causing an appreciable carbon depletion in B stars (supergiants, in particular), while nitrogen remains nearly unchanged.

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Appendix A: Equivalent width measurements

Tables A.1, A.2 and A.3 give the equivalent widths of the spectral lines measured in the spectra of the programme stars of, resp., NGC 3293, NGC 4755, and NGC 6231.

Table A.1. Equivalent widths of the lines in spectra of stars of NGC 3293.

Wave	code	log <i>gf</i>	E_{low} , eV	EW , mÅ					
				8	13	16	18	20	22
4437.55	2.00	-2.03	21.217	99.4	123.5	109.8	119.9	96.0	94.0
3920.68	6.01	-0.21	16.332	131.0	115.0	112.5	100.0		154.0
4074.54	6.01	0.39	24.369				49.0		
4374.28	6.01	0.66	24.652		41.0	22.0	34.8		
5139.17	6.01	-0.74	20.703		23.0				
5143.49	6.01	-0.24	20.703		35.8		37.0		
5145.16	6.01	0.16	20.709	69.0	67.8	57.0	49.0		
5151.09	6.01	-0.20	20.709	55.0	49.8	31.4	30.8		
3955.85	7.01	-0.78	18.465	38.0	39.8	29.0	47.0		
3995.00	7.01	0.28	18.496	135.0	95.5	125.6	110.0	161.0	158.0
4035.08	7.01	0.62	23.123	50.0	51.0	75.0	64.0		58.0
4041.31	7.01	0.85	23.141	52.0	51.9	75.0	48.0		58.0
4043.53	7.01	0.74	23.131	25.0	28.7	44.0	35.0		30.5
4082.27	7.01	0.15	23.131		11.5				
4171.59	7.01	0.28	23.195		7.8		14.2		
4176.16	7.01	0.60	23.195		27.0				
4199.98	7.01	0.03	23.245		19.7				
4227.74	7.01	-0.07	21.598	27.0	28.0	29.0	23.0		
4237.05	7.01	0.55	23.241	49.0	50.2	42.9	48.0		
4241.79	7.01	0.71	23.245	55.0	51.0		62.0		
4432.74	7.01	0.58	23.414		24.1	18.6	30.0		
4442.02	7.01	0.31	23.420		9.2				
4447.03	7.01	0.29	20.408	75.5	65.0	72.5	55.0		45.0
4530.41	7.01	0.67	23.473	43.0	34.0	44.0	35.0		
4607.15	7.01	-0.49	18.461	45.0	57.0	47.8	38.2		44.0
4613.87	7.01	-0.73	18.465	50.0	50.0	62.0	50.0	48.5	
4621.39	7.01	-0.54	18.465	40.0	33.8		17.2		
4630.54	7.01	0.13	18.482	127.0	96.2	110.0	90.9	151.0	113.0
4643.09	7.01	-0.34	18.482	59.0	61.0		55.0		
4788.14	7.01	-0.38	20.652			21.0			
4803.29	7.01	-0.12	20.664	50.0		54.7		63.4	
5002.70	7.01	-1.02	18.461		40.0				
5005.15	7.01	0.61	20.664	102.0	79.0	81.5	87.0		97.0
5007.33	7.01	0.16	20.939	39.0	34.0		46.7		30.0
5010.62	7.01	-0.52	18.465	41.0	33.5	26.3	49.0		37.0
5025.66	7.01	-0.44	20.664	16.0					
5045.10	7.01	-0.33	18.482	59.0	58.0	60.0	56.0		58.0
3945.04	8.01	-0.69	23.418	104.0		109.0	76.0	191.0	104.0
3982.72	8.01	-0.67	23.440	116.0		109.0	77.0	168.0	122.0
4048.24	8.01	-0.37	28.692	24.0	17.0				
4062.91	8.01	-0.09	28.704	43.0	22.1		31.6		
4071.24	8.01	-0.09	28.692		25.5				
4072.16	8.01	0.53	25.648	227.0	90.5	171.0	133.0		213.0
4075.87	8.01	0.69	25.664	267.0	118.0		160.0		
