Discovery of PAHs in the halo of NGC 5907

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ABSTRACT

We have used sensitive archival data from the Infrared Space Observatory (ISO) to make maps of the edge-on low SFR galaxy, NGC 5907, in 6 different MIR bands: LW2, LW5, LW6, LW7, LW8, and LW10, covering the spectrum from 6.5 to 15.0 µm and including several narrow bands that isolate the infrared aromatic spectral features commonly referred to as PAHs. Most of the MIR emission is dominated by PAHs and it is likely that emission from VSGs contribute only negligibly except in the broad IRAS-equivalent band. The flux ratios are typical of galaxies with low SFRs or quiescent regions within galaxies (e.g. M 83) and a very high PAH/continuum ratio is observed. The PAH emission follows the CO distribution and also shows some correlation within the disk with the J=850 µm distribution. However, the PAH emission also reaches larger galactocentric radii than the CO and other correlations suggest that the PAHs are also more widespread. A significant new discovery is the presence of PAHs in the halo of the galaxy. In the narrow bands that isolate single PAH features, the emission shows structure similar to high latitude features seen in other galaxies in other tracers. The features extend as far as 6.5 kpc from the plane but scale heights of 3.5 kpc are more typical. The 11.3/4.7.7 ratio also appears to increase with distance from the major axis. To our knowledge, this is the first time PAHs have been seen in the halo of an external galaxy. Just as significantly, they are seen in a low SFR galaxy, suggesting that strong SNe and winds are not necessary for these large molecules to reach high latitudes.

Key words. galaxies: individual: NGC 5907 – galaxies: halos – galaxies: ISM

1. Introduction

The galactic disk-halo interface is a critical region connecting two very different environments over a relatively small vertical distance. The prevalence, nature, and underlying physical drivers of features which span this region and extend into the halo are particularly important since they play a role in the energy balance as well as the chemical evolution of a galaxy. Moreover, knowledge of the physical conditions of outflowing gas in nearby galaxies can provide important constraints on galaxy formation scenarios (e.g. Dekel & Silk 1986; Navarro & White 1993; Mo & Mao 2004; Scannapieco et al. 2002; Kay et al. 2002; Zhao et al. 2003a,b; Sommer-Larsen et al. 2003; Monaco 2004; Marri & White 2003). Nearby edge-on galaxies therefore, provide an important local laboratory within which such flows can be examined in some detail.

Evidence for outflows includes optical emission line cones from nuclear starbursts (e.g. Ohyama et al. 2002; Veilleux et al. 2003, and others), Hz disk-halo filaments (Rossa & Dettmar 2003), X-ray halos (Wang, Chaves, & Irwin 2003; Strickland et al. 2004a,b), and supershells/bubbles in H I (e.g. Spekkens et al. 2004). Dust in halos has been observed in absorption against the background stellar continuum (Sofue 1987; Ichikawa et al. 1994; Alton et al. 1999; Howk & Savage 1999; Thompson et al. 2004), although this only probes halo dust fairly close to the disk. More recently, however, IR, mm, and sub-mm emission from dust has also been detected in halos (Alton et al. 1998; Neininger & Dumke 1999; Brar et al. 2003; Popescu et al. 2004; Tufts & Popescu 2005). The physical conditions in such halos are, however, not well understood, especially those which allow the survival of dust in hot rarefied halo environments. Indeed, the conditions of dust survival in halos may be similar to those of elliptical galaxies (Tsai & Mathews 1995; Xilouris et al. 2004).

The mid-infrared (MIR) wavelength regime has the potential to characterise the dust properties in halos. The sensitive ISOCAM detector (Cesarsky et al. 1996) on board the Infrared Space Observatory (ISO) satellite (Kessler et al. 1996) provides necessary spatial resolution and, with an appropriate selection from the vast assortment of filters (from 4 to 18 µm), it is possible to discriminate between contributing sources of dust. By now, almost a decade after the launch of ISO, we are beginning to understand the properties of the various components in the MIR regime and what they are tracing in galaxies. ISOCAM surveys of galaxies (Rigopoulou et al. 1999; Dale et al 2000; Laurent et al. 2000; Roussel et al. 2001a; Sturm et al. 2002) and detailed studies of individual galaxies (e.g. Galliano et al. 2003; Vogler et al. 2005) have brought deep insight into the

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extragalactic MIR picture. Among the most prominent of the MIR features in galaxies are the carriers of the unidentified infrared bands, one model for these being the PAHs (polycyclic aromatic hydrocarbons)\(^1\) and small hot (typically a few hundred K, see Sect. 3) grains.

In this study, we use the very sensitive previously unpublished observations of NGC 5907 in the ISO database. This has allowed us to take advantage of an extraordinary number of filters that were used to observe both along the disk as well as into the halo region, where we make clear detections far from the plane. Combining these data, we are able to delve into the MIR properties of the halo. This paper is organized as follows. In Sect. 1.1 we review the previous observations of NGC 5907, in Sect. 2 we discuss the observations and data reduction, Sect. 3 reviews the various contributors to the MIR observing bands, Sect. 4 presents the results for both the disk and halo regions, Sect. 5 provides the discussion and the conclusions are given in Sect. 6.

1.1. NGC 5907

NGC 5907 (Table 1), at a distance of 11 Mpc (1 arcsec = 53.3 pc), is one of the largest edge-on systems in the sky. Its nucleus is an HII region type and no nuclear radio source has been reported (Ulvestad & Ho 2002) in spite of some efforts to find one. Non-circular molecular gas motions near the nucleus have been interpreted in terms of a bar (Garcia-Burillo et al. 1997; Garcia-Burillo & Guelin 1997); however, there is no evidence for a bar in the near infra-red (Jarrett et al. 2003). The CO distribution is centrally peaked (Dumke et al. 1997) with a distribution suggestive of rings or spiral arms (Sofue 1994; Dumke et al. 1997).

NGC 5907 has been an important target in searches for faint stellar halos since it has no appreciable bulge and was originally thought to be isolated. Sackett et al. (1994) first reported the detection of an R-band halo to a distance of 6 kpc around this galaxy, following which and I band extraplanar light was detected by Lequeux et al. (1996), as well as J and K band halos by James & Casali (1998) and Rudy et al. (1997).

NGC 5907 is, in fact, a member of the 396th Lyon Group of Galaxies (LGG 396) of which there are 4 identified members: NGC 5866, NGC 5879, NGC 5907, and UGC 9776 (Garcia 1993). These members are at large separations (Table 2) and are unlikely to be interacting with NGC 5907 now. However, the more recent detection of a nearby companion, PGC 54419, at a projected distance of only 36.9 kpc and velocity separation of \(\Delta V = 45\) km s\(^{-1}\) (Shang et al. 1998) indicate that NGC 5907 is not as isolated as previously thought. Moreover, their additional discovery of a large optical ring around the galaxy indicates that some interaction has indeed occurred. A pronounced HI warp (Sancisi 1976) and some evidence for an optical warp (Morrison et al. 1999) are consistent with this. The ring and nearby companion are faint in comparison to NGC 5907; the luminosity of the ring is \(\lesssim 1.2\%\) of NGC 5907 and the mass of the second companion is 0.5\% of the mass of NGC 5907 (Shang et al. 1998). Dynamical modelling of the ring indicates that it was formed at least 0.8 Gyr ago via the destruction of a satellite of mass \(2 \times 10^8 M_\odot\) (Johnston et al. 2001, who adopted a slightly larger distance to the galaxy). Follow-up observations showed that the ring is highly asymmetric and suggests that NGC 5907 does not after all have a faint extended stellar halo as originally envisioned (Zheng et al. 1999). A near-infrared (3.5–5 \(\mu m\)) search also failed to show evidence for a halo associated with NGC 5907 (Yost et al. 1999) and HST observations find fewer bright giants in the halo region than would be expected from a halo with standard dwarf-to-giant ratios (Zepf et al. 2000). Recent 2MASS K\(_s\) band observations of NGC 5907 (Bizyaev & Mitronova 2002) provide a stellar scale height of \(z_0 = 0.49\) kpc with an isothermal assumption [i.e. \(\rho(z) \propto \exp(-z/z_0)\)].

As for a gaseous halo, no high latitude ionized gas has yet been detected (Rand 1996) nor did Howk & Savage (1999) detect any extraplanar dust in optical absorption. Dumke et al. (1997) find a CO FWHM of 400 pc and an HI FWHM of 1.5 kpc, although the warp makes this latter value uncertain. The most extended component appears to be the radio continuum (Dumke et al. 2000), for which an exponential thick disk component is found with a scale height of 1.5 kpc. An 850 \(\mu m\) map has recently been published by Alton et al. (2004) who find an extrapolan exponential scale height of 0.11 kpc.

