

The bulk of the cosmic infrared background resolved by ISOCAM

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Abstract. Deep extragalactic surveys with ISOCAM revealed the presence of a large density of faint mid-infrared (MIR) sources. We have computed the $15\ \mu\text{m}$ integrated galaxy light produced by these galaxies above a sensitivity limit of $50\ \mu\text{Jy}$. It sets a lower limit to the $15\ \mu\text{m}$ extragalactic background light of $(2.4 \pm 0.5)\ \text{nW m}^{-2}\ \text{sr}^{-1}$. The redshift distribution of the ISOCAM galaxies is inferred from the spectroscopically complete sample of galaxies in the Hubble Deep Field North (HDFN). It peaks around $z \sim 0.8$ in agreement with studies in other fields. The rest-frame $15\ \mu\text{m}$ and bolometric infrared ($8\text{--}1000\ \mu\text{m}$) luminosities of ISOCAM galaxies are computed using the correlations that we establish between the 6.75 , 12 , $15\ \mu\text{m}$ and infrared (IR) luminosities of local galaxies. The resulting IR luminosities were double-checked using radio ($1.4\ \text{GHz}$) flux densities from the ultra-deep VLA and WSRT surveys of the HDFN on a sample of 24 galaxies as well as on a sample of 109 local galaxies in common between ISOCAM and the NRAO VLA Sky Survey (NVSS). This comparison shows for the first time that MIR and radio luminosities correlate up to $z \sim 1$. This result validates the bolometric IR luminosities derived from MIR luminosities unless both the radio-far infrared (FIR) and the MIR-FIR correlations become invalid around $z \sim 1$. The fraction of IR light produced by active nuclei was computed from the cross-correlation with the deepest X-ray surveys from the Chandra and XMM-Newton observatories in the HDFN and Lockman Hole respectively. We find that at most 20% of the $15\ \mu\text{m}$ integrated galaxy light is due to active galactic nuclei (AGNs) unless a large population of AGNs was missed by Chandra and XMM-Newton. About 75% of the ISOCAM galaxies are found to belong to the class of luminous infrared galaxies ($L_{\text{IR}} \geq 10^{11}\ L_{\odot}$). They exhibit star formation rates of the order of $\sim 100\ M_{\odot}\ \text{yr}^{-1}$. The comoving density of infrared light due to these luminous IR galaxies was more than 40 times larger at $z \sim 1$ than today. The contribution of ISOCAM galaxies to the peak of the cosmic infrared background (CIRB) at $140\ \mu\text{m}$ was computed from the MIR-FIR correlations for star forming galaxies and from the spectral energy distribution of the Seyfert 2, NGC 1068, for AGNs. We find that the galaxies unveiled by ISOCAM surveys are responsible for the bulk of the CIRB, i.e. $(16 \pm 5)\ \text{nW m}^{-2}\ \text{sr}^{-1}$ as compared to the $(25 \pm 7)\ \text{nW m}^{-2}\ \text{sr}^{-1}$ measured with the COBE satellite, with less than 10% due to AGNs. Since the CIRB contains most of the light radiated over the history of star formation in the universe, this means that a large fraction of present-day stars must have formed during a dusty starburst event similar to those revealed by ISOCAM.

Key words. galaxies: evolution – infrared: galaxies – galaxies: starburst – galaxies: Seyfert

1. Introduction

The extragalactic background light (EBL) is a measurement of the sum of the light produced by all extragalactic sources over cosmic time. When it is integrated over the full spectral range, the so-called cosmic background is a fossil record of the overall activity of all galaxies from

their birth until now. It can be considered as the global energetic budget available for any model aiming at simulating the birth and fate of galaxies during the Hubble time. However the physical origin of this light will remain unknown until we have pinpointed the individual sources responsible for it. The goal of the present paper is to demonstrate that an important new result has come from the combination of a series of deep extragalactic surveys performed in the mid-infrared (MIR) at $15\ \mu\text{m}$ with

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the ISOCAM camera (Cesarsky et al. 1996a) onboard the Infrared Space Observatory (ISO, Kessler et al. 1996): we suggest here that the galaxies detected in these surveys, which median redshift of $z \sim 0.8$ was measured from a sub-sample of ISOCAM galaxies, contribute dominantly to the cosmic infrared background (CIRB), i.e. the EBL integrated over all wavelengths within $\lambda = 5$ to $1000 \mu\text{m}$.

The CIRB was recently detected and measured thanks to the cosmic background explorer (COBE) instruments FIRAS (Far Infrared Absolute Spectrometer) and DIRBE (Diffuse Infrared Background Experiment) (Puget et al. 1996; Fixsen et al. 1998; Lagache et al. 1999, 2000; Hauser et al. 1998; Dwek et al. 1998; Finkbeiner et al. 2000) from $100 \mu\text{m}$ to 1 mm . It peaks around $\lambda_{\text{max}} \simeq 140 \mu\text{m}$ and was found to represent at least half and maybe two thirds of the overall cosmic background (see Gispert et al. 2000). Hence the CIRB reflects the bulk of the star formation that took place over the history of the universe. By resolving it into individual galaxies, we would therefore pinpoint the times and places where most stars seen in the local universe were formed. Two physical processes were considered for its origin: nucleosynthesis, i.e. stellar radiation in star forming galaxies, and accretion around a black hole, i.e. active galactic nuclei. In both cases, the light is not directly coming from its physical source but is reprocessed by dust, i.e. absorbed and re-radiated thermally by the “warm” dust. Both processes are probably related (see Genzel et al. 1998), but energetic considerations, based on the presence of massive black holes and on the amount of heavy elements in local galaxies, suggest that star formation should by far dominate in the CIRB over AGN activity (Madau & Pozzetti 2000; Franceschini et al. 2001). However, until the individual galaxies responsible for the CIRB are found and studied in detail, this result will remain theoretical.

The spectral energy distribution (SED) in the IR of local galaxies peaks above $\sim 60 \mu\text{m}$ and typically around $80 \pm 20 \mu\text{m}$ (see Sanders & Mirabel 1996). As a result, the distant galaxies responsible for the peak of the CIRB detected by COBE around $\lambda_{\text{max}} \sim 140 \mu\text{m}$ should be located below $z \sim 1.3$ and present a redshift distribution peaked around $z \sim 0.8$, if their SEDs do not strongly differ from those of local galaxies. As we will see this is also the redshift range of the galaxies detected at $15 \mu\text{m}$ with ISOCAM.

The ISOCAM extragalactic surveys were performed with two filters, LW2 ($5\text{--}8.5 \mu\text{m}$) and LW3 ($12\text{--}18 \mu\text{m}$), centered at 6.75 and $15 \mu\text{m}$ respectively. The $6.75 \mu\text{m}$ sample of sources is strongly contaminated by galactic stars, whereas stars are rather easily distinguished from galaxies at $15 \mu\text{m}$ using optical-MIR colour-colour plots. As a consequence, we are only concerned here by the $15 \mu\text{m}$ galaxies. Moreover, the observed $6.75 \mu\text{m}$ light is no more produced by dust emission for galaxies more distant than $z \sim 0.4$ because of k-correction (redshifted stellar light dominates the $6.75 \mu\text{m}$ band above this redshift), whereas the observed $15 \mu\text{m}$ light is mostly due to dust emission for galaxies up to $z \sim 2$.

About 1000 galaxies detected in the $15 \mu\text{m}$ surveys were used to produce number counts (i.e. surface density of galaxies as a function of flux density; see Elbaz et al. 1999). The steep slope of the $15 \mu\text{m}$ counts below $\sim 1 \text{ mJy}$ indicates the presence of an excess of faint sources by one order of magnitude in comparison with predictions assuming no evolution of the $15 \mu\text{m}$ luminosity function with redshift. The presence of broad emission features in the MIR spectrum of galaxies alone cannot explain the shape of the number counts and a strong evolution of either the whole luminosity function (Xu 2000; Chary & Elbaz 2001) or preferentially of a sub-population of starburst galaxies evolving both in luminosity and density (Franceschini et al. 2001; Chary & Elbaz 2001; Xu et al. 2001) is required in order to fit the ISOCAM $15 \mu\text{m}$ counts. In the present paper, we suggest that these ISOCAM galaxies are in fact dusty starbursts responsible for the bulk of the CIRB.

In Sect. 2, we compare the sensitivity of different extragalactic surveys in several wavelength ranges to detect the galaxies responsible for the CIRB. It is suggested that MIR is presently the most efficient technique to detect dusty starbursts up to $z \sim 1.3$.

In Sect. 3, we calculate the $15 \mu\text{m}$ integrated galaxy light (IGL) due to ISOCAM galaxies. The $15 \mu\text{m}$ IGL is the sum of the $15 \mu\text{m}$ fluxes from individual galaxies, down to a given sensitivity limit, per unit area. It represents a lower limit to the $15 \mu\text{m}$ EBL, which remains unknown. Once the redshift distribution and SED of these galaxies is determined, it becomes possible to estimate their contribution to the CIRB.

In Sect. 4, we demonstrate that MIR luminosities at 6.75 , 12 and $15 \mu\text{m}$ are strongly correlated with the bolometric IR luminosity (from 8 to $1000 \mu\text{m}$) for local galaxies. The correlations presented in Chary & Elbaz (2001) are confirmed here with a larger sample of galaxies.

Before spectroscopic redshifts are obtained for the full sample of ISOCAM galaxies used to produce these number counts, the redshift distribution of these galaxies can be inferred from a few sub-samples: HDFN (Aussel et al. 1999, 2001), CFRS-14 (Flores et al. 1999, 2002), CFRS-03 (Flores et al. 2002). The ultra-deep ISOCAM survey of the HDFN samples a flux density range where most of the evolution observed in the number counts takes place and where the bulk of the $15 \mu\text{m}$ IGL is produced. This field is complete in spectroscopic redshifts, so it is used to estimate the bolometric IR luminosities and star formation rates of the ISOCAM galaxies in Sect. 5.

This result relies on two assumptions:

- that the main source for the MIR light in ISOCAM galaxies is star formation and not accretion around a black hole;
- that the correlations found in the local universe between the MIR and bolometric IR luminosity of galaxies remain valid up to $z \sim 1$. The first assumption is discussed and justified in Sect. 5.2, where soft and hard X-ray data from the Chandra and XMM-Newton X-ray observatories are combined with ISOCAM data on galaxies in the HDFN and Lockman Hole regions respectively.

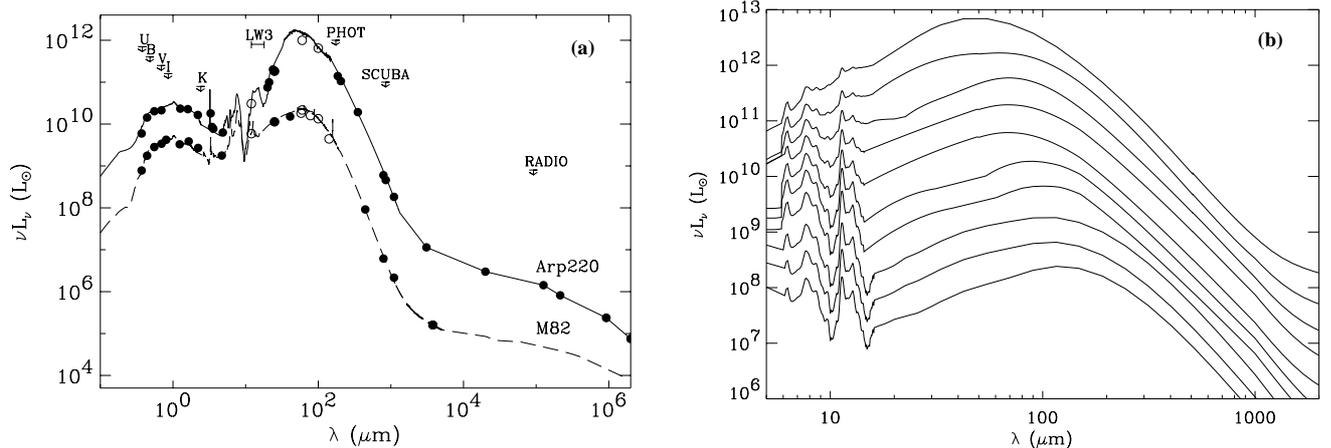


Fig. 1. Spectral energy distribution of galaxies in the infrared. **a)** SEDs of M 82 (dashed line) and Arp 220 (plain line) combining photometric (filled and open dots) and spectroscopic data. Open dots are used when they overlap with spectroscopic data. Radio data for Arp 220 from Anantharamaiah et al. (2000). Sub-millimetric data for M 82 from SCUBA (Hughes et al. 1994). $\lambda = 45\text{--}190\ \mu\text{m}$: ISO-LWS spectra from Fischer et al. (1999). Arp220's $18\ \mu\text{m}$ silicate absorption feature from Smith et al. (1989). $\lambda = 5\text{--}18\ \mu\text{m}$: ISOCAM Circular Variable Filter (CVF) spectra (Arp 220: Charmandaris et al. 1999, M 82: Förster Schreiber et al. 2002) normalized to the IRAS $12\ \mu\text{m}$ flux. For M 82, the CVF spectrum of the central region of the galaxy was complemented outward by adding to it a typical spectrum of a reflection nebula (NGC 7023, Cesarsky et al. 1996b), in order to match the $12\ \mu\text{m}$ IRAS flux. $\lambda = 2.5\text{ to }4\ \mu\text{m}$: for M 82, SWS spectrum of Förster Schreiber et al. (2001) scaled using the NGC 7023 spectrum of Moutou et al. (2000); for Arp 220, fit to the ISOPHOT data (Klaas et al. 1997) combined with the profile from Moutou et al. (2000). Optical to near-infrared: mix of the Coleman et al. (1980) templates to match the observed total magnitudes obtained from NED (Arp 220: 67% Sbc and 33% Ell; M 82: 12% Sbc and 88% Ell). **b)** 10 SEDs with $\log_{10}(L_{\text{IR}}/L_{\odot}) = 8.5$ to 13, with a step of 0.5, from the library of 105 SEDs of Chary & Elbaz (2001).