4078.86	8.01	-0.26	25.637	88.0	39.0	64.0	64.0		87.0
4083.93	8.01	0.15	28.682	62.0	34.0	59.0	60.0		
4085.12	8.01	-0.14	25.648		58.0	105.0	80.0		92.0
4087.20	8.01	0.53	28.675	55.0		48.1	55.0		
4092.93	8.01	-0.25	25.664			61.0	83.0		61.0
4110.78	8.01	-0.89	25.836	52.0	35.0		50.0		
4112.03	8.01	-0.78	25.847	66.0	27.0		48.0		
4119.22	8.01	0.48	25.847	200.0	105.5	212.0	145.0		234.0
4129.31	8.01	-1.12	25.836	37.0	23.0	17.4	26.0		
4132.79	8.01	-0.07	25.830	94.0	49.8	88.0	73.5	131.0	114.0
4153.28	8.01	0.08	25.836	143.5	74.0	134.0	106.0	155.0	127.0
4156.52	8.01	-0.79	25.847	44.0	26.0		38.0		
4185.44	8.01	0.71	28.356	92.0	44.5	85.0	70.5	188.0	88.0
4189.79	8.01	0.80	28.359	112.0	69.0	106.0	84.0	178.0	125.0
4275.53	8.01	0.76	28.856		41.0		66.5		
4291.27	8.01	-0.23	28.820	44.0	37.0	33.0	28.0		
4294.87	8.01	0.36	28.829				49.0		
4307.34	8.01	-0.05	28.837	40.0	20.0	18.5			
4317.15	8.01	-0.32	22.965	188.0	88.0	155.0	135.0	255.0	215.0
4319.64	8.01	-0.32	22.978	191.0	79.1	164.0	125.7	269.0	204.0
4325.76	8.01	-1.06	22.965		45.8	77.0	62.0	138.1	62.0
4366.91	8.01	-0.24	22.998	189.3	73.9	142.6	105.4	200.0	180.0
4369.31	8.01	-0.35	26.224	44.0	31.0	34.0	33.0		
4414.88	8.01	0.17	23.440	298.0	117.9	226.0	157.0		300.0
4416.97	8.01	0.04	23.418	223.0	100.0	178.0	135.0		236.0
4443.04	8.01	0.00	28.356	38.5	21.0	45.0	35.0		
4448.34	8.01	0.13	28.359	52.0	42.0	62.0			
4452.38	8.01	-0.74	23.440	90.0	55.0	90.8	70.3		83.3
4465.44	8.01	0.34	30.423	43.0	42.0		41.0		
4591.01	8.01	0.45	25.660	190.0	77.7	164.9	111.6	214.0	193.5
4596.20	8.01	0.29	25.660	170.0	76.0	139.2	109.0	186.0	163.0
4609.37	8.01	0.67	29.067	76.0	63.0		57.0		
4610.17	8.01	-0.17	29.061		27.0				
4638.86	8.01	-0.30	22.965	198.0	68.9	174.0	131.6	244.0	211.0
4649.14	8.01	0.31	22.998		143.0	301.0	189.0		
4661.64	8.01	-0.17	22.978	217.7	98.7	189.2	121.0	260.0	231.0
4673.75	8.01	-1.07	22.978	63.0					
4676.24	8.01	-0.30	22.998	186.0	82.0	169.0	97.0		
4696.36	8.01	-1.31	22.998	41.0	24.6		38.0		
4699.19	8.01	0.37	28.508	142.0		128.0		143.0	
4701.21	8.01	0.07	28.828	39.0	22.0	39.0	32.0		
4703.26	8.01	0.21	28.511	48.0	35.5	44.0			
4705.32	8.01	0.57	26.248	161.6	71.0	148.0	112.0	159.0	139.5
4709.99	8.01	-0.47	26.224	78.0	47.7	84.0	58.0		50.0
4890.85	8.01	-0.33	26.303	31.5	22.0		41.0		
4906.82	8.01	-0.03	26.303	68.9	33.0	63.0	59.0	80.0	64.0
4941.11	8.01	0.08	26.552	64.0	38.0	59.5	60.0		78.0
4943.00	8.01	0.37	26.560	100.0	46.2	86.4	82.0	140.0	115.0
4955.74	8.01	-0.42	26.560	19.1	22.2	26.3	19.9		

Table A.2. Equivalent widths of the lines in spectra of stars of NGC 4475.