ISO 60, 100, and 180 \(\mu m\) fluxes as well as a 12 \(\mu m\) image have been published by Bendo et al. (2002), and Bendo et al. (2003) also present temperature fits to the data. More recently, Alton et al. (2004) have also fit a two-component dust model to NGC 5907 which includes newer 850 \(\mu m\) data from observations using the Submillimetre Common Use Bolometric Array (SCUBA) on the James Clerk Maxwell Telescope. These temperatures are presented in Table 1. An earlier 2-component dust temperature model of 18 K and 54 K had been reported by Dumke et al. (1997) who observed the 1.2 mm dust component. Note that two-temperature models are simply a rough approximation to the range of temperatures that likely exist in this galaxy. Other parameters for NGC 5907 are also listed in Table 1 including the star formation rate (SFR) which is quite low (cf. SFR = 3.8 M\(_\odot\) yr\(^{-1}\) for the quiescent galaxy, NGC 891, Popescu et al. 2000). In Table 1 we also provide comparative data on the nearby galaxy, M 83, whose spectral energy distribution (SED) has recently been studied in some detail (Vogler et al. 2005) and with which we compare our NGC 5907 results.

2. Observations and data reduction

2.1. Observations

Data on NGC 5907 were available in filters, denoted LW\(_n\), where LW indicates the long wavelength array of the ISOCAM detector and \(n\) specifies the wavelength band. Details of the data acquisition and resulting maps are provided in Table 3. As we prefer to designate the bands according to their wavelength, we use the following nomenclature: LW5, LW6, LW7, LW8, LW2 and LW10 are referred to as 6.8N, 7.7N, 9.6N, 11.3N, 6.7W, and 12W, respectively, where the number represents

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\(^1\) The carriers of the MIR bands are not certain, but in this paper, we adopt the PAH nomenclature for discussion and comparison with other authors.
Table 1. Basic galaxy parameters.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>NGC 5907</th>
<th>M 83</th>
</tr>
</thead>
<tbody>
<tr>
<td>RA (J2000) (h m s)</td>
<td>15 15 53.69</td>
<td></td>
</tr>
<tr>
<td>Dec (J2000) (° ′ ″)</td>
<td>56 19 43.9</td>
<td></td>
</tr>
<tr>
<td>$V_{\text{ls}}$ (km s$^{-1}$)</td>
<td>667</td>
<td></td>
</tr>
<tr>
<td>Distance (Mpc)</td>
<td>11°</td>
<td>4.5°</td>
</tr>
<tr>
<td>Morph. type</td>
<td>SA(s)c: sp HII</td>
<td>SAB(s)c: HII</td>
</tr>
<tr>
<td>Incl. (°)</td>
<td>86.5°</td>
<td>27°</td>
</tr>
<tr>
<td>Major × Minor Axis Diameters (′ × ′)</td>
<td>12.77 × 1.40</td>
<td>12.9 × 1.15</td>
</tr>
<tr>
<td>Major Axis Diameter (kpc)</td>
<td>41</td>
<td>17</td>
</tr>
<tr>
<td>$f_{\text{IR}}$/total</td>
<td>0.24</td>
<td>0.54</td>
</tr>
<tr>
<td>$L_{\text{IR}}$ / (L$_\odot$)</td>
<td>7.1 × 10$^9$</td>
<td>1.4 × 10$^{10}$</td>
</tr>
<tr>
<td>$T_{\text{dust}}$ 1-Comp (1$^{-1}$)</td>
<td>26 K</td>
<td>36</td>
</tr>
<tr>
<td>$T_{\text{dust}}$ 1-Comp (1$^{-2}$)</td>
<td>22 K</td>
<td>29</td>
</tr>
<tr>
<td>$T_{\text{dust}}$ 2-Comp (1$^{-2}$)</td>
<td>13, 26 K</td>
<td></td>
</tr>
<tr>
<td>$T_{\text{dust}}$ 2-Comp (1$^{-15}$)</td>
<td>13, 28 K</td>
<td></td>
</tr>
<tr>
<td>$M_{\text{dust}}$ (M$_\odot$)</td>
<td>6.0 × 10$^7$</td>
<td>3.2 × 10$^{10}$</td>
</tr>
<tr>
<td>$\Sigma_{\text{dust}}$ (M$_\odot$ kpc$^{-2}$)</td>
<td>7.9 × 10$^5$</td>
<td>2.6 × 10$^5$</td>
</tr>
<tr>
<td>SFR (M$_\odot$ yr$^{-1}$)</td>
<td>2.2$^a$, 1.2$^b$</td>
<td>2.4$^a$</td>
</tr>
<tr>
<td>SFR/A (M$_\odot$ yr$^{-1}$ kpc$^{-2}$)</td>
<td>1.7 × 10$^{-3}$</td>
<td>9.0 × 10$^{-4}$</td>
</tr>
</tbody>
</table>

*a* From the NASA Extragalactic Database (NED). This center is the 2MASS center which agrees with the dynamical center (Dumke et al. 1997) to within $\Delta$ RA = 0.8" and $\Delta$ Dec = 0.3".

*b* Sasaki (1987).

*c* Thim et al. (2003).

*d* Garcia-Burillo et al. (1997).

*e* From the major and minor axis diameters and a thin disk assumption.

*f* Flux ratio or FIR luminosity, using the 60 and 100 $\mu$m IRAS fluxes from Sanders et al. 2003. $L_{\text{IR}}$ allows for extrapolations below 40 and above 120 $\mu$m and adjusts to the distance given in this table.

*g* Dust temperature, assuming the emissivity laws shown, for one or two component fits (Bendo et al. 2003). For M 83, a 2-component model does not improve the fit. The second line for the 2-component model gives the fit from Alton et al. (2004).

*h* From Misiriotis et al. (2001).

*i* From Bendo et al. (2003) for fluxes within a 135° aperture, assuming a $\lambda^{-2}$ emissivity and adjusting to the distances in this table.

*j* Dust surface density, using the dust masses in this table (from the adjusted Bendo et al. data) and their aperture sizes.

*k* From $L_{\text{IR}}$ of this table and the calibration of Kennicutt (1998).

*l* SFR per unit area, from values in the immediately preceding row and the optical major axes given in this table.

Table 2. Galaxy group membership.

<table>
<thead>
<tr>
<th>Galaxy Name</th>
<th>Velocity (km s$^{-1}$)</th>
<th>RA (h m s)</th>
<th>Dec (° ′ ″)</th>
<th>Separation$^b$ (Arcmin)</th>
<th>Separation$^b$ (Radii)</th>
</tr>
</thead>
<tbody>
<tr>
<td>NGC 5907</td>
<td>667</td>
<td>15 15 53.69</td>
<td>56 19 43.9</td>
<td>–</td>
<td></td>
</tr>
<tr>
<td>NGC 5866</td>
<td>672</td>
<td>15 06 29.56</td>
<td>55 45 47.9</td>
<td>85.8</td>
<td>13.4</td>
</tr>
<tr>
<td>NGC 5879</td>
<td>772</td>
<td>15 09 46.77</td>
<td>57 00 00.7</td>
<td>64.5</td>
<td>10.1</td>
</tr>
<tr>
<td>UGC 9776</td>
<td>833</td>
<td>15 13 07.44</td>
<td>56 58 07.3</td>
<td>44.7</td>
<td>7.0</td>
</tr>
<tr>
<td>PGC 54419</td>
<td>712</td>
<td>15 14 48.04</td>
<td>56 27 15.4</td>
<td>11.8</td>
<td>1.8</td>
</tr>
</tbody>
</table>

*a* Members of the Galaxy Group LGG 396 (Garcia 1993) as well as PGC companion found by Shang et al. (1998). All values are taken from NED.

*b* “Separation” indicates the distance of the group member from the center of NGC 5907 in units of arcmin and in units of the radius of NGC 5907 (see Table 1).
with a resulting 1′′ pixel spacing. In this case, 12 on-source pointings were made for full coverage of the central part of the galaxy with sufficient sky coverage for subtraction.

The final data set, 6.7W, was obtained on 21 Apr. 1996 in beam switching mode (CAM03). There were 4 positions on-source: one containing the galaxy center but offset to the south-west, another containing the galaxy center offset to the north-east, one at the end of the major axis to the north and one at the end of the major axis at the south. Thus, the final combined image does not provide full coverage of the whole galaxy but rather samples the galaxy at four locations. For each of the 4 pointings, the off-source sky positions were placed at a variety of position angles with respect to that pointing.