The issue of the robustness of the MIR-FIR correlations in the distant universe is addressed in Sect. 5.5, where IR luminosities are also computed from radio (1.4 GHz) flux densities for a sub-sample of 24 distant and 109 local ISOCAM galaxies.

In Sect. 5.6, we compute the cosmic density of IR light due to luminous IR galaxies ($L_{\text{IR}} \geq 10^{11} L_{\odot}$) at $\bar{z} \sim 1$. In Sect. 6, we evaluate the contribution of the ISOCAM galaxies to the CIRB, more precisely to its peak emission around $\lambda_{\text{max}} \sim 140\ \mu\text{m}$. Finally, the nature of ISOCAM galaxies is discussed in the conclusions (Sect. 7).

In the following, we will use the terms ULIG for galaxies with an IR luminosity $L_{\text{IR}} = L[8\text{--}1000\ \mu\text{m}] \geq 10^{12} L_{\odot}$, LIG, when $10^{11} \leq (L_{\text{IR}}/L_{\odot}) < 10^{12}$ and luminous IR galaxies for both ($L_{\text{IR}} \geq 10^{11} L_{\odot}$). Throughout this paper, we will assume $H_0 = 75\ \text{km s}^{-1} \text{Mpc}^{-1}$, $\Omega_{\text{matter}} = 0.3$ and $\Omega_{\Lambda} = 0.7$.

2. How to resolve the CIRB into individual galaxies ?

2.1. Contribution of IRAS galaxies

One of the major results from the IRAS satellite was to demonstrate that the bolometric luminosity of galaxies more luminous than $L_{\text{bol}} \sim 10^{11} L_{\odot}$ was underestimated by more than one order of magnitude before the IR luminosity was accounted for (Soifer et al. 1987). These luminous IR galaxies are not typical of the local galaxy population, since local galaxies radiate only about 30% of their bolometric luminosity in the dust regime from 8 to

1000 μm and only $\sim 2\%$ of the local bolometric luminosity density is due to luminous IR galaxies. If we convert the IR luminosity into a star formation rate (SFR) using the formula of Kennicutt (1998),

$$SFR(M_{\odot}/\text{yr}) = 1.71 \times 10^{-10} L_{\text{IR}}[8\text{--}1000\ \mu\text{m}](L_{\odot}) \quad (1)$$

we find that luminous IR galaxies ($L_{\text{bol}} \sim L_{\text{IR}} \geq 10^{11} L_{\odot}$) form stars at a rate larger than $\sim 20 M_{\odot} \text{yr}^{-1}$. Equation (1) assumes continuous bursts lasting 10–100 Myr, solar abundance and a Salpeter IMF (Kennicutt 1998).

IRAS surveyed about 95% of the sky down to a completeness limit of $\sim 0.5\ \text{mJy}$ at $60\ \mu\text{m}$ ($\sim 1.5\ \text{Jy}$ at $100\ \mu\text{m}$) and the differential counts are well fitted by an Euclidean slope, $dN/dS_{\nu} \sim S_{\nu}^{-2.5}$ (Soifer et al. 1987). The fluxes of all galaxies detected down to this sensitivity limit by IRAS add up to a $60\ \mu\text{m}$ IGL of $\sim 0.15\ \text{nW m}^{-2} \text{sr}^{-1}$. This is less than 1% of the value of the CIRB measured by COBE at $\lambda_{\text{max}} \sim 140\ \mu\text{m}$, $IGL_{140} = (25 \pm 7)\ \text{nW m}^{-2} \text{sr}^{-1}$.

2.2. Contribution of ISOPHOT galaxies

Deeper FIR extragalactic surveys were performed with ISOPHOT (Lemke et al. 1996) onboard ISO at $170\ \mu\text{m}$. At this depth, the differential counts are no more fitted by an Euclidean slope, expected in the case of no evolution of galaxies with redshift, but instead a strong excess of faint sources was found. At the confusion limit of ISOPHOT of about 120 mJy, less than 10% of the value of the CIRB at $170\ \mu\text{m}$ is resolved into individual galaxies

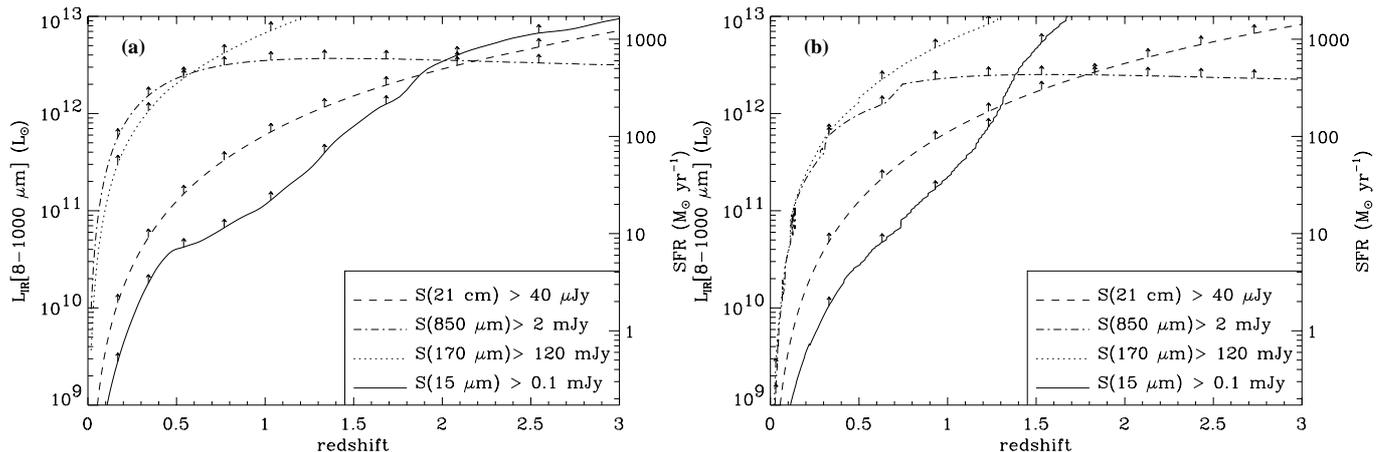


Fig. 2. IR luminosity (left axis) and SFR (right axis) corresponding to the sensitivity limits of ISOCAM (15 μm , plain line), ISOPHOT (170 μm , dotted line), SCUBA (850 μm , dot-dashed line) and VLA/WSRT (21 cm, dashed line) as a function of redshift. *K*-correction from: **a)** the template SED of M 82 (Fig. 1a), normalized to the appropriate IR luminosity, **b)** the library of template SEDs from Chary & Elbaz (2001), plotted in the Fig. 1b.

(Dole et al. 2001). A preliminary follow-up of the ISOPHOT galaxies suggests that a large fraction of these galaxies are located at low redshift.

2.3. Contribution of SCUBA galaxies

In the sub-millimeter range, the SCUBA bolometer array on the James Clerk Maxwell Telescope (JCMT) has produced deep images at 450 and 850 μm . Number counts at both wavelengths present a steep slope compatible with a strong evolution as compared to the local universe (see Smail et al. 2001 and references therein).

At 450 μm , a depth of 10 mJy is reached and the combined fluxes of all SCUBA galaxies produce about 15% of the CIRB measured by COBE-FIRAS at this wavelength (Smail et al. 2001).

At 850 μm , SCUBA is confusion limited at ~ 2 mJy (Hughes et al. 1998; Barger et al. 1999; Eales et al. 2000; Smail et al. 2001), because of its large point spread function (PSF) of 15'' full width half maximum (FWHM). About 20% of the value of the CIRB measured by COBE-FIRAS at 850 μm (the 850 μm EBL) is resolved into galaxies at this depth. However, using gravitational lensing this limit can be lowered to ~ 1 mJy, where $\sim 60\%$ of the CIRB is resolved (Smail et al. 2001; Blain et al. 1999a).

The 850 μm EBL measured by COBE-FIRAS ($\sim 0.5 \pm 0.2 \text{ nW m}^{-2} \text{ sr}^{-1}$) is 50 times lower than the peak value of the CIRB at $\lambda_{\text{max}} \sim 140 \mu\text{m}$. In order to compute the contribution of SCUBA galaxies to the peak of the CIRB it is therefore necessary to determine their redshift distribution and SED. Until now, very few redshifts have been obtained due to the large PSF of SCUBA and optical faintness of these galaxies. Hence the fraction of the CIRB resolved into individual galaxies by SCUBA remains highly uncertain.

2.4. Comparison with the MIR and radio

The deepest ISOCAM surveys reach a completeness limit of $\sim 100 \mu\text{Jy}$ at 15 μm (without lensing, but this limit goes down to $\sim 30\text{--}50 \mu\text{Jy}$ including lensing, see Sect. 3). For a given galaxy SED and redshift, this flux density can be converted into a bolometric IR luminosity ($L_{\text{IR}}[8\text{--}1000 \mu\text{m}]$). Any galaxy more luminous than this L_{IR} will be detected in the survey. In order to compare the sensitivity to distant luminous IR galaxies of ISOCAM, ISOPHOT and SCUBA, we have calculated L_{IR} corresponding to their sensitivity limits for a given redshift between $z = 0$ and 3. On one hand, we used the SED of M 82 (Fig. 1a) for galaxies of all luminosities. On the other hand, we used a library of template SEDs from Chary & Elbaz (2001). In the following, we will call it the “multi-template” technique. This library of 105 SEDs was constructed under the constraint that it reproduces the correlations observed between the 6.75, 12, 15, 60, 100 and 850 μm luminosities of local galaxies. The SEDs were interpolated from a sample of observed SEDs, including MIR spectra obtained with the ISOCAM CVF. They cover the luminosity range $\log_{10}(L_{\text{IR}}/L_{\odot}) = 8.5$ to 13.5 (see Fig. 1b).

We did not use Arp 220’s SED in Fig. 2, because this galaxy presents a large FIR over MIR luminosity ratio, which is not representative of galaxies within the same luminosity range (see Figs. 5c,d).

The radio continuum ($\lambda = 21 \text{ cm}$, $\nu = 1.4 \text{ GHz}$) is also a tracer of L_{IR} because of the tight correlation between both luminosities in the local universe (see Yun et al. 2001 and references therein). If this correlation remains valid in the distant universe then we can translate the sensitivity limit of the deepest radio surveys into a minimum L_{IR} accessible at a given redshift assuming a radio SED. The

correlation between the radio and FIR luminosities is usually described by the “ q ” parameter (Condon et al. 1991):

$$q = \log_{10} \left(\frac{L_{\text{FIR}}(\text{W})}{3.75 \times 10^{12}(\text{Hz})} \times \frac{1}{L_{1.4 \text{ GHz}}(\text{W Hz}^{-1})} \right) \quad (2)$$

where $L_{1.4 \text{ GHz}}$ is the monochromatic luminosity at 1.4 GHz and L_{FIR} is the FIR luminosity between 40 and 120 μm , as defined by Helou et al. (1988),

$$L_{\text{FIR}} = 1.26 \times 10^{-14} (2.58 S_{60} + S_{100}) \times 4\pi d(\text{m})^2 \quad (3)$$

where S_{60} and S_{100} are the IRAS flux densities at 60 and 100 μm in Jy. The relationship between $L_{\text{FIR}}[40\text{--}120 \mu\text{m}]$ and $L_{\text{IR}}[8\text{--}1000 \mu\text{m}]$ was computed with the sample of 300 galaxies from the Bright Galaxy Sample (BGS, Soifer et al. 1987) detected in the four IRAS bands (12, 25, 60 and 100 μm):

$$L_{\text{IR}} = (1.91 \pm 0.17) \times L_{\text{FIR}} \quad (4)$$

L_{IR} is defined as (Sanders & Mirabel 1996):

$$L_{\text{IR}} = [13.48 S_{12} + 5.16 S_{25} + 2.58 S_{60} + S_{100}] \times 1.8 \times 10^{-14} \times 4\pi d(\text{m})^2 \quad (5)$$

where S_{12} and S_{25} are the IRAS flux densities at 12 and 25 μm in Jy. Yun et al. (2001) measured a value of $q = 2.34 \pm 0.01$ from a flux limited ($S_{60} \geq 2$ Jy) sample of 1809 IRAS galaxies.

The deepest radio surveys presently available were performed in the HDFN by Richards (2000) and Garrett et al. (2000) with the VLA and WSRT respectively. Both surveys reach the same sensitivity limit of 40 μJy at 1.4 GHz (5- σ). In order to convert this flux density into a $L_{1.4 \text{ GHz}}$ for a given redshift, we used a power index of $\alpha = 0.8 \pm 0.15$ ($S_\nu \propto \nu^{-\alpha}$), as suggested by Yun et al. (2001) for starburst galaxies.

