

Oscillator strengths between fine structure levels of Fe xxiii^{*}

J. A. Tully¹ and M. C. Chidichimo²

¹ Observatoire de la Côte d'Azur, CNRS Laboratoire Cassini, Bd. de l'Observatoire, BP 4229, 06304 Nice Cedex 4, France

² Department of Applied Mathematics, University of Waterloo, Waterloo, Ontario N2L 3G1, Canada

Received 18 July 2000 / Accepted 4 October 2000

Abstract. We tabulate theoretical line strengths, f -values and transition energies for the beryllium-like ion Fe XXIII. Transitions are between levels $2l_1 2l_2 S' L' J'$ and $2l_3 n l_4 S L J$ with $n = 2, 3, 4$. The calculation uses the well known configuration interaction program CIV3 in which relativistic effects are allowed for by means of the Breit-Pauli approximation. We give a detailed comparison of our oscillator strengths with those which Chen & Ong (1998) obtained using the relativistic Dirac code GRASP2.

Key words. Atomic data

1. Introduction

Numerous papers dealing with the radiative properties of Fe XXIII have appeared over the past few years and this is proof of the interest that this highly ionised Be-like ion continues to arouse. As the last millenium draws to a close two veterans in the field of atomic structure calculations have presented energy levels and oscillator strengths for beryllium-like ions including Fe XXIII (Kingston & Hibbert 2000).

Those wishing to compute collision strengths for electron excitation of an ion need wavefunctions for the ground and several excited states. As part of the IRON Project, Chidichimo et al. (1999) – hereafter CZTB – calculated analytic radial orbitals for Fe XXIII with the program CIV3 (Hibbert 1975) and used these to produce oscillator strengths for transitions within the ground complex: $2l_1 2l_2 \rightarrow 2l_3 2l_4$. These were the allowed transitions for which they computed collision strengths. Here we extend their work by calculating line and oscillator strengths for many more transitions and compare our results with those drawn from the work of several other authors. Of special interest is the detailed comparison we make with the results of Chen & Ong (1998) – hereafter CO – who used the unpublished code called GRASP2, written by F.A. Parpia, I.P. Grant and C. Froese-Fischer. GRASP2 is based on multiconfiguration Dirac-Fock Theory and so makes allowance for relativistic effects in a fuller manner

than we do, our calculation being based on the Breit-Pauli approximation (Hibbert et al. 1991).

2. Atomic orbitals and configurations

Our atomic orbitals $P_{nl}(r)$ are Slater type and have the following analytic form

$$P_{nl}(r) = \sum_{j=1}^k c_{jnl} \frac{(2\zeta_{jnl})^{I_{jnl} + \frac{1}{2}}}{[(2I_{jnl})!]^{\frac{1}{2}}} r^{I_{jnl}} \exp(-\zeta_{jnl} r).$$

We take P_{1s} and P_{2s} from Clementi & Roetti (1974) and calculate P_{3s} , P_{4s} , P_{2p} , P_{3p} , P_{4p} , P_{3d} , P_{4d} , P_{4f} using CIV3. Values of the orbital parameters (i.e. coefficients and exponents) are given in Table 1. Some of the values differ from those found by CZTB. The explanation is simply that in the present work we allowed the number of iterations in the optimisation branch of CIV3 to exceed five, the value Hibbert (1975) recommends. Only the parameters of the 3s, 4s, 4p orbitals are affected to any noticeable extent, but the resulting radial functions are numerically similar to the ones in CZTB.

We use configurations $1s^2 n_a l_a n_b l_b$ in which n_a and n_b can take the values 2, 3, 4. This leads to 295 fine structure states of which 146 (149) have odd (even) parity. Details of these single configuration states are given in Tables 2 and 3. CIV3 calculates the mixing amongst states having the same parity and J value. Each final state is a mixture of configurations and usually, but not always, is dominated by one of them which serves as a convenient label. For further details see CZTB.

Send offprint requests to: J. A. Tully

* Tables 1 to 17 are only available in electronic form at <http://www.edpsciences.org>

Tables 4 through 13 contain length and velocity line strengths S_l and S_v in atomic units; length and velocity absorption oscillator strengths f_l and f_v ; and calculated transition energies E_{ij} in rydberg (Ry) units.

3. Comparison with the work of others

Table 14 is a detailed comparison of our f -values and those of CO who made use of an entirely different method from us. Furthermore, CO used 78 fine structure states compared to our 295 since they restricted the quantum numbers n_a and n_b to 2 and 3. For 28 of the 243 transitions which are listed in Table 14, our f -values and those of CO differ by less than 0.5 per cent. For 144 transitions the differences are between 0.5% and 5%; for 38 transitions between 5% and 10%; for 19 transitions between 10% and 25%. And finally for 13 transitions, the f -values differ by more than 25% but less than 72%; however, 12 of these f -values are quite small. The exception is the transition $2 \rightarrow 34$, namely $2s2p^3P_0^o \rightarrow 2p3p^3S_1$, where $f(\text{CO}) = 0.001846$ and $f(\text{TC}) = 0.001169$.

Table 15 compares the energy levels we obtain from a 78 state calculation using CIV3 with the energies from CZTB and CO. Note that levels with indexes 13 and 14 are almost degenerate in energy; CO find they are ordered differently from CZTB. In other words, their level 13 (14) is level 14 (13) in CZTB. It is interesting to note that when we use 295 states the order of these two levels agrees with that given by CO.

In Table 16 we give f -values for transitions within the ground complex and results to compare with are taken from several sources: Cheng et al. (1979), Nussbaumer & Storey (1979), Bhatia & Mason (1981, 1986), Safronova et al. (1992), Zhang & Sampson (1992). It is possible that Safronova et al. (1992) made a slip when tabulating their results for the transitions $2s^2^1S_0 \rightarrow 2s2p^3P_1^o, ^1P_1^o$ since the numbers they give differ by about a factor of 3 from those obtained by everyone else. Table 17 compares our

line strengths for transitions $2l_1 2l_2 \rightarrow 2l_3 n l_4$, $n = 2, 3$ with those given by Sampson et al. (1980, 1984).

4. A final comment

The present investigation originated as an offshoot to CZTB. Claudio Mendoza and co-workers are calculating radiative data for Fe xxiii with the R-matrix codes and their results will appear in due course in the ‘‘Atomic data from the IRON Project’’ series.

Acknowledgements. Alan Hibbert kindly added an appropriate written statement to his program so making it easier to extract large numbers of line strengths and f -values from CIV3.

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