Radio galaxies and magnetic fields in A514

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Abstract. A514 contains six extended and polarized radio sources located at various projected distances from the cluster center. Here we present a detailed study of these six radio sources in total intensity and polarization using the Very Large Array at 3.6 and 6 cm. Since the radio sources sample different lines of sight across the cluster, an analysis of the Faraday Rotation Measures (\(RMs\)) provides information on the strength and the structure of the cluster magnetic field. These sources show a decreasing Faraday Rotation Measure with increasing distance from the cluster center. We estimate the strength of the magnetic field to be \(\sim 3\,\mu G\) in the cluster center. From the \(R\)M structure across the stronger and more extended sources we estimate the coherence length of the magnetic field to be about 9 kpc at the cluster center.

Key words. galaxies: clusters: general – intergalactic medium – magnetic fields – radio continuum: general

1. Introduction

The study of the magnetic field in clusters of galaxies is relevant to understand the physical conditions and energetics of the intra-cluster medium (ICM), the cluster formation and evolution, the origin of the intergalactic magnetic field itself (primordial, Olinto 1997; injected into the ICM from galactic winds or AGNs, Kronberg et al. 1999; Völk & Atoyan 1999), and its consequences on primordial star formation (Pudritz & Silk 1989).

The presence of a synchrotron radio halo in a cluster of galaxies is a direct evidence for the presence of substantial magnetic fields in the cluster. Radio halos are produced by diffuse, non-thermal emission associated with the ICM rather than with a particular galaxy of the cluster, and have sizes of the order of 1 Mpc (e.g. Coma C, Giovannini et al. 1993).

So far, the studies of the magnetic fields in clusters of galaxies, obtained from different methods, give somewhat discrepant measurements for the field strengths.

Assuming minimum energy content, (i.e., nearly energy equipartition between the magnetic field and the relativistic particles), measurements of the strength of the magnetic fields in some radio halos have been made, and reveal magnetic field strengths of about 0.1–1 \(\mu G\) (Feretti & Giovannini 1996). In the Coma cluster, the equipartition magnetic field is \(\sim 0.4\,\mu G\) (Giovannini et al. 1993).

Estimates by Fusco-Femiano et al. (1999), obtained by comparing the diffuse radio emission with the hard X-ray emission in the Coma cluster (assumed to arise from Inverse Compton scattering of the cosmic microwave background of the relativistic electron responsible of the radio emission) give a field of 0.15 \(\mu G\).

Measurements of the magnetic field strength can also be determined in conjunction with X-ray observations of the hot gas, through the study of the Faraday Rotation Measure (\(RM\)) of radio sources located inside or behind the cluster.

Strong magnetic fields, from \(\sim 5\) up to the value of 30 \(\mu G\) have been found in “cooling flow” clusters (e.g. 3C 295, Perley & Taylor 1991, Allen et al. 2001; Hydra A, Taylor & Perley 1993) where extremely large Faraday rotations have been revealed, suggesting that the generation of very strong ICM magnetic fields may be connected with the cooling flow process (Soker & Sarazin 1990; Godon et al. 1998).

On the other hand, significant magnetic fields have also been detected in clusters without a cooling flow: the measurements of Faraday rotation of polarized radiation through the hot ICM leads to a magnetic field of 2–6 \(\mu G\) in the Coma cluster (Kim et al. 1990; Feretti et al. 1995). Similarly, the presence of a magnetic field was detected in the cluster A119 (Feretti et al. 1999) through the analysis...
of the Faraday Rotation in three extended radio galaxies located at different distances from the cluster center. We found that the RM decreases with the distance from the cluster center, and it is consistent with a field strength of about 5–10 \( \mu G \) at the cluster center, depending on the tangling scale. These values are quite similar to the field strength of about 5 \( \mu G \) found in the cluster containing the radio source 3C129 (Taylor et al. 2001) by similar analysis. From a statistical study of several clusters, Clarke et al. (2001) obtained magnetic fields of 4–8 \( \mu G \), confirming the previous findings.

By studying the spatial distribution of the RMs we can also estimate the coherence length of the magnetic field. The ICM magnetic field can be tangled on scales much smaller than the typical galaxy size. Crusius-Wätzl et al. (1990), studying the depolarization in 5 strong double sources, found tangling on scales of 1–4 kpc. This is confirmed also by the results on the Coma, A119 and 3C129 clusters (Feretti et al. 1995; Feretti et al. 1999; Taylor et al. 2001).

We note that the field strength measured from the RM may be influenced by the presence of filamentary structures (Eilek 1999) and by local turbulences in the ICM induced by the host galaxy motion within the cluster. Therefore the value estimated by RM arguments is generally higher than the equipartition magnetic field, which refers to the “average”, uniformly distributed field permeating the cluster (see also Goldsmith & Rephaeli 1993).