Figure 2 presents the sensitivity of the deepest MIR, FIR, sub-millimeter and radio surveys in the form of L_{IR} (or SFR) as a function of redshift. The sensitivity limits used for ISOCAM, ISOPHOT and SCUBA are taken from the deepest existing surveys at these wavelengths, as previously described. Figure 2a (with M 82) and Fig. 2b (“multi-template” technique) both clearly show that the faintest IR galaxies are best detected at 15 μm up to $z \sim 1.3$. The right axis of the plots shows the corresponding minimum SFR that a galaxy must harbor to be detected at a given redshift in the surveys. At $z \sim 1$, ISOCAM (plain line) detects all LIGs while SCUBA (dot-dashed line) detects only ULIGs. The difference in sensitivity between ISOCAM and the deepest radio surveys is about a factor 2. ISOPHOT (dotted line) is only sensitive to ULIGs above $z \sim 0.5$. In Fig. 2b, where we used the library of template SEDs from Chary & Elbaz (2001), the plain line corresponding to ISOCAM rises faster above $z \sim 1$ than in Fig. 2a, where we used M 82’s SED. This different behavior comes from the fact that at these redshifts ISOCAM measures fluxes at about 7 μm in the rest-frame of the galaxy and we will show that $L_{\text{IR}}/L_{7 \mu\text{m}}$ increases with L_{IR} (Sect. 4).

Above $z \sim 2$, the sub-millimeter becomes the most efficient technique, although only galaxies more luminous than a few $10^{12} L_\odot$ can be detected above a sensitivity limit of 2 mJy at 850 μm . Hence the unlensed ISOCAM and SCUBA surveys are not sampling the same redshift and luminosity ranges. This statement is confirmed observationally in the HDFN itself where none of the ISOCAM sources are detected by SCUBA (Hughes et al. 1998) and over a larger scale in the Canada France Redshift Survey 14 (CFRS-14, Eales et al. 2000). In this latter field of $\sim 50''^2$, only the two brightest 15 μm ISOCAM sources are detected at 850 μm , among a sample of 50 ISOCAM (Flores et al. 1999) and 19 SCUBA sources. This confirms that both instruments detect different sets of objects. More common objects between ISOCAM and SCUBA are expected if they are gravitationally lensed, since the sensitivity limits of both instruments are then decreased by a factor of about two. Indeed in the clusters A 370 and A 2390, three of the four lensed SCUBA galaxies (Smail et al. 2001) are also detected with ISOCAM (Altieri et al. 1999; Metcalfe 2000).

3. The 15 μm extragalactic background light

Above the Earth’s atmosphere, the 15 μm light is strongly dominated by the zodiacal emission from interplanetary dust and it has not yet been possible to make a direct measurement of the 15 μm background, or EBL. Individual galaxies contribute to this background and a lower limit to the 15 μm EBL can be obtained by adding up the fluxes of all ISOCAM galaxies detected per unit area down to a given flux limit. The resulting value is called the 15 μm integrated galaxy light (IGL).

The statistical reliability with which the surface density of MIR sources is computed depends on their flux: large and shallow fields are required for bright sources which are less numerous, whereas faint sources require smaller and deeper fields. The optimization is obtained by observing similar volumes of the sky but in pencil beams for faint sources and in wide surveys for bright sources, instead of one large and deep survey which would be unnecessarily time consuming. Together with the need to account for cosmic variance, this explains why we have combined a set of 10 surveys located in 7 regions of the northern and southern hemispheres, in order to estimate the 15 μm IGL. These surveys cover five orders of magnitude in flux density from $\sim 50 \mu\text{Jy}$ to ~ 4 Jy (including IRAS counts at 12 μm). The number counts presented in Elbaz et al. (1999) are converted here into a differential contribution to the 15 μm IGL as a function of flux density, estimated from the following formula:

$$\frac{dIGL}{dS} = \frac{dN}{dS} \times \left(\frac{S_{15}}{10^{20}} \right) \times \nu_{15} \quad (6)$$

where $dN(\text{sr}^{-1})$ is the surface density of sources with a flux density $S_\nu[15 \mu\text{m}] = S_{15}$ (mJy) over a flux density bin $dS(\text{mJy})$ ($1 \text{ mJy} = 10^{-20} \text{ nW m}^{-2} \text{ Hz}^{-1}$) and $\nu_{15}(\text{Hz})$ is the frequency of the 15 μm photons.

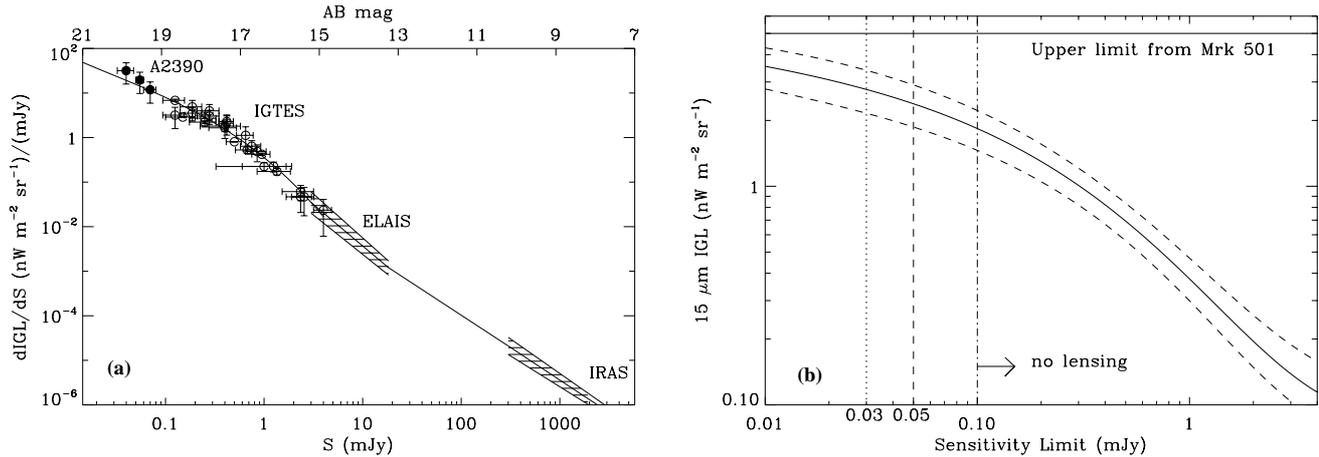


Fig. 3. **a)** Differential contribution to the $15\ \mu\text{m}$ Integrated Galaxy Light as a function of flux density and AB magnitude. The plain line is a fit to the data: Abell 2390 (Altieri et al. 1999), the ISOCAM Guaranteed Time Extragalactic Surveys (IGTES, Elbaz et al. 1999), the European Large Area Infrared Survey (ELAIS, Serjeant et al. 2000) and the IRAS all sky survey (Rush et al. 1993). **b)** Contribution of ISOCAM galaxies to the $15\ \mu\text{m}$ extragalactic background light (EBL), i.e. $15\ \mu\text{m}$ Integrated Galaxy Light (IGL), as a function of sensitivity or AB magnitude ($AB = -2.5 \log(S_{\text{mJy}}) + 16.4$). The plain line is the integral of the fit to $dIGL/dS$ (Fig. 3a). The dashed lines correspond to $1\text{-}\sigma$ error bars obtained by fitting the $1\text{-}\sigma$ upper and lower limits of $dIGL/dS$.

Below $S_{15} \sim 3\ \text{mJy}$, about 600 galaxies were used to produce the points with errors bars in Fig. 3a. Figure 3b shows the $15\ \mu\text{m}$ IGL as a function of depth. It corresponds to the integral of Fig. 3a, where the data below $3\ \text{mJy}$ were fitted with a polynomial of degree 3 and the $1\text{-}\sigma$ error bars on $dIGL/dS$ were obtained from the polynomial fit to the upper and lower limits of the data points. The $15\ \mu\text{m}$ IGL does not converge above a sensitivity limit of $S_{15} \sim 50\ \mu\text{Jy}$, but the flattening of the curve below $S_{15} \sim 0.4\ \text{mJy}$ suggests that most of the $15\ \mu\text{m}$ EBL should arise from the galaxies already unveiled by ISOCAM.

Above the completeness limit of $S_{15} \sim 50\ \mu\text{Jy}$, we computed an IGL of:

$$IGL_{15}(S_{15} \geq 50\ \mu\text{Jy}) = 2.4 \pm 0.5\ \text{nW m}^{-2}\ \text{sr}^{-1} \quad (7)$$

where the error bar corresponds to the 68% confidence level (i.e. $1\text{-}\sigma$). The error bar combines the uncertainty on the measurements for each individual survey with cosmic variance. Each flux bin is covered by two to three independent surveys. ISOCAM number counts are only complete above $S_{15} \sim 100\ \mu\text{Jy}$, but Aussel et al. (1999) computed number counts down to $S_{15} \sim 50\ \mu\text{Jy}$ after correcting for incompleteness and a consistent result is obtained by Altieri et al. (1999) after correcting for lensing magnification a sample of galaxies in the field of the galaxy cluster Abell 2390.

Only 9 sources are detected below $S_{15} \sim 50\ \mu\text{Jy}$ by Altieri et al. (1999), where the flux determination and correction for completeness are less robust and are not confirmed by another survey. A tentative value of $IGL_{15} \sim 3.3 \pm 1.3\ \text{nW m}^{-2}\ \text{sr}^{-1}$ was quoted in Altieri et al. (1999), above $S_{15} = 30\ \mu\text{Jy}$. This value may be slightly overestimated if we consider that the models of Franceschini et al. (2001) and Chary & Elbaz (2001), which reproduce the number counts from ISOCAM at $15\ \mu\text{m}$, from ISOPHOT

at 90 and $170\ \mu\text{m}$ and from SCUBA at $850\ \mu\text{m}$, as well as the shape of the CIRB from 100 to $1000\ \mu\text{m}$, consistently predict a $15\ \mu\text{m}$ EBL of:

$$EBL^{\text{models}}(15\ \mu\text{m}) \sim 3.3\ \text{nW m}^{-2}\ \text{sr}^{-1}. \quad (8)$$

If this prediction from the models is correct then about $73 \pm 15\%$ of the $15\ \mu\text{m}$ EBL is resolved into individual galaxies by the ISOCAM surveys.

This result is consistent with the upper limit on the $15\ \mu\text{m}$ EBL estimated by Stanev & Franceschini (1998) of:

$$EBL^{\text{max}}(15\ \mu\text{m}) \sim 5\ \text{nW m}^{-2}\ \text{sr}^{-1}. \quad (9)$$

This upper limit was computed from the 1997 γ -ray outburst of the blazar Mrk 501 ($z = 0.034$) as a result of the opacity of MIR photons to γ -ray photons, which annihilate with them through e^+e^- pair production. It was recently confirmed by Renault et al. (2001), who found an upper limit of $4.7\ \text{nW m}^{-2}\ \text{sr}^{-1}$ from 5 to $15\ \mu\text{m}$.

Finally, the $15\ \mu\text{m}$ background (EBL) must be contained between the following limits:

$$2.4 \pm 0.5 \leq \frac{EBL(15\ \mu\text{m})}{(\text{nW m}^{-2}\ \text{sr}^{-1})} \leq 5. \quad (10)$$

In Table 1, we give the relative contribution to the $15\ \mu\text{m}$ IGL from sources within a given flux density range. Nearly all the IGL is produced by sources fainter than $3\ \text{mJy}$ (94%) and about 70% by sources fainter than $0.5\ \text{mJy}$. This means that the nature and redshift distribution of the galaxies producing the bulk of the $15\ \mu\text{m}$ IGL can be determined by studying these faint galaxies only. We will therefore assume in the following that the spectroscopically complete sample of $15\ \mu\text{m}$ galaxies detected in the HDFN below $0.5\ \text{mJy}$ (Aussel et al. 1999, 2001) is representative of the population of dusty galaxies producing the bulk of the $15\ \mu\text{m}$ IGL.

Table 1. Relative contributions to the $15\ \mu\text{m}$ IGL as a function of flux density. Ref.: (1) Altieri et al. (1999), Metcalfe (2000), (2) Aussel et al. (1999, 2001), (3) Elbaz et al. (1999), (4) Serjeant et al. (2000), renormalized by a factor 0.5 as in Genzel & Cesarsky (2000), following suggestion from S. Serjeant for photometric calibration, (5) Interpolation in a flux density range where we do not have enough statistics. We used the same slope as in the IRAS counts, (6) IRAS counts from Rush et al. (1993) at $12\ \mu\text{m}$.

Field	$F(\text{mJy})$	%	Ref.
A2390,HDFN	[0.05, 0.1]	23	(1, 2)
IGTES	[0.1, 0.5]	48	(2, 3)
IGTES	[0.5, 3]	23	(3)
ELAIS	[3, 32]	4.5	(4)
Interpolation	[32,300]	1	(5)
IRAS	[300,4000]	0.5	(6)

4. MIR, as a tracer of the bolometric IR luminosity

4.1. SED of galaxies in the IR

Above a bolometric luminosity of a few $10^{10} L_{\odot}$, the SED of a galaxy is dominated by dust radiation with respect to direct stellar emission. Stellar light is absorbed by dust and re-radiated in the IR regime above $\lambda \sim 5\ \mu\text{m}$. The full SEDs, from the UV to the radio, of the starburst galaxy, M 82, and of the ULIG, Arp 220, are plotted in the Fig. 1a. A zoom on the 5–2000 μm wavelength range of the SEDs is shown in the Fig. 4, together with three other galaxies: a starburst, Arp 244, and two Seyfert 2 galaxies, NGC 1068 and NGC 6240. Arp 244, also known as the “Antennae” galaxy, is a system of two interacting spiral galaxies (Vigroux et al. 1996; Mirabel et al. 1998). The optical, MIR and FIR luminosities of these five galaxies are given in the Table 2. Note the increase of the MIR luminosity together with L_{FIR} , or L_{IR} , from M 82 to Arp 244 to NGC 6240. This correlation will be statistically established in Sect. 4.2.2.