### 2.2. Reductions

All data were reduced using the CAM Interactive Reduction Package (CIR, Chanial 2003). First, the dark current was subtracted following the method of Biviano et al. (1998) which includes a dark correction, a second-order dark correction depending on detector temperature and time of observation, and a short-drift correction. Next, high glitches due to cosmic ray (CR) hits were removed automatically via a multi-resolution filtering technique (Starck et al. 1999) at a 6σ level. The effects of memory on a pixel (transient effects) were then corrected (de-glitched) using the Fouks-Schubert method (Coulais & Abergel 2000). Each data set was then examined carefully, frame by frame, and further bad pixels were removed manually.
as required. Most of these occurred in regions immediately adjacent to pixels which had been automatically de-glitched.

For the raster mode observations (CAM01: 6.8N, 7.7N, 9.6N, 11.3N, and 12W), the edges were then trimmed (blanked), the frames were corrected for jitter, and averaged. The results at each pointing were then corrected for the flat field according to the on-line library of calibration flat fields (see Roussel et al. 2001a). The various pointings on-source were then projected onto a common grid in sky coordinates and corrected for field lens distortion. At this point, the data were calibrated to mJy/pixel units (Blommaert et al. 2003). Finally, the sky background/foreground was subtracted by first isolating a region of sky-only in the field and subtracting the median of that region from the field.

For the beam switching observations (CAM03: 6.7W) we carried out an additional stage of manual deglitching to account for the fact that in the off-source positions, the CCD contained a fading memory of the galaxy itself from when the array pointed at the galaxy. Thus, the region of the array which had previously contained the galaxy was blanked in approximately 10 sky frames immediately after moving from the galaxy to the sky. This blanking covered about 20% of the array and left approximately 20 frames where the image had faded and blanking was no longer required. This additional deglitching was required only for the center-most two fields, and not for the two fields at the ends of the galaxy’s major axis, given the weakness of the signal there. Beam switching allowed for sufficient sky coverage off source that the sky positions provided both the flat field and its distortions as well as the sky brightness values. Thus, after the average on-source image was created, it was corrected using an image that was formed from an average of the various off-source positions. The result was then calibrated to mJy/pixel units. This mode of operation provides the best sky subtraction because it does this pixel by pixel rather than using a single value (the median of the sky values) for all pixels, as above. Finally the registration and combination of the 4 pointings, as well as an interpolation over one bad column on the array was carried out outside of the CIR package. In the region of the galaxy in which there was overlap (see center region of Fig. 6a), the flux levels between the two overlapping fields differed by \(-4.5\%\).

All images were interpolated/rotated into RA/Dec coordinates with 1 arcsec square pixels for consistency. The resulting maps are shown in Fig. 1a through 6a. The lowest contour, in each case, is set to \(2\sigma_{\text{map}}\), where \(\sigma_{\text{map}}\) is the rms noise level as measured from the emission map itself.

The field of view is accurately shown by the error maps (next section). The results are discussed in Sect. 4.

2.3. Error analysis

Attention has been paid to the error budget in the final images. The resultant error is the quadratic sum of individual errors, consisting of a) the readout and photon noises added quadratically and calculated by the software, denoted \(\sigma_{\text{rms}}\) in Table 3, b) an error on the background level which was subtracted; this is taken as the \(1\sigma\) dispersion in the histogram of values used to determine the sky value to be subtracted and includes uncertainties due to residual dark current, flat field, glitches, remnant images, and long-term drifts, denoted “1\(\sigma\) Sky Error” in Table 3, and c) the effects of the non-stabilization of the signal (memory) as a function of time over a temporal frame-by-frame block at any pointing; this error has been estimated from variations in the signal that are greater than \(3\sigma\) within a block (assuming that such larger variations are caused by only this error) and applying a factor which accounts for the fact that the error is worse near the beginning of a block (see Galliano 2004; Roussel et al. 2001a), denoted “Transient rms” in Table 3. The quadratic sum of these errors is referred to as the random error (“\(\sigma\)” in the table) and has been computed “per pixel”.
Fig. 2. As in Fig. 1, but for 7.7N and with contours at 0.006 \( (2\sigma_{\text{map}}) \), 0.01, 0.018, 0.05, 0.10, 0.30, 0.50, and 0.65 mJy arcsec\(^{-2}\). The inset shows the emission after an estimate for the stellar contribution has been subtracted (see Sect. 3.1).

Fig. 3. As in Fig. 1, but for 9.6N with contours at 0.015 \( (2\sigma_{\text{map}}) \), 0.025, 0.035, 0.05, 0.10, 0.18, and 0.25 mJy arcsec\(^{-2}\).

where the pixel is the originally sized pixel given in the table, and also has been computed “per square arcsec”, after interpolating onto a 1 arcsec pixel grid. We formed a total Random Error map (showing \( \sigma_i \) for each map point, \( i \), from the quadratic sum of the \( \sigma_{\text{rms}} \) map, the Transient rms map, plus a constant 1\( \sigma \) Sky Error value. These maps are shown in Figs. 1b through 5b. However, for the beam-switching observations, since we also have the sky background errors for each pixel (rather than the same sky value for every pixel), all three sources of error were calculated in one step (thus the missing intermediate step values in Table 3) and the final Random Error map is shown in Fig. 6b.

The calibration error (see Table 3) has not been included in Figs. 1b through 5b since this would globally raise or lower the flux level of the map and these maps are mainly meant to show variations across the fields. However, it is included in any calculations of flux density. A remaining error, that of variations in the flux density calibration according to orbit \( (\approx 5\%, \text{ Blommaert et al. 2003}) \) has not been included. Since 6.8N, 7.7N, 9.6N, and 11.3N observations were all taken at the same orbital point, this error does not enter into a field by field comparison of these 4 bands. However, it will be a factor in a comparison between these bands and the remaining two.
We note that the true errors cannot be determined rigorously, for example, the memory effects of the camera are not accurately known (Roussel et al. 2001a), although they can be tested for, as described below. Thus the quoted errors, though calculated rigorously, should be considered indicative. An example is seen in Fig. 3a in which there remain three residual artifacts to the west of the galaxy which correspond to the edges of the three fields. We have also, therefore, shown the lowest contour level in Figs. 1a through 6a as twice the rms level on the emission map itself, the latter being an indication of the pixel-to-pixel variation across a field in the region of the sky.

Since we are reporting the discovery of halo PAHs in NGC 5907 in bands, 7.7N, 11.3N, and 6.7W (see Sect. 4.2), we performed some additional tests which were designed to check whether this high latitude emission might have been produced artificially. These tests are as follows, focusing on the 7.7N band.

We first examined the point spread function (PSF) whose FWHM values are 7.2″ in these bands (Table 3). The PSFs at different locations on the image in these bands have been carefully examined and characterized by fitting 2D Gaussians (see Galliano 2004) and are, to a high accuracy, symmetric.
Thus, the wings of the PSF should be completely negligible 7.2″ beyond the projection of the disk. To check this, we took the 7.7N image (Fig. 2), measured its semi-major axis (245″), and created an inclined (86.5″) thin disk model (see Irwin et al. 1999 for similar modelling). The resulting projected half-width perpendicular to the plane at the center of the galaxy is 15″ and less elsewhere. This model was then used as a template to blank all emission outside of the inclined thin disk region from the Fig. 2 image. The resulting blanked image then contained only emission from the disk region. The blanked image was then convolved with a 7.2″ Gaussian PSF to create an image consisting of disk plus extended emission due to the PSF only. This resulting map was then compared with the original image by taking the ratio of the two and determine the percentage change per pixel that results. We find that, outside of the projection of the thin disk, all ratios are <1% and the mean is <0.05%, verifying that the PSF contributes negligibly to the halo emission.

Secondly, we compared results between using the library flat field (as described in Sect. 2.2) and using a sky flat, formed from a portion of the image in which the sky is visible. The mean relative error between the two results over comparable regions is 4.5% and high latitude structure, most of which is in the region in which the sky flat could be formed, remains present. Moreover, in both images, artifacts, such as those at the edge of the field in Fig. 3a, cannot be seen.

Finally, we checked further into issues related to memory. Although the resulting map already has memory corrections in place including blanking the first frames of any scan, we applied a more aggressive scheme. This involved substantial blanking of the first time period when a new scan began, determining the median of the last 5 frames of the scan, pixel by pixel, and then fitting a slope to every pixel independently to align the flux of the other frames with this median. Some differences in structure were noted at a low level, however the high latitude emission remained.

3. The MIR spectrum and contributors to the observed bands

To interpret the emission in the observing bands, it is first necessary to understand what the various contributors are to the MIR emission. For this purpose, we provide ISOCAM CVF (circularly variable filter) spectra of regions in two other galaxies for comparison (see Fig. 7). The main spectrum is of the nucleus of the nearby galaxy, M 83, and the inset is of a quiescent region away from the starburst in M 82. Various features are labelled as well as the bands used in our observations.