Here we present the analysis of the magnetic field in the cluster Abell 514, which is characterized by a clumpy and elongated X-ray surface brightness typically found in clusters associated with a merging process.

No radio halo is reported in the literature for this cluster. We analyzed the image taken from the NRAO VLA Sky Survey (NVSS, Condon et al. 1998), with a resolution of 45'' and a sensitivity of 0.45 mJy/beam (see contours in Fig. 1), in a region of about 5’ around the cluster center. No evidence for extended-diffuse low surface brightness emission at the noise level was found. Furthermore, we measured the total flux density in the aforementioned region and then we subtracted the contribution of each individual source. The resulting difference is consistent with the absence of diffuse emission.

However A514 contains six moderately strong, extended, and polarized sources located inside or behind the cluster and at different projected distances from the cluster center. Therefore this cluster is well suited for the analysis of the magnetic field through the study of the Faraday Rotation Measure.

We give a brief description of the cluster in Sect. 2. In Sect. 3 we describe the radio observations. In Sect. 4 we present the total intensity and polarization images for all the sources, and in Sect. 5 the X-ray image and the important X-ray parameters are analyzed. Finally in Sects. 6 and 7 we discuss the Rotation Measure results and the presence of magnetic fields. We assume \( H_0 = 50 \text{ km s}^{-1} \text{ Mpc}^{-1} \) and \( q_0 = 0.5 \) throughout the paper. At the distance of A514 (luminosity distance \( D_L \approx 435.5 \text{ Mpc} \) ) 1'' corresponds to 1.84 kpc.

2. Abell 514

A514 (\( z = 0.0714 \) Fadda et al. 1996) is of richness class 1 (Abell et al. 1989). It is classified as BM type II-III (Bautz & Morgan 1970) and as Rood-Sastry type F (Rood & Sastry 1971).

Figure 1 shows the radio image at 1.4 GHz obtained from the NVSS, overlaid on the X-ray image obtained by us from the ROSAT PSPC archive. The reference point of the image (indicated with a cross in the figure) was chosen to be the optical center of the cluster indicated by Abell et al. (1989) at RA(J2000) = 04\textdegree 47\textquoteright 40\textsec, DEC(J2000) = \( -20\textdegree 25.7\textsec \). The X-ray surface brightness is rather smooth, without any outstanding peak. We adopted as X-ray center of the cluster the highest X-ray brightness peak not related to any point-like emission, at position RA(J2000) = 04\textdegree 48\textquoteright 13\textsec, DEC(J2000) = \( -20\textdegree 27.18\textsec \). The X-ray center of the cluster appears shifted to the East with respect the optical center of about 8’. In Fig. 1 the X-ray center of the cluster is indicated with a star. In Sect. 5 we will discuss in greater detail the X-ray emission of A514 and its implications on the parameter estimates.

Here we examine multi-wavelength Very Large Array (VLA\(^1\)) observations of six radio sources within or behind cluster Abell 514 indicated by labels in Fig. 1. The two radio galaxies which lie close to the cluster center (B1, B2) show a narrow-angle-tailed structure, the radio galaxy D is a tailed, while the other two radio galaxies (A, C) are FRII type radio galaxies. The source E located in the cluster periphery is classified as a quasar.

We derive the Rotation Measures (RM) of the radio sources, which, together with the X-ray data, provide information about the magnetic field in the cluster.

3. Radio observations

The radio sources were observed with the VLA at two frequencies within the 6 cm band and at two frequencies within the 3.6 cm band with a bandwidth of 50 MHz. In Table 1 we provide details of the observations. We note that given the southern declination of A514, the VLA was used in hybrid configurations (BnA, CnB) in order to have roughly circular restoring beams.

The phase calibration was performed using the secondary calibrator 0457+234. The flux-density scale and the absolute polarization position angle were calibrated by observing 3C 48. The instrumental polarization of the antennas was corrected using 0457+234 and also 0530+135 in CnB configuration observed over a wide range of parallactic angles.

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Fig. 1. The 1.4 GHz radio image obtained from the NVSS (contour) superimposed to the X-ray image of Abell 514 taken from ROSAT PSPC archive (grey scale). Radio contour levels are 2, 5, 10, 20, 40, and 80 mJy/beam. The radio image has an angular resolution of 45\arcsec. The ROSAT image has an angular resolution of 25\arcsec (FWHM) and has been smoothed with a Gaussian of $\sigma = 30\arcsec$. The reference point of the image (indicated with a cross in the figure) was chosen to be the optical center of the cluster indicated by Abell et al. (1989). The X-ray center of the cluster is indicated with a star.