The FIR over MIR luminosity ratio of Arp 220 is larger than for galaxies of equivalent luminosity (see Fig. 5 in Sect. 4.2.2). The flat FIR over MIR ratio for NGC 1068 is typical of the hot dust continuum from AGNs. NGC 1068 is the closest Seyfert 2 ($L_{\text{IR}} = 1.7 \times 10^{11} L_{\odot}$). On one hand, its MIR spectrum is completely dominated by a continuum due to the hot dust heated by the central AGN, which produces about 75% of the MIR luminosity (Le Floc’h et al. 2001). On the other hand, the bulk of its FIR luminosity originates from the diffuse region surrounding the nucleus as shown by the SCUBA image at $450\ \mu\text{m}$, which is associated with star formation (Le Floc’h et al. 2001). As a result, the FIR over MIR luminosity ratio is much lower for NGC 1068 than for starbursts (see Fig. 4). This galaxy is a perfect illustration of the close connection between star formation and AGN activity.

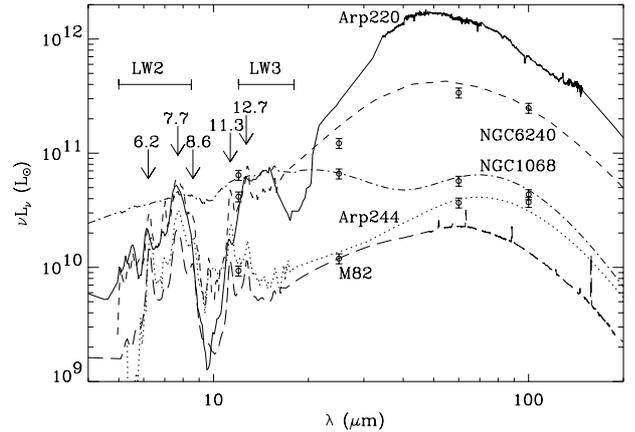


Fig. 4. Spectral energy distributions in the infrared (5–200 μm) of 3 starbursts (M 82, Arp 244, Arp 220) and 2 Seyfert 2’s (NGC 1068, NGC 6240). MIR spectra from ISOCAM CVF: NGC 1068 (Laurent et al. 2000; Le Floc’h et al. 2001), Arp 244 (Vigroux et al. 1996; Mirabel et al. 1998), NGC 6240 (Charmandaris et al. 1999). FIR: fit of the IRAS flux densities at 12, 25, 60 and 100 μm using the template SEDs from Chary & Elbaz (2001) for Arp 244 and NGC 6240. For NGC 1068, the IRAS flux densities were fitted by the sum of two modified black bodies of temperatures 30 and 170 K as in Telesco et al. (1984). Open dots with error bars: IRAS data (from NED) for Arp244, NGC 6240 and NGC 1068.

The MIR (5–40 μm) and FIR (40–1000 μm) emission of a galaxy combine three major components (see Fig. 4):

- broad emission features and their associated underlying continuum, which dominate in the $\lambda \sim 5\text{--}10\ \mu\text{m}$ domain, and up to $\sim 20\ \mu\text{m}$ for galaxies with moderate star formation. These bands located at 6.2, 7.7, 8.6, 11.3 and 12.7 μm , are usually quoted as PAHs (polycyclic aromatic hydrocarbons, Léger & Puget 1984; Puget & Léger 1989; Allamandola et al. 1989), although their exact nature remains uncertain, e.g. Jones & d’Hendecourt (2000).
- the “warm” dust continuum produced by very small dust grains (VSGs, smaller than $0.01\ \mu\text{m}$, Andriessse 1978, Désert et al. 1990) heated at temperatures larger than 200 K without reaching thermal equilibrium. This component dominates the SED between 10–20 to 40 μm for luminous galaxies.
- the “cold” dust continuum produced by big dust grains (“big grains”, of size $>0.01\ \mu\text{m}$) in thermal equilibrium at cool temperatures (below 60–70 K typically). This component is responsible for the bulk of the FIR light where the SED peaks.

4.2. The MIR-FIR correlation

4.2.1. The database

We have compiled a database containing MIR and FIR fluxes for a sample of 157 galaxies detected with ISOCAM and IRAS: 50 galaxies from Boselli et al. (1998), 44 galaxies from Roussel et al. (2001), 19 galaxies from Laurent et al. (1999, 2000), 41 galaxies from Dale et al. (2001), plus M 82, Arp 244 and NGC 1068 (from Fig. 4). Figure 5

Table 2. Luminosities of 3 starbursts (M 82, Arp 244, Arp 220) and 2 Seyfert 2's (NGC 1068, NGC 6240). L_B is the B -band monochromatic luminosity (νL_ν at $0.44 \mu\text{m}$) from RC3 (de Vaucouleurs et al. 1991). L_{IR} and L_{FIR} are the IR and FIR luminosities integrated over 8–1000 and 40–500 μm respectively from the spectra shown in the Fig. 4. L_{LW3} and L_{LW2} are the monochromatic luminosities (νL_ν) at the central wavelengths 15 and 6.75 μm of the ISOCAM LW3 and LW2 filters respectively. At $z \sim 1$, the LW3 filter (12–18 μm) matches the rest-frame LW2 bandpass (5–8.5 μm).

	M 82	Arp 244	Arp 220	N 1068	N 6240
d (Mpc)	3.25	22.0	78.4	14.4	99.6
$\log(L_B)$	10.2	10.2	10.6	10.5	10.5
$\log(L_{\text{IR}})$	10.6	10.8	12.2	11.2	11.8
$\log(L_{\text{FIR}})$	10.4	10.7	12.1	10.8	11.7
$\log(L_{\text{LW3}})$	9.8	10.0	10.6	10.8	10.7
$\log(L_{\text{LW2}})$	9.9	10.1	10.3	10.6	10.5
L_{IR}/L_B	2	3	39	5	18
$L_{\text{IR}}/L_{\text{LW3}}$	15	7	41	5	10
$L_{\text{IR}}/L_{\text{LW2}}$	4	4	87	4	22

shows that the 6.75, 12, 15 and bolometric IR (8–1000 μm) luminosities (L_{IR}) of these galaxies are tightly correlated with each other. L_{IR} is calculated following Eq. (5). At 15 μm (LW3 filter), we completed our sample with the NEPR sample of IRAS galaxies observed with ISOCAM-LW3 (Aussel et al. 2000). We converted the 60 μm luminosities of the NEPR galaxies into L_{IR} by fitting the IR versus 60 μm correlation obtained with the BGS (Soifer et al. 1987). In order to complete our sample in the high luminosity range at 6.75 μm (LW2 filter), we used a sample of 40 ULIGs observed with the spectrometer of ISOPHOT, PHOT-S (from Rigopoulou et al. 1999: the conversion into LW2 flux densities was computed as in Chary & Elbaz 2001). We also used four ULIGs observed with the ISOCAM CVF from Tran et al. (2001). We conservatively selected these star forming galaxies with respect to AGNs from their line over continuum ratio as suggested in Tran et al. (2001). The ISOCAM-CVF spectra were converted into LW2 flux densities using the transmission curve of the LW2 filter.

4.2.2. Correlations

The 6.75, 12 and 15 μm luminosities of local galaxies are tightly correlated with each other (see Figs. 5a,b) following the formulae:

$$\begin{aligned} \nu L_\nu[15 \mu\text{m}] &= 0.43 \times \nu L_\nu[6.75 \mu\text{m}] \\ &\quad \text{for } \nu L_\nu[6.75 \mu\text{m}] < 5 \times 10^9 L_\odot \\ &= 3.8 \times 10^{-7} \times (\nu L_\nu[6.75 \mu\text{m}])^{1.62} \\ &\quad \text{for } \nu L_\nu[6.75 \mu\text{m}] \geq 5 \times 10^9 L_\odot \end{aligned} \quad (11)$$

$$\nu L_\nu[15 \mu\text{m}] = 0.042 \times (\nu L_\nu[12 \mu\text{m}])^{1.12}. \quad (12)$$

These correlations can be used to estimate the rest-frame 15 μm luminosity for galaxies located around $z \sim 0.3$ and 1. At these redshifts, 15 μm observed corresponds to about 12 and 7 μm in the rest-frame.

The MIR luminosity of galaxies is a good tracer of L_{IR} as seen in Figs. 5c,d, where L_{IR} is correlated with νL_ν at 15 and 6.75 μm following the formulae:

$$L_{\text{IR}} = 11.1_{-3.7}^{+5.5} \times (\nu L_\nu[15 \mu\text{m}])^{0.998} \quad (13)$$

$$\begin{aligned} L_{\text{IR}} &= 4.78_{-1.59}^{+2.37} \times (\nu L_\nu[6.75 \mu\text{m}])^{0.998} \\ &\quad \text{for } \nu L_\nu[6.75 \mu\text{m}] < 5 \times 10^9 L_\odot \\ &= 4.37_{-2.13}^{+2.35} \times 10^{-6} \times (\nu L_\nu[6.75 \mu\text{m}])^{1.62} \\ &\quad \text{for } \nu L_\nu[6.75 \mu\text{m}] \geq 5 \times 10^9 L_\odot. \end{aligned} \quad (14)$$

The formulae used in these four correlations are consistent with Chary & Elbaz (2001), although slightly more data were used here. The change of slope in Fig. 5d and Eq. (14) from a slope of 1 to a slope larger than 1 statistically establishes what was suggested by Fig. 4: $L_{\text{IR}}/\nu L_\nu[6.75 \mu\text{m}]$ increases with L_{IR} for galaxies with $L_{\text{IR}} > 2 \times 10^{10} L_\odot$.

The position of the five templates from Fig. 4 are indicated with open squares in Fig. 5. Two galaxies do not fit the correlations: Arp 220 and NGC 1068. Arp 220 presents a very large FIR over MIR ratio, which was interpreted as being due to extinction in the MIR by Haas et al. (2001), while the IR SED of the Seyfert 2 NGC 1068 is flat, which is typical of AGNs. In this sense NGC 6240 is quite atypical, since it falls perfectly on the correlation followed by star forming galaxies. The presence of an active nucleus in NGC 6240 was revealed by its strong X-ray luminosity and Ikebe et al. (2000) suggested that a large fraction of the IR luminosity of this galaxy could originate from its AGN. However, the presence of strong MIR features (Genzel et al. 1998), the absence of a strong continuum emission below $\lambda \sim 6 \mu\text{m}$ (Laurent et al. 2000), together with a FIR over MIR ratio typical of star forming galaxies, suggest that the bulk of its IR luminosity may still be due to star formation.

5. IR luminosities and star formation rates of ISOCAM galaxies

In this section, we compute the IR luminosity and star formation rate of the galaxies with $S_{15} \geq 0.1$ mJy located in the Hubble Deep Field North plus its flanking fields, hereafter called ISOCAM-HDFN galaxies (Aussel et al. 1999, 2001). This field presents the advantage of being spectroscopically complete and deep enough to cover a flux density range ($S_{15} \leq 0.5$ mJy) where more than 70% of the 15 μm IGL is produced. Moreover, the redshift distribution of the ISOCAM-HDFN galaxies ($\bar{z} \sim 0.8$) is similar to the one measured in the CFRS fields CFRS-14 (Flores et al. 1999, 2002) and CFRS-03 (Flores et al. 2002). Together, these fields of $100'^2$ each, possess 81 galaxies with $S_{15} \geq 0.35$ mJy. The redshift distribution of 70% of the CFRS galaxies, for which a spectroscopic redshift was measured, is very similar to the one presented below for the ISOCAM-HDFN galaxies and peaks around $z \sim 0.7$.