There are two main contributors to the MIR spectrum: 1) strong spectral peaks due to very small grains or large molecules (PAHs) which, due to their small sizes (typically ∼0.001 μm) and thus, low heat capacity, are stochastically heated, undergoing wide temperature fluctuations covering 100 s of K from absorption of a single photon (Leger & Puget 1984; Sellgren et al. 1990; Allamandola et al. 1989). The primary emission bands, relevant to the wavelength range observed for NGC 5907, are located at 6.2, 7.7, 8.6, 11.3 and 12.7 μm and can be seen as strong emission peaks in the spectrum. 2) a continuum composed of very small (of order ∼0.01 μm) grains (VSGs) as modeled in our galaxy by Désert et al. (1990). Depending on their size and the local radiation field, these grains could be stochastically heated or in thermal equilibrium. A modeled continuum is shown in the figure.
inset, increasing in strength to longer wavelength. Finally, not shown in the figure, is a possible weak stellar contribution that, if present, would contribute at the low wavelength end of the spectrum, decreasing at higher wavelengths. Since we are able to make numerical estimates of this contribution from previous observations of NGC 5907 itself, we deal with this separately in Sect. 3.1.

Without doubt, the most overwhelming emission in the MIR band is due to PAH features. These features are ubiquitous in normal spirals and account for almost all of the MIR energy (e.g. Vogler 2005; Genzel & Cesarsky 2000, and references therein). It is also now established that most normal spiral galaxies show qualitatively little difference in their spectra in the MIR and that the spectral shape is largely independent of star formation (Lu et al 2003). From higher resolution Galactic observations, we know that the VSGs (producing the dust continuum) peak in the nebular regions, while the PAHs peak outside the HII regions, in the photodissociation regions (PDRs) around molecular cores (e.g. Cesarsky et al 1996; Verstraete et al 1996). While the PAHs emit profusely in the disks of galaxies, excited predominantly by UV photons, they are also observed in the more diffuse ISM (e.g. Chan et al. 2001; Mattila et al. 1996) where optical photons are thought to be the source of excitation. Optical photons are also an explanation for the presence of PAHs in elliptical galaxies (Athey et al 2002; Xilouris et al 2004), where evolved stars are the primary stellar population. Moreover, in a study of 5 galaxies over a range of star forming activity, Haas et al. (2002) find a spatial correlation between PAHs and the cold dust distribution which is more widespread than around SFRs alone.

We thus expect the relative importance of the continuum to diminish in galaxies of lower star formation rate or in quiescent regions of galaxies. This is indeed observed in M 83 (the increasing continuum longward of $\lambda 14 \mu m$ at the nucleus shown in Fig. 7 becomes much less prominent in the interarm regions) as well as other known quiescent star forming regions (Fig. 7 inset). Thorough modeling of ISOCAM CVF spectra using Lorentzian profiles further supports the dominance of the PAH bands, greatly reducing the need for a strong continuum. (Boulanger et al. 1998). Indeed the three PAH features at $\lambda 6.2, 7.7, \text{ and } 8.6 \mu m$ in M 82 can be fit without any significant underlying continuum at all; a continuum contribution is fit with a simple linear curve, increasing with wavelength (Laurent et al. 2000). Vogler et al. (2005) also find that a continuum contributes only $\sim 5\%$ for these three PAH bands. As this is less than a typical error bar on the observations, we consider the continuum contribution below these PAH features to be negligible.

Similar arguments apply to the ionic emission lines. In nuclear regions, a [NeII] emission line at $12.8 \mu m$ has been observed and is blended with the $12.7 \mu m$ PAH feature. However, in a sample of 69 nearby normal spiral galaxies with ISOCAM spectra, Roussel et al. (2001c) found that the contribution from the [NeII] line is negligible. The $\lambda 7.0 \mu m$ [ArII] line, visible at the nucleus of M 83 (Fig. 7) is no longer seen in the interarm regions of this galaxy (Vogler et al. 2005).

The final possible MIR feature to be remarked upon is a broad silicate absorption feature at $9 \mu m$ which has been detected in narrow central regions of very dusty starbursts and AGNs (e.g. Laurent et al. 2000; Tran et al. 2001; Sturm et al. 2000). This feature is not likely to be important toward the central regions of normal starbursts or in quiescent disks (e.g. Roussel et al 2001a, 2001b; Vogler et al 2005, Dale et al. 2001, also arguments in Spoon et al. 2004), even in the case of edge-on spiral galaxies (Mattila et al. 1999, for NGC891).

Therefore, the MIR spectrum of a quiescent galaxy like NGC 5907 can be expected to be dominated by the PAH emission bands, with a weak underlying continuum which becomes stronger at longer wavelengths. We thus interpret the emission in the various observed bands as follows:

- **6.8N**: Lorentzian wings of the $\lambda 6.2 \mu m$ and $\lambda 7.7 \mu m$ PAH bands.
- **6.7W**: Lorentzian wings of the $\lambda 6.2 \mu m$ and $\lambda 7.7 \mu m$ PAH bands.
- **6.8N**: Lorentzian wings of the $\lambda 7.7 \mu m$ PAH bands.
- **12W**: Lorentzian wings of the $\lambda 11.3 \mu m$ PAH bands.
- **9.6N**: Lorentzian wings of the $\lambda 7.7 \mu m$ PAH bands.
- **12.7W**: Lorentzian wings of the $\lambda 11.9 \mu m$ PAH bands.
- **7.7W**: Lorentzian wings of the $\lambda 7.7 \mu m$ PAH bands.
7.7N: λ7.7 μm PAH band.
9.6N: Lorentzian wings of the λ7.7 μm, 8.6 μm, and λ11.3 μm PAH bands plus possible minor contribution from a continuum.
11.3N: λ11.3 μm PAH band plus possible continuum.
6.7W: Combined λ6.2, 7.7, and part of the λ8.6 μm PAH band emission.
12W: Broad band emission, equivalent to the IRAS λ 12 μm band, dominated by PAH features with likely continuum.

3.1. The stellar contribution

The stellar contribution to the MIR band is illustrated beautifully in Lu et al. (2003) for both a reddened and de-reddened case (their Fig. 6). These authors have obtained complete ISO spectra in the 2.4 to 5.9 μm and 5.8 to 11.6 μm bands for 45 galaxies finding, as have other authors, that the characteristic shape of the spectrum in these bands is extremely consistent, galaxy to galaxy. The stellar contribution can be approximated by a modified black body with a λ−2 emissivity law. This behaviour for the stellar contribution also accurately describes the elliptical galaxies in their sample. Thus, we expect the stellar contribution in both the disk and halo regions of NGC 5907 to have a similar λ dependence.

We have obtained the Ke band (λ 2.159 μm) image for NGC 5907, from the 2MASS Large Galaxy Atlas (LGA, Jarrett et al. 2003), which is the band most likely to show only stars with negligible dust obscuration. This image was interpolated onto the same grid as our MIR images and then smoothed to 7.2″ resolution. The “data numbers” (dn), in the FITS file were converted to Jy using the transformation, 
f = f(0) \times 10^{-6} \text{dn} \text{Jy}^2.

As a check on our processing steps, we measured the total flux density of the source in the regridded, smoothed image, finding f = 1.322 Jy which corresponds to mK = 6.76. This agrees with the magnitude given in the LGA for NGC 5907 of mK = 6.76 ± 0.02. The resulting image has an rms noise of \sigma = 1.6 \text{mJy arcsec}^{-2}. From the emissivity law (following Lu et al.), we expect the stellar contribution, extrapolated into the MIR bands in which PAHs dominate to be 7.9% in the 7.7N band, 3.7% at 11.3N, and 10% at 6.7W of whatever value is in the Ke band at the relevant location. With this extrapolation, we created 3 maps from the Ke band image for the 3 bands dominated by PAH emission and subtracted each from the total emission map. These results are shown in the insets to Figs. 2, 4, and 6 with the same contour levels as the main maps.

We then measured the global flux of the stellar-subtracted maps and compared this to the global flux of the main maps finding a difference of 5%, 2%, and 2% for the 7.7 μm, 11 μm, and 6.7 μm maps, respectively (ignoring emission from the two outer frames in the latter case). Since the flux corrections are well within other error bars (Sect. 2) and the appearance of the maps is also close to the original non-subtracted maps, we consider the stellar contribution to be negligible and continue our discussion and analysis with respect to the unsubtracted maps.

4. Results

All maps along with the best estimated error maps are shown in Figs. 1 to 6 as described in Sect. 2.2. The error maps delineate the final field of view. Emission has been detected in every observing band as shown in these figures.