Images in all Stokes parameters were produced with the NRAO AIPS package following the standard procedures. Self calibration was applied to minimize the effect of amplitude and phase uncertainties of atmospheric and instrumental origin. The $(u, v)$ data at the same frequencies but from different configurations were first handled separately and then added together.

The images of the polarized intensity $P = (Q^2 + U^2)^{1/2}$, the degree of polarization $m = P/I$ and the position angle of polarization $\Psi = 0.5 \tan^{-1}(U/Q)$ were derived from the $I$, $Q$ and $U$ images. Total intensity images have been produced by averaging the two frequencies in the same band (4535/4885 MHz; 8085/8465 MHz) while $U$ and $Q$ images have been obtained for each frequency separately.

4. Total intensity and polarization images

In Table 2 we give the parameters for the total intensity images at 1.6\arcmin ($\approx$3 kpc) resolution. The total flux density was estimated after the primary beam correction. In all the radio images presented in this work contours are total intensity while vectors represent the orientation of the projected $E$-field and their length is proportional to the fractional polarization. In the polarization images, points with error in fractional polarization greater than 10% were clipped. In the following we present the individual sources.


In the Digitized Sky Survey (DSS) image from the UK Schmidt Telescope there is a very faint optical object coincident with the radio source. No redshift is available, but the object appears fainter by about 2 magnitudes than the other cluster members. Its peripheral position and optical properties suggest that it is a background source.

This radio source was observed by Owen et al. (1996) in a VLA Survey of radio source in Abell clusters. They excluded this object from their complete sample of radio galaxies in Abell Cluster because the source is at a projected distance larger than 0.3 Abell radii. Their radio observations were made with the VLA at 1.4 GHz in B and C array. They measured a total flux density of 270 mJy obtained by fitting the unresolved source with a Gaussian model.
### Table 1. Pointing and VLA observations of A514.

<table>
<thead>
<tr>
<th>Name</th>
<th>Other Name</th>
<th>Label</th>
<th>RA (J2000)</th>
<th>DEC (J2000)</th>
<th>Frequency (MHz)</th>
<th>Config.</th>
<th>Date</th>
<th>Duration (Hours)</th>
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<td>8085/8465</td>
<td>BnA, CnB</td>
<td>Oct.99, Feb.00</td>
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<td>BnA, CnB</td>
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<td>8085/8465</td>
<td>BnA, CnB</td>
<td>Oct.99, Feb.00</td>
<td>1.2</td>
</tr>
</tbody>
</table>

Column 1: source name; Col. 2: other name; Col. 3: label name; Cols. 4, 5: pointing position (RA, DEC); Col. 6: observation frequencies; Col. 7: VLA configuration; Col. 8: dates of observation; Col. 9: duration of observation.

**Fig. 2.** Left: image of J0447–2014 (A514A) at 3.6 cm with an angular resolution of 1.6″ × 1.6″. The rms noise level is 0.027 mJy/beam. The peak is 29.6 mJy/beam. Right: image of J0447–2014 (A514A) at 6 cm with an angular resolution of 1.6″ × 1.6″. The rms noise level is 0.028 mJy/beam. The peak is 47.2 mJy/beam. In both images contour levels are: −0.1, 0.1, 0.2, 0.4, 0.8, 1.6, 3.2, 6.4, 12.8, and 25.6 mJy/beam. The lines represent the orientation of the electric vector (E-field) and are proportional in length to the fractional polarization (1″ ≈ 11%).

Our observations resolve the radio source. In Fig. 2 we show the radio images at 3.6 cm and 6 cm with an angular resolution of 1.6″. The maximum projected angular size of the source is about 12″. At both frequencies, the radio source shows a double structure. An extension in the North-West side is visible in the images.

The source is polarized at 3.6 cm and 6 cm, with similar values of the polarization percentage. The fractional polarization is ≈5% in the northern lobe and ≈4% in the southern lobe. The extended emission in the North-West is polarized at 6 cm (about 10%) while at 3.6 cm the sensitivity is not enough to reveal any significant polarized emission and we only set an upper limit of about 13% to the fractional polarization.

In Fig. 3 we present the same source at 3.6 cm with the highest angular resolution. The image shows the unresolved nucleus at position: RA(2000) = 04h47m13.15s, DEC(2000) = −20°14'03". For this component we estimate a spectral index α ≈ 0.7 (where α is defined as $S(\nu) \propto \nu^{-\alpha}$); at this resolution, this quite steep spectrum in the compact component could be due to the contamination of other structures in the source. We found a spectral index α ≈ 0.9 and α ≈ 0.8 respectively in the northern and in the southern lobe.
Table 2. Parameters of total intensity radio images, with 1.6\" resolution.