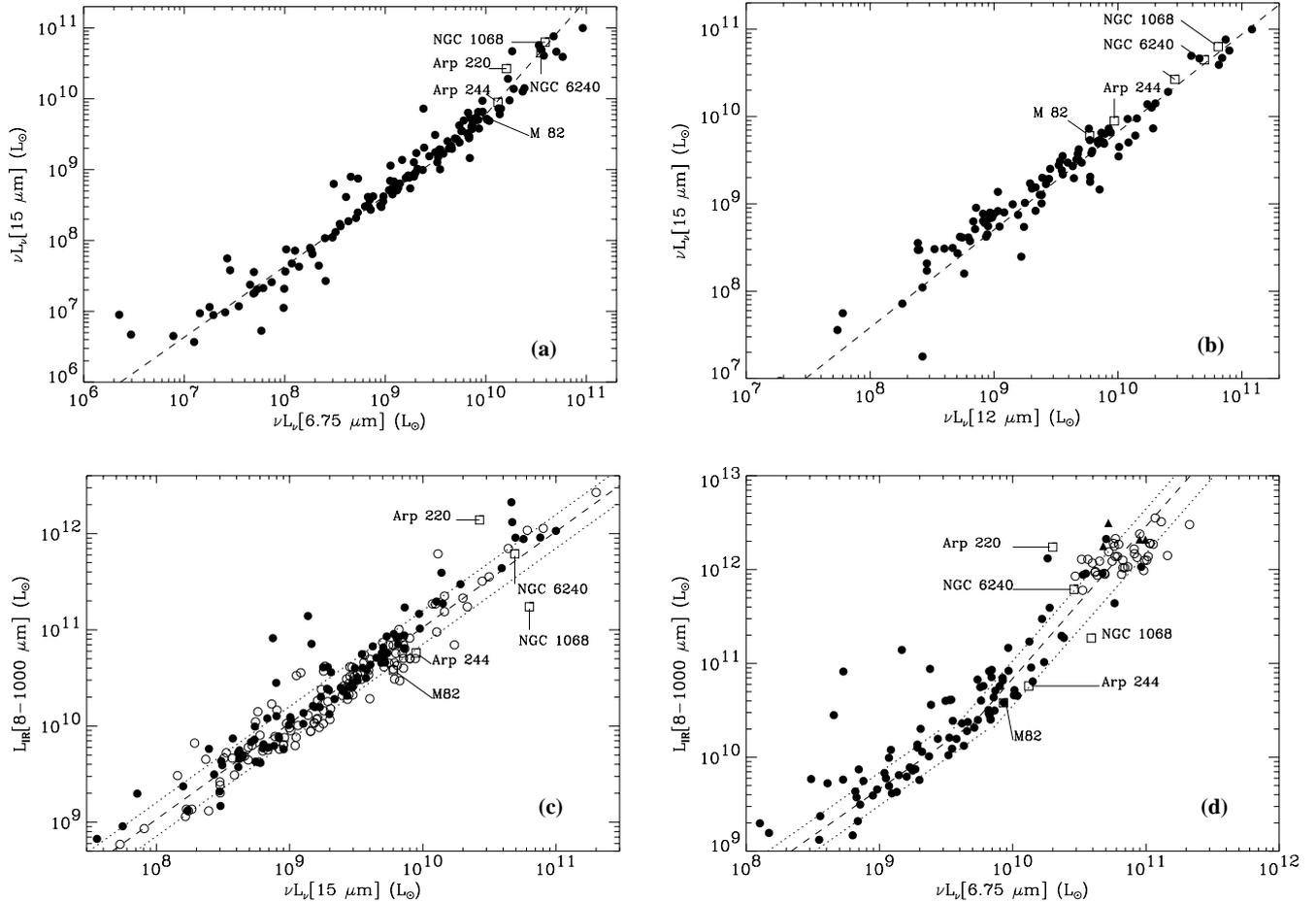


Fig. 5. IR luminosity correlations for local galaxies (see text for the origin of the sample). The five galaxies from Fig. 4 and Table 2 are marked with open squares. **a)** ISOCAM-LW3 (15 μm) versus ISOCAM-LW2 (6.75 μm) luminosities (νL_ν) (56 galaxies). **b)** ISOCAM-LW3 (15 μm) versus IRAS-12 μm luminosities (45 galaxies). **c)** $L_{\text{IR}}[8-1000 \mu\text{m}]$ versus ISOCAM-LW3 (15 μm) luminosity (120 galaxies). **d)** $L_{\text{IR}}[8-1000 \mu\text{m}]$ versus LW2-6.75 μm luminosities (91 galaxies). Filled dots: galaxies from the ISOCAM guaranteed time (47 galaxies including the open squares). Open dots: 40 galaxies from Rigopoulou et al. (1999). Filled triangles: 4 galaxies from Tran et al. (2001). Galaxies below $L_{\text{IR}} \sim 10^{10} L_\odot$ present a flatter slope and have $L_{\text{IR}}/L_B < 1$.

5.1. Redshift distribution of the ISOCAM-HDFN galaxies

The list of ISOCAM-HDFN sources used in the following comes from a revised version of the list published in Aussel et al. (1999). In this version (Aussel et al. 2001), 4 ISOCAM sources were deblended using PSF deconvolution based on the position of the optical counterparts and the flux density of 10 other galaxies were re-evaluated. With a PSF $FWHM$ of 4.6'' at 15 μm (Okumura 1998), optical counterparts are easily identified for ISOCAM galaxies and all of them are found within 3'' of the PSF centroid. This can also be seen in the crowded field of the nearby galaxy cluster Abell 1689 (Fadda et al. 2000). More than 90% of the ISOCAM galaxies present an optical counterpart brighter than $I(Kron-Cousins) = 22.5$, which is natural since ISOCAM is mostly sensitive to galaxies closer than $z \sim 1.3$ (see Fig. 2). It results from a combination of sensitivity limit and k-correction: dust emission becomes negligible below $\sim 6 \mu\text{m}$ (except for bright AGNs which can be detected further away).

A total of 41 ISOCAM-HDFN galaxies have 15 μm flux densities between $S_{15} = 0.1$ and 0.5 mJy (see Fig. 6a) over a field of 25.8 l^2 (HDFN+FF). Monte-Carlo simulations give a completeness of $\sim 90\%$ in this flux density range (Aussel et al. 1999, 2001). All but one galaxy (HDF_PS3_6a) possess a spectroscopic redshift (98% complete). Their mean and median redshift is $\bar{z} \sim 0.8$. Except for one galaxy located at $z = 2.01$ (HDF_PM3_5, classified as an AGN, see Sect. 5.2) and four galaxies closer than $z = 0.15$, all ISOCAM galaxies have a redshift between $z \sim 0.4$ and 1.3 (see Fig. 6b), as expected from the Fig. 2.

5.2. Fraction of AGN galaxies

Cohen et al. (2000) classified the optical spectra of 2/3 of the ISOCAM galaxies and found only two galaxies presenting broad emission lines typical of AGNs:

HDF_PM3_6 (C36367_1346, $z = 0.960$),

HDF_PM3_20 (C36463_1404, $z = 0.962$).

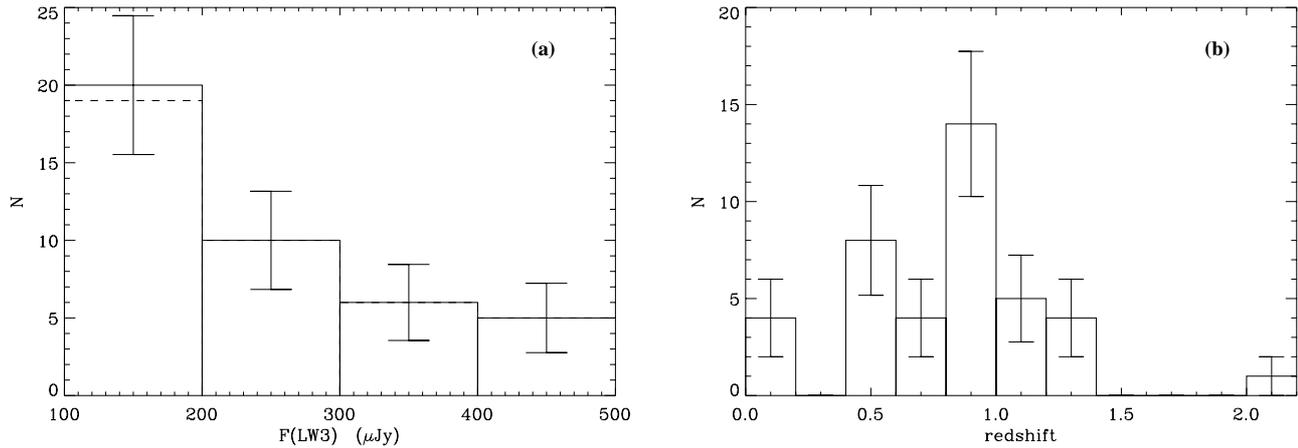


Fig. 6. a) Histogram of the flux densities of the 41 ISOCAM-HDFN galaxies ($S_{15} \geq 0.1$ mJy, Aussel et al. 1999, 2001) with Poissonian error bars. Only one galaxy has no spectroscopic redshift (dashed line). **b)** Redshift distribution of the 40 galaxies with spectroscopic redshift (98% complete in redshift, Poissonian error bars).

However, dusty AGNs may not be easily identified in the optical.

The best technique at present for the identification of dusty AGNs is offered by hard X-ray observations, which are much less affected by extinction. The deepest soft plus hard X-ray observation currently available was performed with the Chandra X-ray Observatory in the HDFN by Hornschemeier et al. (2001) and Brandt et al. (2001, 2002). The Chandra and ISOCAM catalogs were cross-correlated by Fadda et al. (2002). A total of 16 ISOCAM-HDFN ($S_{15} \geq 0.1$ mJy) are detected in the 1 Ms Chandra image at the depth of $S[0.5-2 \text{ keV}] \simeq 3 \times 10^{-17} \text{ erg cm}^{-2} \text{ s}^{-1}$ and $S[2-8 \text{ keV}] \simeq 2 \times 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1}$. They represent 39% of the ISOCAM-HDFN galaxies. However, the Chandra image is so deep that it is able to detect starburst galaxies up to large redshifts and among the 16 ISOCAM galaxies in common with Chandra sources, only 5 are classified as being AGNs or AGN dominated by Fadda et al. (2002). These 5 sources include the two galaxies previously identified as AGNs from their optical spectra and the following 3 galaxies:

HDF_PM3_5 (J123635.6+621424, $z = 2.010$),
HDF_PS3_10 (J123642.1+621545, $z = 0.857$),
HDF_PM3_21 (J123646.4+621529, $z = 0.851$).

Galaxies detected in the ultra-hard X-ray band (4–8 keV) or with $L_X \geq 10^{43} \text{ erg s}^{-1}$ were classified as AGNs. Galaxies with $L_X \leq 10^{40} \text{ erg s}^{-1}$ or which could be fitted with the SED of M 82 or Arp 220 were classified in the starburst category.

Finally only $(12 \pm 5)\%$ (5/41 galaxies, Poissonian error bar) of the ISOCAM-HDFN galaxies are found to be powered by an active nucleus. These AGNs contribute to $(18 \pm 7)\%$ of the sum of the flux densities (the IGL) of all 41 ISOCAM-HDFN galaxies. Since the active nuclei in these galaxies may not be responsible for the totality of the IR emission, this result suggests that at most $(18 \pm 7)\%$ of the $15 \mu\text{m}$ IGL due to galaxies with $S_{15} = 0.1-0.5$ mJy is produced by AGN activity.

The $S_{15} = 0.5-3$ mJy flux density range is covered by the ISOCAM image of the Lockman Hole, where Fadda et al. (2002) cross-correlated ISOCAM with the XMM-Newton X-ray observatory catalog of sources brighter than $S[2-10 \text{ keV}] \simeq 1.4 \times 10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}$ and $S[5-10 \text{ keV}] \simeq 2.4 \times 10^{-15} \text{ erg cm}^{-2} \text{ s}^{-1}$ (Hasinger et al. 2001). Among 103 ISOCAM sources with a flux density ranging from 0.5 to 3 mJy, 13 sources were found to be AGN dominated by Fadda et al. (2002), i.e. $(13 \pm 4)\%$. Such galaxies could be fitted by the SEDs of either type-I or type-II AGNs by Franceschini et al. (2002). These AGNs contribute to $(15 \pm 5)\%$ of the $15 \mu\text{m}$ IGL due to galaxies in the 0.5–3 mJy flux range.

The combination of the HDFN and Lockman Hole cover a flux density range of 0.1–3 mJy where 70% of the $15 \mu\text{m}$ IGL is measured. In this range, these two fields suggest that AGNs contribute to at most $(17 \pm 6)\%$. It must be noted that the XMM-Newton and Chandra experiments were limited to energies below 10 keV. Some harder X-ray AGNs may not have been detected by these instruments yet and may be common with our ISOCAM source list. The presence of such objects will remain highly uncertain until the launch of the next generation X-ray experiments (XEUS, Constellation X). We will assume in the following that the AGN fraction computed in the HDFN (which is larger than in the Lockman Hole) applies for the $15 \mu\text{m}$ IGL as a whole.

5.3. Rest-frame $15 \mu\text{m}$ luminosity of ISOCAM galaxies

In this section, we compare three different techniques to estimate the rest-frame $15 \mu\text{m}$ luminosity of ISOCAM galaxies. The first and simplest technique consists in using a single template SED to compute the k -correction. We chose the SED of M 82 (Fig. 1a) which falls on the correlations plotted in the Fig. 5 (crosses in Fig. 7b). As in Sect. 2, we decided not to use the SED of Arp 220 (Fig. 1a) because it is well outside the correlations found

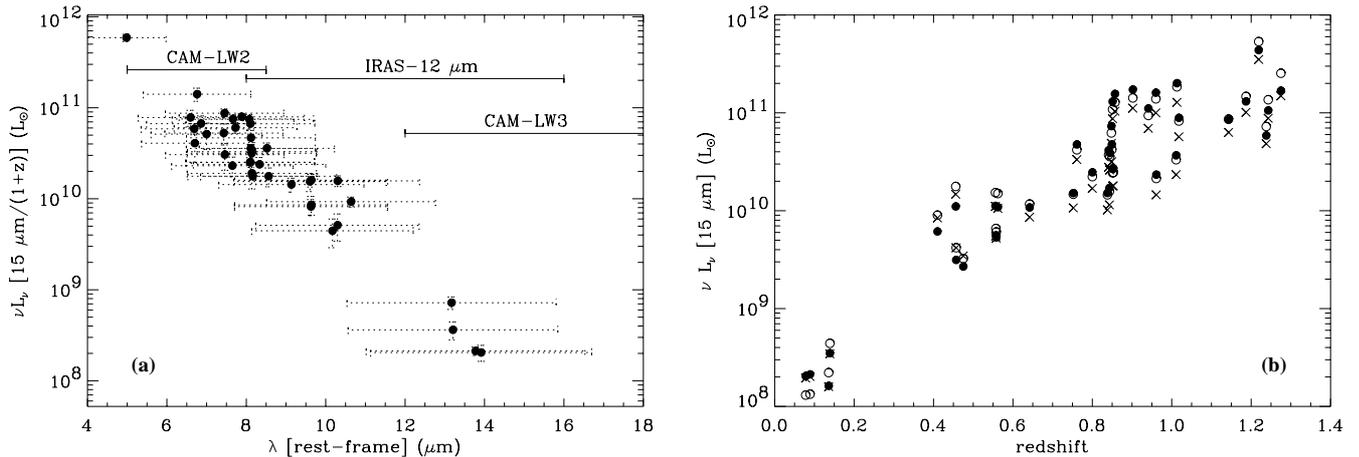


Fig. 7. Rest-frame MIR luminosity of ISOCAM-HDFN galaxies. **a)** luminosity of ISOCAM galaxies at $\lambda_{\text{rf}} = 15/(1+z) \mu\text{m}$ as a function of λ_{rf} . The dotted lines represent the redshifted LW3 (12–18 μm) band. **b)** rest-frame 15 μm luminosity (νL_{ν}) versus redshift estimated with three different techniques: M 82 SED (crosses), “closest filter” (filled dots) and “multi-template” (open dots).

for galaxies of similar luminosities with its large FIR over MIR luminosity ratio.