4.1. Disk emission

Emission is clearly detected along the disk of NGC 5907. In no individual band do we have full coverage of the entire galaxy. However, for the narrow bands, 6.8N, 7.7N, 9.6N, and 11.3N, we have full coverage of the south-east major axis and for 6.7W, small segments at both ends of the galaxy have been observed. In all cases, the emission extends along the optical disk with detections to the “end” of the optical disk in some cases.

The extent of the PAH emission, specifically, is illustrated by Fig. 8a in which we show the smoothed 7.7N emission (λ 7.7 μm PAHs) over the DSS optical image. It is clear from this overlay that the PAH emission extends all along the optical disk of the galaxy. It is further illustrated in Fig. 9 in which we show the smoothed, high sensitivity 6.7W emission (λ6.2, 7.7, and 8.6 μm PAHs) superimposed on the HI total intensity map from Shang et al. (1998). PAH emission is seen very far out along the galaxy disk where the HI emission is also strong. The total extent of the PAH emission along the major axis is 11.0 arcmin (~35 kpc), as measured to the outermost detectable contiguous features on the smoothed map (cf. 41 kpc for the optical disk, Table 1). However, there is also some evidence for real features farther out. For example, several emission peaks can be seen at the edges of the 6.7W field of view on the south-east end of the major axis at RA = 15h16m12. Dec = 56°13′09″. A channel-by-channel comparison (in velocity) of the HI distribution with the PAH emission, however, fails to show any clear correlation.

The lack of global correlation with HI is further illustrated in Fig. 10a in which we show major axis slices of the HI total intensity distribution, the CO, and the smoothed 7.7N emission. It is clear from this plot that the PAH emission follows the molecular gas distribution and not the atomic gas distribution. This was also found by Dumke et al. (1997) for the λ1.2 mm dust distribution as well as by Alton et al. (2004) for the λ850 μm emission. Figure 10a shows two sets of peaks in the λ7.7 μm distribution on either side of the nucleus, one at ~±70″ and the second at ~±100″. The CO comparison also show two peaks, interpreted by Dumke et al. (1997) to represent rings or spiral arms, but at somewhat offset radii of ~±60″ and ~±120″.

In Fig. 10b, we show a comparison between major axis cuts for our ISO λ7.7 μm data and the λ850 μm SCUBA data, both averaged over a minor axis extent of 16″ and at the same spatial resolution. The shape of the PAH emission curve and the λ850 μm emission curve, the latter which traces cool dust, are remarkably similar showing the same number and approximate, but not exact, positions of peaks on either side of the

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2 This transformation uses the calibration, mK = ZP − 2.5 log (dn) and f = f(0) 10^{-6}dn, where ZP = 20.0704994 is the zero point provided in the FITS header, and f(0) = 666.7 Jy, the latter from http://ssc.spitzer.caltech.edu/documents/cookbook/html/cookbook.html
Fig. 8. a) The 7.7N band image (λ7.7 µm PAH band), smoothed to 12” resolution superimposed on the optical DSS image. Contours are at 4.6 (2σmap), 7.5, 13, 25, 50, 75, 200, and 380 × 10⁻³ mJy arcsec⁻². b) The 7.7N band image as in a) superimposed on a greyscale λ850 SCUBA µm map from Alton et al. (2004) with a resolution of 16”. Note that the SCUBA field of view ends at a Declination of 56°17’.

Fig. 9. The 6.7W band image, smoothed to 12” resolution superimposed on the HI total intensity image (data from Shang et al. 1998). Contours are at 0.0004 (2σmap), 0.0008, 0.0015, 0.003, 0.01, 0.04, and 0.10 mJy arcsec⁻².

nucleus. Positional offsets between the two different slices for 5 peaks range from zero to 19” with a typical value of ~13”. Given that the resolution is 16” and relative positional uncertainties between maps is of order 7” (see caption), the positions of the peaks are likely consistent, though we cannot rule out the possibility of some displacement which might be detected with higher resolution observations. The line ratios from these curves (calculated though not plotted) are also consistent with the values of Haas et al. (2002) who found λ7.7 µm/λ850 µm = 2 (1.8 to 2.2 for 5 galaxies). Their denominator is from SCUBA observations but the numerator represents the peak λ7.7 µm flux rather than the average in the 7.7N band as we have plotted. All of the emission in our 7.7N band is from the λ7.7 µm PAH line but the peak, if we could measure it directly, would be higher than the average. From Fig. 3, the increase should be of order a factor of ≈1.5 which would put our ratio in even closer agreement with Haas et al. Therefore, overall, our results are consistent with a possible relationship between PAHs and the 850 µm cool dust distribution.

Finally, there is excess λ7.7 µm emission between radii of ~150” and ~250” (Fig. 10a) in comparison to the CO distribution. (Note that we cannot discern this in Fig. 10b because of the truncated field of view of the SCUBA map.) We believe that this excess is significant. Indeed, Dumke et al. (1997) also found an excess in the λ1.2 mm emission in comparison to CO over the same radii but on both the north and south sides of the nucleus. This shows that there is PAH emission as well as cool dust where there is little molecular gas. We will return to this point in Sect. 5.

In Fig. 11 we show the broad band 12W emission (equivalent to the IRAS 12 µm band) superimposed on the DSS optical image. This MIR band should contain the highest contribution of continuum emission from VSGs compared to the other bands (Fig. 7, inset). The 12W band has the highest spatial resolution of all observations (1”) and reveals that the strongest ridge of MIR emission lies along the optical dust lane (note that the field of view does not cover the entire disk), indicating a good correlation between VSGs and the optically obscuring dust.

4.2. High latitude emission

A particularly striking result is the evidence for high latitude emission, especially in the 7.7N band (Fig. 2), the 11.3N band (Fig. 4), and the high-sensitivity beam-switched 6.7W band (Fig. 6) which all trace PAH emission. The 11.3N band in principle, can contain PAHs plus continuum, but the continuum
Fig. 10. Total intensity cuts along the major axis of NGC 5907. Positional uncertainties are of order ±7′.

(a) The HI total intensity (faint extended curve) from data of Shang et al. (1998), the CO distribution (faint solid plus dashed curves showing data points and error bars) from Dumke et al. (1997), and the λ7.7 µm from this work (darkest curve), the latter smoothed to the same resolution as the HI data (18.6′′ × 17.8′′ at PA = −38°). The CO resolution is 21′′ and the center point has been placed approximately at the CO center of RA(J2000) = 15h 15m 53.8s, Dec(J2000) = 56° 19′ 37″.

(b) The λ7.7 µm PAH emission from the 7.7N band, smoothed to 16′′ resolution shown with the λ850 µm SCUBA emission at the same resolution. Curves and intensities represent an average over a 16″ swath parallel to the major axis. The arrow marks the edge of the SCUBA field of view.

should be negligible so far from star forming regions. Thus it is clear that PAHs exist far from the plane of NGC 5907, in the case of 6.7W, as far as 6.5 kpc (2′) from the plane, though ~3 kpc is more typical. Even the 6.8N band (Fig. 1), which should consist only of emission from PAH band wings, shows some evidence for features away from the plane. Given the galaxy’s edge-on orientation (86.5°, Table 1), the correction for galaxy inclination is negligible. For example, for a PAH disk radius of 250″ (Fig. 10), the semi-minor axis is only 15.3″ in a thin disk model. Moreover, the CO semi-minor axis is only ~4 arcsec if we use the HWHM of 200 pc modeled by Dumke et al. (1997). These features resemble the extra-planar emission seen in other galaxies, (see, e.g. the radio continuum features in NGC 5775, Lee et al. 2001) and some are similarly arc-like in appearance (e.g. the east side of the λ7.7 µm emission in Fig. 2). Recall also (Sect. 1.1) that the stellar scale height has been measured to be 0.49 kpc (9′′). Note that the narrow bands, 7.7N and 11.3N, show more structure than the wide band, 6.7W which may result from the ability of these bands to isolate a single PAH line. Although structural details may vary, both narrow bands show high latitude features at approximately the same locations (see also Sect. 2.3).

Given that the beam-switched 6.7W band observations are our most sensitive and also the fact that this band is dominated by several PAH emission features, we have taken the original, unsmoothed 6.7W data shown in the central regions of Fig. 6 and averaged the emission in strips parallel to the major axis. The resulting minor axis profile is shown in Fig. 12. The wings on this averaged profile approach the noise between 100″ and 150″ (5.3 kpc to 8.0 kpc) and the line half-width at the 3σ level is ~95″ (5 kpc). We fit various functional forms to the profile up to 3 components in total, including Gaussians, Lorentzians, Voigt, combinations of these forms and a combination of Gaussian + exponential. The best fit results were achieved by 3 Gaussians or 1 Gaussian plus 1 Voigt profile, which resulted in equivalent residual rms values. However, since the Voigt profile does not have easily identifiable parameters, we show the fits (Fig. 12) for the Gaussian combinations only and provide tabular information (Table 4) for these combinations as well as the Gaussian + exponential for comparative purposes. The broadest Gaussian scale length is...
is a measure also, since the high latitude $\mu/\lambda$ mJy arcsec $^2$ represent $\mu/\sigma$.