<table>
<thead>
<tr>
<th>Source</th>
<th>$\lambda$ (cm)</th>
<th>$\sigma(I)$ (mJy/beam)</th>
<th>Peak brightness (mJy/beam)</th>
<th>Flux density (mJy)</th>
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<td>29.6</td>
<td>79.2</td>
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<tr>
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<td>$2.8 \times 10^{-2}$</td>
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<td>131.0</td>
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<td>$1.9 \times 10^{-2}$</td>
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<td>6.1</td>
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<td>51.4</td>
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<td>0.9</td>
<td>7.6</td>
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<tr>
<td></td>
<td>6</td>
<td>$2.4 \times 10^{-2}$</td>
<td>1.6</td>
<td>14.5</td>
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<tr>
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<td>7.0</td>
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<td>48.9</td>
<td>70.6</td>
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<td>6</td>
<td>$2.6 \times 10^{-2}$</td>
<td>65.3</td>
<td>105.9</td>
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</table>

Column 1: source name; Col. 2: observation wavelength; Col. 3: rms noise; Col. 4: peak brightness; Col. 5: flux density.

Fig. 3. Image of J0447–2014 (A514A) at 3.6 cm with an angular resolution of 0.8\" x 0.6\". Contour levels are: -0.07, 0.07, 0.1, 0.14, 0.20, 0.28, 0.40, 0.56, 0.79, 1.12, 1.58, 2.24, 3.17, 4.48, 6.34 mJy/beam. The rms noise level is 0.023 mJy/beam. The peak is 19.1 mJy/beam.

4.2. J0448–2026 (A514B1)

A 1.4 GHz image is presented in Owen et al. (1993) who measured a total flux density of 95 mJy. They considered this galaxy to be a member of the cluster.

In Fig. 4 we show the total intensity and polarization images at 3.6 cm and at 6 cm with an angular resolution of 1.6\". The source shows a Narrow-Angle Tail (NAT) structure and in projection it is located near the cluster center. In our 6 cm image the source has a size of 66\" (\pm 120 kpc).

At 3.6 cm the head is readily visible, while most of the low brightness emission in the tail is below the sensitivity of the present observations. The mean spectral index is $\alpha \approx 0.9$ in the head while it is steeper ($\alpha \approx 1.8$) in the tail. At both wavelengths the head is $\sim 4\%$ polarized, while the tail is polarized ($\sim 5\%$) at 6 cm but not detected in polarized flux at 3.6 cm. In the tail, at 3.6 cm, we quote an upper limit of 26\% to the fractional polarization. Due to its low polarization this source was not used to study the rotation measure.

4.3. J0448–2025 (A514B2)

Slee et al. (1994) identified the host galaxy as a D galaxy. The radio source morphology is that of a head tail source. The source was also studied by Owen et al. (1993, 1997) to be a member of the cluster. A 1.46 GHz image is shown in Owen et al. (1993) who calculated a total flux density of 119 mJy at this frequency.

In Fig. 5 we show the total intensity and polarization images obtained at 3.6 cm and at 6 cm. In both images the morphology of the radio galaxy is well defined as tailed elongated to the West. At 6 cm the total intensity radio emission reaches an extension of 50\", while at the same wavelength the polarized emission is detected only out to 25\" from the core. At both wavelengths the percentage of polarization is $\sim 4\%$ in the head. The polarization along the tail is $\sim 8\%$ at 6 cm. The mean spectral index is $\alpha \approx 0.5$ in the head while it becomes steeper in the brightest part of the tail, $\alpha \approx 1.2$.

4.4. J0448–2005 (A514C)

No reference for this radio source has been found in the literature and no optical identification is seen on the DSS. Since the host galaxies of radio sources are typically bright
Fig. 4. Left: image of J0448–2026 (A514B1) at 3.6 cm with an angular resolution of 1.6′′ × 1.6′′. The rms noise level is 0.019 mJy/beam. The peak is 4.1 mJy/beam. Right: image of J0448–2026 (A514B1) at 6 cm with an angular resolution of 1.6′′ × 1.6′′. The rms noise level is 0.018 mJy/beam. The peak is 8.2 mJy/beam. In both images contour levels are: −0.06, 0.06, 0.12, 0.48, 0.96, 1.92, 3.84, and 7.68 mJy/beam. Superimposed lines represent the orientation of the electric vector (E-field) and are proportional in length to the fractional polarization (1′′ ≈ 9%).