Wave	code	log <i>gf</i>	<i>E</i> _{low} , eV	<i>EW</i> , mÅ							
				2	3	5	6	7	18	106	414
4437.55	2.00	-2.03	21.217	112.0	133.0	114.0	117.0	115.3	104.0	104.0	134.6
3918.97	6.01	-0.51	16.331						100.0		
3920.68	6.01	-0.21	16.332	184.5	199.0			117.0	90.5	155.0	102.0
4374.28	6.01	0.66	24.652					29.3			
5122.09	6.01	-0.53	20.844		24.3					24.5	
5143.49	6.01	-0.24	20.703							32.0	
5145.16	6.01	0.16	20.709					43.0	36.3	64.0	
5151.09	6.01	-0.20	20.709				70.0		26.0		
3955.85	7.01	-0.78	18.465	75.0	76.0			79.0			
3995.00	7.01	0.28	18.496	212.0	241.0	126.9	133.0		69.6	294.0	
4035.08	7.01	0.62	23.123	41.8	111.0	79.0	102.0	75.0	45.0	89.0	
4041.31	7.01	0.85	23.141	81.1				69.7	32.8	114.0	50.4
4043.53	7.01	0.74	23.131	52.0						61.0	
4082.27	7.01	0.15	23.131							20.4	
4171.59	7.01	0.28	23.195		18.6						
4199.98	7.01	0.03	23.245	19.6							
4227.74	7.01	-0.07	21.598		37.0					58.0	
4237.05	7.01	0.55	23.241		68.0		66.5	78.0	31.0		
4241.79	7.01	0.71	23.245		102.0					39.0	
4432.74	7.01	0.58	23.414							12.5	
4447.03	7.01	0.29	20.408	96.5	103.0			112.0	33.5	143.0	
4530.41	7.01	0.67	23.473		74.0			63.0	32.0		
4601.48	7.01	-0.37	18.465		120.0			123.0	40.0	150.0	
4607.15	7.01	-0.49	18.461		121.5			79.5	30.3	149.0	33.6
4613.87	7.01	-0.73	18.465	74.0	86.5			93.0	31.0	129.0	
4621.39	7.01	-0.54	18.465	54.0	82.5			86.0	42.0	124.0	
4630.54	7.01	0.13	18.482	199.0	240.1	201.0	140.0	126.0	57.0	318.0	67.2
4643.09	7.01	-0.34	18.482	123.0	158.5			91.0	36.4		
4779.72	7.01	-0.56	20.645				50.0	30.0		62.0	
4788.14	7.01	-0.38	20.652	52.7	46.4			38.7	12.4		
4803.29	7.01	-0.12	20.664		86.0					100.0	
5005.15	7.01	0.61	20.664	120.0	155.0	150.0			46.0	217.0	75.0
5007.33	7.01	0.16	20.939		74.0			53.4	30.0	93.0	
5010.62	7.01	-0.52	18.465		113.0				24.0	142.0	35.0
5045.10	7.01	-0.33	18.482	110.0	141.0			85.0	30.0	196.0	
3945.04	8.01	-0.69	23.418			162.0		73.5	50.0	93.0	
3982.72	8.01	-0.67	23.440	40.0					27.5	94.0	
4072.16	8.01	0.53	25.648	69.7	82.0			100.0	44.0		
4075.87	8.01	0.69	25.664						75.0		
4078.86	8.01	-0.26	25.637		26.1			65.0		60.0	
4087.20	8.01	0.53	28.675							25.1	35.0
4092.93	8.01	-0.25	25.664					64.0		41.0	
4119.22	8.01	0.48	25.847		91.0			168.0	66.0	170.0	
4132.79	8.01	-0.07	25.830	34.0					32.0	92.0	
4153.28	8.01	0.08	25.836	57.0	73.0	134.0	104.0		47.4	101.5	
4185.44	8.01	0.71	28.356	26.0	43.0	97.0	105.0	54.4	30.8	69.0	48.0
4189.79	8.01	0.80	28.359	27.0			106.0	70.0	36.0	88.9	
4317.15	8.01	-0.32	22.965	72.0	91.0	207.0	160.0	104.4	53.0	173.0	78.0
4319.64	8.01	-0.32	22.978	74.0	97.0	192.0		103.3	44.0	173.0	72.7
4325.76	8.01	-1.06	22.965	24.1	36.8			57.0	30.0	53.0	
4366.91	8.01	-0.24	22.998		59.0	162.0		99.0	49.0	183.0	
4414.88	8.01	0.17	23.440	88.0	127.0		213.0	136.0	61.5	257.0	77.7
4416.97	8.01	0.04	23.418	86.0	111.0			128.0	57.8	194.0	74.2
4452.38	8.01	-0.74	23.440	46.3	37.1	74.0		73.0	28.0	72.5	
4591.01	8.01	0.45	25.660	70.0	75.0	179.0		111.0	41.4	163.0	61.0
4596.20	8.01	0.29	25.660		62.0			104.0	50.0	128.7	
4609.37	8.01	0.67	29.067					95.0	29.0		46.3
4638.86	8.01	-0.30	22.965	78.0	87.5	235.0		104.0	49.5	185.0	70.0
4649.14	8.01	0.31	22.998	143.0	206.5			199.0	79.0		
4650.85	8.01	-0.36	22.965						55.0		
4661.64	8.01	-0.17	22.978	94.0	94.0	210.0	133.0	102.0	48.9	209.0	71.0
4673.75	8.01	-1.07	22.978						13.8	48.0	
4676.24	8.01	-0.30	22.998			217.0	143.0	105.0	31.9	164.0	
4696.36	8.01	-1.31	22.998					41.0			
4699.19	8.01	0.37	28.508	32.4		97.0				78.0	48.6
4703.26	8.01	0.21	28.511					45.0	20.0		
4705.32	8.01	0.57	26.248	57.0	57.5	119.0		102.0	47.8	104.0	
4709.99	8.01	-0.47	26.224			58.0		59.0	22.8	44.0	
4906.82	8.01	-0.03	26.303			96.0	82.0	52.0	16.3	64.0	
4941.11	8.01	0.08	26.552					54.0	23.5		
4943.00	8.01	0.37	26.560		45.0			63.0	30.5	68.0	