Plot of the data over the full range of intensity. b) Plot of the data and two of the models over only the lowest 1.7% of the total intensity. The dark solid curve shows the data with the noise, at an rms level of $2 \times 10^{-4}$ mJy arcsec$^{-2}$, now obvious. The grey curve shows the best fit 2-component Gaussian model and the narrower dark solid curves shows the best fit 3-component Gaussian model. See Table 4 for the parameters of the Gaussian fits.

Fig. 12. Average profile of the PAH emission as a function of $z$ from the high sensitivity 6.7W data (Fig. 6). Negative values of $z$ represent the western side of the major axis. a) Plot of the data over the full range of intensity. b) Plot of the data and two of the models over only the lowest 1.7% of the total intensity. The dark solid curve shows the data with the noise, at an rms level of $2 \times 10^{-4}$ mJy arcsec$^{-2}$, now obvious. The grey curve shows the best fit 2-component Gaussian model and the narrower dark solid curves shows the best fit 3-component Gaussian model. See Table 4 for the parameters of the Gaussian fits.

65.3" (3.5 kpc). Increasing the number of components would reduce the residual rms further but such a situation is likely not physical. The profile, in general, consists of a central Gaussian component plus broad wings. Also, since the high latitude emission is highly structured (Fig. 6), the scale length of the broad component will vary with position. Nevertheless, it is clear that an extended halo or thick disk of PAH emission is present in NGC 5907 with a characteristic scale height of 3.5 to 5 kpc.

Given the apparent correlation between the ISO $\lambda 7.7 \mu m$ emission and the SCUBA $\lambda 850 \mu m$ emission in the disk (Sect. 4.1), it is of interest to see if this correlation extends to the halo region. Figure 8b shows a comparison between these bands. Note that the $\lambda 850 \mu m$ field of view is truncated at a Declination of 56°17′ and the 7.7N map is truncated on the eastern side. Note also that the greyscale scaling apparent emission away from the plane in the SCUBA data is all below the 3σ level (Alton et al. 2004, their Fig. 1) with the exception of the “detached” feature at RA = 15h15m52s, Dec = 56°17′45″. Indeed, Alton et al. find an exponential scaleheight of only 0.11 kpc (2″) for the $\lambda 850 \mu m$ map. If the $\lambda 850 \mu m$ emission existed in the halo at the same fraction of the in-disk emission as the PAHs, then the $\lambda 850 \mu m$ emission should have been detected in the halo above the noise of the Alton et al. map. This suggests that either the large grains are under-represented in the halo in comparison to the PAHs, or are under-emitting.

The broad band 12W emission shows very little vertically extended emission (Figs. 5, 11) in comparison to the PAH bands (cf. Figs. 2 and 4). This is likely due to the fact that the dynamic range in the 12W band is only 10 to 17% as great as in the 7.7N and 11.3N bands (Table 3). For example, the extended vertical emission seen in Figs. 2 and 4, if present in Fig. 5 at the same relative brightness with respect to the maximum disk emission, would mostly fall at the level of the noise or lower. One obvious feature is visible at the edge of the field of view on the north, extending to the east. (The optical feature within this extension appears to be a foreground star.) The reality of the large scale feature is not clear but its “footprints” close to the disk have counterparts in the PAH bands (e.g. Fig. 4) suggesting that the PAH contribution within this band may be responsible for it.

4.3. Flux densities and band ratios

In no case do we have complete coverage of the galaxy so global flux densities cannot be obtained. However, we have determined the flux density for each of the fields displayed. These are listed in Table 5. Since the 6.8N through 11.3N bands have the same field of view, their fluxes can be compared directly. There is one common field for all maps, however, which is outlined in Fig. 6b. Thus, we also calculate fluxes for this common region (Table 5) and compare these results for all bands. The resulting flux ratios are provided in Table 6 with some comparative values, where known, also given.

The 6.7W:12W ratio is interesting since it is significantly higher than the values given by Dale et al. (2001). This ratio represents a comparison between a broad ISO/CAM band that contains mainly 3 PAH bands (Sect. 3) with the broader 12 μm IRAS-equivalent band which may contain both PAHs and continuum. Dale et al. (2001) determined average values of this ratio, using the IRAS 12 μm flux itself, for normal galaxies in different $f_{60}/f_{100}$ bins. The IRAS flux ratio, $f_{60}/f_{100}$ is a measure of the heating of classical grains in temperature equilibrium. As such, it is widely used as a measure of star formation activity, with an increasing ratio implying greater dust heating (Helou 1986; Dale et al. 2000, 2001). As the SFR increases ($f_{60}/f_{100}$ increases), the 6.7W/(broad-band 12 μm) ratio decreases, reflecting the fact that a) the contribution of a hot dust continuum increases with SFR and b) PAHs are expected to be destroyed in regions of high heating intensity such as active star forming regions (Dale et al. 2001 and references therein). The $f_{60}/f_{100}$ ratio for NGC 5907 (0.24, Table 1) places it at a SFR that falls below all of the galaxies of the Dale et al. sample. Moreover, the $f_{60}/f_{100}$ ratio for NGC 5907 could be even lower than that quoted in Table 1, i.e. 0.21 from earlier ISO data, Bendo et al. (2002). Thus, the high 6.7W:12W ratio for NGC 5907 (Table 6) is likely consistent with the data from other galaxies in that it would represent an extension to lower SFR.

The ratios for the bands 6.8N through 11.3N (Table 6) all have the same field of view and so can be compared directly. (Note also, that all results for the “field” and for the “common” regions agree within errors.) For these bands, we make a comparison to the galaxy, M 83, for which similar observations have been made. NGC 5907 is not perfectly identical to M 83 in its global properties (Table 1). The two galaxies are of similar morphological types and have similar global star formation rates, infrared luminosities, and modelled temperatures. However, M 83 is physically smaller than NGC 5907. Thus, NGC 5907 contains an order of magnitude more dust but M 83 is the more active galaxy when SFR per unit area is determined. Nevertheless, the shapes of the different spectra in M 83 (see Vogler et al. 2005 and Fig. 7) are remarkably constant as a function of location in that galaxy, i.e. as a function of differing local SFR, the main differences being increased.
Table 4. 6.7W z profile.

<table>
<thead>
<tr>
<th>Data or type of fit</th>
<th>Component No.</th>
<th>Peak intensity (mJy arcsec$^{-2}$)</th>
<th>FWHM$^*$ (arcsec)</th>
<th>Residual rms (mJy arcsec$^{-2}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Data</td>
<td>1-component Gaussian</td>
<td>0.315</td>
<td>18.3</td>
<td>3.96 × 10$^{-3}$</td>
</tr>
<tr>
<td></td>
<td>2-component Gaussian</td>
<td>0.287</td>
<td>16.2</td>
<td>3.96 × 10$^{-3}$</td>
</tr>
<tr>
<td></td>
<td>Gaussian + Exponential</td>
<td>0.065</td>
<td>16.5</td>
<td>1.11 × 10$^{-3}$</td>
</tr>
<tr>
<td></td>
<td>3-component Gaussian</td>
<td>0.0564</td>
<td>27.2</td>
<td>0.86 × 10$^{-3}$</td>
</tr>
</tbody>
</table>

$^a$ Unless otherwise indicated.
$^b$ Exponential scale length.

Table 5. Flux densities.

<table>
<thead>
<tr>
<th>Band</th>
<th>Flux density (field)$^a$ (mJy)</th>
<th>Flux Density (Common)$^b$ (mJy)</th>
</tr>
</thead>
<tbody>
<tr>
<td>6.8N</td>
<td>1170 ± 97</td>
<td>519 ± 42</td>
</tr>
<tr>
<td>7.7N</td>
<td>2170 ± 141</td>
<td>918 ± 60</td>
</tr>
<tr>
<td>9.6N</td>
<td>919 ± 66</td>
<td>360 ± 26</td>
</tr>
<tr>
<td>11.3N</td>
<td>2071 ± 116</td>
<td>828 ± 46</td>
</tr>
<tr>
<td>12W</td>
<td>940 ± 38</td>
<td>557 ± 57</td>
</tr>
<tr>
<td>6.7W</td>
<td>666 ± 22</td>
<td>591 ± 20</td>
</tr>
</tbody>
</table>

$^a$ Flux density determined over the entire field shown in Fig. 1a through Fig. 6a. The error (here and throughout) is calculated via $\sqrt{\sum i \sigma^2}$ over the map pixels, $i$, where $\sigma_i$ is the quadratic sum at each pixel of $\sigma_{rms}$, the Transient rms, the 1σ sky error and the calibration error (see Table 3).