4.5. J0448–2032 (A514D)

This radio galaxy, belonging to A514 (Owen et al. 1993), is located in the South-East corner of the cluster, and appears as an FRI in the high resolution image, but is an extended tailed radio source (see Fig. 1). In Fig. 8 we show the images of J0448–2032 at 3.6 cm and at 6 cm. In both images the synthesized restoring beam is a circular Gaussian with a FWHM of 1.6′′. The radio images reveal a bright unresolved core and two oppositely directed jets. The jet to the North is well collimated up to about 20′′ where it flares, widening and bending slightly before diffusing into the Northern lobe. The jet to the South is much shorter (a few arcsec), and its emission blends rapidly into the Southern lobe. The structure of the southern lobe suggests a smooth bending at a small angle along the line of sight. The surface brightness of this lobe, and possibly the Northern lobe as well, fades gradually. The maximum projected angular size is about 85′′, and in the radio images presented here we do not detect the low brightness southern tail clearly visible in the NVSS image (see Fig. 1) and in the image at 1.46 GHz of Owen et al. (1993). At this frequency, the total flux density calculated by Owen et al. (1993) is 502 mJy.
Fig. 5. Top: image of J0448–2025 (A514B2) at 3.6 cm with an angular resolution of 1.6\arcsec × 1.6\arcsec. The peak is 9.1 mJy/beam. Bottom: image of J0448–2025 (A514B2) at 6 cm with an angular resolution of 1.6\arcsec × 1.6\arcsec. The peak is 11.4 mJy/beam. In both images contour levels are: −0.06, 0.06, 0.12, 0.24, 0.48, 0.96, 1.92, 3.84, and 7.68 mJy/beam and the rms noise level is 0.017 mJy/beam. Superimposed lines represent the orientation of the electric vector (E-field) and are proportional in length to the fractional polarization (1\arcsec ≈ 10\%).

In our images no polarized emission is detected in the core region. At both wavelengths the jet on the North side has an average polarization percentage of \(~7\%\) and the percentage of polarization increases along the jet with distance from the core. Although the southern jet is very short and extends only for about two resolution elements, the jet has an average polarization percentage of \(~5\%). This last measure is, however, uncertain given that the jet is rather short and is embedded in the southern lobe emission. Both lobes are strongly polarized at both 6 and 3.6 cm. At both wavelengths the lobe on the North side has a polarization percentage of about 12\%, and the lobe on the South has a polarization percentage of about 16\%.

We found an inverted spectrum ($\alpha \approx −0.1$) in the core and a mean spectral index $\alpha \approx 1.1$ in both brighter parts of the lobes.

4.6. J0448–2044 (A514E)

This source has been identified with a Quasar at a redshift 1.894 (Osmer et al. 1994). In Fig. 9 we show the images at 3.6 cm and at 6 cm. The maximum angular size of the source is \(55\arcsec\) corresponding to a linear size of about 450 kpc. A collimated knotty jet starts from the core toward the West. At a distance of about 25\arcsec from the core the jet bends to the North. The hot-spots on the West and
East sides are slightly misaligned with respect to the core at angles of 21° and 18° respectively if we assume that the source axis is defined by the jet. At the location of the knot in the jet the polarization at 6 cm increases. The jet on the East side is not visible.

Polarized emission at the level of ~2% is detected in the core. The polarization along the jet is detected at 6 cm and not at 3.6 cm, probably because of the sensitivity limit. At 3.6 cm we quote an upper limit of about 20% to the fractional polarization. At the jet bend the percentage of polarization is about 11% at both frequencies. At both wavelengths the hot-spot on the West side has a polarization percentage of about 18%, while the hot-spot on the East has a polarization percentage of about 7%.

We note that the source structure is characterized by the core (α ≈ 0.5) and two compact hot-spots (α ≈ 1.2 and 1.3 in the Western and Eastern hot-spots respectively), while the radio lobes are below our detection limit.

5. X-ray properties of Abell 514

From observations with the Einstein Imaging Proportional Counter, Jones & Forman (1999), show that the cluster A514 has at least three mass condensations with different X-ray luminosities. A514 was also studied by Ebeling et al. (1996) in a X-ray flux-limited sample of 242 Abell cluster of galaxies compiled from the ROSAT All Sky Survey (RASS) data in the soft X-ray energy band (0.1–2.4 KeV). They published a X-ray luminosity in the energy band 0.1–2.4 KeV of $1.44 \times 10^{44}$ erg/s and adopted the $kT - L_X$ relation to derive an estimated gas temperature of 3.6 keV.

The cluster A514 has been the target of X-ray observations with the ROSAT PSPC, for a total exposure time of 18,000 s. In Fig. 1 (grey scale) and in Fig. 10 (contours) we show the X-ray images in the 0.5–2 keV band obtained from the ROSAT public archive by binning the photon event table in pixels of 15′′ and by smoothing the image with a Gaussian of $σ = 30′′$. As already pointed out by Bliton et al. (1998), the X-ray brightness distribution of A514 is clumpy and elongated to the South-East direction.