Table A.3. Equivalent widths of the lines in spectra of stars of NGC 6231.

Wave	code	log <i>gf</i>	$E_{\text{low,eV}}$	$EW, \text{m}\text{\AA}$						
				161	232	286	287	292	293	323
4009.26	2.00	-1.47	21.217							151.0
4143.76	2.00	-1.19	21.217							218.0
4437.55	2.00	-2.03	21.217	33.0	105.0	94.0	78.0	69.0	30.0	
4713.14	2.00	-1.23	20.963		249.3					
3920.68	6.01	-0.21	16.332	44.0	71.0	124.5	76.0	46.0		
4374.28	6.01	0.66	24.652				29.0			
5137.26	6.01	-0.93	20.700				7.8			
5139.17	6.01	-0.74	20.703				4.5			
5143.49	6.01	-0.24	20.703			41.0	37.0			
5145.16	6.01	0.16	20.709		63.0	63.0	41.0			
5151.09	6.01	-0.20	20.709		45.2	48.0	29.0			
3955.85	7.01	-0.78	18.465			27.1	38.0			
3995.00	7.01	0.28	18.496		111.0	99.0	75.3	47.0		
4035.08	7.01	0.62	23.123		77.0		50.0			
4041.31	7.01	0.85	23.141		49.0	51.0	53.0			
4043.53	7.01	0.74	23.131				31.0			
4073.05	7.01	-0.16	23.123				17.0			
4082.27	7.01	0.15	23.131				17.0			
4171.59	7.01	0.28	23.195				15.0			
4176.16	7.01	0.60	23.195				19.1			
4199.98	7.01	0.03	23.245				25.0			297.0
4227.74	7.01	-0.07	21.598				15.0			
4237.05	7.01	0.55	23.241		59.0	60.0	41.0			
4241.79	7.01	0.71	23.245				53.0			
4432.74	7.01	0.58	23.414			36.0	26.0			
4447.03	7.01	0.29	20.408			57.0	46.5			
4530.41	7.01	0.67	23.473				41.0			
4601.48	7.01	-0.37	18.465				37.0			
4607.15	7.01	-0.49	18.461		42.3	49.9	33.3			
4613.87	7.01	-0.73	18.465		90.0	73.0	44.0			
4621.39	7.01	-0.54	18.465		58.0	85.0	45.0			
4630.54	7.01	0.13	18.482	116.0		89.0	63.0	183.0		
4643.09	7.01	-0.34	18.482				45.0			
4779.72	7.01	-0.56	20.645		32.0	46.0				
4788.14	7.01	-0.38	20.652		30.8	34.0	18.0			
4803.29	7.01	-0.12	20.664			76.0	27.0			
5002.70	7.01	-1.02	18.461			40.0	23.0			
5005.15	7.01	0.61	20.664		90.0	70.0	58.0			
5007.33	7.01	0.16	20.939			31.0	39.0			
5010.62	7.01	-0.52	18.465				38.0			
5025.66	7.01	-0.44	20.664				6.9			
5045.10	7.01	-0.33	18.482		57.0	60.0	46.0			
3945.04	8.01	-0.69	23.418		61.0	72.0	56.8	50.0		
3982.72	8.01	-0.67	23.440		100.0	45.5	68.0	80.0		
4062.91	8.01	-0.09	28.704				25.0			
4071.24	8.01	-0.09	28.692				34.0			