$^b$ Excluding the 3 small regions with residual errors (see Sect. 2.3).

contributions of continuum, especially longwands of 14 µm (which we sample only marginally in NGC 5907 in the 12W band). For M 83, the main differences between spectra can be seen as an increased continuum contribution from hot grains which become more prominent towards the nucleus. Thus, we provide the 6.8N through 11.3N ratios for M 83 for the inner, hotter, 1′ radius region as well as the outer, cooler, region between radii of 1′ and 3′ in Table 6. Given the estimated errors on the various ratios, there is agreement between all values for M 83 and those of NGC 5907, although the most variation is seen for 9.6N. The band, 9.6N, measures a mixture of PAH wings plus possible hot continuum. The 9.6N:7.7N ratio for NGC 5907 is somewhat low in comparison to the nucleus of M 83 but agrees with the result for the interarm region of M 83. This is again quite consistent with the low SFR for NGC 5907.

The Dale et al. comparison values for 6.7W:12W are globally determined values. Our data allow us to further examine this ratio as a function of projected distance from the major axis, as shown in Fig. 13. In Fig. 13a, we show the increasing 6.7W:12W ratio with projected distance from the major axis. Given that this ratio can only be determined over a restricted area of the “common” region due to signal-to-noise considerations, this may reflect either an increase in this ratio with $z$ height or an increase with distance from the center of the galaxy since the region shown does not extend beyond the projected optical disk. In one case, this is likely due to an increase in the distance from star forming regions and therefore a decrease in the contribution of the hot continuum that would be expected in the 12W band. In Fig. 13b, we show the ratio of the emission in two PAH bands: 11.3N/7.7N. In this case, the ratio can be determined to much higher latitude. The maximum change in ratio is of order 2.4 with the in-disk value lower. Given that more continuum, if present, is expected in the 11.3N band than the 7.7N band in the disk, then the in-disk ratio could be lower than shown. As indicated above, the global ratio of these two bands is quite consistent with what is seen in other galaxies so we may be observing a real variation in the 11.3/7.7 µm PAH ratio between the changing environments of the disk and halo. The nature of this variation is unclear and the exact magnitude will require the acquisition of spectra and modelling. However, we note that the variation seen here is not inconsistent with the variations between PAH feature strengths observed for regions in our own Milky Way, i.e. of order a factor of 5 (see Peeters et al. 2002).

4.4. Radio continuum correlation?

Vogler et al. (2005) find a good spatial correlation between the 16 cm radio continuum emission and the 6.7W (largely PAH) MIR emission in M 83, and Brar et al. (2003) also find a good correlation between the 6850 µm emission and the 617 MHz radio continuum emission in the edge-on galaxy, NGC 5775. These results, which link the MIR or sub-mm with the radio continuum, are reminiscent of the well-known FIR-radio correlation whose origin is not fully understood but is thought to relate to star formation in galaxies (see Groves et al. 2003, and references therein). Thus, a comparison of our ISO results with radio data is of interest. The best available radio data sets are those of the NRAO VLA Sky Survey (NVSS) at 1420 MHz (Condon et al. 1998) and the Westerbork Northern Sky Survey (WENSS) at 327 MHz (Rengelink et al. 1997). The radio emission in the in the VLA Faint Images of the Radio
Table 6. Flux density ratios.

<table>
<thead>
<tr>
<th>Bands</th>
<th>Ratio or relation</th>
<th>Galaxy or region</th>
<th>Ref.</th>
</tr>
</thead>
<tbody>
<tr>
<td>6.7W:12(IRAS)</td>
<td>$[0.52^{+0.27}_{-0.25}]:1^f$</td>
<td>8 galaxies</td>
<td>Dale et al. (2001)</td>
</tr>
<tr>
<td>6.7W:12W</td>
<td>$[1.06 \pm 0.16]:1$</td>
<td>Common$^d$</td>
<td>This work</td>
</tr>
<tr>
<td>6.8N:7.7N:9.6N:11.3N</td>
<td>$[0.57 \pm 0.08]:1:[0.39 \pm 0.05]:[0.90 \pm 0.11]$, Common$^d$ This work</td>
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<td>$0.49:1:0.59:1.12^c$, M 83 ($R \leq 1'$)$^f$ Vogler et al. (2005)</td>
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<tr>
<td>6.8N:7.7N:9.6N:11.3N</td>
<td>$0.55:1:0.48:0.95^c$, M 83 ($1' \leq R \leq 3'$)$^f$ Vogler et al. (2005)</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

$^a$ The center wavelengths of the bands is given below each sequence that is the same.
$^b$ Average of 8 galaxies in a flux ratio, $f_{60}/f_{100}$, bin from 0.28 to 0.35 (Dale et al. 2001).
$^c$ This error includes an additional 5% error to account for variations in orbit (Sect. 2.3).
$^d$ Common or Field region as quoted in Table 5.
$^e$ Errors (estimates only) are believed to be of order 20%.
$^f$ $R$ is the projected radial distance from the center of the galaxy.

Fig. 13. a) Profile of the 6.7W/12W ratio from the common region only (see Fig. 6b) as a function of projected distance from the major axis. The values are averages over 71’’ strips parallel to the major axis.

b) The $z$ profile of the 11.3N/7.7N ratio for the fields shown in Figs. 2 and 4, but averaged over 320’’ strips. Note that the $z$ extent is much larger than in a).

Sky at Twenty-Centimeters survey (FIRST) is too faint for comparison (but see note below). In Fig. 14, we show the 7.7N ($\lambda 7.7 \mu m$ PAH) emission, which is unlikely to contain any significant continuum component, in comparison to these two radio images. The strong radio source at the end of the south-eastern major axis (actually a double) is believed to be from a background source (Dumke et al. 2000). There are significant differences between the fields of view and resolutions of the ISO and radio data sets which make comparison difficult, and smoothing the 7.7N emission results in confusion and truncations near the field edges. In addition, details in the structure of the 7.7N map at low emission levels may be approximate. Thus any apparent correlations are tentative and are presented so that higher resolution data in the future might be compared to these results. In particular, a radio continuum extension on the west side of the major axis at RA $\approx 15h 15m 45s$, Dec $\approx 56' 20''$ shows extended PAH emission (though the ISO field of view also truncates here). South of this feature is another at RA $\approx 15h 15m 50s$, Dec $\approx 56' 18''$. On the east side of the major axis, there is a feature at RA $\approx 15h 16m 00s$ Dec $\approx 56' 20''$. These features in the radio are more easily seen in the 327 MHz image and at least the northernmost extension on the west side can also be seen in a 20 cm image of Dumke et al. (2000). We are not able to determine correlations along the disk due to the limited ISO field of view and resolution and also because of confusion with the strong background source.

Finally, we note that there is a previously unreported point source in the FIRST data with coordinates, RA(J2000) = 15h15m49.10s, Dec(J2000) = 56°20’’27.1’’, which is displaced 38’’ to the west and 43’’ to the north of the nucleus (2.0 and 2.3 kpc, respectively, in projection) with a flux density of 3.1 mJy. This flux density corresponds to a spectral power of $4.5 \times 10^{19}$ Watts/Hz which is intermediate between that of Cas A and the brightest supernova remnant in M 82 (see Irwin et al. 2000) and is therefore likely a supernova remnant or remnants.

5. Discussion

The mid-IR emission in NGC 5907 is consistent with NGC 5907 as a low SFR galaxy. As described in Sect. 4.3, this galaxy is “PAH-rich” and “hot-dust-poor”. Our flux ratios are consistent with those measured in other galaxies at the low end of the range of SFR and are also consistent with those of Vogler et al. (2005) for M 83 for the regions between spiral arms. For M 83, the continuum contributes no more than $\approx$10% of the total MIR 4 to 18 $\mu m$ emission with $\approx$90% due to PAHs in these
regions. Although we do not have CVF data for NGC 5907, all other indicators suggest that the situation is likely comparable.

It is quite clear that the PAH distribution follows that of the CO (Fig. 10). Both the $\lambda 850 \mu m$ and $\lambda 1.2 \ mm$ cool dust also follow the CO distribution, and the PAH and $\lambda 850 \mu m$ are spatially well correlated within the disk of NGC 5907. This suggests that both the PAHs and larger grains are associated with the molecular gas component. PAHs are well known to trace classical photodissociation regions (PDRs) of molecular clouds which are excited by far-UV photons from a hot stellar population in our own Galaxy (see Sect. 3). However there is also evidence that the PAHs are extended radially in comparison to CO both from our results as well as the cool 1.2 mm dust distribution of Dumke et al. (1997). The cool dust and PAH emission both extend to radii of $\sim 250''$ in comparison to $\sim 175''$ for CO. While there may be some weak CO present, the PAH and cool dust distributions show an excess with respect to the CO at large radii in comparison to smaller radii. These results suggest that, while PAHs are strongly associated with the molecular gas, a more widespread distribution, likely associated with cool HI, is also present but at a much weaker level of emission.