In a second technique, that we will call “closest filter”, we use directly the correlations described in Sect. 4 (Eqs. (11) and (12)) between the rest-frame 15 μm and 6.75 or 12 μm luminosities. We first calculate the rest-frame wavelength $\lambda_{\text{rf}} = 15/(1+z) \mu\text{m}$ and find which of the three filters ISOCAM-LW3, ISOCAM-LW2 and IRAS-12 has its central wavelength the closest to λ_{rf} (see Fig. 7a). Then we apply the formula given in the Sect. 4.2.2 for this “closest filter” to compute the rest-frame 15 μm luminosity (filled dots in Fig. 7b). Since most ISOCAM-HDFN galaxies are located around the median redshift $\bar{z} \sim 0.8$, the correlation between LW3 and LW2 (Fig. 5d, Eq. (11)) is most often used.

The third technique is the multi-template technique already used in Sect. 2.4, where we use the library of 105 template SEDs from Chary & Elbaz (2001). For a given galaxy of flux density S_{15} located at a given redshift z , we have computed S'_{15} for each one of the 105 template SED if it were located at the same redshift z . The template for which the computed S'_{15} was the closest to the observed S_{15} was then used to compute the rest-frame 15 μm luminosity, $\nu L_{\nu}[15 \mu\text{m}]$ (open circles in Fig. 7b), after a normalization by a factor S_{15}/S'_{15} .

The rest-frame 15 μm luminosities estimated using the three techniques for the ISOCAM-HDFN galaxies are consistent within 20% (1- σ , see Fig. 7b). Above $z \sim 0.6$, the first technique, which makes use of the SED of M 82, systematically underestimates $\nu L_{\nu}[15 \mu\text{m}]$, as compared to the other two techniques. This was to be expected since M 82 is a moderately luminous IR galaxy with a lower FIR over 6.75 μm luminosity ratio than more luminous galaxies (see Fig. 5d).

In the following, we will use the multi-template technique to estimate the IR luminosity and SFR of galaxies, because it is consistent with the correlations between MIR

and FIR luminosities and provides a continuous interpolation between the 6.75, 12, 15 μm and FIR luminosities based on observed spectra. Since the templates fit the correlation between L_{IR} and $\nu L_{\nu}[15 \mu\text{m}]$, it is equivalent to compute $\nu L_{\nu}[15 \mu\text{m}]$ first and then use Eq. (13) which links L_{IR} and $\nu L_{\nu}[15 \mu\text{m}]$ or to estimate directly L_{IR} with the best template SED.

5.4. Luminosities and star formation rates

In this section, we use the multi-template technique to directly estimate the L_{IR} and SFR of the ISOCAM-HDFN galaxies. The results are shown in the Fig. 8 and Table 3. The 1- σ error bars in Fig. 8 were computed from the dispersion measured in the correlations of the Fig. 5. For example, for a galaxy located at $z \simeq 1$, since λ_{rf} is close to 7 μm , the error bar on L_{IR} is the quadratic sum of the error bar in Eq. (11), which gives L_{15} from L_7 , and of the error bar in Eq. (13), which gives L_{IR} from L_{15} .

The dashed line in Fig. 8 corresponds to the sensitivity limit of $S_{15} = 0.1$ mJy converted into a $L_{\text{IR}}(z)$ as in the Fig. 2. Table 3 summarizes the properties of the ISOCAM galaxies in each of these four galaxy categories plus the AGN category.

About 75% of the galaxies dominated by star formation are either LIGs or ULIGs, while the remaining $\sim 25\%$ are nearly equally distributed among either “starbursts” ($10^{11} > (L_{\text{IR}}/L_{\odot}) \geq 10^{10}$) or “normal” galaxies ($L_{\text{IR}} < 10^{10} L_{\odot}$). The contribution of each galaxy type to the 15 μm IGL (IGL_{15}) is given in the last column of the Table 3. Luminous IR galaxies are responsible for about 60% of IGL_{15} while AGNs contribute to about 20%, the remaining 20% being due to the normal and starburst galaxy types. This repartition will be used in Sect. 6 to compute the 140 μm IGL (IGL_{140}) from ISOCAM galaxies.

Table 3. Properties of the ISOCAM galaxies detected above 0.1 mJy in the HDFN+FF (40 galaxies with spectroscopic redshift within a total of 41 galaxies). Column 1: the galaxies have been divided into 5 categories (see text). Column 2: median redshift. Column 3: logarithm of the median $15\ \mu\text{m}$ luminosity ($L_{15} = \nu L_\nu[15\ \mu\text{m}]$). Column 4: logarithm of the median IR luminosity (8–1000 μm), calculated from the multi-template technique. Column 5: median star formation rate, calculated from L_{IR} using Eq. (1). The $1\text{-}\sigma$ error bars in Cols. 3–5 are the median of the $1\text{-}\sigma$ error bars for each individual galaxy. They characterize the uncertainty on L_{15} , L_{IR} and $\langle SFR \rangle$ for individual galaxies, not their dispersion around the median value. Column 6: number of galaxies in each category. Column 7: relative contribution of each galaxy type to the $15\ \mu\text{m}$ IGL (IGL_{15}) produced by the sum of all 40 galaxies. The mean and median values given in the last two lines concern only the 35 star forming galaxies.

Type	\bar{z}	$\log(L_{15})$ (L_\odot)	$\log(L_{\text{IR}})$ (L_\odot)	$\langle SFR \rangle$ (M_\odot/yr)	Nb	% IGL_{15}
ULIGs	1.2	11.3 ± 0.1	12.3 ± 0.2	332 ± 148	6	17
LIGs	0.8	10.5 ± 0.1	11.5 ± 0.2	56 ± 26	20	44
Starbursts	0.5	9.8 ± 0.1	10.7 ± 0.2	9 ± 4	5	9
Normal	0.1	8.3 ± 0.1	9.4 ± 0.2	0.4 ± 0.2	4	11
AGNs	0.9	11.1 ± 0.7	11.5 ± 0.7	-	5	19
Median	0.8	10.4 ± 0.1	11.4 ± 0.2	41 ± 19	35	100
Mean	0.8	10.8 ± 0.1	11.8 ± 0.2	109 ± 51	35	100

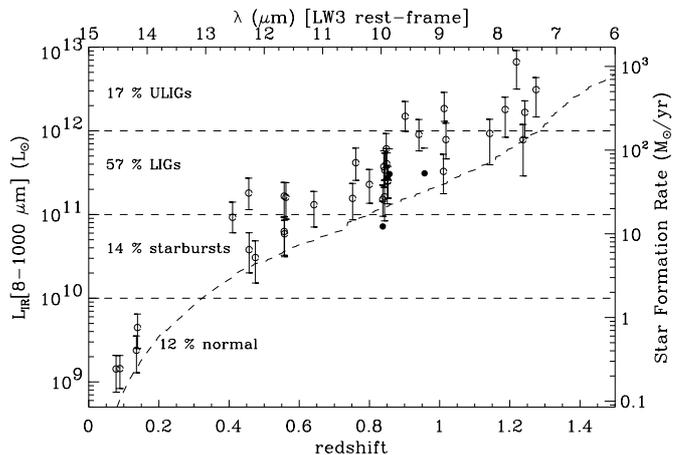


Fig. 8. $L_{\text{IR}}[8\text{--}1000\ \mu\text{m}]$ and star formation rate (SFR) versus redshift and λ_{rf} for the ISOCAM-HDFN galaxies ($S_{15} \geq 0.1$ mJy). The $15\ \mu\text{m}$ completeness limit of 0.1 mJy is materialized with a dashed line. Only the left axis is meaningful for the five AGNs (filled dots, Sect. 5.2). HDFN_PS3_6a (no spectroscopic redshift) is not included.

The mean IR luminosity of the ISOCAM-HDFN galaxies is $L_{\text{IR}} \simeq 6 \times 10^{11} L_\odot$. It corresponds to a $SFR \simeq 100 M_\odot \text{yr}^{-1}$ (from Eq. (1)). In the following section, we evaluate the robustness of this result with an independent estimator of the IR luminosity: the radio luminosity.

5.5. The radio-MIR correlation

The radio continuum is also a tracer of the bolometric IR luminosity of star forming galaxies as a result of the tight correlation between IR and 1.4 GHz radio continuum in local galaxies (Condon 1992; Yun et al. 2001; see Sect. 2.4). The origin of this correlation remains unclear, but it is generally assumed that massive stars are both responsible for the UV photons that heat dust before it

radiates in the IR and for the synchrotron acceleration of electrons, producing the radio continuum, when they explode as supernovae.

In Fig. 9a, we have plotted the 1.4 GHz and $15\ \mu\text{m}$ luminosities of the 109 local galaxies (small filled dots) from the ISOCAM sample described in Sect. 4.2.1, which were detected in the NRAO VLA Sky Survey (NVSS, Condon et al. 1998). As expected both luminosities are correlated with each other since both the 1.4 GHz and $15\ \mu\text{m}$ luminosities are correlated with L_{IR} .

Half of the 35 star forming HDFN-ISOCAM galaxies with spectroscopic redshift (AGNs excluded) present flux densities larger than $40\ \mu\text{Jy}$ at 1.4 GHz ($5\text{-}\sigma$) in the VLA and WSRT catalogs from Richards (2000) and Garrett et al. (2000) respectively. Their rest-frame $15\ \mu\text{m}$ luminosities were computed in the Sect. 5.3. Their rest-frame 1.4 GHz luminosities were computed assuming a power-law as in Sect. 2.4: $S_\nu \propto \nu^{-\alpha}$, where $\alpha = 0.8 \pm 0.15$ as suggested for star forming galaxies in Yun et al. (2001). The rest-frame $15\ \mu\text{m}$ and 1.4 GHz luminosities of these 17 HDFN-ISOCAM galaxies are plotted as filled dots with error bars in the Fig. 9a. The error bars on $\nu L_\nu[1.4\ \text{GHz}]$ were computed from the quadratic sum of the error bars on the measurement of the radio flux densities plus the error bar on α . Nine galaxies common in both catalogs from Richards (2000) and Garrett et al. (2000) present up to 40% differences in their radio flux densities. For these galaxies, we used the mean value and included the difference between both measurements in the error bars.

We have also included 7 galaxies from the CFRS-14 field (Flores et al. 1999), which rest-frame luminosities were computed with the same techniques (open dots).

The MIR and radio luminosities of this sample of 24 distant dusty galaxies ($\bar{z} \simeq 0.8$) are also strongly correlated with each other. This is the first time that such correlation is found at these redshifts. The plain line in the Fig. 9a is a power-law fit to this correlation. For comparison, we have plotted in dashed line the correlation that

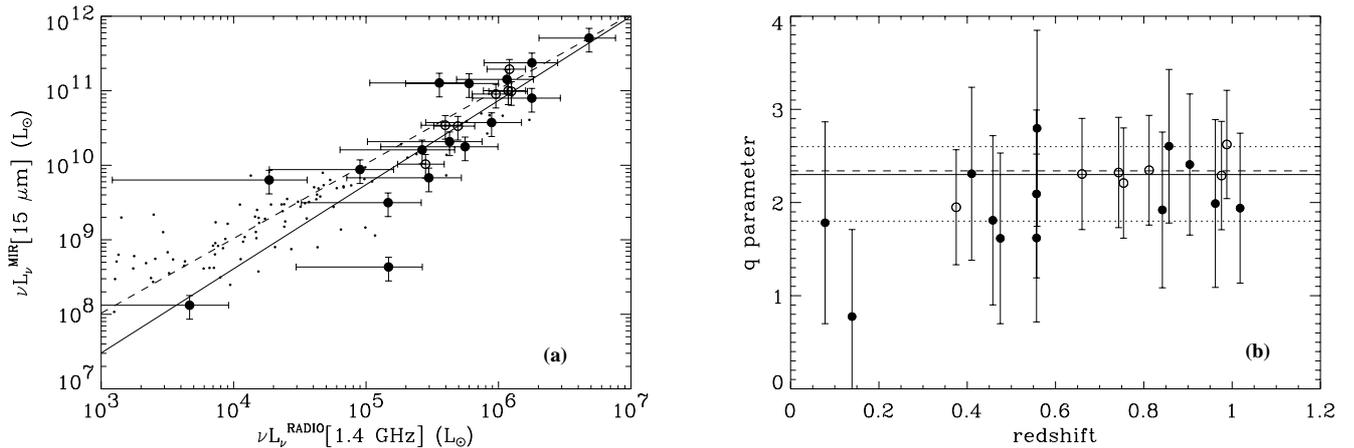


Fig. 9. **a)** 15 μm versus radio continuum (1.4 GHz) luminosities. Small filled dots: sample of 109 local galaxies from ISOCAM and NVSS. Filled dots with error bars: 17 HDFN galaxies in common between ISOCAM and VLA or WSRT (see text). Open dots with error bars: 7 CFRS-14 galaxies in common between ISOCAM and VLA (Flores et al. 1999). **b)** “ q ” parameter as a function of redshift for the HDFN and CFRS-14 galaxies. Plain line: median value of “ q ” ($2.3^{+0.3}_{-0.5}$). Dashed line: measured value from local galaxies ($q = 2.34 \pm 0.01$; Yun et al. 2001). Dotted line: 1- σ error bar on “ q ” for the 24 galaxies.

one would expect from the combination of Eq. (13) (L_{IR} from L_{15}) and the radio “ q ” parameter (see Sect. 2.4). The slope of the correlation observed for distant galaxies is marginally steeper than this dashed line. It is also slightly steeper than the correlation found for local galaxies (small filled dots), but the difference is too marginal with respect to the number of galaxies at high luminosities for detailed interpretation.