Given its irregular and asymmetric X-ray emission, the cluster brightness distribution cannot be described by a hydrostatic isothermal model. However, in order to estimate the gas parameters necessary in the calculation of the magnetic field we attempted a fit to a portion of the cluster which looks approximately regular.

Therefore, part of the X-ray brightness image was fitted with a hydrostatic isothermal model (Cavaliere & Fusco-Femiano 1981):

$$ S(r) = S_0 (1 + r^2/r_c^2)^{-3\beta+0.5} + S_b $$

where $S_0$ is the central surface brightness, $r_c$ is the core radius, and $β$ is the ratio between the galaxy and the gas temperatures.
Fig. 8. Left: image of J0448-2032 (A514D) at 3.6 cm with an angular resolution of 1.6′ × 1.6′. The rms noise level is 0.018 mJy/beam. The peak is 7.0 mJy/beam. Right: image of J0448-2032 (A514D) at 6 cm with an angular resolution of 1.6′ × 1.6′. The rms noise level is 0.019 mJy/beam. The peak is 6.4 mJy/beam. In both images contour levels are: −0.06, 0.06, 0.12, 0.24, 0.48, 0.96, 1.92, and 3.84 mJy/beam. Superimposed lines represent the orientation of the electric vector (E-field) and are proportional in length to the fractional polarization (1σ ≈ 10%).

It is important to note that due to the cluster irregularity, the obtained β-model parameters are strongly dependent on the choice of the centroid and the portion of the X-ray surface brightness used in the fit. Different attempts were carried out to test reasonable values of the X-ray parameters. In Fig. 11 we show the radial profile of the X-ray surface brightness obtained by integrating the PSPC counts over concentric rings of 15″ in radius, after subtracting discrete X-ray sources and considering only the X-ray emission from North to South on the West side of the centroid. In this way we have excluded from the calculation all the irregular X-ray extension in the South-East side. As discussed in Sect. 2, the centroid of the X-ray emission was taken at the position of the X-ray peak RA(2000) = 04 h 48 m 13 s, DEC(2000) = −20°27′18″. This is approximately coincident with one of the X-ray peaks given by Jones & Forman (1999).

The best fit, reported in Fig. 11 as a solid line, only provides a rough approximation to the data (reduced χ² ≈ 3.9). From this fit we obtain a core radius \( r_c = 570^{+200}_{-130} \) kpc and \( \beta = 0.6^{+0.12}_{-0.07} \). The large uncertainty (at one sigma level of significance) obtained for the parameters reflects the irregular cluster structure.

Assuming a cluster temperature of 3.6 keV and varying the β-model parameters within their interval of confidence we obtain a central density in the range 0.44–0.50 \( 10^{-3} \) cm\(^{-3}\).

6. Rotation measure in A514

Polarized radiation from cluster and background radio galaxies may be rotated by Faraday effect if magnetic fields are present in the Intra Cluster Medium (ICM) together with substantial ionized material. In this case, the observed polarization angle (\( \Psi_\Lambda \)) is related to the intrinsic polarization angle (\( \Psi_{\text{int}} \)) through:

\[
\Psi_\Lambda = \Psi_{\text{int}} + RM \lambda^2
\]

where the Rotation Measure (\( RM \)) is related to the electron density, \( n_e \) and magnetic field along the line of sight \( H_\parallel \) through the cluster according to:

\[
RM = 812 \int_0^L n_e H_\parallel dl \ \text{rad/m}^2
\]

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where $L$ in kpc, is the path-length through the ICM, $H_{\parallel}$ is in $\mu G$ and $n_e$ in cm$^{-3}$.

The observed $RM$ in a radio galaxy is the sum of the contributions to the $RM$ from all magneto-ionic components along the line of sight. The $RM$ distribution of radio sources can therefore be used to derive information on the magnetic field along line of sights crossing different regions of the cluster.

We derived the images of the rotation measure using the position angles obtained at the 4 frequencies 4535, 4885, 8085 and 8465 MHz with a resolution of 1.6\". Following the definition in Eq. (2), the Faraday $RM$s were obtained by performing a least-squares fit of the polarization angle images at each pixel as a function of $\lambda^2$. The pixels in which the uncertainty in the polarization angle exceeds 10\" were blanked. For each target, except A514B1 where the polarized emission is too weak, an image of the $RM$ was obtained.

In Fig. 12 we give examples of $RM$ fits in the source A514D. The images and the histograms of the $RM$ distribution are presented in Figs. 14–18. In each caption we give a brief description of the $RM$ distribution.