Haas et al. (2002) also find a correlation between PAHs and cold dust, the latter as measured by the SCUBA $\lambda 850 \mu m$ continuum emission, in galaxies. These authors suggest that PAH carriers are widespread and correlate with large dust grains and neutral molecular clouds, but require a minimal radiation field in order to be excited, thus explaining global correlations between SFR and PAH emission as well as the fact that PAHs are distributed more widely than the SF regions alone. This is also consistent with Peeters et al. (2004) who suggest that PAHs may be better tracers of B stars, which dominate the Galactic stellar energy budget, rather than as a tracer of massive star formation. As indicated in Sect. 3, the PAHs in our Galaxy peak outside of HII regions and are also diffuse emitters in the diffuse ISM of galaxy disks, where the excitation may be from optical photons from an older population. These studies therefore support the association of PAHs with more extended neutral atomic gas. Tacconi-Garman et al. (2005) find that the global distribution of PAHs in NGC 253 and NGC 1808 follows the starburst activity, but on small scales, there is little correspondence (correlation or anti-correlation) with star forming regions. It may well be that PAHs are widespread and excited by both UV photons from a hot stellar population as well as optical photons from the more widespread stellar population from which the interstellar radiation field (ISRF) originates. The fact that they correlate globally with the CO distribution would then be a statement that there are more UV photons available at the appropriate wavelength and also that there is a greater concentration of dust and PAHs in regions of higher density molecular clouds. In regions (or galaxies) with very high SFRs the PAHs will be destroyed in comparison to the VSGs; this is supported by observations of M 83 (Vogler et al. 2005) and of starbursts and AGNs (e.g. Sturm et al. 2000).

An important result is the presence of PAH emission to very high latitude, with scale heights of order 3.5 to 5 kpc and up to 6.5 kpc in some locations. By comparison, the CO FWHM in the $z$ direction is only $\sim 8'' = 426 $ pc after deconvolution with the disk (Dumke et al. 1997) and the $\lambda 850 \mu m$ exponential scale height is only 110 pc to the limits of those data. An equally significant result is that such high latitude gas is seen in a galaxy with a low SFR. To our knowledge, this is the first evidence for high latitude PAH emission in any external galaxy. A $\lambda 3.3 \mu m$ PAH spur of $z$ extent $< 120$ pc has now been detected by Tacconi-Garman et al. (2005) in the outflowing wind of the starburst galaxy, NGC 253, lending support to the concept of “mass-loaded” nuclear starburst winds and the possibility that some PAHs can survive such blasts, the intense UV radiation and shocks that are expected to accompany them.

The situation is quite different in NGC 5907, however, which does not show strong starbursting or outflowing winds. There is, however, a feature with anomalous velocities in high resolution CO data that cannot be explained by a bar or rotation model (Garcia-Burillo et al. 1997, their Feature F). A CO loop-like feature of size $\sim 20''$ associated with the anomalous structure in the CO maps is suggestive of a small-scale wind outflow. The radio point-source described in Sect. 4.4 (not at the same location as the CO feature) indicates that some SF activity is indeed occurring now in NGC 5907, although at a low level in comparison to other SF galaxies. The main independent evidence for high latitude emission in NGC 5907 is from the radio continuum (Dumke et al. 2000 and Fig. 14) which shows a $\sim 20$ cm thick disk scale length of 1.5 kpc and a large scale field of $\sim 1 \mu G$. This indicates that cosmic rays and magnetic fields do exist in the halo, although a detailed search for a correlation has not yet been possible (Sect. 4.4).
A mechanism for depositing PAHs in the halo of a low SFR galaxy is not yet clear. Since SF and magnetic fields are both clearly present in NGC 5907, the PAHs may be deposited into the halo initially from SNe and stellar winds from SF regions and then assisted in reaching high latitudes via coupling to the magnetic field. Excitation could be from UV photons that leak from the disk into the halo. An interesting possibility is that the 3-D VSGs are themselves converted into planar PAH molecules (Duley et al. 2000). This would occur in regions of shocks and high UV excitation in the disk. The smaller PAHs may then be more easily "levitated" above the plane. Jones et al. (1996) have also noted that PAHs and smaller particles can build up at the expense of larger ones in shocks or winds and that wind velocities as low as tens of km s\(^{-1}\) can modify the grain properties. The apparent dearth of \(\lambda 850\) \(\mu\)m emission in the halo (Sect. 4.2) supports this view. It would be interesting to obtain more sensitive \(\lambda 850\) \(\mu\)m and \(\lambda 1.2\) mm data in order to compare the results with starburst galaxies, such as NGC 5775 in which \(\lambda 850\) \(\mu\)m emission has been detected in the halo (Brar et al. 2003). Determining the admixture of different dust components as a function of height would provide important clues as to what mechanisms may be involved. Conditions in galaxy halos are quite different from those in the disk and may more closely resemble those seen in elliptical galaxies.

Although there is a large optical ring around NGC 5907, as described in Sect. 1.1, there is no evidence that the high latitude PAH emission is directly related to it. The intersection of the optical ring with the disk of NGC 5907 occurs at two locations, one of which (to the north) is not in the field of view of our maps and the other is near the southern tip of the major axis (coordinates given in Fig. 11). We see no evidence for enhanced emission or extensions at the latter location. There does appear to be a disturbance in the Hz map of Rand et al. (1996) approximately at this location, however.

6. Conclusions

We have mapped the MIR emission of the low SFR galaxy, NGC 5907, in the 6.7W, 6.8N, 7.7N, 9.6N, 11.3N, and 12W (LW2, LW5, LW6, LW7, LW8, and LW10, respectively) ISO bands, at high sensitivity and high resolution. Although we do not have spectra for NGC 5907, we have compared the line strengths and line ratios with those of other galaxies, especially M 83 in quiescent regions between the spiral arms. These comparisons indicate that the excitation conditions are quite typical of regions and/or galaxies with a low SFR. The contribution from hot VSGs is likely no more than \(\sim 10\%\) across these bands and is likely negligible in all bands except the broad, IRAS-equivalent 12W band. The MIR emission is dominated by PAH emission.

In the disk, the high spatial resolution (1\(\arcsec\)) broad band 12W emission is peaked towards the optical dust lane. The bands dominated by PAH emission (at 6\(\arcsec\) resolution) are observed all along the major axis of the galaxy and extend almost to the end of the optical disk (cf. \(\sim 35\) kpc in comparison to 41 kpc) and possibly farther. The distribution of PAHs in the disk roughly follows that of the CO and not the broad HI distribution. The PAH distribution shows an apparent correlation with the SCUBA \(\lambda 850\) \(\mu\)m distribution to within the positional errors of the data. The \(\lambda 1.2\) mm emission also follows that of the molecular gas. At the same time, there is some PAH emission (and also \(\lambda 1.2\) mm emission) at larger galactocentric radii than the CO, suggesting that PAH emission is more widespread than the molecular gas alone. It is likely that greater PAH emission is seen where there is more molecular gas because there are more UV photons available from SF in these regions and because of the greater concentration of grains and PAHs in higher density gas. A smaller fraction of the PAH emission will be coming from the more extended HI, excited by the general diffuse ISRF.

An important finding of this work is the detection of PAH emission at high latitudes. The emission features show structure in the narrow bands that isolate single PAH features, appearing similar to features seen in H\alpha, radio continuum, and other bands in other edge-on galaxies (see, e.g. Lee et al. 2002). Individual features extend to latitudes of up to 6.5 kpc but a more typical scale height is 3.5 kpc. Previous high latitude emission on kpc scales has been seen only in the radio continuum component in this galaxy. Thus, cosmic rays and magnetic fields are known to exist in the halo, but other components have not (yet) been observed. Indeed, the PAHs appear to be selectively represented in the halo in comparison to large grains. We note also an increase in the \(\lambda 1.13/\lambda 7.7\) PAH ratio with distance from the major axis. Although NGC 5907 is a low SFR galaxy, there is some SF activity occurring, such as has been seen via some anomalous CO emission as well as a previously unreported off-nucleus radio continuum point source which is likely a supernova remnant or remnants. Thus, the halo PAHs may be transported to high latitudes with the help of coupling to magnetic fields (likely with gas as well), from regions of SF.

Although NGC 5907 is surrounded by a faint stellar ring, there is no indication in these data that interaction with the ring plays a direct role in the presence of the high latitude emission.

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