In Fig. 9b, the “ q ” parameter for the 24 ISOCAM-HDFN and CFRS-14 galaxies is plotted as a function of redshift. The “ q ” parameter is the logarithm of the ratio of L_{FIR} over $L(1.4 \text{ GHz})$ (see Eq. (2)). L_{IR} was estimated as described in Sect. 5.4, then converted into L_{FIR} from Eq. (4). The median value of the “ q ” parameter for the 24 galaxies is: $q = 2.3^{+0.3}_{-0.5}$ (plain and dotted lines in Fig. 9b), in perfect agreement with the local value of $q = 2.34 \pm 0.01$ (Yun et al. 2001, dashed line in Fig. 9b).

This study shows that the IR luminosities estimated from the ISOCAM 15 μm flux densities using the MIR-FIR correlations are perfectly consistent with those estimated from the radio. Although it is not clear whether the radio versus IR correlation also applies up to $z \sim 1$, this result independently validates our estimate of the bolometric IR luminosity of the ISOCAM galaxies.

5.6. Evolution of the infrared volume emissivity with redshift

We have seen that the excess of faint MIR sources in number counts was due to the presence of distant ($z \sim 1$) luminous IR galaxies. As a consequence, the amount of star formation per comoving volume hidden by dust must have rapidly decreased from $z \sim 1$ to 0. Indeed, the ISOCAM-HDFN galaxies with $0.6 \leq z \leq 1.3$ and $L_{\text{IR}} \geq 10^{11} L_{\odot}$ (LIGs and ULIGs; AGNs excluded) produce a 15 μm luminosity density of $\mathcal{L}_{15}(\bar{z} \sim 1) \simeq (60 \pm 12) \times 10^6 L_{\odot} \text{ Mpc}^{-3}$, while in the local universe luminous IR galaxies only make

$\mathcal{L}_{15}(z \sim 0) \simeq (11 \pm 3) \times 10^5 L_{\odot} \text{ Mpc}^{-3}$ (see Table 4). The comoving luminosity density produced by luminous IR galaxies at 15 μm was therefore about 55 times larger at $z \sim 1$ than in the local universe. The 15 μm luminosity densities at $z \sim 0$ were calculated using the local luminosity function (LLF) from Xu et al. (1998). In order to estimate the contribution of luminous IR galaxies, we used the Eq. (13) to convert the IR luminosity cut-off of $L_{\text{IR}} = 10^{11} L_{\odot}$ into a 15 μm luminosity cut-off of $L_{15} \sim 10^{10} L_{\odot}$. The 15 μm luminosity density that we have computed is consistent with the one measured at 12 μm from IRAS, $\mathcal{L}_{12}^{z \sim 0} \sim (1.7 \pm 0.5) \times 10^6 L_{\odot} \text{ Mpc}^{-3}$, converted to 15 μm using Eq. (12).

If we now consider the bolometric IR luminosity (from 8 to 1000 μm) of ISOCAM galaxies, as estimated in the Sect. 5, we find that the comoving density of IR luminosity radiated by dusty starbursts was about (70 ± 35) times larger at $\bar{z} \sim 1$ than today (computed from the LLF of Soifer et al. 1987). Since the IR luminosity is directly proportional to the extincted star formation rate of a galaxy, this means that the comoving density of star formation taking place in luminous IR galaxies was about (70 ± 35) times larger at $z \sim 1$ than today. In case of a pure density evolution proportional to $(1+z)^n$, this would translate into a value of $n \simeq 6$. However, we want to emphasize that we are only considering here the galaxies detected by ISOCAM, not the full luminosity function. For comparison, the B -band (0.44 μm) luminosity density was only about three times larger at $z \sim 1$ than today.

The redshift evolution of the comoving density of IR luminosity is compared to the UV (2800 Å) one in Fig. 10a. Both wavelengths exhibit similar luminosity densities at low redshift but the IR rises faster than the UV and reaches a larger value at $z \sim 1$, implying that a much larger fraction of star formation was hidden by dust at $z \sim 1$ than today.

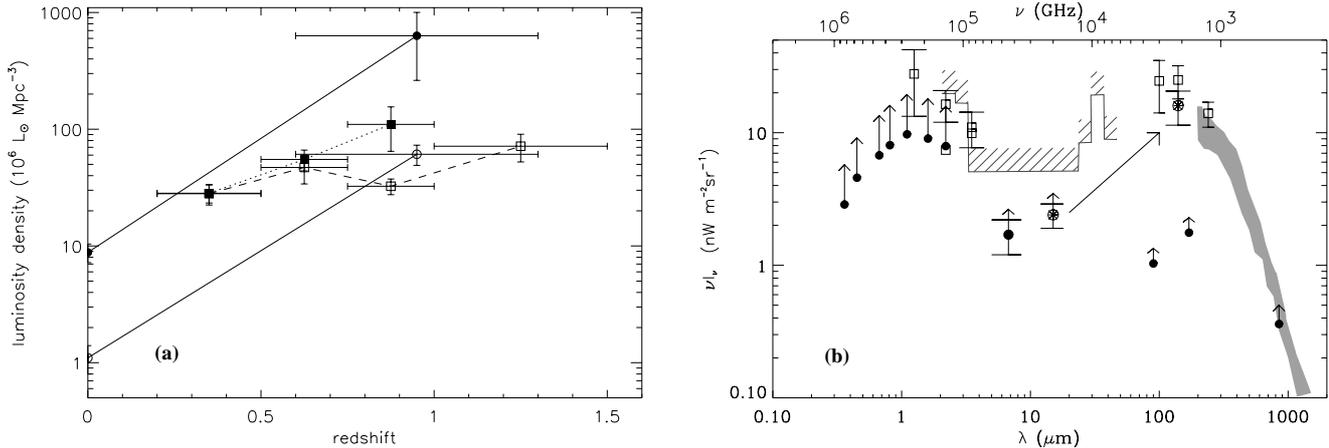


Fig. 10. **a)** FIR (filled circles), MIR-15 μm (open circles) and UV-2800 \AA luminosity density (in $L_{\odot} \text{ Mpc}^{-3}$) as a function of redshift. UV-2800 \AA : open squares from Cowie et al. (1999); filled squares from Lilly et al. (1996). **b)** Integrated Galaxy Light (IGL, filled dots) and Extragalactic Background Light (EBL, open squares, grey area) from the UV to sub-millimeter (see Table 5 for the origin of the data). EBL measurements from COBE: 200-1500 μm EBL from COBE-FIRAS (grey area, Lagache et al. 1999), 1.25, 2.2, 3.5, 100, 140 μm EBL from COBE-DIRBE (open squares, see Table. 5). IGL in the U, B, V, I, J, H, K bands from Madau & Pozzetti (2000). The upper end of the arrows indicate the revised values suggested by Bernstein et al. (2001, factor two higher). Our estimate of the 15 μm IGL ($2.4 \pm 0.5 \text{ nW m}^{-2} \text{sr}^{-1}$) is marked with a surrounded star. 6.75 μm (ISOCAM-LW2 filter) IGL from Altieri et al. (1999, filled dot). Hatched upper limit from Mrk 501 (Stanev & Franceschini 1998). The arrow from 15 to 140 μm indicates our computation of the 140 μm IGL due to ISOCAM galaxies.

Finally, we note that the projected density of galaxies detected in the B -band in the HDFN (529 galaxies/ μ^2 with $B_{\text{AB}} \leq 29$, Pozzetti et al. 1998) is 330 times greater than the projected density of ISOCAM galaxies (1.6 sources/ μ^2 , with $S_{15} \geq 0.1 \text{ mJy}$ or $AB(15 \mu\text{m}) \leq 18.9$). However the ISOCAM galaxies produce a 15 μm IGL which is only twice lower than the B -band IGL ($IGL_B \sim 4.57^{+0.73}_{-0.47} \text{ nW m}^{-2} \text{sr}^{-1}$, Madau & Pozzetti 2000). This confirms that the very luminous galaxies detected in the MIR radiate mostly in the IR.

6. Contribution of ISOCAM galaxies to the cosmic infrared background

Figure 10b summarizes the current estimates of the cosmic background for $\lambda = 0.1$ to $1500 \mu\text{m}$. Lower limits to the EBL (filled dots) come from integrating galaxy fluxes (IGL), while direct measurements of the EBL come from the COBE experiments DIRBE (open squares) and FIRAS (grey area, $\lambda > 200 \mu\text{m}$, Lagache et al. 1999). Table 5 gives the origin of the data used in Fig. 10b. Three independent studies obtain the same FIR to sub-millimeter EBL above $\lambda \sim 140 \mu\text{m}$. We did not consider here the EBL at $\lambda = 60 \mu\text{m}$ (Finkbeiner et al. 2000) because it remains controversial (see Puget & Lagache 2001).

Below $\lambda \sim 5 \mu\text{m}$, the cosmic background originates from direct stellar light, while above this wavelength, it comes from either stellar or AGN light reprocessed by dust. In the dust domain, i.e. from $\lambda \sim 5$ to $1500 \mu\text{m}$, the cosmic background peaks around $\lambda \sim 140 \mu\text{m}$. It is not directly measured over this whole wavelength range, but below $\lambda \sim 40 \mu\text{m}$ and down to about $2 \mu\text{m}$, an

Table 4. Comoving luminosity density in units of $10^6 L_{\odot} \text{ Mpc}^{-3}$ in the local universe and at $z \sim 1$ ($0.6 \leq z \leq 1.3$). Columns “All”: contribution of galaxies of all luminosities are considered. Columns “LIGs”: only galaxies with $L_{\text{IR}} \geq 10^{11} L_{\odot}$. AGNs were excluded for the computation of the IR values. The B -band ($0.44 \mu\text{m}$) values are from Lilly et al. (1996) at $z \sim 0$ and Connolly et al. (1997) at $z \sim 1$ (mean of the 0.5–1 and 1–1.5 redshift bins).

	15 μm		8–1000 μm		0.44 μm
	All	LIGs	All	LIGs	All
$z \sim 0$	7.5 ± 1.6	1.1 ± 0.3	120 ± 30	8.8 ± 1.6	82 ± 2
$z \sim 1$	-	61 ± 12	-	632 ± 291	240 ± 100

upper limit is set by the observations of the TeV outburst of Mrk 501 (see Sect. 3).

The 15 μm IGL ($IGL_{15} = (2.4 \pm 0.5) \text{ nW m}^{-2} \text{sr}^{-1}$, see Sect. 3) is marked with a surrounded star in Fig. 10b. It is about ten times lower than the peak of the CIRB at $\lambda_{\text{max}} \sim 140 \mu\text{m}$. At the median redshift of the ISOCAM galaxies, $\bar{z} \sim 0.8$ (Sect. 5.1), the rest-frame wavelengths corresponding to the observed 12–18 μm (ISOCAM-LW3 bandpass) and 140 μm wavelengths, are 6.5–9.8 μm and $\sim 80 \mu\text{m}$. These rest-frame wavelengths are close to the ISOCAM-LW2 (5–8.5 μm) and IRAS (60, 100 μm) bands. We have seen in the Fig. 5d that the luminosities of local galaxies in both bands correlate with each other. If this correlation remains valid up to $z \sim 1$, then it implies that we can compute the 140 μm IGL due to ISOCAM galaxies. There is no obvious reason why the SEDs of galaxies at $z \sim 1$ would exhibit very different IR SEDs. For example, the SED of the extremely red object HR 10 ($z \sim 1.44$) is very similar to the one of Arp 220 normalized by a factor of ~ 4 (Elbaz et al. 2002).

We propose here to use the multi-template technique, as in Sect. 5, to compute the contribution of ISOCAM galaxies to the $140\mu\text{m}$ background. We will first separate the relative contributions to IGL_{15} from ULIGs, LIGs, SBs, normal galaxies and AGNs as estimated in the Sect. 5. For each galaxy type, a median redshift, z_{med} , and the ratio $\mathcal{R} = \nu S_{\nu}[140/(1+z_{\text{med}})]/\nu S_{\nu}[15/(1+z_{\text{med}})]$ are estimated based on the HDFN sample (see Table 1). Although a larger sample would improve the statistical reliability of this computation, this choice is justified by the fact that the bulk of the $15\mu\text{m}$ IGL is produced by galaxies with flux densities in the range $S_{15} \simeq 0.35\text{--}1\text{ mJy}$ as the HDFN and CFRS-14 galaxies, which share similar redshifts distributions. It is important to separate the contribution from galaxies of different luminosities since the IR versus $6.75\mu\text{m}$ luminosity correlation presents a slope larger than one (Fig. 5d). The contribution from AGNs is computed assuming that they all share the SED of NGC 1068. This is a conservative hypothesis since NGC 1068 exhibits the lowest FIR over MIR ratio that we know.

The results are summarized in the Table 6. We first separate the contribution of each type of galaxy to the $15\mu\text{m}$ IGL (Cols. 3, 4). The $140\mu\text{m}$ IGL is then computed by converting each contribution to the $15\mu\text{m}$ IGL into a contribution at $140\mu\text{m}$, by multiplying by \mathcal{R} each time. We find a total of $IGL_{140} \simeq (16 \pm 5)\text{ nW m}^{-2}\text{ sr}^{-1}$ (see Fig. 10b), which corresponds to $(65 \pm 35)\%$ of the observed value from COBE-DIRBE ($EBL_{140} = (25 \pm 7)\text{ nW m}^{-2}\text{ sr}^{-1}$, Hauser et al. 1998; Lagache et al. 2000; Finkbeiner et al. 2000).