In galactic coordinates A514 is located at $l = 219^\circ$ and $b = -36^\circ$. The Galactic contribution to the $RM$ in the region of A514 is expected to be about $-16$ rad/m$^2$ based on the average of the $RM$ galactic contribution published by Simard-Normandin et al. (1981) for sources near the cluster. The aforementioned value is consistent with the $RM$ we derive for the sources located in the cluster periphery. Hereafter we do not apply the correction of the galactic $RM$ since we are most interested in the differential $RM$ at various cluster locations.

In the Table 3 we report the maximum absolute value, the mean, and the $\sigma$ of the $RM$ values for the sources ordered according to an increasing projected distance from the X-ray cluster center. We expect a contribution to the $\sigma_{RM}$ from the noise in the measurements at a level of $\sim 10$ rad/m$^2$.

Table 3 shows that the innermost sources have the largest $\sigma_{RM}$ and highest absolute $RM$s, and that these values fall with increasing projected distance from the cluster center. This result is consistent with the interpretation that the external Faraday screen can be the same for all 5 sources, i.e., the radial profile of $\sigma_{RM}$ in
Table 3. Rotation measure.

<table>
<thead>
<tr>
<th>Name</th>
<th>Dist. cluster center (Mpc)</th>
<th></th>
<th>RM max</th>
<th>&lt;RM&gt;</th>
<th>(\sigma_{RM})</th>
<th>Cluster member</th>
</tr>
</thead>
<tbody>
<tr>
<td>A514B2</td>
<td>0.30</td>
<td></td>
<td>154</td>
<td>104</td>
<td>63</td>
<td>y</td>
</tr>
<tr>
<td>A514D (North)</td>
<td>0.73</td>
<td></td>
<td>120</td>
<td>25</td>
<td>54</td>
<td>y</td>
</tr>
<tr>
<td>A514D (South)</td>
<td>0.81</td>
<td></td>
<td>126</td>
<td>56</td>
<td>38</td>
<td>y</td>
</tr>
<tr>
<td>A514A</td>
<td>2.10</td>
<td></td>
<td>105</td>
<td>-19</td>
<td>48</td>
<td>n</td>
</tr>
<tr>
<td>A514E (right lobe)</td>
<td>2.22</td>
<td></td>
<td>34</td>
<td>6</td>
<td>35</td>
<td>n</td>
</tr>
<tr>
<td>A514E (core)</td>
<td>2.30</td>
<td></td>
<td>68</td>
<td>58</td>
<td>31</td>
<td>n</td>
</tr>
<tr>
<td>A514E (left lobe)</td>
<td>2.36</td>
<td></td>
<td>26</td>
<td>-31</td>
<td>27</td>
<td>n</td>
</tr>
<tr>
<td>A514C</td>
<td>2.55</td>
<td></td>
<td>15</td>
<td>-19</td>
<td>24</td>
<td>n</td>
</tr>
</tbody>
</table>

Column 1: source name; Col. 2: distance from the cluster center (arcmin); Col. 3: maximum absolute value of RM; Col. 4: average value of RM; Col. 5: RM dispersion; Col. 6: cluster member.

Fig. 10. X-ray PSPC image of the cluster A514. The image has been smoothed with a Gaussian with \(\sigma = 30\)'. Contour levels are: 0.96, 1.36, 1.92, 2.72, 3.84, 5.43, 7.72, and 10.86 counts/pixel (1 pixel = 15'\times15').

Fig. 13 is due to the intracluster medium in A514, whose differential contribution depends on how much magnetized medium is crossed by the polarized emission. Thus, the data show good evidence for the existence of magnetic fields associated with the intracluster medium.

The RM images show fluctuations on scales of about 5' (9 kpc at the cluster redshift), therefore we can consider this value as the coherence length of the magnetic field. This scale can be primarily evaluated from the RM structure across the stronger and more extended source J0448-2032 (A514D) and this value is consistent with coherence length of the magnetic field found in previous works in the literature see e.g. Coma (Feretti et al. 1995), A119 (Feretti et al. 1999) and 3C129 (Taylor et al. 2001).

7. Magnetic field in A514

The determination of the strength of the magnetic field in a cluster depends on several assumptions, including...
the model for the X-ray emitting gas distribution and the magnetic field structure.