4072.16	8.01	0.53	25.648			98.0	86.0			
4075.87	8.01	0.69	25.664			124.0	103.0			
4078.86	8.01	-0.26	25.637			49.7	44.5	52.0		
4083.93	8.01	0.15	28.682				30.0			
4085.12	8.01	-0.14	25.648			44.0	59.0			
4087.20	8.01	0.53	28.675		53.0	41.0	43.0			
4092.93	8.01	-0.25	25.664			83.0				
4119.22	8.01	0.48	25.847		181.0	168.0	117.0			
4129.31	8.01	-1.12	25.836				18.0			
4132.79	8.01	-0.07	25.830			71.0	53.0	77.0	27.4	
4153.28	8.01	0.08	25.836		120.0	105.5	65.0	128.0		
4156.52	8.01	-0.79	25.847				28.5			
4185.44	8.01	0.71	28.356			46.0	51.5	232.0		
4189.79	8.01	0.80	28.359		121.0		73.0	108.0		
4275.53	8.01	0.76	28.856	76.0			55.0			
4294.87	8.01	0.36	28.829				52.0			
4313.42	8.01	0.46	28.882				33.0			
4317.15	8.01	-0.32	22.965		135.0	101.0	85.0	154.0		
4319.64	8.01	-0.32	22.978		126.0	86.0	84.5	123.0		
4325.76	8.01	-1.06	22.965		62.0	36.0	45.0			
4366.91	8.01	-0.24	22.998		136.0	105.0	78.0	116.0	71.0	
4369.31	8.01	-0.35	26.224				29.0			
4414.88	8.01	0.17	23.440	79.0	143.0	96.9	109.0	212.0		
4416.97	8.01	0.04	23.418	85.0	129.0	78.6	105.0			
4443.04	8.01	0.00	28.356				29.5			
4448.34	8.01	0.13	28.359			41.0	45.0			
4452.38	8.01	-0.74	23.440		67.0	42.0	51.0	48.0		
4591.01	8.01	0.45	25.660		118.2	85.0	86.5	138.0		43.0
4596.20	8.01	0.29	25.660		125.6	90.1	85.0	99.0		
4602.06	8.01	0.51	29.061				40.0			
4609.37	8.01	0.67	29.067				51.0			
4610.17	8.01	-0.17	29.061				24.0			
4638.86	8.01	-0.30	22.965		110.0	95.0	80.0	151.0		
4649.14	8.01	0.31	22.998			160.0	125.0			
4661.64	8.01	-0.17	22.978	60.0	131.0	79.8	78.0		62.0	
4673.75	8.01	-1.07	22.978				20.0			
4676.24	8.01	-0.30	22.998	36.0	102.0	77.0	68.0			
4696.36	8.01	-1.31	22.998				28.3			
4699.19	8.01	0.37	28.508					94.0		
4703.26	8.01	0.21	28.511			28.2	35.0			
4705.32	8.01	0.57	26.248		93.0	64.3	77.2	107.0		33.9
4709.99	8.01	-0.47	26.224	92.0						
4906.82	8.01	-0.03	26.303		64.0	33.3	39.0	66.0		
4941.11	8.01	0.08	26.552			36.8	42.0			
4943.00	8.01	0.37	26.560		63.0	54.7	55.0			
4955.74	8.01	-0.42	26.560				20.0			

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