Hence the galaxies detected in ISOCAM deep surveys are found to be responsible for the bulk of the CIRB. About half of this $140\mu\text{m}$ IGL is due to LIGs and about one third from ULIGs. The contribution of AGNs is estimated to be as low as $\sim 4\%$ of the $140\mu\text{m}$ IGL. It could be as high as $\sim 20\%$ if instead of NGC 1068, we had used the SED of NGC 6240. In this case, ISOCAM galaxies would produce a $IGL_{140} \simeq (19 \pm 5)\text{ nW m}^{-2}\text{ sr}^{-1}$. If we had included the contribution from galaxies with flux densities between 30 and $50\mu\text{Jy}$ only detected through lensing magnification (see Sect. 3), IGL_{140} would be larger by an extra $\sim 4\text{ nW m}^{-2}\text{ sr}^{-1}$.

At the mean redshift of $z \simeq 0.6$ where the bulk of the B -band IGL is produced (Pozzetti & Madau 2001), light was radiated in the UV in order to be observed at $0.44\mu\text{m}$. As we have seen, the bulk of the $140\mu\text{m}$ is also produced in the same redshift range. Hence the $140\mu\text{m}$ and B -band probe the UV emission from stars with and without dust extinction. The ratio of IGL_{140} over IGL_B should therefore give a rough approximation of the ratio of extincted over non extincted star formation around $z \sim 0.5\text{--}1$. This ratio is close to 5, which suggests again that the bulk of the UV photons radiated by young stars in this redshift range was strongly affected by dust extinction. This value is consistent with the one obtained by Chary & Elbaz (2001) or Franceschini et al. (2001).

Table 5. Integrated Galaxy Light (IGL , upper half of the table) and Extragalactic Background Light (EBL, lower half of the table), νI_{ν} , in units of $(*)\text{ nW m}^{-2}\text{ sr}^{-1}$. (a) $1\text{-}\sigma$ error bars. (b),(c) include the systematic error quoted by the authors. Ref: (1) Madau & Pozzetti (2000), (2) Altieri et al. (1999), (3) Matsuhara et al. (2000), (4) Puget et al. (1999), (5) Barger, Cowie & Sanders (1999), (6) Blain et al. (1999b), lensed galaxies (l), (7) Wright (2000), (8) Cambr sy et al. (2001), (9) Dwek & Arendt (1998), (10) Gorjian et al. (2000), see also Wright & Reese (2000), Wright (2000), (11) Finkbeiner et al. (2000) and Hauser et al. (1998) for the 140 and $240\mu\text{m}$ values, (12) Lagache et al. (2000), (13) Lagache et al. (1999), Fixsen et al. (1998).

Filter	λ μm	$\nu I_{\nu}^{(*)}$	$[\sigma^+, \sigma^-]^{(a)}$	S_{min} mJy	AB mag	Ref.
<i>U</i>	0.36	2.87	[0.58,0.42]	0.023	28	1
<i>B</i>	0.45	4.57	[0.73,0.47]	0.009	29	1
<i>V</i>	0.67	6.74	[1.25,0.94]	0.002	30.5	1
<i>I</i>	0.81	8.04	[1.62,0.92]	0.009	29	1
<i>J</i>	1.1	9.71	[3.00,1.90]	0.009	29	1
<i>H</i>	1.6	9.02	[2.62,1.68]	0.009	29	1
<i>K</i>	2.2	7.92	[2.04,1.21]	0.23	25.5	1
<i>LW2</i>	6.75	1.7	0.5	0.03	20.2	2
<i>LW3</i>	15	2.4	0.5	0.05	19.7	
PHOT	90	1.03		1.5e3	11	3
PHOT	170	0.88		1.5e3	11	3
PHOT	170	1.76		1.2e3	11	4
SCUBA	850	0.36		2	15.6	5
SCUBA ^(l)	850	0.5	0.2	0.5	17.2	6
	1.25	28.9	16.3			7
	1.25	54.0	16.8			8
	2.2	27.8	6.7			8
<i>K</i>	2.2	7.4	6.9			9
<i>K</i>	2.2	22.4	6.0			10
	3.5	7.4	6.9			9
	3.5	11.0	3.3			10
DIRBE	60	28.1	8.8 ^(b)			11
DIRBE	100	24.6	10.5 ^(c)			11
DIRBE	100	23.4	6.3			12
DIRBE	140	25.0	7.0			11
DIRBE	140	24.2	11.6			12
DIRBE	240	14.0	3.0			11
DIRBE	240	11.0	6.9			12
FIRAS	850	0.5	0.2			13

7. Conclusions

The cosmic IR background is a fossil record of the light radiated by galaxies since their formation. It peaks around $\lambda_{\text{max}} \simeq 140\mu\text{m}$ whereas the spectral energy distribution of galaxies peaks above $60\mu\text{m}$. This suggests that the bulk of the cosmic IR background is due to galaxies located below $z \sim 1.3$. We have shown that the best technique currently available to unveil dusty galaxies up to $z \sim 1.3$ is provided by the ISOCAM MIR extragalactic surveys.

We have computed the contribution of ISOCAM galaxies to the $15\mu\text{m}$ background, the $15\mu\text{m}$ integrated galaxy light, and found a value of $IGL_{15} \simeq (2.4 \pm 0.5)\text{ nW m}^{-2}\text{ sr}^{-1}$. This is about ten times below the

Table 6. Contribution of the ISOCAM galaxies to the $140\ \mu\text{m}$ extragalactic background light. Column 1 galaxy type: AGNs and star forming galaxies above and below $\log(L_{\text{IR}}/L_{\odot}) = 10$, 11 and 12 (see Sect. 5.4). \bar{z} , in Col. 2, is the median redshift for each galaxy type. “%15”, in Col. 3, is the fractional contribution of each galaxy type to IGL_{15} (Eq. (7)). Column 4, IGL_{15} for each galaxy type in $\text{nW m}^{-2}\ \text{sr}^{-1}$. Column 5, $\mathcal{R} = \nu S_{\nu}[140\ \mu\text{m}]/\nu S_{\nu}[15\ \mu\text{m}]$. For AGNs, we used the SED of NGC 1068. IGL_{140} , in Col. 4, is the $140\ \mu\text{m}$ IGL produced by each galaxy type in $\text{nW m}^{-2}\ \text{sr}^{-1}$ and %140, in Col. 5, is the fractional contribution to the $140\ \mu\text{m}$ IGL for each galaxy type.

Type	\bar{z}	%15	IGL_{15}	\mathcal{R}	IGL_{140}	%140
ULIGs	1.2	17	0.4 ± 0.1	14 ± 6	5.7 ± 2.8	36
LIGs	0.8	44	1.1 ± 0.2	7 ± 3	7.7 ± 3.6	48
SBs	0.5	9	0.2 ± 0.1	4 ± 1	1.0 ± 0.4	6
normal	0.1	11	0.3 ± 0.1	4 ± 1	0.9 ± 0.5	6
AGNs	1.0	19	0.5 ± 0.1	1.4	0.7 ± 0.1	4
Total	0.8	100	2.4 ± 0.5	7 ± 3	16.0 ± 4.6	100

cosmic background measured by COBE at $\lambda_{\text{max}} \simeq 140\ \mu\text{m}$: $EBL_{140} \simeq (25 \pm 7)\ \text{nW m}^{-2}\ \text{sr}^{-1}$.

We have demonstrated that the MIR luminosities at 6.75 , 12 and $15\ \mu\text{m}$ were correlated with each other and with the bolometric IR luminosity for local galaxies. This suggests that the contribution of ISOCAM galaxies to the CIRB can be computed from IGL_{15} , unless distant galaxies SEDs strongly differ from local ones.

The redshift distribution of ISOCAM galaxies was measured from the spectroscopically complete sample of galaxies in the region of the HDFN. This redshift distribution is consistent with the twice larger sample of ISOCAM galaxies detected in the CFRS fields CFRS-14 and CFRS-03 (Flores et al. 1999, 2002). At the median redshift of $\bar{z} \simeq 0.8$, the observed 15 and $140\ \mu\text{m}$ wavelengths correspond to about $7\ \mu\text{m}$ (ISOCAM-LW2 filter) and $80\ \mu\text{m}$ (IRAS bands) in the rest-frame of the galaxies. Luminosities at both wavelengths are correlated (see Fig. 5d). If the correlations between MIR and FIR luminosities remain valid up to $z \sim 1$, then they can be used to compute IGL_{140} . We have checked with a sample of galaxies detected both in the MIR with ISOCAM and in the radio with the VLA and WSRT, that the MIR-FIR and radio-FIR correlations are consistent up to $z \sim 1$. This comparison independently validates our estimate of the bolometric IR luminosity of the ISOCAM galaxies, although it is not clear whether the radio-FIR correlation works also up to $z \sim 1$.

The fraction of active nuclei responsible for the $15\ \mu\text{m}$ luminosity of ISOCAM galaxies was estimated from the deepest soft and hard X-ray surveys available at present by the XMM-Newton and Chandra X-ray observatories in the Lockman Hole and HDFN respectively (Fadda et al. 2002). It was found that about $(12 \pm 5)\%$ of the ISOCAM galaxies are powered by an AGN and that the AGN contribution to IGL_{15} was about $(17 \pm 6)\%$. The AGN contribution to IGL_{140} was found to be as low as $\sim 4\%$ assuming that they

all share the SED of the local Seyfert 2, NGC 1068. This is a conservative choice since NGC 1068 presents the flattest IR SED that we know. However, we note that the cosmic X-ray background (CXB) peaks around $30\ \text{keV}$ (see Fig. 1 of Wilman et al. 2000), while both XMM-Newton and Chandra were limited to energies below $10\ \text{keV}$. It is therefore possible that a population of hard X-ray AGNs was missed by these surveys. But as claimed by the authors of these deep X-ray surveys, the bulk of the CXB had been resolved into individual galaxies in the Lockman Hole and HDFN images. Moreover, using estimates of the present comoving density of black holes, Madau & Pozzetti (2000) calculated that less than 20% of the CIRB could be due to dusty AGNs.

For the remaining star forming galaxies, we used a library of template SEDs, reproducing the MIR-FIR correlations, to compute their IR luminosity and contribution to IGL_{140} . We find that LIGs and ULIGs produce about 60% of IGL_{15} . The comoving density of IR luminosity produced by these luminous IR galaxies was about (70 ± 35) times larger at $z \sim 1$ than today, while in the same redshift interval the B -band or UV luminosity densities only decreased by a factor ~ 3 . Since the IR luminosity measures the dusty star formation rate of a galaxy, this also implies that the comoving density of star formation, due to luminous IR galaxies, decreased by a similar factor in this redshift range, i.e. much more than expected by studies at other wavelengths affected by dust extinction. This indicates that a large fraction of present day stars were formed during a dusty starburst event.

We estimate a contribution of ISOCAM galaxies to the peak of the CIRB at $\lambda_{\text{max}} \simeq 140\ \mu\text{m}$ of $(16 \pm 5)\ \text{nW m}^{-2}\ \text{sr}^{-1}$ as compared to the measured value of $(25 \pm 7)\ \text{nW m}^{-2}\ \text{sr}^{-1}$ from COBE. This study therefore suggests that the ISOCAM galaxies are responsible for the bulk of the CIRB.

We have started a systematic spectroscopic follow-up of these galaxies with the aim of studying their physical properties and the origin of their large SFR. Franceschini et al. (2001) estimated their baryonic masses to be of the order of $\langle M \rangle \sim 10^{11}\ M_{\odot}$ by fitting their optical and near-IR luminosities with template SEDs (from Silva et al. 1998) and assuming a Salpeter initial mass function (from 0.15 to $100\ M_{\odot}$). Their colors are similar to field galaxies of similar magnitudes (Cohen 2001), hence they could not have been selected on the basis of their optical colors. The technique which consists in using the spectral slope in the UV domain to correct luminosities from extinction (Meurer et al. 1999) fails to differentiate the luminous dusty galaxies detected with ISOCAM from other field galaxies in the HDFN+FF (Cohen 2001). This was to be expected since this technique only works for galaxies with $L_{\text{IR}} \lesssim 4 \times 10^{11}\ L_{\odot}$ (Meurer et al. 2000) while most ISOCAM galaxies are more luminous than this threshold. A property of the ISOCAM galaxies that may give a hint on their origin is their association with small groups of galaxies. A preliminary study of 22 ISOCAM-HDFN galaxies by Cohen et al. (2000) found that nearly

all ISOCAM galaxies belonged to small groups, while the fraction of field galaxies with similar optical magnitudes belonging to such groups was 68%. The study of the full sample of ISOCAM-HDFN galaxies by Aussel et al. (2001) shows that all of them belong to physical groups, hence suggests that dynamical effects such as merging and tidal interactions are responsible for their large SFR.

The Space IR Telescope Facility (SIRTF) will soon provide a powerful insight on the FIR emission of distant dusty starbursts and its $24\ \mu\text{m}$ band, less affected by confusion, should prolongate their detection up to $z \sim 2.5$. Extending the redshift range surveyed by ISOCAM to such redshifts is particularly important to measure the direct IR emission of the distant population of Lyman break galaxies, whose dust extinction remains highly uncertain (Steidel et al. 1999). However, the direct measurement of the MIR and FIR emission of distant galaxies will only become possible with the launch of the Next Generation Space Telescope (NGST) and Herschel (FIRST) satellite scheduled for 2009 and 2007.

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