The effect of a Faraday screen with different gas density distributions and with a tangled magnetic field has been analyzed by several authors (Lawler & Dennison 1982; Tribble 1991; Felten 1996). Assuming a randomly oriented magnetic field in cells of uniform size and strength, and a gas density distribution given by the hydrostatic isothermal $\beta$ model (Cavaliere & Fusco-Femiano 1981) i.e.:

$$n_e(r) = n_0 (1 + r^2/r_c^2)^{-\beta/2}$$

(4)

where $n_0$ is the central (electron) density, and $r_c$ is the core radius of the gas distribution, the $RM$ dispersion at different projected distances from the cluster center was evaluated by Feretti et al. (1995) and Felten (1996) by solving the integral of Eq. (3):

$$\sigma_{RM} = \frac{441Hn_0r_c^{1/2}}{(1 + r^2/r_c^2)^{0.5}} \sqrt{\Gamma(3\beta - 0.5)}$$

(5)

where $\Gamma$ is the Gamma function, $r_c$ is the core radius in kpc, and $l$ is the size of each cell in kpc, the central gas density $n_0$ is in cm$^{-3}$ and $H$ is in $\mu$G; and it is assumed that the source is as distant from the observer as the cluster center.

For the cluster galaxies J0448–2025 (A514B2) and J0448–2032 (A514D), using the gas parameters calculated from the $\beta$-model fit, the values of $\sigma_{RM}$ fit fairly well in the model described above and are consistent with a magnetic field strength in the range 3–7 $\mu$G for a cell size of about 9 kpc. The uncertainty in the magnetic field strength is dominated by errors in the gas parameters calculated from the $\beta$-model fit. Assuming the above magnetic field, the cell size, and the $\beta$-model fit calculated for A514, we derived the radial profile of $\sigma_{RM}$ expected from Eq. (5). This is plotted as a solid line in Fig. 13, up to a distance corresponding to the extent of the cluster X-ray emission above the background.

The calculated $\sigma_{RM}$ of the background sources (A514A, A514C, A514E) is generally higher than that expected from the absence of cluster contribution at these large projected distances from the cluster center. This may be due to effects internal to the radio sources, to local effects of the host galaxies or to some other effects along the line of sight.

The model we used to derive the magnetic field strength is oversimplified and has the following limitations: (1) the radio galaxies belonging to the cluster can be at different locations along the line of sight; (2) the cluster shows an irregular X-ray emission, therefore the described King density model is only a rough approximation; and (3) the magnetic field structure is likely to be more complicated than assumed.

Even allowing for the uncertainties related to the previous computation, the observational evidence favors the existence of a strong magnetic field in the intergalactic medium of A514, over a wide scale of about 1.4 Mpc in diameter.

At the cluster center the energy density in the magnetic field results about 4–24% of the thermal energy density.
8. Conclusions

We have obtained indirect evidence for the presence of a magnetic field in the cluster A514. Six radio sources located at different projected distances from the cluster center of A514 have been studied in total intensity and polarization. In conjunction with hot gas density estimates based on the X-ray properties of the cluster, observations of Faraday rotation measures in the radio sources can most reasonably be explained by the presence of cluster magnetic fields with a strength of ~3–7 μG spread throughout the central 1.4 Mpc of the cluster. This magnetic field is...
Fig. 16. Left: the image of the rotation measure in A514C (J0448-2005), computed using the polarized angle images at the frequencies 4535, 4885, 8085 and 8465 MHz with a resolution of 1.6″. The contours refer to the total intensity image at 3.6 cm. Right: the histogram of the rotation measure for all significant pixels. The values of \( RM \) range between \(-80\) rad/m\(^2\) and 50 rad/m\(^2\). The \( <RM> \) is \(-19\) rad/m\(^2\) and the \( \sigma_{RM} \) is 24 rad/m\(^2\).

Fig. 17. Left: the image of the rotation measure in A514D (J0448-2032), computed using the polarized angle images at the frequencies 4535, 4885, 8085 and 8465 MHz with a resolution of 1.6″. The contours refer to the total intensity image at 3.6 cm. Right: the histograms of the rotation measure for all significant pixels for North and the South lobe respectively. The \( <RM> \) and \( \sigma_{RM} \) for the North and the South lobes are reported in Table 3, while considering all the significant pixels, the values of \( RM \) range between \(-200\) rad/m\(^2\) and \(+240\) rad/m\(^2\).
Fig. 18. Top: the image of the rotation measure in A514E (J0448-2044), computed using the polarized angle images at the frequencies 4535, 4885, 8085 and 8465 MHz with a resolution of 1.6″. The contours refer to the total intensity image at 3.6 cm. Bottom: the histograms of the rotation measure for all significant pixels for the East lobe, core and the West lobe respectively. The $<RM>$ and $\sigma_{RM}$ for the East lobe, core and the West lobe are reported in Table 3, while considering all the significant pixels, the values of $RM$ range between $-150$ rad/m$^2$ and $+150$ rad/m$^2$.

consistent with the magnetic field calculated using similar methods in other clusters of galaxies without a cooling flow (see e.g. Coma, A119, 3C 129).

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